# LEGENDRIAN GRONWALL CONJECTURE 

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#### Abstract

The Gronwall conjecture states that a planar 3-web of foliations which admits more than one distinct linearizations is locally equivalent to an algebraic web. We propose an analogue of the Gronwall conjecture for the 3-web of foliations by Legendrian curves in a contact three manifold. The Legendrian Gronwall conjecture states that a Legendrian 3-web admits at most one distinct local linearization, with the only exception when it is locally equivalent to the dual linear Legendrian 3-web of the Legendrian twisted cubic in $\mathbb{P}^{3}$. We give a partial answer to the conjecture in the affirmative for the class of Legendrian 3-webs of maximum rank. We also show that a linear Legendrian 3-web which is sufficiently flat at a reference point is rigid under local linear Legendrian deformation.


## 1. Introduction

Let $M$ be a connected contact three manifold. A Legendrian $d$-web on $M$ is by definition a set of $d$ pairwise transversal foliations of $M$ by Legendrian curves. The Legendrian web was introduced in [Wa2] for a second order generalization of the classical planar web.

Abelian relations and rank, two of the central concepts in web geometry, are analogously defined for the Legendrian web. The main result of [Wa2] was that the rank of a Legendrian $d$-web admits the optimal bound $\rho_{d}=\frac{(d-1)(d-2)(2 d+3)}{6}$. We also gave an analytic characterization of the Legendrian 3-webs of maximum rank three.

In this paper we study the linearization problem for the Legendrian 3-webs with a more algebro-geometric perspective. As described in [Wa2], the algebraic model for the Legendrian web theory is provided by the projective duality associated with the simple Lie group $\mathrm{Sp}_{2} \mathbb{C}$, Figure [1.1, see Section 1.1 [ $\overline{\mathrm{Br}}$ ]. By the standard dual construction, a null degree $d$ surface $\Sigma \subset \mathbb{Q}^{3}$ induces the dual $d$-web of Legendrian lines on a generic small open subset of $\mathbb{P}^{3}$ Generalizing this, a linearization of a Legendrian $d$-web on a contact three manifold $M$ is defined as a contactomorphism $M \hookrightarrow \mathbb{P}^{3}$ for which the image of each leaf of the Legendrian foliations is mapped to a Legendrian line. A natural question arises as to which Legendrian $d$-webs are linearizable, and how unique such a linearization is. The linearizability problem for a planar web can be traced back to Balschke, [BB AGL, GL, PP Pir] and the references therein. See [Wa1] for the references on the Gronwall conjecture.

It is the uniqueness part of the linearization problem that we are interested in. In planar web geometry, the Gronwall conjecture states that a planar 3 -web admits more than one distinct linearizations (uniqueness fails) whenever the 3 -web is locally equivalent to an algebraic web. In other words, the conjecture claims that the failure of unique linearization, or equivalently the linear deformability, implies that the planar 3-web is essentially algebraic.

One of the motivations for the present work was therefore the idea that the condition for linear deformability may lead to a special, possibly algebraic, class of linear Legendrian 3-webs. Through these examples one might hope to gain an insight on the analogues of Abel's theorem and its converse in our setting, Section 4 .

For the Legendrian $d$-webs, $d \geq 4$, the uniqueness of linearization follows from the local normal form for the projectively flat third order ODE's, [SY] Corollary 4.2]. The case of Legendrian 3-webs on the other hand is not

[^0]

Figure 1.1. Projective duality associated with $\mathrm{Sp}_{2} \mathbb{C}$
obvious. We shall employ the method of moving frames and the local differential analysis to analyze the linear Legendrian deformation and rigidity of the linear Legendrian 3-webs on an open subset of $\mathbb{P}^{3}$.

## Main results.

1. The number of distinct local linearizations of a Legendrian 3-web is uniformly bounded, Theorem 3.14 (the bound is far from being optimal).
2. A linear Legendrian 3-web on a connected open subset of $\mathbb{P}^{3}$ is rigid under linear Legendrian deformation when it is sufficiently flat at a single point, Theorem 3.16.
3. We propose the Legendrian analogue of the Gronwall conjecture, Conjecture 5.1 The conjecture claims that a Legendrian 3-web on a connected contact three manifold admits at most one distinct linearization, with the only exception when the 3-web is locally equivalent to the dual linear Legendrian 3-web of the Legendrian twisted cubic curve in $\mathbb{P}^{3}$ (in which case it has exactly two distinct local linearizations). We verify the conjecture for the class of Legendrian 3-webs of maximum rank three, Theorem 4.1

Let us give a description of the deformable Legendrian 3-web in Main results 3. Let $\gamma \subset \mathbb{P}^{3}$ be the Legendrian twisted cubic. Take a generic point $\mathrm{x} \in \mathbb{P}^{3}$. Let $\mathbb{P}_{\mathrm{x}}^{2}$ be the contact 2-plane at x , which is the union of the Legendrian lines through x. Since $\gamma$ is a curve of degree $3, \mathbb{P}_{\mathrm{x}}^{2}$ intersects $\gamma$ at 3 points. By definition of $\mathbb{P}_{\mathrm{x}}^{2}$ these points determine the corresponding 3 Legendrian lines through x . It is clear that as x varies this construction defines a linear Legendrian 3 -web $\mathcal{W}_{\gamma}$ on a generic small open subset of $\mathbb{P}^{3}$. Our analysis shows that $\mathcal{W}_{\gamma}$ admits another (exactly one more) distinct linearization, and that it is the unique such linear Legendrian 3-web of maximum rank.

The three Abelian relations of $\mathcal{W}_{\gamma}$ admit the following geometrical interpretation as a generalized addition law for the Legendrian twisted cubic. Take a set of 3 generic points $\mathrm{p}_{i}, i=1,2,3$, on $\gamma$. Then the 3 planes $\mathbb{P}_{\mathrm{p}_{i}}^{2}$ intersect at a single point, and one consequently gets a set of 3 concurrent Legendrian lines. The choice of three points on a given curve $\gamma$ depends on three one-dimensional parameters. Once these points are fixed, the choice of a Legendrian line through each of these points also depends on three one-dimensional parameters. ${ }^{2}$ It is evident from the above incidence relation that, roughly speaking, the latter three parameters are functions of the former three parameters. The three Abelian relations of $\mathcal{W}_{\gamma}$ implies that in fact there exist three additive form of functional relations among the total six one-dimensional parameters.

Considering the obvious analogue of the converse of Abel's theorem, Section 4 the relevant point here would be that the incidence geometry described above manifests itself through the additive functional equations among the related geometric quantities which by construction are the first integrals of the foliations of $\mathcal{W}_{\gamma}$. As is the case with the planar web geometry, perhaps this additivity provides a concrete evidence to infer, and in turn it would imply via the converse of Abel's theorem, the underlying algebraic structure.

There is another algebro-geometric implication of the three Abelian relations of $\mathcal{W}_{\gamma}$ on the degeneration of K3 surfaces in $\mathbb{P}^{4}$ in the context of the presumed Legendrian analogue of Abel's theorem and its converse. Consider the surface $\Sigma_{\gamma}=\pi_{1} \circ \pi_{0}^{-1}(\gamma) \subset \mathbb{Q}^{3}$. Since $\gamma$ is Legendrian, $\Sigma_{\gamma}$ is the tangent developable of the Gauss map of $\gamma$, which is a null rational normal curve in $\mathbb{Q}^{3}$. Via the converse of Legendrian Abel's theorem, the three Abelian relations of $\mathcal{W}_{\gamma}$ would then imply that $\Sigma_{\gamma}$ supports three generalized closed holomorphic 1-forms. On the other hand, $\Sigma_{\gamma}$ lies in the intersection of $\mathbb{Q}^{3}$ with a cubic hypersurface of $\mathbb{P}^{4}$. A smooth complete intersection of type $(2,3)$ in $\mathbb{P}^{4}$ is a K3 surface, which has no nonzero holomorphic 1-forms. The analytic surface $\Sigma_{\gamma}$ represents in

[^1]this way a degeneration of K 3 surfaces where the dimension of the space of closed holomorphic 1 -forms jumps by three.

Let us give an outline of the paper. In Section 1.1 we record the generalities on the duality between the projective space $\mathbb{P}^{3}$ and the 3 -quadric $\mathbb{Q}^{3} \subset \mathbb{P}^{4}$. In Section 2 the moving frame method is applied to determine a normalized frame bundle associated with a linear Legendrian 3-web on an open subset of $\mathbb{P}^{3}$. The normalized frame bundle serves as a base point for the linear Legendrian deformation of the 3-web. In Section 3 the standard deformation analysis in terms of the deformed Maurer-Cartan equation leads to the closed structure equation for the deformation parameters. The differential compatibility conditions among these parameters imply a set of polynomial equations that they must satisfy. The rigidity results Theorem 3.14 and Theorem 3.16 are the immediate consequences of the analysis of the root structure of the first few low degree polynomial compatibility equations. In Section 4 we implement the procedure established in Section 3 to the class of maximum rank Legendrian 3webs. Drawing from [Wa2] the analytic characterization of such 3-webs, we first observe that a Legendrian 3-web of maximum rank always admits a linearization as the dual 3-web of a union of three hyperplane sections in $\mathbb{Q}^{3}$. A further analysis of the polynomial compatibility equations shows that there exists a unique linear Legendrian 3web of this kind which is contact equivalent (and not projectively equivalent) to the dual 3-web of the Legendrian twisted cubic described above, and that otherwise these 3 -webs are rigid under linear Legendrian deformation. Based on this analysis, we propose the Legendrian analogue of the converse of Abel's theorem. In Section 55, we state the Legendrian Gronwall conjecture, and present a few remarks on the related problems.

