

ON THE 3D STEADY FLOW OF A SECOND GRADE FLUID PAST AN OBSTACLE

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ABSTRACT. We study steady flow of a second grade fluid past an obstacle in three space dimensions. We prove existence of solution in weighted Lebesgue spaces with anisotropic weights and thus existence of the wake region behind the obstacle. We use properties of the fundamental Oseen tensor together with results achieved in [6] and properties of solutions to steady transport equation to get up to arbitrarily small ε the same decay as the Oseen fundamental solution.

1. INTRODUCTION

The flow of a second grade fluid is governed by the following system of equations.

$$\left. \begin{aligned} \rho \frac{\partial \mathbf{v}}{\partial t} + \rho \mathbf{v} \cdot \nabla \mathbf{v} + \nabla p &= \operatorname{div} \mathbf{T} + \rho \mathbf{f} \\ \operatorname{div} \mathbf{v} &= 0 \end{aligned} \right\} \text{in } (0, T) \times \Omega, \quad (1)$$

where \mathbf{v} denotes the fluid velocity, p is the pressure, ρ is the constant density of the fluid, \mathbf{f} stands for the external force and \mathbf{T} is the Cauchy stress tensor which for the second grade fluid is given by (see e.g. [13])

$$\mathbf{T} = 2\mu \mathbf{D} + 2\alpha_1 \mathbf{A}_1 + 4\alpha_2 \mathbf{D}^2. \quad (2)$$

Here μ is a constant viscosity, $\mathbf{D} = \frac{1}{2}(\nabla \mathbf{v} + (\nabla \mathbf{v})^T)$ is the symmetric part of the velocity gradient, $\alpha_1 > 0$ and α_2 are the stress moduli and \mathbf{A}_1 is given by

$$\mathbf{A}_1 = \frac{\partial}{\partial t} \mathbf{D} + \mathbf{v} \cdot \nabla \mathbf{D} + (\nabla \mathbf{v})^T \mathbf{D} + \mathbf{D} \nabla \mathbf{v}. \quad (3)$$

The condition of thermodynamical stability yields $\alpha_1 + \alpha_2 = 0$, see [2].

Remark 1. *The question of signs and values of the stress moduli α_1 , α_2 and especially of $\alpha_1 + \alpha_2$ in this model is not clear. In [1] the authors show that the constraint $\alpha_1 + \alpha_2 = 0$ is not necessary for the mathematical problem being well set. In [5] the authors show that for $\alpha_1 < 0$ the rest state of flow of second grade fluid in exterior domain is unstable. Our results can be easily adapted also for the case $\alpha_1 + \alpha_2 \neq 0$, however we keep this thermodynamic constraint for simplicity.*

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We consider a steady flow past an obstacle, that is Ω is an exterior domain $\mathbb{R}^3 \setminus \mathcal{B}$, where \mathcal{B} is a simply connected compact set and we assume that $B_{\kappa L}(\mathbf{0}) \subset \mathcal{B} \subset B_L(\mathbf{0})$ for some $\kappa > 0$ and $L > 0$. Plugging (1) - (3) together we get

$$\left. \begin{aligned} -\mu\Delta\mathbf{v} - \alpha_1(\mathbf{v} \cdot \nabla)\Delta\mathbf{v} + \nabla p &= -\rho(\mathbf{v} \cdot \nabla)\mathbf{v} + \rho\mathbf{f} + \\ &+ \alpha_1\operatorname{div}[(\nabla\mathbf{v})^T(\nabla\mathbf{v} + (\nabla\mathbf{v})^T)] \\ \operatorname{div}\mathbf{v} &= 0 \end{aligned} \right\} \text{in } \Omega \quad (4)$$

$$\begin{aligned} \mathbf{v} &= \mathbf{0} && \text{on } \partial\Omega = \partial\mathcal{B} \\ \mathbf{v} &\rightarrow \mathbf{v}_\infty && \text{as } |\mathbf{x}| \rightarrow \infty, \end{aligned}$$

where \mathbf{v}_∞ is the prescribed constant velocity at infinity. Assuming $\mathbf{v}_\infty \neq \mathbf{0}$ we can rotate the coordinate system in such a way that $\mathbf{v}_\infty = \beta\mathbf{e}_1 = (\beta, 0, 0)$ and denoting $\mathbf{u} = \mathbf{v} - \mathbf{v}_\infty$ we get from (4)

