

压电体椭圆孔边的力学分析

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摘要 基于复变函数的方法, 以 PZT-4 材料为例, 分别采用精确电边界条件和非导电电边界条件进行了远场均匀载荷作用下的横观各向同性压电体椭圆孔的力学分析并与相关结果进行对比. 结果表明当椭圆孔退化为圆孔时, 无论在远场作用力载荷或电载荷, 两种电边界条件下的结果均能完全吻合. 随着椭圆孔的愈加尖锐化, 非导电电边界条件逐渐不能适用.

关键词 横观各向同性, 压电体, 椭圆孔, 圆孔, 电边界

引 言

近十多年来, 压电陶瓷材料在较大的领域内日渐为人所瞩目, 以含有各种缺陷的无限大压电体为研究对象的论文纷纷呈现. 运用多种方法, 得到了大量结果, 其中的一些结果还有待实验进一步验证.

在这一领域中, 由于三维问题的复杂性, 当前的的工作大都集中在二维尺度内. 其中裂纹问题备受关注. 在 1980 年, Deeg 分析了压电体中的位错、裂纹和夹杂问题, 提出裂纹上下表面电位移的径向分量可以认为等于零. 为了验证这一假设的合理性, Pak^[1] 详细讨论了忽略裂纹内电位移时的情况. Sosa 和 Pak^[2] 研究了一个很一般的裂纹尖端问题. Pak^[3] 和 Suo^[4,5] 分析了有限裂纹的弹性场和电场. 此外还有大量工作^[6~8] 分析研究了裂纹与界面裂纹问题. 另外, 由于椭圆孔问题所具有的重要意义, 有较多工作研究了压电体的椭圆孔问题. Sosa^[9] 和 Sosa, Khutoryansky^[10] 基于复变函数法采用不同的电边界条件研究了含椭圆孔的横观各向同性压电介质的平面应变问题, Gao 等^[11] 和 Zhang 等^[12] 也对此问题进行了分析, Chung 和 Ting^[13] 基于 Stroh 公式法分析了压电介质椭圆孔问题. 在这些文献中, 多数并没有全面具体地对孔边力学量的变化进行计算和分析. 最近, Deng 和 Wang^[14] 基于 Stroh 公式法进行了该问题的分析.

本文在以上工作的基础上, 运用 Lekhnistkii 复变法基于两种不同的运用广泛的电边界条件对横观各向同性压电体椭圆孔边力学量的变化进行了分析计算, 并与 Sosa^[9] 及 Deng 等^[14] 的结果进行了比

较, 得出了一些有益的结论.

1 问题的复变描述

考虑含一个受有远场均匀力电载荷的含椭圆孔 (长轴 x 半长 a , 短轴 z 半长 b) 的横观各向同性压电体. 为了在目前情况下能够解析求解, 仅考虑平面应变问题. 以下将全部以 x_1, x_2, x_3 分别替代 x, z, y ; 令 $z_k = x_1 + \mu_k x_2$, 则参照文献 [9] 可推出 x_1-x_2 平面内的场解可表示为

$$\left. \begin{aligned} \langle \sigma_{11}, \sigma_{22}, \sigma_{12} \rangle &= 2\text{Re} \sum_{k=1}^3 \langle \mu_k^2, 1, -\mu_k \rangle \varphi'_k(z_k) \\ \langle D_1, D_2 \rangle &= 2\text{Re} \sum_{k=1}^3 \langle -\lambda_k \mu_k, \lambda_k \rangle \varphi'_k(z_k) \\ \langle E_1, E_2 \rangle &= -2\text{Re} \sum_{k=1}^3 \langle d_k, d_k \mu_k \rangle \varphi'_k(z_k) \\ \phi &= 2\text{Re} \sum_{k=1}^3 d_k \varphi_k(z_k) \\ \langle u, w \rangle &= 2\text{Re} \sum_{k=1}^3 \langle p_k, q_k \rangle \varphi_k(z_k) \end{aligned} \right\} \quad (1)$$

其中 ϕ 为电势; $\mu_k, \lambda_k, d_k, p_k, q_k$ 可由二维电弹性的控制方程、变形协调方程、几何方程及本构方程求出. 其边界条件可表示为

$$\left. \begin{aligned} \frac{dU}{dx_1} &= - \int_l t_{2s} ds = - \int_l \sigma_{12} dx_2 - \sigma_{22} dx_1 \\ \frac{dU}{dx_2} &= \int_l t_{1s} ds = \int_l \sigma_{11} dx_2 - \sigma_{12} dx_1 \\ - \int_l D_n ds &= - \int_l D_1 dx_2 - D_2 dx_1 \end{aligned} \right\} \quad (2)$$