We assume the complex analytic category. The analysis and results are valid within the real smooth category with only minor modifications. The moving frame method, and the over-determined PDE machinery are used throughout the paper. We refer the reader to [BCG3][IL] for the standard references.
1.1. Projective duality. Let $V=\mathbb{C}^{4}$ be the four dimensional complex vector space. Let $\varpi$ be the standard symplectic 2-form on $V$. Let $\mathbb{P}^{3}=\mathbb{P}(V)$ be the projectivization equipped with the induced contact structure. The contact 2-plane field $\mathcal{H}$ on $\mathbb{P}^{3}$ is defined by

$$
\begin{equation*}
\left.\mathcal{H}_{\mathrm{x}}=\mathbb{P}((\hat{\mathrm{x}}\lrcorner \varpi)^{\perp}\right), \text { for } \mathrm{x} \in \mathbb{P}^{3} \tag{1.1}
\end{equation*}
$$

where $\hat{x} \in V$ is any de-projectivization of x . Since $\varpi$ is non-degenerate, $(\hat{\mathrm{x}}\lrcorner \varpi)^{\perp} \subset V$ is a codimension one subspace containing $\hat{x}$, and its projectivization $\left.\mathbb{P}((\hat{\mathrm{x}}\lrcorner \varpi)^{\perp}\right) \subset \mathbb{P}^{3}$ is a hyperplane at x . The symplectic group $\operatorname{Sp}_{2} \mathbb{C}$ acts transitively on $\mathbb{P}^{3}$ as a group of contact transformation.

Let $\bigwedge_{0}^{2} V=\langle\varpi\rangle^{\perp} \subset \bigwedge^{2} V$ be the five dimensional subspace of isotropic 2-vectors. Let $\operatorname{Lag}(V)$ be the set of two dimensional Lagrangian subspaces of $V . \operatorname{Lag}(V)$ is identified with the 3-quadric $\mathbb{Q}^{3} \subset \mathbb{P}^{4}=\mathbb{P}\left(\bigwedge_{0}^{2} V\right)$ via the Plücker embedding. $\mathrm{Sp}_{2} \mathbb{C}$ acts transitively on $\mathbb{Q}^{3}$, and $\mathbb{Q}^{3}$ inherits the $\mathrm{Sp}_{2} \mathbb{C}$ invariant non-degenerate conformal structure.

Let $\mathrm{Z}=\mathbb{P}(\mathcal{H}) \rightarrow \mathbb{P}^{3}$ be the bundle of Legendrian ( $\mathcal{H}$-horizontal) line elements. It can be defined as the incidence space

$$
\mathrm{Z}=\left\{(\mathrm{x}, \mathrm{~L}) \in \mathbb{P}^{3} \times \mathbb{Q}^{3} \mid \hat{\mathrm{x}} \wedge \hat{\mathrm{~L}}=0\right\}
$$

The projective duality in Figure 1.1 represents the following $\mathrm{Sp}_{2} \mathbb{C}$ equivariant incidence double fibration, Figure 1.2 By definition of Z , it is clear that both fibers of $\pi_{0}, \pi_{1}$ are isomorphic to $\mathbb{P}^{1}$. A fiber of $\pi_{0}$, the set of Legendrian lines through a point in $\mathbb{P}^{3}$, projects to a null line of $\mathbb{Q}^{3}$, and dually a fiber of $\pi_{1}$, the set of null lines through a point in $\mathbb{Q}^{3}$, projects to a Legendrian line of $\mathbb{P}^{3}$. The duality correspondence can be summarized by;
$\mathbb{P}^{3}$ is the space of null lines in $\mathbb{Q}^{3}$, and dually $\mathbb{Q}^{3}$ is the space of Legendrian lines in $\mathbb{P}^{3}$.
There exists an immediate application of the duality principle. Let $\gamma \subset \mathbb{P}^{3}$ be a Legendrian ( $\mathcal{H}$-horizontal) curve. The Gauss map $\hat{\gamma} \subset \mathbb{Q}^{3}$ is defined as the tangent map of $\gamma$. By duality, $\hat{\gamma}$ is the envelope of a one parameter family of null lines. Hence $\hat{\gamma}$ itself is a null curve in $\mathbb{Q}^{3}$.

To fix the notation, let us define the projection maps $\pi, \pi_{0}$, and $\pi_{1}$ explicitly. Let $Z=\left(Z_{0}, Z_{1}, Z_{2}, Z_{3}\right)$ denote the $\mathrm{Sp}_{2} \mathbb{C} \subset \mathrm{SL}_{4} \mathbb{C}$ frame of $V$ such that the 2-vector $\varpi_{\mathrm{b}}=Z_{0} \wedge Z_{2}+Z_{1} \wedge Z_{3}$ is dual to the symplectic


Figure 1.2. Incidence double fibration
form $\varpi$. Define

$$
\begin{align*}
\pi(Z) & =\left(\left[Z_{0}\right],\left[Z_{0} \wedge Z_{1}\right]\right),  \tag{1.2}\\
\pi_{0}\left(\left[Z_{0}\right],\left[Z_{0} \wedge Z_{1}\right]\right) & =\left[Z_{0}\right], \\
\pi_{1}\left(\left[Z_{0}\right],\left[Z_{0} \wedge Z_{1}\right]\right) & =\left[Z_{0} \wedge Z_{1}\right] .
\end{align*}
$$

In this formulation, the stabilizer subgroup P in Figure 1.2 is of the form

$$
\mathrm{P}=\left\{\left(\begin{array}{cc}
A & B  \tag{1.3}\\
\cdot & \left(A^{t}\right)^{-1}
\end{array}\right)\right\},
$$

where $\left(A^{-1} B\right)^{t}=A^{-1} B$, and

$$
A=\left\{\left(\begin{array}{ll}
* & * \\
\cdot & *
\end{array}\right)\right\}
$$

Here ' '' denotes 0 and ' ${ }^{\prime}$ ' is arbitrary.
The $\mathrm{Sp}_{2} \mathbb{C}$-frame $Z$ satisfies the structure equation

$$
\begin{equation*}
d Z=Z \phi \tag{1.4}
\end{equation*}
$$

for the Maurer-Cartan form $\phi$ of $\mathrm{Sp}_{2} \mathbb{C}$. $\phi$ satisfies the structure equation

$$
\begin{equation*}
d \phi+\phi \wedge \phi=0 \tag{1.5}
\end{equation*}
$$

The components of $\phi$ are denoted by

$$
\phi=\left(\begin{array}{cc}
\alpha & \gamma  \tag{1.6}\\
\beta & -\alpha^{t}
\end{array}\right)=\left(\begin{array}{cccc}
\alpha_{0}^{0} & \alpha_{1}^{0} & \gamma_{0}^{0} & \gamma_{1}^{0} \\
\alpha_{0}^{1} & \alpha_{1}^{1} & \gamma_{0}^{1} & \gamma_{1}^{1} \\
\beta_{0}^{0} & \beta_{1}^{0} & -\alpha_{0}^{0} & -\alpha_{0}^{1} \\
\beta_{0}^{1} & \beta_{1}^{1} & -\alpha_{1}^{0} & -\alpha_{1}^{1}
\end{array}\right)
$$

where $\{\alpha, \beta, \gamma\}$ are 2-by-2 matrix 1-forms such that $\beta^{t}=\beta, \gamma^{t}=\gamma$.

## 2. Linear Legendrian 3-webs

Let $\mathcal{W}$ be a Legendrian $d$-web on a contact three manifold $M$, Wa2]. A linearization of $\mathcal{W}$ is a contactomorphism $M \hookrightarrow \mathbb{P}^{3}$ such that each leaf of the foliations is mapped to a Legendrian line. Two linearizations are equivalent when they are isomorphic up to projective transformation by $\mathrm{Sp}_{2} \mathbb{C}$, and otherwise distinct. More specifically;

Definition 2.1. Let $\mathbb{P}^{3}$ be equipped with the standard $\mathrm{Sp}_{2} \mathbb{C}$ invariant homogeneous contact structure, Section 1.1 Let $M \subset \mathbb{P}^{3}$ be a connected open subset. A linear Legendrian $d$-web on $M$ is a set of $d$ pairwise transversal foliations of $M$ by Legendrian lines.

There exists a distinguished class of linear Legendrian webs. Recall that an analytic surface $\Sigma \subset \mathbb{Q}^{3}$ has null degree $d$ when it intersects a generic null line of $\mathbb{Q}^{3}$ at $d$ points.

Definition 2.2. Let $\mathcal{W}$ be a linear Legendrian $d$-web on an open subset of $\mathbb{P}^{3} . \mathcal{W}$ is algebraic, and is denoted by $\mathcal{W}_{\Sigma}$, when it is induced from a null degree d analytic surface $\Sigma \subset \mathbb{Q}^{3}$ by the standard dual construction.

Let us give a description of a particular class of algebraic Legendrian webs induced from a curve in $\mathbb{P}^{3}$. Let $\gamma \subset \mathbb{P}^{3}$ be a degree $d$-curve. Take a generic point $\mathrm{x} \in \mathbb{P}^{3}$. The contact 2-plane $\mathcal{H}_{\mathrm{x}}$ intersects $\gamma$ at $d$ points $\mathrm{p}_{i}(\mathrm{x}), i=1,2, \ldots d$. By definition of $\mathcal{H}_{\mathrm{x}}$, (1.1), one gets a set of $d$ Legendrian lines through x . It is clear then that this incidence construction defines a linear Legendrian $d$-web $\mathcal{W}_{\gamma}$ on a generic small open subset of $\mathbb{P}^{3}$. Note from Figure 1.2 that $\mathcal{W}_{\gamma}=\mathcal{W}_{\pi_{1} \circ \pi_{0}^{-1}(\gamma)}$.

The dual Legendrian web $\mathcal{W}_{\gamma}$ obtained in this way inherits a set of Abelian relations from the holomorphic 1-forms on $\gamma$. Let $\Omega \in H^{0}\left(\gamma, \Omega^{1}\right)$ be a holomorphic 1-form. By Hartogs' theorem, the trace of $\Omega$,

$$
\operatorname{Tr} \Omega=\sum \mathrm{p}_{i}^{*} \Omega
$$

is a holomorphic 1 -form on $\mathbb{P}^{3}$, and hence must vanish identically 3 Each 1 -form $p_{i}^{*} \Omega$ trivially vanishes on the Legendrian line determined by $\mathrm{p}_{i}$. This implies that there exists a linear map from $H^{0}\left(\gamma, \Omega^{1}\right)$ to the space of Abelian relations of $\mathcal{W}_{\gamma}$.