$$\left. \begin{aligned} -\mu\Delta\mathbf{u} - \alpha_1(\mathbf{u} \cdot \nabla)\Delta\mathbf{u} - \alpha_1\beta\Delta\frac{\partial\mathbf{u}}{\partial x_1} + \rho\beta\frac{\partial\mathbf{u}}{\partial x_1} + \nabla p &= \\ -\rho(\mathbf{u} \cdot \nabla)\mathbf{u} + \rho\mathbf{f} + \alpha_1\operatorname{div}[(\nabla\mathbf{u})^T(\nabla\mathbf{u} + (\nabla\mathbf{u})^T)] \\ \operatorname{div}\mathbf{u} &= 0 \end{aligned} \right\} \text{in } \Omega \quad (5)$$

$$\begin{aligned} \mathbf{u} &= -\mathbf{v}_\infty = -\beta\mathbf{e}_1 && \text{on } \partial\Omega = \partial\mathcal{B} \\ \mathbf{u} &\rightarrow \mathbf{0} && \text{as } |\mathbf{x}| \rightarrow \infty. \end{aligned}$$

Next we rewrite the equations in dimensionless form, i.e. we introduce new velocity $\mathbf{U} = \mathbf{u}/\beta$ and new independent variable $\mathbf{X} = \mathbf{x}/L$, where L is the diameter of the obstacle. We renormalize the pressure $P = \frac{p}{\rho\beta^2}$ and the external force $\mathbf{F} = \frac{\mathbf{f}L}{\beta^2}$ and introduce the Reynolds number $\mathcal{R} = \frac{\rho\beta L}{\mu}$ and the Weisenberg number $\mathcal{W} = \frac{\alpha_1\beta}{L\mu}$. However, for the sake of transparency, we keep writing small letters instead of capital letters. After renormalization we end up with

$$\left. \begin{aligned} -\Delta\mathbf{u} - \mathcal{W}(\mathbf{u} \cdot \nabla)\Delta\mathbf{u} - \mathcal{W}\Delta\frac{\partial\mathbf{u}}{\partial x_1} + \mathcal{R}\frac{\partial\mathbf{u}}{\partial x_1} + \mathcal{R}\nabla p &= \\ -\mathcal{R}(\mathbf{u} \cdot \nabla)\mathbf{u} + \mathcal{R}\mathbf{f} + \mathcal{W}\operatorname{div}[(\nabla\mathbf{u})^T(\nabla\mathbf{u} + (\nabla\mathbf{u})^T)] \\ \operatorname{div}\mathbf{u} &= 0 \end{aligned} \right\} \text{in } \Omega \quad (6)$$

$$\begin{aligned} \mathbf{u} &= -\mathbf{e}_1 && \text{on } \partial\Omega \\ \mathbf{u} &\rightarrow \mathbf{0} && \text{as } |\mathbf{x}| \rightarrow \infty, \end{aligned}$$

where the renormalized domain $\Omega = \mathbb{R}^3 \setminus \mathcal{D}$ and $B_\kappa(\mathbf{0}) \subset \mathcal{D} \subset B_1(\mathbf{0})$. Finally we follow the decomposition procedure proposed in [8]. We introduce new pressure q as a solution to

$$q + \mathcal{W}[(\mathbf{u} + \mathbf{e}_1) \cdot \nabla]q = \mathcal{R}p \quad (7)$$

and we denote

$$-\Delta\mathbf{u} + \mathcal{R}\frac{\partial\mathbf{u}}{\partial x_1} + \nabla q =: \mathbf{z}. \quad (8)$$

Then \mathbf{z} satisfies

$$\begin{aligned} \mathbf{z} + \mathcal{W}[(\mathbf{u} + \mathbf{e}_1) \cdot \nabla]\mathbf{z} &= \mathcal{R}\mathbf{f} - \mathcal{R}(\mathbf{u} \cdot \nabla)\mathbf{u} + \mathcal{W}\operatorname{div}[(\nabla\mathbf{u})^T(\nabla\mathbf{u} + (\nabla\mathbf{u})^T)] - \\ &- \mathcal{W}(\nabla\mathbf{u})^T\nabla q + \mathcal{R}\mathcal{W}(\mathbf{u} \cdot \nabla)\frac{\partial\mathbf{u}}{\partial x_1} + \mathcal{R}\mathcal{W}\frac{\partial^2\mathbf{u}}{\partial x_1^2}. \end{aligned} \quad (9)$$

