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广义经典完整非保守力学系统 Lie 对称性及其守恒量

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摘要:研究了广义完整非保守力学系统的 Lie 对称性及其守恒量. 建立了系统的运动微分方程, 给出其确定方程、结构方程和守恒量, 得到了系统的 Lie 对称性定理和逆定理, 最后举例说明结果的应用.

关键词:广义力学系统; Lie 对称性; 守恒量

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Lie symmetries and conserved quantities of holonomic non-conservative dynamical systems in generalized classical mechanics

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Abstract: Lie symmetries and conserved quantities of holonomic non-conservative dynamical systems in generalized classical mechanics were studied. First, the determining equations of Lie symmetries are established by using Lie's invariance method of infinitesimal transformations of differential equations. Second, the structure equations and the conserved quantities were obtained. Third, the inverse Lie symmetries system problems were discussed. Finally, an example to illustrate the application of the results was given.

Key words: generalized mechanics system; Lie symmetries; conserved quantities

0 引言

描述动力学系统的 Lagrange 量含广义坐标对时间高阶微商情形(广义力学系统)的研究,在数学、力学和物理学等诸多领域广为应用. Ostrogradsky 和 Jacobi 最早对其进行了开创性研究^[1],文献[2]用现代数学方法描述了广义经典力学系统和场论. 近年来,对于此类系统的研究更加受到人们的关注,并取得许多研究成果^[3-11].

动力学系统的守恒量在力学研究中具有十分重要的作用,它不仅能使复杂的运动微分方程易于求解,而且反映着深刻的物理本质. 对称性和守恒量之间的联系在经典理论中通常由 Noether 定理与 Lie 对称性定理给出^[12-17],文献[18]研究了广义经典完整非保守力学系统的 Noether 定理,本文研究系统的 Lie 对称性,建立系统逆变代数形式的运动方程,给出系统的确定方程、结构方程和守恒量,得到 Lie 对称性定理及其逆定理,最后举例说明结果的应用.

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1 系统的运动微分方程

研究广义完整非保守力学系统. 设系统的位形由 n 个广义坐标 $q_s (s = 1, \dots, n)$ 确定, 其 Lagrange 函数 L 是时间 t , 广义坐标 $q_s (s = 1, \dots, n)$ 及其一阶到 k 阶导数的函数, 系统受广义非势力 $Q_s = Q_s(q_s, \dot{q}_s, \dots, q_s^{(k)}, t)$, 则广义 Euler-Lagrange 方程为

$$\sum_{j=0}^k (-1)^j \frac{d^j}{dt^j} \frac{\partial L}{\partial q_s^{(j)}} + Q_s(q_s, \dot{q}_s, \dots, q_s^{(k)}, t) = 0 \quad (1)$$

引入广义动量和哈密顿函数

$$p_{s/m} = \sum_{j=0}^{k-m} (-1)^j \frac{d^j}{dt^j} \frac{\partial L}{\partial q_s^{(j+m)}}, \quad (s = 1, \dots, n; m = 1, \dots, k) \quad (2)$$

$$H(q_s^{(m-1)}, p_{s/m}, t) = \sum_{s=1}^n \sum_{m=1}^k p_{s/m} q_s^{(m)} - L, \quad (3)$$

则方程(1)可写成正则形式

$$\frac{d}{dt} q_s^{(m-1)} = \frac{\partial H}{\partial p_{s/m}}, \quad \frac{d}{dt} p_{s/m} = - \frac{\partial H}{\partial q_s^{(m-1)}} + \tilde{Q}_s(q_s^{(m-1)}, p_{s/m}, t). \quad (s = 1, \dots, n; m = 1, \dots, k) \quad (4)$$

式中

$$\tilde{Q}_s(q_s^{(m-1)}, p_{s/m}, t) = Q_s(q_s^{(m-1)}, \dot{q}_s^{(m-1)}(q_s^{(m-1)}, p_{s/m}, t), t). \quad (s = 1, \dots, n; m = 1, \dots, k) \quad (5)$$

令

$$\tilde{Q}_s = \Omega_{s/m} \frac{\partial H}{\partial p_{s/m}}, \quad (6)$$

$$a^\mu = \begin{cases} q_s^{(m-1)}, & (\mu = (2m-2)n + s; m = 1, \dots, k; s = 1, \dots, n) \\ p_{s/m}, & (\mu = (2m-1)n + s; m = 1, \dots, k; s = 1, \dots, n) \end{cases} \quad (7)$$

$$R_\mu = \begin{cases} p_{s/m}, & (\mu = (2m-2)n + s; m = 1, \dots, k; s = 1, \dots, n) \\ 0, & (\mu = (2m-1)n + s; m = 1, \dots, k; s = 1, \dots, n) \end{cases} \quad (8)$$

$$\Lambda_\mu = \begin{cases} \Omega_{s/m} \frac{\partial H}{\partial p_{s/m}}, & (\mu = (2m-2)n + s; m = 1, \dots, k; s = 1, \dots, n) \\ 0, & (\mu = (2m-1)n + s; m = 1, \dots, k; s = 1, \dots, n) \end{cases} \quad (9)$$