Consider the special case when $\gamma$ is itself a Legendrian curve. The associated surface $\pi_{1} \circ \pi_{0}^{-1}(\gamma)$ becomes the tangent developable of the Gauss map of $\gamma$. We shall see that in this case there are Abelian relations of $\mathcal{W}_{\gamma}$ which do not come from $H^{0}\left(\gamma, \Omega^{1}\right)$, Section 4
2.1. Equivalence problem. Let $\mathcal{W}$ be a linear Legendrian 3-web. For definiteness, we assume the three Legendrian foliations of $\mathcal{W}$ are ordered, and denote them by $\mathcal{W}=\cup_{i=1}^{3} \mathcal{F}^{i}$. Note by duality Figure 1.1 that each foliation $\mathcal{F}^{i}$ corresponds to a (possibly singular) surface $\Sigma^{i} \subset \mathbb{Q}^{3}$. An immersed surface in the three manifold $\mathbb{Q}^{3}$ is locally described as a graph of one scalar function of two variables. One may argue that the local moduli space of linear Legendrian 3-webs in $\mathbb{P}^{3}$ depends on three arbitrary scalar functions of two variables.

In this section, we apply the method of moving frames and determine the $\mathrm{Sp}_{2} \mathbb{C}$ invariant structure equation for a linear Legendrian 3-web. The analysis will result in a principal bundle $\mathcal{B} \rightarrow M$ equipped with a normalized $\mathfrak{s p}_{2} \mathbb{C}$-valued Maurer-Cartan form $\phi$. The functional relations among the coefficients of the components of $\phi$ are the basic local invariants of a linear Legendrian 3-web. Since the data $(\mathcal{B}, \phi)$ are canonically associated with the given linear Legendrian 3-web, they will serve as a reference point for the problem of deformation and rigidity to be discussed in Section 3 ,

Let $\mathcal{W}$ be a linear Legendrian 3-web on an open subset $M \subset \mathbb{P}^{3}$. Let $\mathrm{B} \subset \mathrm{Sp}_{2} \mathbb{C} \rightarrow M$ denote the induced principal right P -bundle. We continue the analysis from Section 1.1

Step 0. From (1.6), set $\alpha_{0}^{1}=\omega^{1}, \beta_{0}^{1}=\omega^{2} ; \beta_{0}^{0}=2 \theta ; \alpha_{0}^{0}=\rho_{0}$. The Maurer-Cartan form $\phi$ is written as

$$
\phi=\left(\begin{array}{cccc}
\rho_{0} & \alpha_{1}^{0} & \gamma_{0}^{0} & \gamma_{1}^{0}  \tag{2.3}\\
\omega^{1} & \alpha_{1}^{1} & \gamma_{0}^{1} & \gamma_{1}^{1} \\
2 \theta & \omega^{2} & -\rho_{0} & -\omega^{1} \\
\omega^{2} & \beta_{1}^{1} & -\alpha_{1}^{0} & -\alpha_{1}^{1}
\end{array}\right) .
$$

From the general theory of moving frames, one may apply the fiber group action by $\mathrm{P} \subset \mathrm{Sp}_{2} \mathbb{C}$ to arrange so that the three linear Legendrian foliations are defined by

$$
\begin{equation*}
\mathcal{F}^{i}=\left\langle\omega^{i}, \theta\right\rangle^{\perp}, i=1,2,3, \tag{2.4}
\end{equation*}
$$

where $\omega^{3}=-\left(\omega^{1}+\omega^{2}\right)$.

[^2]Let $\mathcal{B}_{0} \subset \mathrm{~B} \rightarrow M$ be the sub-bundle defined by (2.4). Assuming that the foliations are ordered, the structure group $P_{0} \subset \mathrm{P}$ of $\mathcal{B}_{0}$ is reduced to

$$
P_{0}=\left\{\left(\begin{array}{cccc}
* & * & * & * \\
\cdot & \pm 1 & * & \cdot \\
\cdot & \cdot & * & \cdot \\
\cdot & \cdot & * & \pm 1
\end{array}\right)\right\}
$$

where '. ' denotes 0 . On $\mathcal{B}_{0} \rightarrow M$, the 1 -forms $\beta_{1}^{1}, \gamma_{1}^{1} ; \alpha_{1}^{1}$ are semi-basic, and one may write

$$
\begin{align*}
& \beta_{1}^{1}=\epsilon_{1} \omega^{1}+\epsilon_{2} \omega^{2}+A_{1} \theta,  \tag{2.5}\\
& \gamma_{1}^{1}=\epsilon_{3} \omega^{1}+\epsilon_{4} \omega^{2}+A_{2} \theta, \\
& \alpha_{1}^{1}=\epsilon_{5} \omega^{1}+\epsilon_{6} \omega^{2}+A_{9} \theta,
\end{align*}
$$

for the coefficients $\epsilon_{j} ; A_{k}$.
Step 1. The condition that $\mathcal{W}$ is linear imposes a set of relations on $\epsilon_{j}$ 's. By (1.4), (2.4), one must have

$$
\begin{aligned}
d Z_{0}, d Z_{1} & \equiv 0, \quad \bmod \quad Z_{0}, Z_{1} ; \quad \omega^{2}, \theta, \\
d Z_{0}, d Z_{3} & \equiv 0, \quad \bmod \quad Z_{0}, Z_{3} ; \quad \omega^{1}, \theta \\
d Z_{0}, d\left(Z_{1}-Z_{3}\right) & \equiv 0, \quad \bmod \quad Z_{0}, Z_{1}-Z_{3} ; \quad \omega^{3}, \theta .
\end{aligned}
$$

A computation shows that this implies in (2.5)

$$
\begin{equation*}
\epsilon_{1}=0, \quad \epsilon_{4}=0, \quad 2\left(\epsilon_{5}-\epsilon_{6}\right)-\epsilon_{2}-\epsilon_{3}=0 . \tag{2.6}
\end{equation*}
$$

Step 2. $\phi$ satisfies the structure equation $d \phi+\phi \wedge \phi=0$. For a notational purpose, denote $\Phi=d \phi+\phi \wedge \phi$, which must vanish identically. For instance, $\Phi_{0}^{1}, \Phi_{0}^{3} ; \Phi_{0}^{2}$ give the formulae for the exterior derivatives $d \omega^{1}, d \omega^{2} ; d \theta$, etc.
$\Phi_{3}^{1} \wedge \omega^{2} \wedge \theta, \Phi_{1}^{3} \wedge \omega^{1} \wedge \theta$ show that

$$
\begin{array}{lll}
d \epsilon_{2} \equiv 2 \alpha_{1}^{0}, & \bmod \quad \omega^{1}, \omega^{2}, \theta, \rho_{0} \\
d \epsilon_{3} \equiv 2 \gamma_{1}^{0}, & \bmod \quad \omega^{1}, \omega^{2}, \theta, \rho_{0}
\end{array}
$$

One may apply the fiber group action by $P_{0}$ to translate so that

$$
\begin{equation*}
\epsilon_{2}, \epsilon_{3}=0 \tag{2.7}
\end{equation*}
$$

From (2.6), denote $\epsilon_{5}=\epsilon_{6}=A_{0}$. The exterior derivative of $A_{0}$ is written as

$$
d A_{0}=-A_{0} \rho_{0}+A_{0,1} \omega^{1}+A_{0,2} \omega^{2}+A_{0,0} \theta
$$

We shall adopt the similar notation for the covariant derivative of a coefficient for the rest of the paper.
Let $\mathcal{B}_{1} \subset \mathcal{B}_{0}$ be the sub-bundle defined by (2.7). The structure group $P_{1} \subset P_{0}$ of $\mathcal{B}_{1}$ is reduced to the form

$$
P_{1}=\left\{\left(\begin{array}{cccc}
* & \cdot & * & \cdot \\
\cdot & \pm 1 & \cdot & \cdot \\
\cdot & \cdot & * & \cdot \\
\cdot & \cdot & \cdot & \pm 1
\end{array}\right)\right\}
$$

On $\mathcal{B}_{1} \rightarrow M$, the 1 -forms $\alpha_{1}^{0}, \gamma_{1}^{0}$ are semi-basic. $\Phi_{3}^{1} \wedge \theta, \Phi_{1}^{3} \wedge \theta$ show that one may write

$$
\begin{align*}
& \alpha_{1}^{0}=\frac{A_{1}}{2} \omega^{1}+\epsilon_{7} \omega^{2}+B_{1} \theta,  \tag{2.8}\\
& \gamma_{1}^{0}=\epsilon_{8} \omega^{1}-\frac{A_{2}}{2} \omega^{2}+B_{2} \theta,
\end{align*}
$$

for the coefficients $\epsilon_{7}, \epsilon_{8} ; B_{1}, B_{2}$.
$\Phi_{1}^{1} \wedge \theta$ with this relation gives $A_{9}=-\epsilon_{7}-\epsilon_{8}-A_{0,1}+A_{0,2}-2 A_{0}^{2}$.