Note that we still have the conditions

$$\begin{aligned} \operatorname{div} \mathbf{u} &= 0 & \text{in } \Omega \\ \mathbf{u} &= -\mathbf{e}_1 & \text{on } \partial\Omega \\ \mathbf{u} &\rightarrow \mathbf{0} & \text{as } |\mathbf{x}| \rightarrow \infty, \end{aligned} \quad (10)$$

Our main result is the following

Theorem 1. *Let $\mathbf{f} = \operatorname{div} \mathbf{H}$, $\mathbf{H} \in W^{k,2}(\Omega)$, $k \geq 3$. Let $\Omega \in C^{k+1}$ be an exterior domain in \mathbb{R}^3 and let $\mathcal{R}_0, \mathcal{W}_0$ be sufficiently small. Then for any $\mathcal{R} \in (0, \mathcal{R}_0)$, $\mathcal{W} \in (0, \mathcal{W}_0)$ there exists a unique solution (\mathbf{u}, q) to the problem (7) - (10) for which the following estimates hold*

$$\mathcal{R}^{\frac{1}{4}} \|\mathbf{u}\|_4 + \|\nabla \mathbf{u}\|_{k,2} + \|q\|_{k,2} \leq K. \quad (11)$$

If in addition $\mathbf{f}, \mathbf{H} \in L^p(\Omega, \mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, 2\omega}(\cdot, \mathcal{R}))$ for some $p > 6$ and \mathcal{R} and \mathcal{W} are sufficiently small, the previously obtained solution (\mathbf{u}, q) has the following properties

$$\begin{aligned} \mathbf{u} &\in L^p(\Omega, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R})) \\ \nabla \mathbf{u}, \nabla^2 \mathbf{u} &\in L^p(\Omega, \mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R})) \\ q, \nabla q &\in L^p(\Omega, \mu_{\frac{1}{2}-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R})). \end{aligned} \quad (12)$$

In particular

$$\mathbf{u} \in L^\infty(\Omega, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R})). \quad (13)$$

Remark 2. *The weights $\mu_B^{A,\omega}$ are defined in (32). As the power p can be chosen arbitrarily large, we get almost the same asymptotic structure as for the fundamental solution \mathcal{O} of the Oseen system.*

Remark 3. *Note that for $\mathbf{f} \in L^\infty(\Omega, \eta_1^{\frac{3}{2}}(\mathbf{x}))$ it holds*

$$\|\mathbf{f}\|_{L^p(\Omega, \mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, 2\omega}(\cdot, \mathcal{R}))} \leq C \mathcal{R}^{-2\omega-\frac{3}{p}} \|\mathbf{f}\|_{L^\infty(\Omega, \eta_1^{\frac{3}{2}}(\cdot))}. \quad (14)$$

Throughout this paper we shall assume \mathcal{W} and \mathcal{R} small. We introduce operator

$$\mathcal{M} : (\mathbf{w}, s) \mapsto \mathbf{z} \mapsto (\mathbf{u}, q), \quad (15)$$

where for given (\mathbf{w}, s) , \mathbf{z} is the solution to the transport equation

$$\begin{aligned} \mathbf{z} + \mathcal{W}[(\mathbf{w} + \mathbf{e}_1) \cdot \nabla] \mathbf{z} &= \mathcal{R} \mathbf{f} - \mathcal{R}(\mathbf{w} \cdot \nabla) \mathbf{w} + \mathcal{W} \operatorname{div} [(\nabla \mathbf{w})^T (\nabla \mathbf{w} + (\nabla \mathbf{w})^T)] - \\ &- \mathcal{W}(\nabla \mathbf{w})^T \nabla s + \mathcal{R} \mathcal{W}(\mathbf{w} \cdot \nabla) \frac{\partial \mathbf{w}}{\partial x_1} + \mathcal{R} \mathcal{W} \frac{\partial^2 \mathbf{w}}{\partial x_1^2} =: \mathbf{B}(\mathbf{f}, \mathbf{w}, s) \quad \text{in } \Omega \end{aligned} \quad (16)$$

and (\mathbf{u}, q) is the solution to the Oseen problem

$$\begin{aligned} -\Delta \mathbf{u} + \mathcal{R} \frac{\partial \mathbf{u}}{\partial x_1} + \nabla q &= \mathbf{z} & \text{in } \Omega \\ \operatorname{div} \mathbf{u} &= 0 & \text{in } \Omega \\ \mathbf{u} &= -\mathbf{e}_1 & \text{on } \partial\Omega \\ \mathbf{u} &\rightarrow \mathbf{0} & \text{as } |\mathbf{x}| \rightarrow \infty. \end{aligned} \quad (17)$$