则方程(4)可进一步表示为

$$\left(\frac{\partial R_\nu}{\partial a^\mu} - \frac{\partial R_\mu}{\partial a^\nu} \right) \dot{a}^\nu - \frac{\partial H}{\partial a^\mu} = - \Lambda_\mu. \quad (\mu, \nu = 1, \dots, 2kn) \quad (10)$$

与式(10)相对应的逆变代数方程为^[19]

$$\dot{a}^\mu - s^{\mu\nu} \frac{\partial H}{\partial a^\mu} = 0. \quad (\mu, \nu = 1, \dots, 2kn) \quad (11)$$

此处

$$s^{\mu\nu} = \omega^{\mu\nu} + T^{\mu\nu}. \quad (12)$$

其中

$$\omega^{\mu\nu} = \frac{\partial R_\nu}{\partial a^\mu} - \frac{\partial R_\mu}{\partial a^\nu}, \quad (\omega_{\mu\nu}) = (\omega^{\mu\nu})^{-1}, \quad (13)$$

且

$$\begin{cases} (\omega^{\mu\nu}) = \begin{bmatrix} M_1 & & \\ & \ddots & \\ & & M_k \end{bmatrix}_{2nk \times 2nk}, \\ M_i = \begin{bmatrix} 0_{n \times n} & I_{n \times n} \\ -I_{n \times n} & 0_{n \times n} \end{bmatrix}. \quad (i = 1, \dots, k) \end{cases} \quad (14)$$

显然,系统存在守恒量式(22).

3 系统 Lie 对称性的逆定理

假设已知广义完整非保守力学系统(1)有初积分

$$I = I(t, a) = \text{const}. \quad (23)$$

则

$$\frac{dI}{dt} = \frac{\partial I}{\partial t} + \sum_{\mu=1}^{2kn} \frac{\partial I}{\partial a^\mu} \dot{a}^\mu = 0. \quad (24)$$

将系统运动方程(10)的两端乘以 $\bar{\xi}_\mu^\alpha = \xi_\mu^\alpha - \dot{a}^\mu \xi_0^\alpha$, 并对 s 求和,得

$$\sum_{\mu=1}^{2kn} \bar{\xi}_\mu^\alpha \left(\sum_{\nu=1}^{2kn} \omega^{\nu\alpha} \dot{a}^\nu - \frac{\partial H}{\partial a^\mu} + \Lambda_\mu \right) = 0. \quad (\alpha = 1, \dots, r) \quad (25)$$

将式(24)与(25)相加后,分离出含 \dot{a}^ν 的项,并令其系数为0,得到

$$\frac{\partial I}{\partial a^\nu} + \sum_{\mu=1}^{2kn} \omega_{\nu\mu} \xi_\mu^\alpha + \left(\frac{\partial H}{\partial a^\mu} - \Lambda_\nu \right) \xi_0^\alpha = 0. \quad (\nu = 1, \dots, 2kn, \alpha = 1, \dots, r) \quad (26)$$

考虑到式(13),有

$$\xi_\mu^\alpha = \sum_{\nu=1}^{2kn} \omega^{\nu\alpha} \left[\frac{\partial I}{\partial a^\nu} + \left(\frac{\partial H}{\partial a^\nu} - \Lambda_\nu \right) \xi_0^\alpha \right]. \quad (\mu = 1, \dots, 2kn, \alpha = 1, \dots, r) \quad (27)$$

其中

$$\sum_{\nu=1}^{2kn} \omega^{\nu\alpha} \omega_{\nu\rho} = \delta_{\mu\rho}. \quad (\mu, \rho = 1, \dots, 2kn) \quad (28)$$

为使式(17)为 Lie 对称变换,令初积分式(23)等于守恒量式(22),得

$$\sum_{\mu=1}^{2kn} R_\mu \xi_\mu^\alpha - H \xi_0^\alpha + G^\alpha = I. \quad (\alpha = 1, \dots, r) \quad (29)$$

定义 2 如果式(27)和(29)确定的无限小生成元 $\xi_0^\alpha, \xi_\mu^\alpha$ 满足确定方程(20). 则称相应的变换为 Lie 对称变换.

4 举例

设力学系统的 Lagrange 函数为

$$L = \frac{1}{2} \dot{q}_1^2 + \frac{1}{2} q_1^2 \dot{q}_2^2 + \frac{1}{2} \ddot{q}_1^2. \quad (30)$$

系统受有非势广义力为 $Q = -\dot{q}_1$, 试研究系统的对称性与守恒量.