Step 3. $\omega^{1} \wedge \Phi_{1}^{0}+\Phi_{3}^{0} \wedge \omega^{2}$ gives

$$
d\left(\epsilon_{7}-\epsilon_{8}\right) \equiv-2 \gamma_{0}^{0}, \quad \bmod \quad \omega^{1}, \omega^{2}, \theta, \rho_{0}
$$

One may translate, and denote

$$
\begin{equation*}
\epsilon_{7}=\epsilon_{8}=B_{0} . \tag{2.9}
\end{equation*}
$$

Let $\mathcal{B} \subset \mathcal{B}_{1}$ be the sub-bundle defined by (2.9). The structure group $P \subset P_{1}$ of $\mathcal{B}$ is reduced to

$$
\left.P=\left\{\begin{array}{cccc}
* & \cdot & \cdot & \cdot  \tag{2.10}\\
\cdot & \pm 1 & \cdot & \cdot \\
\cdot & \cdot & * & \cdot \\
\cdot & \cdot & \cdot & \pm 1
\end{array}\right)\right\}
$$

On $\mathcal{B} \rightarrow M$, the 1 -form $\gamma_{0}^{0}$ is semi-basic. One may write

$$
\begin{equation*}
\gamma_{0}^{0}=C_{1} \omega^{1}+C_{2} \omega^{2}+C_{9} \theta \tag{2.11}
\end{equation*}
$$

for the coefficients $C_{1}, C_{2}, C_{9} 4^{4}$
Step 4. Differentiating (2.5), (2.8), (2.11), and examining the rest of the components of $\Phi$, one obtains the following structure equations 5

$$
\begin{align*}
& d A_{0}=-A_{0} \rho_{0}+A_{0,1} \omega^{1}+A_{0,2} \omega^{2}+A_{0,0} \theta,  \tag{2.12}\\
& d A_{0,1} \equiv-2 A_{0,1} \rho_{0}+A_{0,1,1} \omega^{1}+\left(A_{0,1,1}+6 A_{0} A_{0,1}+4 A_{0} B_{0}+2 A_{0}^{3}+2 A_{0,0}-2 A_{1} A_{0}-5 B_{1}-2 C_{1}\right) \omega^{2}, \\
& d A_{0,2} \equiv-2 A_{0,2} \rho_{0}+\left(A_{0,1,1}-2 C_{1}+4 A_{0} B_{0}+5 A_{0} A_{0,1}-A_{0} A_{0,2}-2 A_{1} A_{0}+2 A_{0}^{3}+A_{0,0}-5 B_{1}\right) \omega^{1} \\
& +\left(A_{0,1,1}+5 B_{2}+2 C_{2}+5 A_{0} A_{0,2}-2 A_{0} A_{2}-2 C_{1}-5 B_{1}+3 A_{0,0}-2 A_{1} A_{0}+5 A_{0} A_{0,1}\right) \omega^{2}, \\
& d A_{0,0} \equiv-3 A_{0,0} \rho_{0}+\left(A_{0,1,0}+A_{0} B_{1}-2 A_{0} C_{1}-2 A_{0,2} A_{1}+4 A_{0,1} B_{0}-A_{0,1} A_{0,2}+2 A_{0,1} A_{0}^{2}+A_{0,1}^{2}\right) \omega^{1} \\
& +\left(A_{0,2,0}-4 A_{0,2} B_{0}-2 A_{0,1} A_{2}+A_{0} B_{2}-2 A_{0} C_{2}-A_{0,1} A_{0,2}-2 A_{0,2} A_{0}^{2}+A_{0,2}^{2}\right) \omega^{2}, \\
& d A_{1} \equiv-2 A_{1} \rho_{0}+2 A_{1} A_{0} \omega^{1}+\left(-2 B_{1}+2 A_{1} A_{0}\right) \omega^{2}, \\
& d A_{2} \equiv-2 A_{2} \rho_{0}+\left(-2 A_{0} A_{2}-2 B_{2}\right) \omega^{1}-2 A_{0} A_{2} \omega^{2}, \\
& d B_{0} \equiv-2 B_{0} \rho_{0}+\left(-C_{1}-2 B_{1}\right) \omega^{1}+\left(C_{2}+2 B_{2}\right) \omega^{2}, \\
& d B_{1} \equiv-3 B_{1} \rho_{0}+\left(-A_{0,2} A_{1}+A_{1} A_{0,1}+2 A_{1} A_{0}^{2}+A_{0} B_{1}+\frac{1}{2} A_{1,0}\right) \omega^{1} \\
& +\left(-4 B_{0}^{2}+2 C_{9}+A_{0} B_{2}+B_{2,1}-A_{1} A_{2}+A_{0} B_{1}\right) \omega^{2}, \\
& d B_{2} \equiv-3 B_{2} \rho_{0}+B_{2,1} \omega^{1}+\left(2 A_{0}^{2} A_{2}-A_{2} A_{0,2}-A_{0} B_{2}-\frac{1}{2} A_{2,0}+A_{0,1} A_{2}\right) \omega^{2}, \\
& d C_{1} \equiv-3 C_{1} \rho_{0}+C_{1,1} \omega^{1}+\left(-\frac{5}{2} B_{2,1}+A_{0} C_{1}+A_{1} A_{2}-\frac{5}{2} A_{0} B_{2}-2 C_{9}+4 B_{0}^{2}\right) \omega^{2}, \\
& d C_{2} \equiv-3 C_{2} \rho_{0}+\left(-\frac{5}{2} B_{2,1}-A_{0} C_{2}-\frac{5}{2} A_{0} B_{2}-3 C_{9}+6 B_{0}^{2}+\frac{3}{2} A_{1} A_{2}\right) \omega^{1}+C_{2,2} \omega^{2} \text {, } \\
& d C_{9} \equiv-4 C_{9} \rho_{0}+\left(-B_{2} A_{1}+2 B_{1} B_{0}+C_{1,0}+4 C_{1} B_{0}+C_{1} A_{0,1}-C_{1} A_{0,2}+2 C_{1} A_{0}^{2}-2 C_{2} A_{1}\right) \omega^{1} \\
& +\left(-2 B_{2} B_{0}-B_{1} A_{2}+C_{2,0}-2 C_{1} A_{2}-4 C_{2} B_{0}-A_{0,1} C_{2}+C_{2} A_{0,2}-2 C_{2} A_{0}^{2}\right) \omega^{2}, \\
& d B_{2,1} \equiv-4 B_{2,1} \rho_{0}+\left(-2 B_{1,0}-4 B_{1} A_{0}^{2}+2 B_{1} A_{0,2}-2 C_{1,0}-4 C_{1} A_{0}^{2}+2 C_{1} A_{0,2}-2 B_{1} A_{0,1}-2 C_{1} A_{0,1}\right. \\
& \left.-16 C_{1} B_{0}-16 B_{1} B_{0}-4 B_{2} A_{1}-A_{0,1} B_{2}-A_{0} B_{2,1}\right) \omega^{1}+\left(B_{2} A_{0,2}-2 A_{0,1} B_{2}+\frac{1}{2} A_{0} A_{2,0}-3 A_{0}^{2} B_{2}\right. \\
& \left.+4 B_{2} B_{0}+2 B_{2,0}-2 A_{0}^{3} A_{2}+4 C_{1} A_{2}+4 B_{1} A_{2}-A_{0} A_{0,1} A_{2}+A_{0} A_{0,2} A_{2}\right) \omega^{2}, \quad \bmod \theta .
\end{align*}
$$

The $\theta$-derivative terms, e.g., $A_{0,1,0}, A_{1,0}, B_{1,0}, \ldots$, are all independent with the one exception that

$$
B_{0,0}=C_{9}-2 B_{0}^{2}+A_{0} B_{2}+B_{2,1}-\frac{1}{2} A_{1} A_{2} .
$$

Remark 2.13. In the language of the theory of differential systems, this set of structure equations is involutive and a general analytic solution (linear Legendrian 3-web with the given structure equations) depends on three arbitrary functions of two variables as expected, see the remark at the beginning of this section. See [ $[\overline{\mathrm{BCG} 3]}$ for the details.

[^3]Proposition 2.14. Let $\mathcal{W}$ be a linear Legendrian 3-web on a connected open subset $M \subset \mathbb{P}^{3}$. There exists a canonically associated principal bundle $\mathcal{B} \subset \mathrm{Sp}_{2} \mathbb{C} \rightarrow M$ with the structure group (2.10). The induced $\mathfrak{s p}_{2} \mathbb{C}$ valued Maurer-Cartan form $\phi$, (2.3), is normalized on $\mathcal{B}$ such that

$$
\phi=\left[\begin{array}{cccc}
\rho_{0} & \frac{1}{2} A_{1} \omega^{1}+B_{0} \omega^{2}+B_{1} \theta & C_{1} \omega^{1}+C_{2} \omega^{2}+C_{9} \theta & B_{0} \omega^{1}-\frac{1}{2} A_{2} \omega^{2}+B_{2} \theta  \tag{2.15}\\
\omega^{1} & \phi_{1}^{1} & B_{0} \omega^{1}-\frac{1}{2} A_{2} \omega^{2}+B_{2} \theta & A_{2} \theta \\
2 \theta & \omega^{2} & -\rho_{0} & -\omega^{1} \\
\omega^{2} & A_{1} \theta & -\frac{1}{2} A_{1} \omega^{1}-B_{0} \omega^{2}-B_{1} \theta & -\phi_{1}^{1}
\end{array}\right],
$$

where $\phi_{1}^{1}=-A_{0} \omega^{3}+\left(A_{0,2}-2 B_{0}-A_{0,1}-2 A_{0}^{2}\right) \theta$. The structure coefficients $A_{i}, B_{j}, C_{k}$ and their derivatives satisfy (2.12).

Two linear Legendrian 3-webs $\mathcal{W}, \mathcal{W}^{\prime}$ are congruent up to $\mathrm{Sp}_{2} \mathbb{C}$ motion whenever the corresponding data $(\mathcal{B}, \phi)$ and $\left(\mathcal{B}^{\prime}, \phi^{\prime}\right)$ are isomorphic.

Let us rephrase the argument at the beginning of this section with a view to applying Proposition 2.14, Given a linear Legendrian 3 -web $\mathcal{W}$ on $M \subset \mathbb{P}^{3}$, it determines 3 sections $M^{i} \subset \mathrm{Z}, i=1,2,3$, by definition of the duality in Figure 1.2. The linearity of $\mathcal{W}$ implies that each $M^{i}$ is tangent to the fibers of $\pi_{1}$. Under the projection by $\pi_{1}, M^{i}$ is mapped to a surface $\Sigma^{i} \subset \mathbb{Q}^{3}$. The image $\pi_{1} \circ \pi_{0}^{-1}(\mathrm{x})$ is the dual null line of $\mathrm{x} \in M$ which intersects $\Sigma=\cup_{i=1}^{3} \Sigma^{i}$ at 3 points. The local geometry of a linear Legendrian 3-web in this way corresponds to the semiglobal geometry of a union of 3 pieces of surfaces in $\mathbb{Q}^{3}$. Before we proceed to the problem of deformation, let us consider an example where this dual interpretation allows a simple description of a class of linear Legendrian 3-webs.