We have decomposed the original problem into the Oseen problem (17) and the steady transport equation (16). Due to the presence of the Oseen problem we expect the structure of solutions to correspond to the structure of the Oseen fundamental solution, especially the existence of the wake region behind the obstacle (compare with [3], [12] for incompressible Navier-Stokes equations and [10] for viscoelastic fluid). Denoting $s(\mathbf{x}) = |\mathbf{x}| - x_1$ one might expect the solution \mathbf{u} to satisfy

$$|\mathbf{u}(\mathbf{x})| \leq C |\mathbf{x}|^{-1} (1 + s(\mathbf{x}))^{-1} \quad (18)$$

for $|\mathbf{x}|$ sufficiently large. However we are only able to prove

$$|\mathbf{u}(\mathbf{x})| \leq C |\mathbf{x}|^{-1+\varepsilon} (1 + s(\mathbf{x}))^{-1+\varepsilon} \quad (19)$$

for arbitrarily small ε . This is due to the presence of the linear term $\mathcal{RW} \frac{\partial^2 \mathbf{u}}{\partial x_1^2}$ on the right-hand side of the transport equation (16). This implies that the solution to the transport equation has the same decay as $\nabla^2 \mathbf{w}$ (quadratic terms decay faster). Moreover, for the Oseen system previously available estimates for the second gradient (which were using techniques of fundamental solution) lose logarithmic factor in the weight in the L^∞ norm, L^p estimates lose ε in the weight, see [11]. Thus fixed point theorem argument would not work. Fortunately due to recent results of Koch [6] we have at least L^p estimates without mentioned ε loss in the weight and therefore fixed point argument works.

It is worth mentioning that in [10], where the model of viscoelastic fluid is considered, authors are able to overcome these problems by introducing modified Oseen problem with the problematic term $\frac{\partial^2 \mathbf{u}}{\partial x_1^2}$ being included in the Oseen operator. Then all terms on the right hand side are quadratic and thus with better decay. In our problem this cannot be repeated since this linear term appears in the transport equation.

The drawback of using L^p estimates is that in order to get L^∞ estimate we have to use embedding theorems and thus we are able to prove

$$|\mathbf{u}(\mathbf{x})| \leq C |\mathbf{x}|^{-1+\varepsilon} (1 + s(\mathbf{x}))^{-1+\varepsilon}$$

for arbitrarily small ε . Details will be specified in the proper part of the proof.

2. PRELIMINARIES

Throughout this paper we will use standard notation for the Lebesgue spaces $L^p(\Omega)$ with the norm $\|\cdot\|_p$, the Sobolev spaces $W^{k,p}(\Omega)$ with the norm $\|\cdot\|_{k,p}$ and the homogenous Sobolev spaces $D^{k,p}(\Omega)$ with the norm $|\cdot|_{k,p}$. Let $g \in L^1_{loc}(\Omega)$ be a nonnegative weight. Then $L^p(\Omega, g)$ denotes the weighted L^p space with the norm

$$\|u\|_{p,(g)} = \|ug\|_p$$

for any $p \in [1, \infty]$. Similarly, $W^{k,p}(\Omega, g)$ denotes the weighted Sobolev space with the norm

$$\|u\|_{k,p,(g)} = \|ug\|_{k,p}.$$

Note that if there is no confusion we sometimes omit writing the domain and instead of $L^p(\Omega, g)$ we write simply $L^p(g)$.

As we decomposed the original problem into an Oseen problem and a steady transport equation, we shall mention several classical results about these problems in three dimensional exterior domains. Let us start with the Oseen problem (17).