其中 U 是平面问题的应力函数. 对于本问题而言, 远场受均匀力电载荷. 在孔边, 可将孔内真空视作刚度为零的夹杂, 并如相关文献一样假设孔内电场为常电场, 则孔边条件可由上式表示, 并要求两边电势在孔端相等. 此时可将问题分解为两部分, 一部分是没有孔的一个无限大板在受有远场载荷时的应力场, 另一部分是出现孔后在孔边施加相关载荷时的应力场. 两部分叠加后, 远场处的应力应等于所加载荷, 孔边的应力场则应使之保持为自由面, 从而使该问题得解.

2 结果与讨论

下面我们将对椭圆孔边力学量的变化进行具体分析, 为便于与相关文献进行比较, 将以 PZT-4 陶瓷为例, 并采用 Sosa^[9] 给出的压电材料参数和真空介电系数. 本文结果经 Mathematic 整理、计算, 因而具有很高的计算精度. 需说明的是, 精确电边界条件在上面已给出来了, 非导电边界条件则是公式 (2) 的第 3 式所表示的压电体在孔边的值等于零 (即认为孔内真空的电位移为零^[11]) 而不再要求孔端电势相等, 其他两式不变. 由此就可以求出这两个不同电边界的解. 下面试作一些讨论:

当孔为中心圆孔时, 由图 1 至图 9 可以看出无论远场作用力或电载荷, 两个电边界的解都能完全吻合, 并且比较本文结果与 Deng 等^[14] 中的相应结果可知二者的两个电边界解的吻合程度都是令人满意的. 在远场只作用力载荷 $\sigma_{22}^{\infty} = \sigma_0$ 的情况下, Sosa^[9] 的结果与本文非导电边界的结果是一样的. 但在远场只作用电载荷 $D_2^{\infty} = D_0$ 时, Sosa^[9] 的部分结果与本文非导电边界解是完全不同的, 由于 Sosa^[9]

也是运用非导电边界进行求解的, 所以他的计算也许存在错误.

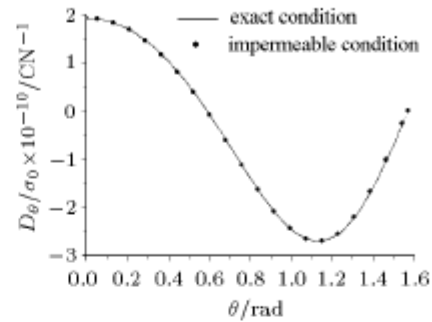


图 2 无穷远处力载荷下两种电边界条件的 D_{θ} 圆孔边分布
Fig.2 The normalized electric displacement D_{θ} variations on the rim of a circular hole subjected to remote tension loading $\sigma_{22}^{\infty} = \sigma_0$

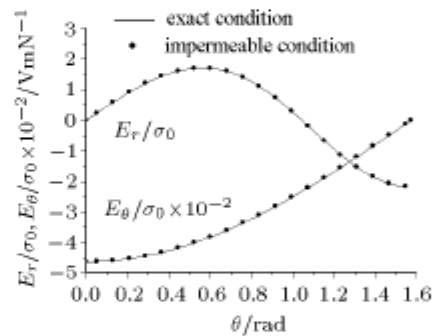


图 3 无穷远处力载荷下两种电边界条件的 E_r, E_{θ} 圆孔边分布
Fig.3 The normalized electric fields E_r, E_{θ} variations on the rim of a circular hole subjected to remote tension loading $\sigma_{22}^{\infty} = \sigma_0$

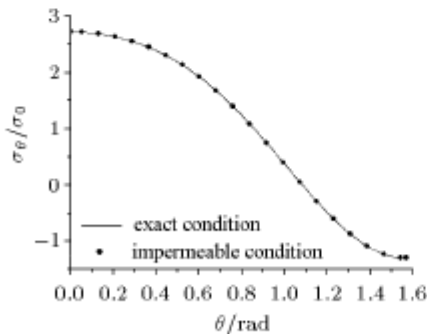


图 1 无穷远处力载荷下两种电边界条件的 σ_{θ} 圆孔边分布
Fig.1 The non-dimensional normal stress σ_{θ} variations on the rim of a circular hole subjected to remote tension loading $\sigma_{22}^{\infty} = \sigma_0$