Suppose for a linear Legendrian 3-web the relative invariant $A_{0}$ vanishes identically,

$$
A_{0} \equiv 0
$$

An analysis shows that in this case the Maurer-Cartan form $\phi$ reduces to

$$
\phi=\left[\begin{array}{cccc}
\rho_{0} & \frac{1}{2} A_{1} \omega^{1}+B_{0} \omega^{2} & \left(2 B_{0}^{2}+\frac{1}{2} A_{1} A_{2}\right) \theta & B_{0} \omega^{1}-\frac{1}{2} A_{2} \omega^{2}  \tag{2.16}\\
\omega^{1} & -2 B_{0} \theta & B_{0} \omega^{1}-\frac{1}{2} A_{2} \omega^{2} & A_{2} \theta \\
2 \theta & \omega^{2} & -\rho_{0} & -\omega^{1} \\
\omega^{2} & A_{1} \theta & -\frac{1}{2} A_{1} \omega^{1}-B_{0} \omega^{2} & 2 B_{0} \theta
\end{array}\right]
$$

where $d A_{1}=-2 A_{1} \rho_{0}, d A_{2}=-2 A_{2} \rho_{0}, d B_{0}=-2 B_{0} \rho_{0}$.
Choose the Legendrian foliation $\mathcal{F}^{1}$ defined by $\left\langle\omega^{2}, \theta\right\rangle^{\perp}$. The corresponding surface $\Sigma^{1} \subset \mathbb{Q}^{3} \subset \mathbb{P}\left(\bigwedge_{0}^{2} \mathbb{C}^{4}\right)$ is described by $\left[Z_{0} \wedge Z_{1}\right]$ (here we follow the notation from Section 1.1). A direct computation by successively differentiating $\left[Z_{0} \wedge Z_{1}\right]$ shows that $\Sigma^{1}$ is a part of a hyperplane section $H^{1} \subset \mathbb{Q}^{3}$. From the similar analysis for the foliations $\mathcal{F}^{2}, \mathcal{F}^{3}$, one concludes that;

Let $\mathcal{W}$ be a linear Legendrian 3-web with vanishing relative invariant $A_{0}$. Then $\mathcal{W}$ is a part of the algebraic Legendrian 3-web $\mathcal{W}_{\Sigma}$ induced by $\Sigma=\cup_{i=1}^{3} H^{i}$, a union of 3 hyperplane sections in $\mathbb{Q}^{3}$.

We shall see in Section 4 that this class of Legendrian 3-webs account for all the linear Legendrian 3-webs of maximum rank, with the only exception of the dual 3-web of the Legendrian twisted cubic curve.

## 3. DEFORMATION, AND RIGIDITY

In this section, we establish the fundamental structure equation for the linear Legendrian deformation of a linear Legendrian 3-web. A variant of the moving frame method is applied, and the analysis leads to the closed structure equation for the three deformation parameters. The differential compatibility conditions of this structure equation generate a sequence of polynomial equations that the deformation parameters must satisfy.

We currently have a partial understanding of the root structure of these polynomial equations. An elementary examination of the first few polynomials shows that; 1) the number of distinct linearizations of a Legendrian 3 -web
is uniformly bounded, Theorem 3.14, 2) if the Legendrian 3-web is sufficiently flat at a point, it admits at most one distinct local linearization, Theorem 3.16

Let $\mathcal{W}$ be a linear Legendrian 3 -web on a connected open subset $M \subset \mathbb{P}^{3}$. Let $\mathcal{B} \rightarrow M$ be the associated adapted bundle with the $\mathfrak{s p}_{2} \mathbb{C}$-valued normalized Maurer-Cartan form $\phi$, Proposition 2.14 Let $M \hookrightarrow \mathbb{P}^{3}$ be another distinct linearization of $\mathcal{W}$. Let $\mathrm{B}^{\prime} \rightarrow M$ be the associated $\mathrm{P}^{\prime} \simeq \mathrm{P}$-bundle with the Maurer-Cartan form $\phi^{\prime}$. We employ the method of moving frames to normalize the frame bundle $\mathrm{B}^{\prime}$ based at $\mathcal{B}$. The structure equation for the difference $\phi^{\prime}-\phi$ then gives the aforementioned polynomial compatibility equations.

Set

$$
\begin{equation*}
\phi^{\prime}=\phi+\delta \phi \tag{3.1}
\end{equation*}
$$

The components of $\delta \phi$ are denoted by

$$
\delta \phi=\left(\begin{array}{cc}
\delta \alpha & \delta \gamma \\
\delta \beta & -\delta \alpha^{t}
\end{array}\right)
$$

For the notational purpose, set

$$
\begin{equation*}
\Delta=d(\delta \phi)+\delta \phi \wedge \phi+\phi \wedge \delta \phi+\delta \phi \wedge \delta \phi \tag{3.2}
\end{equation*}
$$

Maurer-Cartan equations for $\phi^{\prime}$ and $\phi$ imply that $\Delta$ must vanish identically.
Step $\mathbf{0}^{\prime}$. Applying the fiber group action by $\mathrm{P}^{\prime} \subset \mathrm{Sp}_{2} \mathbb{C}$ as in Step 0 of Section 2 , one may translate

$$
\begin{equation*}
\delta \alpha_{0}^{1}=0, \delta \beta_{0}^{1}=0 ; \delta \beta_{0}^{0}=0 \tag{3.3}
\end{equation*}
$$

Under these relations, the structure group is reduced to

$$
\mathrm{P}_{0}^{\prime}=\left\{ \pm\left(\begin{array}{cccc}
1 & \cdot & * & \cdot \\
\cdot & 1 & \cdot & \cdot \\
\cdot & \cdot & 1 & \cdot \\
\cdot & \cdot & \cdot & 1
\end{array}\right)\right\}
$$

$\Delta_{0}^{2}$ gives $\delta \alpha_{0}^{0} \wedge \theta=0$. Applying the fiber group action by $\mathrm{P}_{0}^{\prime}$, one may translate

$$
\begin{equation*}
\delta \alpha_{0}^{0}=0 \tag{3.4}
\end{equation*}
$$

Under this relation, the structure group is reduced to the center of $\mathrm{Sp}_{2} \mathbb{C}, Z\left(\mathrm{Sp}_{2} \mathbb{C}\right)=\left\{ \pm I_{4}\right\}$. The $\mathfrak{s p}_{2} \mathbb{C}$-valued deformation 1-form $\delta \phi$ becomes

$$
\delta \phi=\left(\begin{array}{cccc}
\cdot & * & * & * \\
\cdot & * & * & * \\
\cdot & \cdot & \cdot & \cdot \\
\cdot & * & * & *
\end{array}\right)
$$

Starting from this initial state, by successive applications of (3.2), we intend to find the compatibility equations for the linear Legendrian deformation.

Differentiating (3.3), (3.4), one gets

$$
\left(\begin{array}{ccc}
\delta \alpha_{1}^{1} & \delta \gamma_{1}^{1} & 2 \delta \gamma_{1}^{0} \\
-\delta \beta_{1}^{1} & \delta \alpha_{1}^{1} & 2 \delta \alpha_{1}^{0} \\
\delta \alpha_{1}^{0} & \delta \gamma_{1}^{0} & 2 \delta \gamma_{0}^{0}
\end{array}\right) \wedge\left(\begin{array}{c}
\omega^{1} \\
\omega^{2} \\
\theta
\end{array}\right)=0 .
$$

By the Cartan's lemma, one may write

$$
\left(\begin{array}{c}
\delta \alpha_{1}^{1} \\
\delta \gamma_{1}^{1} \\
\delta \beta_{1}^{1} \\
\delta \alpha_{1}^{0} \\
\delta \gamma_{1}^{0} \\
\delta \gamma_{0}^{0}
\end{array}\right)=\left(\begin{array}{ccc}
\mu_{1} & \mu_{2} & 2 \mu_{9} \\
\mu_{2} & \mu_{4} & 2 \mu_{10} \\
\mu_{5} & -\mu_{1} & -2 \mu_{7} \\
\mu_{7} & \mu_{9} & 2 \mu_{11} \\
\mu_{9} & \mu_{10} & 2 \mu_{12} \\
\mu_{11} & \mu_{12} & \mu_{11,0}
\end{array}\right)\left(\begin{array}{c}
\omega^{1} \\
\omega^{2} \\
\theta
\end{array}\right),
$$

for the coefficients $\mu_{l}$.
Step 1'. The condition that it is a linear Legendrian deformation imposes a set of relations on $\mu_{k}$ 's. A computation similar as in Step 1 of Section 2 shows that $\mu_{4}=0, \mu_{5}=0, \mu_{1}-\mu_{2}=0$. We set accordingly

$$
\left(\begin{array}{l}
\delta \alpha_{1}^{1}  \tag{3.5}\\
\delta \gamma_{1}^{1} \\
\delta \beta_{1}^{1} \\
\delta \alpha_{1}^{0} \\
\delta \gamma_{1}^{0} \\
\delta \gamma_{0}^{0}
\end{array}\right)=\left(\begin{array}{ccc}
a_{0} & a_{0} & 2 b_{0} \\
a_{0} & \cdot & 2 b_{3} \\
\cdot & -a_{0} & -2 b_{1} \\
b_{1} & b_{0} & 2 c_{1} \\
b_{0} & b_{3} & 2 c_{2} \\
c_{1} & c_{2} & c_{9}
\end{array}\right)\left(\begin{array}{c}
\omega^{1} \\
\omega^{2} \\
\theta
\end{array}\right) .
$$

The derivatives of these coefficients will be denoted similarly as before, e.g., $d a_{0}=-a_{0} \rho_{0}+a_{0,1} \omega^{1}+a_{0,2} \omega^{2}+$ $a_{0,0} \theta$.