2.1. Oseen problem. Existence and properties. We denote by $(\mathcal{O}, \mathbf{e})$ the fundamental solution to the Oseen problem. It can be shown (see for example [11]) that

$$\mathbf{e}(\mathbf{x}) = \nabla \mathcal{E}(\mathbf{x}), \quad (20)$$

where $\mathcal{E}(\mathbf{x})$ is the fundamental solution to the Laplace equation. The tensor $\mathcal{O}(\mathbf{x}, \mathcal{R})$ (here \mathcal{R} denotes the constant standing in front of $\frac{\partial \mathbf{u}}{\partial x_1}$ in the equation) satisfies the following property

$$\mathcal{O}(\mathbf{x}, \mathcal{R}) = \mathcal{R} \mathcal{O}(\mathcal{R}\mathbf{x}, 1) \quad (21)$$

and therefore it is sufficient to study the tensor $\mathcal{O}(\mathbf{x}, 1)$. For $|\mathbf{x}| \rightarrow \infty$ we have

$$\begin{aligned} \mathcal{O}(\mathbf{x}, 1) &\sim |\mathbf{x}|^{-1} (1 + s(\mathbf{x}))^{-1} \\ D^\alpha \mathcal{O}(\mathbf{x}, 1) &\sim |\mathbf{x}|^{-1 - \frac{|\alpha|}{2}} (1 + s(\mathbf{x}))^{-1 - \frac{|\alpha|}{2}} \\ D^\alpha \mathcal{O}(\mathbf{x}, 1) &\sim |\mathbf{x}|^{-1 - \alpha_1 - \frac{|\alpha| - \alpha_1}{2}} (1 + s(\mathbf{x}))^{-1 - \frac{|\alpha| - \alpha_1}{2}}, \end{aligned} \quad (22)$$

i.e. the derivatives with respect to the first variable decay faster.

Next, we present some results for the general Oseen problem

$$\begin{aligned} -\Delta \mathbf{u} + \mathcal{R} \frac{\partial \mathbf{u}}{\partial x_1} + \nabla P &= \mathbf{f} = \operatorname{div} \mathcal{G} \quad \text{in } \Omega \\ \operatorname{div} \mathbf{u} &= 0 \quad \text{in } \Omega \\ \mathbf{u} &= \mathbf{u}_* \quad \text{on } \partial\Omega \\ \mathbf{u} &\rightarrow \mathbf{0} \quad \text{as } |\mathbf{x}| \rightarrow \infty, \end{aligned} \quad (23)$$

where Ω is an exterior domain.

The proof of the following classical theorem can be found in [4] or in [11].

Theorem 2. *Let $\Omega \subset \mathbb{R}^3$ be an exterior domain of class C^{k+2} . Let $\mathbf{f} \in D_0^{-1,q}(\Omega) \cap W^{k,2}(\Omega)$, $\mathbf{u}_* \in W^{k+\frac{3}{2},2}(\partial\Omega)$, $q \in (\frac{3}{2}, 3)$, $k \geq 0$. Then there exists exactly one q -weak solution (i.e. weak solution such that $\mathbf{u} \in W^{1,q}(\Omega)$) to (23). Moreover*

$$\mathbf{u} \in L^{\frac{4q}{4-q}}(\Omega) \quad \text{and} \quad \nabla \mathbf{u}, P \in L^q(\Omega) \cap W^{k+1,2}(\Omega) \quad (24)$$

and

$$\begin{aligned} a_2 \|\mathbf{u}\|_{\frac{4q}{4-q}} + \|\mathbf{u}\|_{1,q} + \|\nabla \mathbf{u}\|_{k+1,2} + \|P\|_q + \|P\|_{k+1,2} &\leq \\ &\leq C(\|\mathbf{f}\|_{-1,q} + \|\mathbf{f}\|_{k,2} + \|\mathbf{u}_*\|_{k+\frac{3}{2},2,\partial\Omega}), \end{aligned} \quad (25)$$

where for $\mathcal{R} \in (0, \mathcal{R}_0]$ the constant $C = C(k, q, \Omega, \mathcal{R}_0)$ and $a_2 = \min\{1, \mathcal{R}^{\frac{1}{4}}\}$.

We need the following integral representation of solutions to (23) to obtain weighted estimates.

Let us denote

$$\begin{aligned} T_{ij}(\mathbf{u}, P) &= \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - P \delta_{ij} \\ \mathcal{T}_{ij}(\mathbf{e}) &= \frac{\partial e_i}{\partial x_j} + \frac{\partial e_j}{\partial x_i} + \mathcal{R} e_1 \delta_{ij} \end{aligned} \quad (26)$$