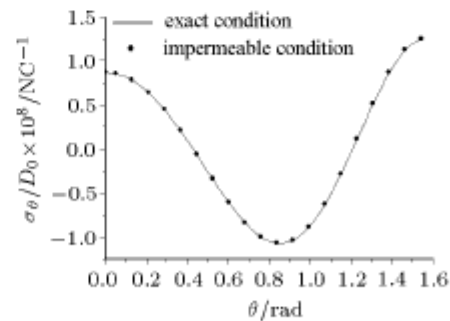


图 4 无穷远处电载荷下两种电边界条件的 σ_{θ} 圆孔边分布
Fig.4 The normalized stress σ_{θ} variations on the rim of a circular hole subjected to remote electric displacement loading $D_2^{\infty} = D_0$

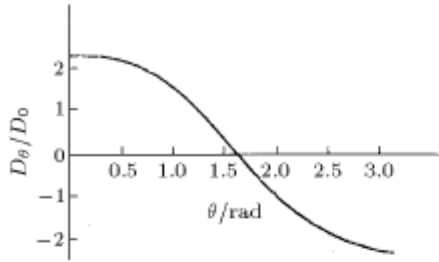


图 5 无穷远处电荷载荷下, Sosa 的 σ_θ 圆孔边分布
Fig.5 The normalized stress σ_θ variations on the rim of a circular hole subjected to remote electric displacement loading $D_2^\infty = D_0$. The results by Sosa

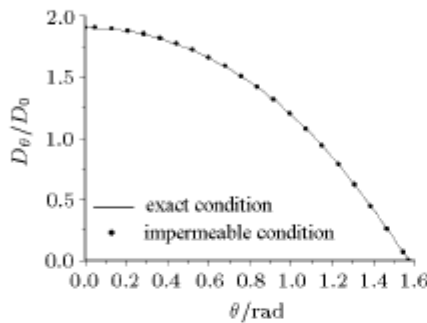


图 6 无穷远处电荷载荷下两种电边界条件的 D_θ 圆孔边分布
Fig.6 The non-dimensional normal electric displacement D_θ variations on the rim of a circular hole subjected to remote electric displacement loading $D_2^\infty = D_0$

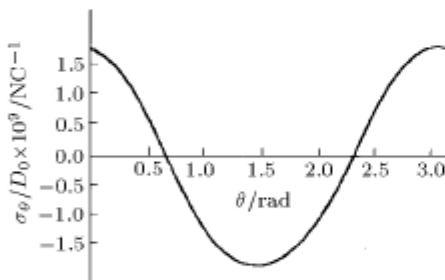


图 7 无穷远处电荷载荷下 Sosa 的 D_θ 圆孔边分布
Fig.7 The non-dimensional normal electric displacement D_θ variations on the rim of a circular hole subjected to remote electric displacement loading $D_2^\infty = D_0$. The results by Sosa

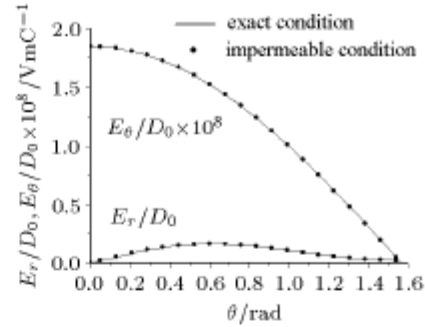


图 8 无穷远处电荷载荷下两种电边界条件的 E_r, E_θ 圆孔边分布
Fig.8 The non-dimensional normal electric fields E_r, E_θ variations on the rim of a circular hole subjected to remote electric displacement loading $D_2^\infty = D_0$

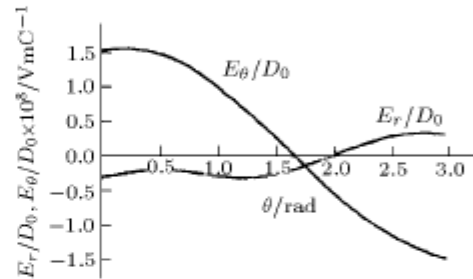


图 9 无穷远处电荷载荷下 Sosa 的 E_r, E_θ 圆孔边分布
Fig.9 The non-dimensional normal electric fields E_r, E_θ variations on the rim of a circular hole subjected to remote electric displacement loading $D_2^\infty = D_0$. The results by Sosa

当椭圆孔为不同的形状时, 两种载荷下两个电边界条件的解在表 1, 表 2 中均已给出, 从中可以看出, 随着孔的尖锐化, 非导电边界条件将越来越不适用. 另外, 这些解与 Deng 等^[14] 的解在变化趋势上是一致的, 但在数字上逐渐有了一些不同, 这