解 首先研究 Lie 对称性正问题

第一步 列写系统的运动微分方程

由式(2)求得

$$\begin{aligned} P_{1/1} &= \frac{\partial L}{\partial \dot{q}_1} - \frac{d}{dt} \frac{\partial L}{\partial \ddot{q}_1} = \dot{q}_1 - \ddot{q}_1, \\ P_{1/2} &= \frac{\partial L}{\partial \dot{q}_2} = \dot{q}_1, \\ P_{2/1} &= \frac{\partial L}{\partial \dot{q}_2} - \frac{d}{dt} \frac{\partial L}{\partial \ddot{q}_2} = q_1^2 \dot{q}_2, \\ P_{2/2} &= \frac{\partial L}{\partial \ddot{q}_2} = 0. \end{aligned} \quad (31)$$

式(3)给出

$$H = P_{1/1} \dot{q}_1 + P_{1/2} \dot{q}_1 + P_{2/1} \dot{q}_2 - \frac{1}{2} \dot{q}_1^2 - \frac{1}{2} q_1^2 \dot{q}_2^2 - \frac{1}{2} \ddot{q}_1^2, \quad (32)$$

令

$$\begin{aligned}
a^1 &= q_1, & a^2 &= q_2, & a^3 &= P_{1/1}, & a^4 &= p_{2/1}, \\
a^5 &= \dot{q}_1, & a^6 &= \dot{q}_2, & a^7 &= P_{1/2}, & a^8 &= 0, \\
R_1 &= P_{1/1}, & R_2 &= P_{2/1}, & R_3 &= 0, & R_4 &= 0, \\
R_5 &= P_{1/2}, & R_6 &= P_{2/2}, & R_7 &= 0, & R_8 &= 0.
\end{aligned} \tag{33}$$

则有

$$H = a^3 a^5 + \frac{1}{2} (a^7)^2 + a^4 a^6 - \frac{1}{2} (a^5)^2 - \frac{1}{2} (a^1)^2 (a^6)^2. \tag{34}$$

再令

$$\tilde{Q}_1 = Q = -\dot{q}_1, \quad \tilde{Q}_2 = 0.$$

由式(6),(9)和(33),得

$$\Lambda_1 = -a^7, \quad \Lambda_2 = 0, \quad \Lambda_3 = 0, \quad \Lambda_4 = 0, \quad \Lambda_5 = -a^7, \quad \Lambda_6 = 0, \quad \Lambda_7 = 0, \quad \Lambda_8 = 0. \tag{35}$$

式(12)给出

$$\mathbf{S}^{\mu\nu} = \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ -1 & 0 & -a^7/a^5 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & -1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \end{pmatrix}. \tag{36}$$

据式(11),系统的运动微分方程为

$$\begin{aligned}
\dot{a}^1 &= a^5, & \dot{a}^2 &= a^6, & \dot{a}^3 &= a^1 (a^6)^2 - a^7, & \dot{a}^4 &= 0, \\
\dot{a}^5 &= a^7, & \dot{a}^6 &= 0, & \dot{a}^7 &= a^5 - a^4 - a^7, & \dot{a}^8 &= (a^1)^2 a^6 - a^4.
\end{aligned} \tag{37}$$

第二步 建立确定方程并求解

取 $r=1$, 确定方程式(20)给出

$$\begin{cases} \dot{\xi}_1 - a^5 \xi_0 = \xi_5, \\ \dot{\xi}_2 - a^6 \xi_0 = \xi_6, \\ \dot{\xi}_3 + [a^1 (a^6)^2 + a^7] \xi_0 = (a^6)^2 \xi_1 + 2a^1 a^6 \xi_6 + \xi_7, \\ \dot{\xi}_4 = 0, \\ \dot{\xi}_5 - a^7 \xi_0 = \xi_7, \\ \dot{\xi}_6 = 0, \\ \dot{\xi}_7 + (a^3 - a^5 + a^7) \xi_0 = \xi_3 - \xi_5 + \xi_7, \\ \dot{\xi}_8 = \xi_4 - 2a^1 a^6 \xi_1 - (a^1)^2 \xi_6. \end{cases} \tag{38}$$

取无穷小变换生成元

$$\xi_0 = -1, \quad \xi_1 = \xi_2 = \xi_3 = \xi_4 = \xi_5 = \xi_6 = \xi_7 = \xi_8 = 0. \tag{39}$$

$$\xi_0 = 0, \quad \xi_1 = 0, \quad \xi_2 = 1, \quad \xi_3 = \xi_4 = \xi_5 = \xi_6 = \xi_7 = \xi_8 = 0. \tag{40}$$

显然,生成元式(39)和(40)满足确定方程(37),因此,它们对应系统的 Lie 对称,对于生成元(39),结构方程(21)给出

$$a^7 (a^5 + a^7) + \dot{G} = 0. \tag{41}$$

由此找不到规范函数,对于生成元(40),结构方程(21)给出

$$G = 0. \tag{42}$$

将式(40)和(42)代入式(22),得系统的守恒量

$$I = a^4 = \text{const}. \tag{43}$$

其次,研究 Lie 对称性逆问题. 假设系统有守恒量(43),求与其相应的 Lie 对称性.