Theorem 3. Let $\Omega \in C^2$ be an exterior domain, $\mathbf{g} \in C_0^\infty(\overline{\Omega})$ and (\mathbf{u}, P) be the unique solution to (23). Let \mathbf{T} be defined in (26) and (\mathbf{O}, \mathbf{e}) be the fundamental solution to the Oseen problem. Then

$$\begin{aligned} u_j(\mathbf{x}) = & \int_{\Omega} \frac{\partial}{\partial x_k} \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) \mathcal{G}_{ik}(\mathbf{y}) d\mathbf{y} + \\ & + \int_{\partial\Omega} [-\mathcal{R} \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) u_i(\mathbf{y}) \delta_{1k} + u_i(\mathbf{y}) T_{ik}(\mathbf{O}_{\cdot j}, e_j)(\mathbf{x} - \mathbf{y}, \mathcal{R}) + \\ & + \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) T_{ik}(\mathbf{u}, P)(\mathbf{y}) + \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) \mathcal{G}_{ik}(\mathbf{y})] n_k(\mathbf{y}) dS \end{aligned} \quad (27)$$

$$\begin{aligned} D^\alpha u_j(\mathbf{x}) = & - \int_{\Omega} D^\alpha \frac{\partial}{\partial x_k} \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) \mathcal{G}_{ik}(\mathbf{y}) d\mathbf{y} + \\ & + \int_{\partial\Omega} [-\mathcal{R} D^\alpha \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) u_i(\mathbf{y}) \delta_{1k} + u_i(\mathbf{y}) D^\alpha T_{ik}(\mathbf{O}_{\cdot j}, e_j)(\mathbf{x} - \mathbf{y}, \mathcal{R}) + \\ & + D^\alpha \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) T_{ik}(\mathbf{u}, P)(\mathbf{y}) + D^\alpha \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) \mathcal{G}_{ik}(\mathbf{y})] n_k(\mathbf{y}) dS \end{aligned} \quad (28)$$

for $|\alpha| = 1$ and

$$\begin{aligned} D^\alpha u_j(\mathbf{x}) = & v.p. \int_{\Omega} D^\alpha \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) \frac{\partial}{\partial y_k} \mathcal{G}_{ik}(\mathbf{y}) d\mathbf{y} + c_{ij\alpha_1\alpha_2} \frac{\partial \mathcal{G}_{ik}}{\partial x_k}(\mathbf{x}) + \\ & + \int_{\partial\Omega} [-\mathcal{R} D^\alpha \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) u_i(\mathbf{y}) \delta_{1k} + u_i(\mathbf{y}) D^\alpha T_{ik}(\mathbf{O}_{\cdot j}, e_j)(\mathbf{x} - \mathbf{y}, \mathcal{R}) + \\ & + D^\alpha \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) T_{ik}(\mathbf{u}, P)(\mathbf{y})] n_k(\mathbf{y}) dS \end{aligned} \quad (29)$$

for $|\alpha| = 2$.

Remark 4. The integral representation formulas hold for much larger classes of functions. For example it holds for a.a. $\mathbf{x} \in \Omega$ if $\mathbf{u} \in W_{loc}^{2,q}(\overline{\Omega})$ and $P \in W_{loc}^{1,q}(\overline{\Omega})$ for some $q \in (1, \infty)$ and

- (27) if $\mathbf{g} \in L^q(\Omega)$ and $\operatorname{div} \mathbf{g} \in L_{loc}^r(\overline{\Omega})$ for $q \in (1, 4)$, $r \in (1, \infty)$,
- (28) if $\mathbf{g} \in L^q(\Omega)$ and $\operatorname{div} \mathbf{g} \in L_{loc}^r(\overline{\Omega})$ for $q, r \in (1, \infty)$,
- (29) if $\operatorname{div} \mathbf{g} \in L_{loc}^r(\overline{\Omega})$ for $r \in (1, \infty)$.

For the pressure we have also integral representation formulas.

Theorem 4. Let $\Omega \in C^2$ be an exterior domain, $\mathbf{g} \in C_0^\infty(\overline{\Omega})$ and (\mathbf{u}, P) be the unique solution to (23). Let T_{ij} and \mathcal{T}_{ij} be defined in (26). Then

$$\begin{aligned} P(\mathbf{x}) = & v.p. \int_{\Omega} \frac{\partial}{\partial x_k} e_i(\mathbf{x} - \mathbf{y}) \mathcal{G}_{ik}(\mathbf{y}) d\mathbf{y} + c_{ik} \mathcal{G}_{ik}(\mathbf{x}) + \\ & + \int_{\partial\Omega} [-\mathcal{R} e_i(\mathbf{x} - \mathbf{y}) u_i(\mathbf{y}) \delta_{1l} + u_i(\mathbf{y}) \mathcal{T}_{il}(\mathbf{e})(\mathbf{x} - \mathbf{y}) + \\ & + e_i(\mathbf{x} - \mathbf{y}) T_{il}(\mathbf{u}, P)(\mathbf{y}) + e_i(\mathbf{x} - \mathbf{y}) \mathcal{G}_{il}(\mathbf{y})] n_l(\mathbf{y}) dS \end{aligned} \quad (30)$$