Step 2'. $\Delta_{3}^{1} \wedge \theta, \Delta_{1}^{3} \wedge \theta, \Delta_{1}^{1} \wedge \theta$ give

$$
\begin{align*}
& a_{0,1}=-4 b_{1}+a_{0} A_{0}+2 a_{0}^{2}, \\
& a_{0,2}=4 b_{3}-a_{0} A_{0}-2 a_{0}^{2},  \tag{3.6}\\
& b_{0}=b_{1}+b_{3}-a_{0} A_{0}-\frac{3}{4} a_{0}^{2} .
\end{align*}
$$

The remaining equations from $\Delta_{3}^{1}, \Delta_{1}^{3}, \Delta_{1}^{1}$ give

$$
\begin{align*}
b_{1,1}= & 2 A_{0} b_{1}-2 A_{1} a_{0}+2 b_{1} a_{0}, \\
b_{1,2}= & 2 c_{1}+2 A_{0} a_{0}^{2}+\frac{3}{2} a_{0}^{3}-2 a_{0} b_{3}-\frac{1}{2} A_{0,2} a_{0}+2 A_{0} b_{1}+\frac{1}{2} A_{0,1} a_{0}-A_{1} a_{0}+A_{0}^{2} a_{0}+\frac{1}{2} a_{0,0}, \\
b_{3,1}= & -2 c_{2}-2 A_{0} a_{0}^{2}-\frac{3}{2} a_{0}^{3}+2 b_{1} a_{0}-A_{0}^{2} a_{0}-2 A_{0} b_{3}-A_{2} a_{0}-\frac{1}{2} A_{0,1} a_{0}+\frac{1}{2} A_{0,2} a_{0}+\frac{1}{2} a_{0,0}, \\
b_{3,2}= & -2 A_{0} b_{3}-2 A_{2} a_{0}-2 b_{3} a_{0},  \tag{3.7}\\
c_{1}= & -\frac{9}{2} a_{0}^{3}-\frac{11}{2} A_{0} a_{0}^{2}+3 b_{1} a_{0}+\frac{2}{3} A_{0} b_{1}+\frac{10}{3} A_{0} b_{3}+6 a_{0} b_{3}-3 A_{0}^{2} a_{0}-\frac{4}{3} A_{0,1} a_{0} \\
& +\frac{5}{3} A_{0,2} a_{0}-2 B_{0} a_{0}+\frac{1}{2} A_{1} a_{0}, \\
c_{2}= & c_{1}-3 a_{0} b_{3}-\frac{8}{3} A_{0} b_{3}-\frac{1}{2} A_{2} a_{0}-\frac{1}{3} A_{0,1} a_{0}-\frac{1}{3} A_{0,2} a_{0}+3 b_{1} a_{0}+\frac{8}{3} A_{0} b_{1}-\frac{1}{2} A_{1} a_{0} .
\end{align*}
$$

The compatibility equations from $d\left(d\left(a_{0}\right)\right)=0$ give

$$
\begin{align*}
& a_{0,0}=-\frac{1}{3} A_{0,1} a_{0}-\frac{1}{3} A_{0,2} a_{0}+\frac{8}{3} A_{0} b_{1}-\frac{8}{3} A_{0} b_{3}, \\
& b_{1,0}=5 A_{0} a_{0}^{3}+\left(-\frac{2}{3} A_{0,2}+7 A_{0}^{2}+A_{1}+\frac{1}{3} A_{0,1}+2 B_{0}\right) a_{0}^{2}+\mathcal{O}\left(b_{1} a_{0}, b_{3} a_{0} ; a_{0}, b_{1}, b_{3}\right),  \tag{3.8}\\
& b_{3,0}=-5 A_{0} a_{0}^{3}+\left(-\frac{2}{3} A_{0,1}-7 A_{0}^{2}+A_{2}+\frac{1}{3} A_{0,2}-2 B_{0}\right) a_{0}^{2}+\mathcal{O}\left(b_{1} a_{0}, b_{3} a_{0} ; a_{0}, b_{1}, b_{3}\right) .
\end{align*}
$$

Here $\mathcal{O}\left(b_{1} a_{0}, b_{3} a_{0} ; a_{0}, b_{1}, b_{3}\right)$ means the terms that are linear combination of $\left\{b_{1} a_{0}, b_{3} a_{0} ; a_{0}, b_{1}, b_{3}\right\}$ with the coefficients in $A_{i}, B_{j}, C_{k}$ 's and their derivatives.

At this stage, there are three components $\Delta_{1}^{0}, \Delta_{3}^{0} ; \Delta_{2}^{0}$ left to be checked.
Step $3^{\prime} . \Delta_{1}^{0} \wedge \omega^{1}$ finally gives

$$
\begin{align*}
c_{9}= & \frac{81}{8} a_{0}^{4}+19 A_{0} a_{0}^{3}+\left(-26 A_{2}+27 b_{1}-4 A_{1}-105 b_{3}-\frac{5}{2} A_{0,2}-19 B_{0}-\frac{19}{2} A_{0,1}+5 A_{0}^{2}\right) a_{0}^{2}  \tag{3.9}\\
& +2 b_{1}^{2}+50 b_{3}^{2}+26 b_{1} b_{3}+\mathcal{O}\left(b_{1} a_{0}, b_{3} a_{0} ; a_{0}, b_{1}, b_{3}\right) .
\end{align*}
$$

It follows that $\delta \phi \equiv 0, \bmod a_{0}, b_{1}, b_{3}$, and that the structure equations for $\left\{a_{0}, b_{1}, b_{3}\right\}$ are closed, i.e., their derivatives are expressed as the functions of themselves and do not involve any new variables.
Remark 3.10. Note that

$$
d a_{0} \equiv-4\left(b_{1} \omega^{1}-b_{3} \omega^{2}\right), \quad \bmod \theta ; a_{0}
$$

Hence if $a_{0}$ vanishes up to order one at a point on a connected open subset $M$, the uniqueness theorem of ODE implies that $a_{0}, b_{1}, b_{3} \equiv 0$ identically. Hence in this case $\delta \phi=0$ and the deformation is trivial.
Proposition 3.11. Let $\mathcal{W}$ be a linear Legendrian 3-web on a connected open subset $M \subset \mathbb{P}^{3}$. Consider a linear Legendrian deformation of $\mathcal{W}$ represented by the $\mathfrak{s p}_{2} \mathbb{C}$-valued 1-form $\delta \phi$ satisfying the initial conditions (3.3), (3.4). Then the components of $\delta \phi$ are given by, and satisfy the structure equations (3.5), (3.6), (3.7), (3.8), (3.9).

Suppose $\delta \phi \equiv 0, \bmod \theta$ at a point in $M$. Then $\delta \phi$ vanishes identically and the deformation is trivial.

Proof. If $\delta \phi \equiv 0, \bmod \theta$ at a point, (3.5) shows that $a_{0}, b_{1}, b_{3}=0$ at the given point. The rest follows from Remark 3.10 .
3.1. Bound on the number of distinct linearizations. The remaining compatibility conditions from $\Delta_{1}^{0}, \Delta_{3}^{0}, \Delta_{2}^{0} ; d\left(d\left(b_{1}\right)\right)=0, d\left(d\left(b_{3}\right)\right)=0$, impose a set of polynomial equations (and their successive derivatives) on the deformation parameters $\left\{a_{0}, b_{1}, b_{3}\right\}$. The analysis of the root structure of these equations inevitably leads to a variety of case by case analysis problems depending on the relative invariants of the original 3-web $\mathcal{W}$, e.g., the resultants of a set of polynomial compatibility equations for $\left\{a_{0}, b_{1}, b_{3}\right\}$ are expressed in terms of the local invariants of $\mathcal{W}$.

Due to the complexity and size of the algebraic analysis involved, we shall consider the first few lowest order compatibility equations. In this section, we examine these equations without any extra conditions on the local invariants of the original web $\mathcal{W}$, and determine an upper bound on the number of distinct linearizations.

We continue the analysis from Step 3'.
The identities $d\left(d\left(b_{1}\right)\right) \wedge \theta=0, d\left(d\left(b_{3}\right)\right) \wedge \theta=0$ give a set of two compatibility equations which must vanish identically for a linear Legendrian deformation.

$$
\begin{align*}
& E q_{1} \equiv b_{1}^{2}+2 b_{1} b_{3},  \tag{3.12}\\
& E q_{3} \equiv b_{3}^{2}+2 b_{1} b_{3}, \quad \bmod \left\{\text { terms at most linear in } b_{1}, b_{3}\right\}
\end{align*}
$$

$E q_{1}, E q_{3}$ are polynomials of degree 4 in $\left\{a_{0}, b_{1}, b_{3}\right\}$. One computes then that $\Delta_{1}^{0}, \Delta_{3}^{0} \equiv 0$ modulo $E q_{1}, E q_{3}$. $\Delta_{2}^{0}$ modulo $E q_{1}, E q_{3}$ gives another set of two equations

$$
\begin{align*}
& E q_{0} \equiv a_{0}^{5}+\frac{314}{111} A_{0} a_{0}^{4},  \tag{3.13}\\
& E q_{9} \equiv A_{0} a_{0}^{4}, \quad \bmod \left\{\text { terms at most linear in } b_{1}, b_{3}, \text { and of degree } \leq 3 \text { in } a_{0}\right\} .
\end{align*}
$$

$E q_{0}, E q_{9}$ are polynomials of degree 5, 4 in $\left\{a_{0}, b_{1}, b_{3}\right\}$ respectively.
Theorem 3.14. Let $\mathcal{W}$ be a Legendrian 3-web on a connected contact three manifold $M$. Then it admits at most $4 \cdot 4 \cdot 5+1=81$ distinct local linearizations. If the local invariant $A_{0}$ of a linearization of $\mathcal{W}$ is nonzero, it admits at most $4 \cdot 4 \cdot 4+1=65$ distinct local linearizations. ${ }^{6}$

Proof. It is clear that $E q_{1}$ and $E q_{3}$ cannot have a common linear factor. The rest follows by counting the degrees of the polynomials $\left\{E q_{1}, E q_{3}, E q_{0}\right\}$, and $\left\{E q_{1}, E q_{3}, E q_{9}\right\}$.

It is unlikely that this is the optimal bound. We suspect that the optimal bound on the number of distinct local linearizations is two, see Conjecture 5.1 in Section 5. Theorem 3.14 shows that the number of distinct local linearizations of a Legendrian 3-web is uniformly bounded.
3.2. Rigidity of the linear Legendrian 3-webs with a flat point. Remark 3.10 implies that a linear Legendrian 3 -web is rigid under linear Legendrian deformation when the various compatibility equations force the deformation parameters $\left\{a_{0}, b_{1}, b_{3}\right\}$ to vanish at a single point. In this section we examine $\left\{E q_{1}, E q_{3}, E q_{0}\right\}$ in (3.12), (3.13), and show that this occurs in case $\mathcal{W}$ is sufficiently flat at a point.