方程(27),(29)给出

$$\begin{cases} \xi_1 = a^5 \xi_0, \\ \xi_2 = 1 + a^6 \xi_0, \\ \xi_3 = [a^1(a^6)^2 - a^7] \xi_0, \\ \xi_4 = 0, \\ \xi_5 = a^7 \xi_0, \\ \xi_6 = 0, \\ \xi_7 = (a^5 - a^3 - a^7) \xi_0, \\ \xi_8 = [(a^1)^2 a^6 - a^4] \xi_0, \\ (a^3 a^5 + a^4 a^6 + a^5 - a^3 - a^7) \xi_0 - H \xi_0 + G = 0. \end{cases} \quad (44)$$

引入规范函数

$$G = 0 \quad (45)$$

则

$$\xi_0 = 0, \xi_1 = 0, \xi_2 = 1, \xi_3 = \xi_4 = \xi_5 = \xi_6 = \xi_7 = \xi_8 = 0. \quad (46)$$

容易验证,生成元(46)满足确定方程(38),因此,这个对称性是系统的 Lie 对称性.

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参考文献:

- [1] WHITAKER E. T. Analytical dynamics[M]. Cambridge: Cambridge University Press, 1937.
- [2] De Leon M, Redríguez P R. Generalized classical mechanics and field theory[M]. Amsterdam. North-Holland: Elsevier Science Publishers B V, 1985.
- [3] 乔永芬,岳庆文. 广义力学系统的最小作用量原理[J]. 科学通报, 1993, 38(4): 314-318.
- [4] 李子平. 经典和量子约束系统及其对称性质[M]. 北京: 北京工业大学出版社, 1993.
- [5] 梅凤翔. 李群和李代数对约束力学系统的应用[M]. 北京: 科学出版社, 1999.
- [6] ZHANG Yi, SHANG Mei, MEI Fengxiang. Symmetries and conserved quantities for systems of generalized classical mechanics[J]. Chinese Physics, 2000, 9(6): 401-407.
- [7] 乔永芬,赵淑红. 准坐标下广义力学系统的 Lie 对称定理及其逆定理[J]. 物理学报, 2001, 50(1): 1-7.
- [8] 乔永芬,李仁杰,赵淑红. 高维增广相空间中广义力学系统的对称性和不变量[J]. 物理学报, 2001, 50(5): 811-815.
- [9] 张毅. 广义经典力学系统的第一积分与变分方程特解的联系[J]. 物理学报, 2001, 50(11): 2059-2061.
- [10] 乔永芬,赵淑红. 准坐标下广义力学系统的 Lie 对称定理及其逆定理[J]. 物理学报, 2001, 50(1): 1-7.
- [11] ZHANG Yi, SHANG Mei, MEI Fengxiang. Symmetries and conserved quantities for systems of generalized classical mechanics[J]. Chinese Physics, 2000, 9(6): 401-407.
- [12] 刘端. 非完整非保守动力学系统的守恒律[J]. 力学学报, 1989, 21(1): 75-83.
- [13] LIU Duan. Noether's theorem and its inverse of nonholonomic nonconservative dynamical systems[J]. Science in China, Series A, 1990, 34(2): 419-429.
- [14] 梅凤翔. 变质量完整力学系统的 Lie 对称与守恒量[J]. 应用数学和力学, 1999, 20(6): 592-596.
- [15] 赵跃宇,梅凤翔. 力学系统的对称性与守恒量[M]. 北京: 科学出版社, 1999.
- [16] ZHANG Hongbin, GU Shulong. Lie symmetries and conserved quantities of Birkhoff systems with unilateral constraints[J]. Chinese Physics, 2002, 11(8): 765-770.
- [17] 方建会. 二阶非完整力学系统的 Lie 对称性与守恒量[J]. 应用数学和力学, 2002, 23(9): 582-586.
- [18] 乔永芬,岳庆文,董永安. 广义力学中完整非保守系统的 Noether 守恒量[J]. 应用数学和力学, 1994, 15(9): 973-981.
- [19] 董文山. 高维增广相空间中完整非保守力学系统积分不变量的构造[J]. 山东大学学报: 理学版, 2006, 41(5): 108-111.

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