$$\begin{aligned} D^\alpha P(\mathbf{x}) = & v.p. \int_{\Omega} D^\alpha e_i(\mathbf{x} - \mathbf{y}) \frac{\partial}{\partial y_k} \mathcal{G}_{ik}(\mathbf{y}) d\mathbf{y} + c_{ik} \frac{\partial}{\partial x_k} \mathcal{G}_{ik}(\mathbf{x}) + \\ & + \int_{\partial\Omega} [-\mathcal{R} D^\alpha e_i(\mathbf{x} - \mathbf{y}) u_i(\mathbf{y}) \delta_{1l} + u_i(\mathbf{y}) D^\alpha \mathcal{T}_{il}(\mathbf{e})(\mathbf{x} - \mathbf{y}) + \\ & + D^\alpha e_i(\mathbf{x} - \mathbf{y}) T_{il}(\mathbf{u}, P)(\mathbf{y})] n_l(\mathbf{y}) dS \end{aligned} \quad (31)$$

for $|\alpha| = 1$.

Remark 5. *The integral representation formulas for pressure hold also for much larger classes of functions. For example it holds for a.a. $\mathbf{x} \in \Omega$ if $\mathbf{u} \in W_{loc}^{2,q}(\overline{\Omega})$ and $P \in W_{loc}^{1,q}(\overline{\Omega})$ for some $q \in (1, \infty)$ and*

- (30) if $\mathcal{G} \in L^q(\Omega)$ and $\operatorname{div} \mathcal{G} \in L_{loc}^r(\overline{\Omega})$ for $q, r \in (1, \infty)$,
- (31) if $\operatorname{div} \mathcal{G} \in L_{loc}^r(\overline{\Omega})$ for $r \in (1, \infty)$.

The proof of these representation formulas can be found in [4] or in [11].

2.2. Results for weighted spaces. We introduce the following weights which will be useful in studying the asymptotic structure of solutions

$$\begin{aligned}
 \eta_B^A(\mathbf{x}) &= (1 + |\mathbf{x}|)^A (1 + s(\mathbf{x}))^B \\
 \nu_B^A(\mathbf{x}) &= |\mathbf{x}|^A (1 + s(\mathbf{x}))^B \\
 \mu_B^{A,\omega}(\mathbf{x}) &= \eta_B^{A-\omega}(\mathbf{x}) \nu_0^\omega(\mathbf{x}) \\
 \eta_B^A(\mathbf{x}, \mathcal{R}) &= (1 + |\mathcal{R}\mathbf{x}|)^A (1 + s(\mathcal{R}\mathbf{x}))^B \\
 \nu_B^A(\mathbf{x}, \mathcal{R}) &= |\mathbf{x}|^A (1 + s(\mathcal{R}\mathbf{x}))^B \\
 \mu_B^{A,\omega}(\mathbf{x}, \mathcal{R}) &= \eta_B^{A-\omega}(\mathbf{x}, \mathcal{R}) \nu_0^\omega(\mathbf{x}, \mathcal{R})
 \end{aligned} \tag{32}$$

We recall that weight g belongs to the class A_p if there exists a constant C such that

$$\sup_Q \left[\left(\frac{1}{|Q|} \int_Q g^p(\mathbf{x}) d\mathbf{x} \right) \left(\frac{1}{|Q|} \int_Q g^{-\frac{p}{p-1}}(\mathbf{x}) d\mathbf{x} \right)^{p-1} \right] \leq C < \infty, \tag{33}$$

where the supremum is taken over all cubes $Q \subset \mathbb{R}^3$.

A basic property of the weight η_{-b}^{-a} is the following condition

$$\int_{\mathbb{R}^3} \eta_{-b}^{-a}(\mathbf{x}) d\mathbf{x} < \infty \iff a + \min\{1, b\} > 3. \tag{34}$$

For a proof of this and further properties we refer the reader to [11].

One of the main tools we use is the following theorem due to Koch [6]. We should mention that this estimate is an essential improvement of what has been known about weighted estimates of the solutions to the Oseen system. Without this result it was impossible to obtain weighted estimates for the model of the second grade fluid or for Maxwell and Oldroyd-type fluids, see [11].