Fix a reference point $\mathrm{x}_{0} \in M$. Assume all the coefficients $A_{i}, B_{j}, C_{k}$ 's and their derivatives of sufficiently high order vanish at $\mathrm{x}_{0}$. The compatibility equations $E q_{1}, E q_{3}, E q_{0}$ in (3.12), (3.13) evaluated at $\mathrm{x}_{0}$ become (up to scaling by constants)

$$
\begin{align*}
\left.\left(E q_{1}-E q_{3}\right)\right|_{\mathrm{x}_{0}} & =\left(b_{1}-b_{3}\right)\left(4 b_{1}+4 b_{3}-11 a_{0}^{2}\right), \\
\left.E q_{1}\right|_{\mathrm{x}_{0}} & =\left(4 b_{1}+5 a_{0}^{2}\right)\left(2 b_{1}+4 b_{3}-3 a_{0}^{2}\right),  \tag{3.15}\\
\left.E q_{0}\right|_{\mathrm{x}_{0}} & =a_{0}\left(185 a_{0}^{4}-60 a_{0}^{2} b_{3}-220 b_{1} a_{0}^{2}-208 b_{1} b_{3}\right) .
\end{align*}
$$

It is easily checked that the only root to this system of equations for $\left\{a_{0}, b_{1}, b_{3}\right\}$ is $a_{0}, b_{1},\left.b_{3}\right|_{\mathrm{p}_{0}}=0$. By Remark 3.10 in this case $\left\{a_{0}, b_{1}, b_{3}\right\}$ vanish identically and the deformation is trivial.

[^4]The following theorem describes up to which order the coefficients $A_{i}, B_{j}, C_{k}$ 's should vanish at the reference point to put the equations $E q_{1}, E q_{3}, E q_{0}$ into the reduced form (3.15).

Theorem 3.16. Let $\mathcal{W}$ be a linear Legendrian 3-web on a connected open subset $M \subset \mathbb{P}^{3}$. Let $A_{i}, B_{j}, C_{k}$ be the structure coefficients of $\mathcal{W}$, (2.15). Let $\mathrm{x}_{0} \in M$ be a reference point. Suppose

$$
\begin{aligned}
A_{0} & \text { vanishes to order } 3 \text { at } \mathrm{p}_{0}, \\
A_{1}, A_{2}, B_{1}, B_{2}, C_{1}, C_{2} & \text { vanish to order } 1 \text { at } \mathrm{p}_{0}, \\
B_{0}, C_{9} & \text { vanish at } \mathrm{p}_{0} .
\end{aligned}
$$

Then $\mathcal{W}$ is rigid and does not admit any nontrivial local linear Legendrian deformations.
Proof. Examining the equations $E q_{1}, E q_{3}, E q_{0}$, the vanishing conditions on $A_{i}, B_{j}, C_{k}$ are sufficient to imply (3.15). The rest follows from Remark 3.10

This agrees with the partial proof of the Gronwall conjecture for the planar 3-webs obtained in [Wa1].

## 4. LINEARIZATION OF THE LEGENDRIAN 3-WEBS OF MAXIMUM RANK

Theorem 3.16states that a linear Legendrian 3-web is rigid when it is sufficiently flat at a point. A question arises as to which abstract Legendrian 3-webs are such that their linearizations are likely to have the similar property. The first candidates would be the Legendrian 3-webs of maximum rank three, [Wa2]. The following refinement of Theorem 3.16 gives a partial proof of the Legendrian Gronwall Conjecture, Section [5, for this class of Legendrian 3 -webs.

Theorem 4.1. Let $\mathcal{W}$ be a Legendrian 3-web on a connected contact three manifold $M$. Suppose $\mathcal{W}$ has the maximum rank three.
a) $\mathcal{W}$ admits a local linearization $M \hookrightarrow \mathbb{P}^{3}$ as the dual Legendrian 3-web of an analytic surface $\Sigma \subset \mathbb{Q}^{3}$ which is the union of three hyperplane sections $\Sigma=\cup_{i=1}^{3} H^{i}, H^{i} \subset \mathbb{Q}^{2}$. Conversely, for any set of three distinct hyperplane sections $H^{i}$ in $\mathbb{Q}^{3}$, the dual Legendrian 3-web $\mathcal{W}_{\mathrm{U}_{i=1}^{3} H^{i}}$ has the maximum rank.
b) The linearization of $\mathcal{W}$ in a) is unique up to motion by $S p_{2} \mathbb{C}$, with the only exception when the structure invariants of $\mathcal{W}$ in (A-1) satisfy the relation

$$
\begin{equation*}
S=-T=2 R \neq 0 . \tag{4.2}
\end{equation*}
$$

In this case, the Legendrian 3-web admits exactly one more distinct local linearization as the dual Legendrian 3-web of the Legendrian twisted cubic.

Corollary 4.3. A Legendrian 3-web of maximum rank is algebraic.
Proof of Theorem 4.1
a) Let $\mathcal{W}$ be a Legendrian 3-web of maximum rank three. Then the structure invariants of $\mathcal{W}$ in (A-1) satisfy the relations

$$
\begin{aligned}
d R, d S, d T & \equiv 0, \quad \bmod \rho, \\
L=K & =0 .
\end{aligned}
$$

Recall (2.16) in Section 2 for the linear Legendrian 3-webs with the vanishing local invariant $A_{0}$. Substituting

$$
\begin{equation*}
B_{0}=\frac{R}{4}, A_{1}=-\frac{T}{2}, A_{2}=-\frac{S}{2} ; \rho_{0}=-\rho, \tag{4.4}
\end{equation*}
$$

one has $d \phi+\phi \wedge \phi=0$, and it induces a local linearization of $\mathcal{W}$.
b) Given the linearization of $\mathcal{W}$ defined by $\phi$ in (2.16), (4.4), the deformation analysis as in Section 3 shows that the various compatibility equations force $\delta \phi=0$, except when $S=-T=2 R \neq 0$ (this requires a long but
straightforward case by case analysis). In this case $\mathcal{W}$ admits exactly one more distinct linearization induced by the (deformed) Maurer-Cartan form

$$
\phi^{\prime}=\left[\begin{array}{cccc}
-\rho & \frac{3}{2} R \omega^{1}+\frac{3}{4} R \omega^{2} & -\frac{27}{8} R^{2} \theta & \frac{3}{4} R \omega^{1}+\frac{3}{2} R \omega^{2}  \tag{4.5}\\
\omega^{1} & \frac{1}{2} R \theta+a_{0} \omega^{1}+a_{0} \omega^{2} & \frac{3}{4} R \omega^{1}+\frac{3}{2} R \omega^{2} & R \theta+a_{0} \omega^{1} \\
2 \theta & \omega^{2} & \rho & -\omega^{1} \\
\omega^{2} & -R \theta-a_{0} \omega^{2} & -\frac{3}{2} R \omega^{1}-\frac{3}{2} R \omega^{2} & -\frac{1}{2} R \theta-a_{0} \omega^{1}-a_{0} \omega^{2}
\end{array}\right],
$$

where $a_{0}^{2}=2 R$.
We claim that this describes the dual 3-web of the Legendrian twisted cubic. Choose the linear Legendrian foliation $\mathcal{F}^{1}$ defined by $\left\langle\omega^{2}, \theta\right\rangle^{\perp}$ via $\phi^{\prime}$. The corresponding surface $\Sigma^{1} \subset \mathbb{Q}^{3}$ is described by $\left[Z_{0} \wedge Z_{1}\right]$. The conformal structure of $\mathbb{Q}^{3}$ is represented by the quadratic form $\beta_{0}^{0} \beta_{1}^{1}-\left(\beta_{0}^{1}\right)^{2}$, whose restriction on $\Sigma^{1}$ becomes the perfect square $\left(\omega^{2}-a_{0} \theta\right)^{2}$. Differentiating $Z_{0} \wedge Z_{1}$, one gets that the two dimensional subspace spanned by $\left\{Z_{0} \wedge Z_{1},\left(2 Z_{1}-a_{0} Z_{0}\right) \wedge\left(Z_{2}+\frac{a_{0}}{2} Z_{3}\right)\right\}$ is constant along the leaves of the foliation on $\Sigma^{1}$ defined by $\left(\omega^{2}-a_{0} \theta\right)^{\perp}$. By duality, $\Sigma^{1}$ is ruled by the null line dual to $\mathrm{p}_{1}=\left[2 Z_{1}-a_{0} Z_{0}\right] \in \mathbb{P}^{3}$. The similar analysis for the foliations $\mathcal{F}^{2}, \mathcal{F}^{3}$ shows that the corresponding surfaces $\Sigma^{2}, \Sigma^{3}$ are ruled by the null lines dual to $\mathrm{p}_{2}=\left[2 Z_{3}+a_{0} Z_{0}\right], \mathrm{p}_{3}=$ $\left[2\left(Z_{1}-Z_{3}\right)+a_{0} Z_{0}\right]$ respectively. Note that $\cap_{i} \mathbb{P}_{\mathrm{p}_{i}}^{2}=\left[Z_{0}\right]$.