表 1 $\sigma_{22}^\infty = \sigma_0$ 条件下的 $\sigma_\theta(a, 0)$ 和 $D_\theta(a, 0)$

Table 1 The normal stress $\sigma_\theta(a, 0)$ and the normal electric displacement $D_\theta(a, 0)$ under remote tension $\sigma_{22}^\infty = \sigma_0$

b/a	$\theta = 0, r = a$					
	The impermeable electric boundary condition		The exact electric boundary condition		Relative error (%)	
	σ_θ/σ_0	$D_\theta/\sigma_0 \times 10^{-10}$	σ_θ/σ_0	$D_\theta/\sigma_0 \times 10^{-10}$	stress	electric displacement
0.001	1722.46	1924.68	1732.23	2938.19	0.564 013	34.494 37
0.01	173.146	192.468	173.308	209.271	0.093 475	8.029 302
1/30	52.643 8	57.740 5	52.6596	59.371 6	0.030 004	2.747 273
0.1	18.214 6	19.246 8	18.2165	19.443 5	0.010 43	1.011 649
1/3	6.164 38	5.774 05	6.164 59	5.795 87	0.003 407	0.376 475
1	2.721 46	1.924 68	2.721 5	1.928 39	0.001 47	0.192 388
3	1.573 82	0.641 562	1.573 83	0.642 4	0.000 635	0.130 448

表 2 $D_2^\infty = D_0$ 条件下的 $\sigma_\theta(a, 0)$ 和 $D_\theta(a, 0)$
 Table 2 The normal stress $\sigma_\theta(a, 0)$ and the normal electric displacement $D_\theta(a, 0)$
 under remote tension $D_2^\infty = D_0$

b/a	$\theta = 0, r = a$					
	The impermeable electric boundary condition		The exact electric boundary condition		Relative error (%)	
	$\sigma_\theta/D_0 \times 10^8$	D_θ/D_0	$\sigma_\theta/D_0 \times 10^8$	D_θ/D_0	stress	electric displacement
0.001	872.545	906.159	493.051	512.48	76.968 51	76.818 41
0.01	87.254 5	91.515 9	80.959 8	84.985 9	7.775 093	7.683 628
1/30	26.176 4	28.154 8	25.564 5	27.52	2.393 554	2.306 686
0.1	8.725 45	10.051 6	8.651 41	9.974 79	0.855 814	0.770 041
1/3	2.617 64	3.715 48	2.609 35	3.706 88	0.317 704	0.232 001
1	0.872 545	1.905 16	0.871 118	1.903 68	0.163 812	0.077 744
3	0.290 848	1.301 72	0.290 521	1.301 38	0.112 556	0.026 126

应该是由于本文中孔内介质是真空而文献 [14] 中是硅油.

3 结 论

本文基于复变函数的方法, 以 PZT-4 材料为例, 分别采用精确电边界条件和非导通电边界条件进行了远场均匀载荷作用下的横观各向同性压电体椭圆孔的力学分析并与 Sosa^[9] 和 Deng 等^[14] 的结果进行了对比. 结果表明当椭圆孔退化为圆孔时, 无论在远场作用力载荷或电载荷, 两种电边界条件下的结果均能完全吻合. 随着椭圆孔的愈加尖锐化, 非导通电边界条件逐渐不能适用. 同时, 本文也验证了 Deng 等^[14] 得出的 Sosa^[9] 的部分结果有错的结论, 同时也指出这只是 Sosa 的计算错误而并非是由非导通电边界条件所造成的.

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ON THE PROBLEM OF PIEZOELECTRIC SOLID WITH AN ELLIPTIC HOLE

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Abstract An electroelastic analysis is performed on a transversely isotropic material containing elliptic hole, which subject to a uniform stress field σ^∞ and a uniform electric displacement field D^∞ at infinity while the surface of the hole is free of traction and electrically open. Taking PZT-4 ceramic into consideration, the results based on the exact electric boundary conditions and the impermeable electric boundary conditions are all obtained. It can be seen that the related results of these two electric boundary conditions are entirely identical when the hole is a circular hole. The impermeable electric boundary condition is becoming incorrect with the elliptic hole degenerating into a crack. A comparison between these present results and the related results shows that the present results are validity. It also shows that some results given by Sosa are incorrect maybe due to his calculation, not the impermeable electric boundary.

Key words transversely isotropic, piezoelectric, elliptical hole, circular hole, electric boundary