Our claim is that the loci of $\mathrm{p}_{1}, \mathrm{p}_{2}, \mathrm{p}_{3}$, all lie in the same Legendrian twisted cubic. A short computation shows that each $p_{i}$ describes a Legendrian curve. One may then verify by direct computation that the following three quadratic polynomials generate the well defined, covariant constant, three dimensional space of quadratic polynomials in a neighborhood of $\left[Z_{0}\right]$ which vanish simultaneously on $\mathrm{p}_{i}$ 's. Here $W=\left(W_{0}, W_{1}, W_{2}, W_{3}\right)^{t}$ is the dual frame of $Z=\left(Z_{0}, Z_{1}, Z_{2}, Z_{3}\right)$ that satisfies the structure equation $d W=-\phi^{\prime} W$.

$$
\begin{aligned}
& Q_{0}=-R W_{1}^{2}-R W_{1} W_{3}+2 W_{0}^{2}-R W_{3}^{2}+\frac{27}{8} R^{2} W_{2}^{2}, \\
& Q_{1}=a_{0} W_{1}^{2}+2 W_{0} W_{1}-3 R W_{2} W_{3}+2 a_{0} W_{1} W_{3}-\frac{3}{2} R W_{1} W_{2}, \\
& Q_{3}=-2 a_{0} W_{1} W_{3}-a_{0} W_{3}^{2}+2 W_{0} W_{3}+\frac{3}{2} R W_{2} W_{3}+3 R W_{1} W_{2} .
\end{aligned}
$$

It is clear that these polynomials define the Legendrian twisted cubic in $\mathbb{P}^{3}$ that connects the three arcs traversed by $\mathrm{p}_{i}$ 's.

The analysis in the proof of b ) above has the following algebro-geometric implication. Let $\mathcal{W}$ be the dual 3-web of the Legendrian twisted cubic just described. Since the Gauss map of the Legendrian twisted cubic in $\mathbb{P}^{3}$ is the null rational normal curve in $\mathbb{Q}^{3} \subset \mathbb{P}^{4}, \mathcal{W}$ is also the dual 3-web of the analytic surface $\Sigma \subset \mathbb{Q}^{3}$ which is the tangent developable of the null rational normal curve.

Consider the following analogue of the converse of Abel's theorem, [CG, HP].

## Legendrian analogue of the converse of Abel's theorem

Let $\mathrm{x}_{0} \in \mathbb{P}^{3}$ be a reference point. Let $N_{0} \subset \mathbb{Q}^{3}$ be the null line dual to $\mathrm{x}_{0}$. Let $\Sigma^{i} \subset \mathbb{Q}^{3}, i=1,2, \ldots d$, be a set of $d$ distinct pieces of local analytic surfaces each of which intersects $N_{0}$ transversally at a single point $\mathrm{q}_{i}\left(\mathrm{x}_{0}\right)$. Let $\Omega_{i}$ be a meromorphic 1-form on $\Sigma^{i}$ which is regular at $\mathrm{q}_{i}\left(\mathrm{x}_{0}\right)$. Suppose the following local trace vanishes in a neighborhood of $\mathrm{x}_{0}$.

$$
\operatorname{Tr}\left(\Omega_{i}\right)=\sum_{i} \mathrm{q}_{i}^{*} \Omega^{i}=0
$$

Then there exists a null degree $d$ analytic surface $\Sigma$, and a meromorphic 1-form $\Omega$ on $\Sigma$, which analytically extends the given data $\cup_{i}\left(\Sigma^{i}, \Omega_{i}\right)$.

Assuming this is true, the relevant observation is that the tangent developable of the null rational normal curve lies in the intersection of $\mathbb{Q}^{3}$ with a cubic hypersurface of $\mathbb{P}^{4}$ It is known that a smooth complete intersection

[^5]of type $(2,3)$ in $\mathbb{P}^{4}$ is a K 3 surface, which has no nonzero holomorphic 1 -forms. According to Theorem 4.1, when this complete intersection degenerates to $\Sigma$, it supports three closed generalized holomorphic 1 -forms. We currently do not have any purely algebro-geometric explanation of this phenomenon.

## 5. Concluding remarks

1. The Legendrian analogue of the Gronwall conjecture can be stated as follows.

Conjecture 5.1 (Legendrian Gronwall conjecture). Let $\mathcal{W}$ be a Legendrian 3-web on a connected contact three manifold. Then it admits at most one distinct local linearization in $\mathbb{P}^{3}$, with the only exception when $\mathcal{W}$ is locally equivalent to the dual 3 -web of the Legendrian twisted cubic curve in $\mathbb{P}^{3}$ (in which case $\mathcal{W}$ admits exactly two distinct linearizations).

Note the discrepancy when compared with the planar web case. An algebraic planar 3-web admits infinitely many distinct local linearizations.

From the analysis carried out in this paper, one suspects that a generic Legendrian 3-web admits at most one distinct linearization, although we did not write down the condition for rigidity explicitly. Since the other extreme case of the maximum rank Legendrian 3-webs are treated in Section 4 the next step toward the proof of the Legendrian Gronwall conjecture would be to consider the class of Legendrian 3-webs with one, or two independent Abelian relations. One may then attempt to show that a Legendrian 3-web with two distinct linearizations necessarily possesses at least one Abelian relation. It appears to be a difficult problem to directly analyze the root structure of the polynomial integrability equations for the linear Legendrian deformation discussed in Section 3
2. Let $\Sigma \subset \mathbb{Q}^{3}$ be a null degree $d$ surface. We say that $\Sigma$ is extremal when the associated dual linear Legendrian $d$-web $\mathcal{W}_{\Sigma}$ has the maximum rank $\rho_{d}=\frac{(d-1)(d-2)(2 d+3)}{6}, W \mathrm{Wa} 2$. Theorem 4.1 implies that the only null degree 3 extremal surfaces are the union of three hyperplane sections, and the tangent developable to the null rational normal curve. Can one give an algebro-geometric proof of this? The direct differential analysis for the Legendrian $d$-webs of maximum rank for the case $d \geq 4$ is complicated. One may hope to generalize the algebro-geometric classification of the extremal null degree 3 surfaces to the extremal null degree $d$ surfaces in general.

One may start by considering the following sub-problem. Let $\gamma \subset \mathbb{P}^{3}$ be a degree $d$ curve. We say that $\gamma$ is extremal when the associated dual linear Legendrian $d$-web $\mathcal{W}_{\gamma}$ has the maximum rank $\rho_{d}$. Theorem 4.1 shows that it is possible for $\mathcal{W}_{\gamma}$ to have an Abelian relation which is not induced from a holomorphic 1 -form on $\gamma$. We suspect that this kind of auxiliary Abelian relations exist only when $\gamma$ is itself Legendrian. Can one give an algebro-geometric classification of the extremal Legendrian curves? As mentioned in Section 11 such an extremal Legendrian curve is tightly controlled by a large number of generalized addition laws.
3. As discussed briefly in [Wa2], the geometry of a Legendrian 2-web is locally equivalent to the geometry of a single scalar second order ODE up to point transformation. In particular, the structure of a Legendrian 2-web has local invariants, and not every two of them are locally equivalent, [Ca].

A question arises as to the geometric meaning of the linearization of a Legendrian 2-web. Can the linearizability be considered as a local counterpart of the notion of completeness of the associated projective connection discussed in [McK]?

## Appendix

Let $\mathcal{W}$ be a Legendrian 3 -web on a contact three manifold $M$. There exists a sub-bundle $B$ of the $\mathrm{GL}_{3} \mathbb{C}$ principal frame bundle of $M$ on which the tautological 1-forms $\theta ; \omega^{i}, i=1,2,3, \sum \omega^{i}=0$, satisfy the following structure equations.

$$
\begin{align*}
d\left(\begin{array}{c}
\omega^{1} \\
\omega^{2} \\
\theta
\end{array}\right) & =-\left(\begin{array}{ccc}
\rho & \cdot & \cdot \\
\cdot & \rho & \cdot \\
\cdot & \cdot & 2 \rho
\end{array}\right) \wedge\left(\begin{array}{c}
\omega^{1} \\
\omega^{2} \\
\theta
\end{array}\right)+\left(\begin{array}{c}
\theta \wedge\left(R \omega^{1}+S \omega^{2}\right) \\
\theta \wedge\left(T \omega^{1}-R \omega^{2}\right) \\
\omega^{1} \wedge \omega^{2}
\end{array}\right),  \tag{A-1}\\
d \rho & =\theta \wedge\left(L \omega^{1}+K \omega^{2}\right) .
\end{align*}
$$

Here $R, S, T, L, K$ are torsion coefficients. The contact structure on $M$ is defined by $\langle\theta\rangle^{\perp} . \mathcal{W}$ is defined by the three line fields $\left\langle\omega^{i}, \theta\right\rangle^{\perp}$. See [Wa2] for the derivation of this structure equation.

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    Key words and phrases. Legendrian 3-web, linearization, Gronwall conjecture.
    ${ }^{1}$ An analytic surface $\Sigma \subset \mathbb{Q}^{3} \subset \mathbb{P}^{4}$ has null degree $d$ when it intersects a generic null line of $\mathbb{Q}^{3}$ at $d$ points. Hence it has degree $2 d$ as a surface in $\mathbb{P}^{4}$.

[^1]:    ${ }^{2}$ Here we mean a choice of a section of a $\mathbb{P}^{1} \times \mathbb{P}^{1} \times \mathbb{P}^{1}$-bundle over $\gamma \times \gamma \times \gamma$.

[^2]:    ${ }^{3}$ The trace is obtained by the pull back of the map $\mathbb{P}^{3} \backslash \gamma \rightarrow \gamma^{(d)}$.

[^3]:    ${ }^{4}$ We use the subscript '9' in $C_{9}$ in place of ' 0 ' to indicate that $C_{9}$ has the higher scaling weight than $C_{1}, C_{2}$ under the action of the structure group $P$, see 2.12 .

    5 The method of differential analysis used here is referred to as the prolongation, [BCG3]. It is the process of successively adding the derivatives as the new variables, under the condition of contact which indicates that these new variables are the derivatives. It allows one to access the differential relations (not necessarily of higher order) which are possibly hidden and can only be detected by examining the higher order derivatives. The computation was carried out using the Maple.

[^4]:    ${ }^{6}$ In case the invariant $A_{0} \equiv 0$ identically, the linear Legendrian 3-web admits at most two distinct local linearizations, see Section 4

[^5]:    ${ }^{7}$ Let $V_{m}$ be the irreducible $\mathrm{SL}_{2} \mathbb{C}$-module of dimension $m+1$. By Clebsch-Gordan, the symmetric cubic tensor product decomposes into $S^{3}\left(V_{4}\right)=V_{12} \oplus V_{8} \oplus V_{6} \oplus V_{4} \oplus V_{0}$. The $V_{0}$ piece vanishes on the tangent developable of the null rational normal curve.