Theorem 5. (Koch) *Let T be an integral operator with the kernel $\frac{\partial^2}{\partial x_i \partial x_j} \mathcal{O}$ on \mathbb{R}^3 . Then the following estimate hold:*

$$\|Tf\|_{p,(g),\mathbb{R}^3} \leq C \|f\|_{p,(g),\mathbb{R}^3} \tag{35}$$

for $p \in (1, \infty)$ and $g = \eta_B^A(\mathbf{x})$ for A, B satisfying

$$\begin{aligned}
 A, B &\in \left(-\frac{1}{p}, \frac{2(p-1)}{p} \right) \\
 A + B &> -\frac{1}{p} \\
 2A - B, 2B - A &< \frac{2(p-1) + 1}{p}
 \end{aligned} \tag{36}$$

The proof of this theorem can be found in [6]. Note that in [6] the theorem is formulated only for the case $p = 2$, nevertheless the proof is given for general $p \in (1, \infty)$. The same estimate holds also for the case of an exterior domain.

Corollary 1. *For the same operator T as in Theorem 5 we have*

$$\|Tf\|_{p,(g_1),\Omega} \leq C\mathcal{R}^\omega \|f\|_{p,(g_2),\Omega}, \quad (37)$$

where $g_1 = \mu_B^{A,\omega}(\mathbf{x}, \mathcal{R})$, $g_2 = \mu_B^{A,2\omega}(\mathbf{x}, \mathcal{R})$, A, B satisfying (36) and $\omega \in [0, \frac{A}{2})$.

Proof. First we observe that for all $\mathbf{x} \in \Omega$ it holds

$$\mu_B^{A,\omega}(\mathbf{x}) \leq \eta_B^A(\mathbf{x}) \leq (1 + \frac{1}{\kappa})^\omega \mu_B^{A,\omega}(\mathbf{x}), \quad (38)$$

where κ was introduced earlier by condition $B_\kappa(\mathbf{0}) \subset \mathcal{D}$ and $\Omega = \mathbb{R}^3 \setminus \mathcal{D}$. In other words the weights $\eta_B^A(\mathbf{x})$ and $\mu_B^{A,\omega}(\mathbf{x})$ are equivalent in Ω . The presence of the term \mathcal{R}^ω is an easy consequence of a rescaling argument, because $\nabla^2 \mathcal{O}(\mathbf{x}, \mathcal{R}) = \mathcal{R}^3 \nabla^2 \mathcal{O}(\mathcal{R}\mathbf{x}, 1)$. \square

The proofs of the following theorems can be found for example in [7] or in [11].

Theorem 6. *Let T be an integral operator with the kernel $|\nabla \mathcal{O}|$, $T : f \rightarrow |\nabla \mathcal{O}| * f$ and $p \in (1, \infty)$. Then T is a well defined continuous operator:*

$$L^p(\mathbb{R}^3, \eta_B^{A+\frac{1}{2}}(\cdot, \mathcal{R})) \mapsto L^p(\mathbb{R}^3, \eta_B^A(\cdot, \mathcal{R})) \quad (39)$$

for $B \in (0, \frac{3}{2} - \frac{3}{2p})$, $A + B > -\frac{1}{p}$, $A < \frac{3}{2} - \frac{2}{p}$, $A - B < \frac{1}{2} - \frac{1}{p}$. Moreover we have for A, B specified above

$$\| |\nabla \mathcal{O}(\cdot, \mathcal{R})| * f \|_{p, \eta_B^A(\cdot, \mathcal{R}), \mathbb{R}^3} \leq C\mathcal{R}^{-1} \|f\|_{p, \eta_B^{A+\frac{1}{2}}(\cdot, \mathcal{R}), \mathbb{R}^3} \quad (40)$$

Corollary 2. *For the same operator T as in Theorem 6 one has*

$$\|Tf\|_{p,(g_1),\Omega} \leq C\mathcal{R}^{-1+\omega} \|f\|_{p,(g_2),\Omega}, \quad (41)$$

where $g_1 = \mu_B^{A,\omega}(\mathbf{x}, \mathcal{R})$, $g_2 = \mu_B^{A,2\omega}(\mathbf{x}, \mathcal{R})$, A, B satisfy the assertions of Theorem 6 and $\omega \in [0, \frac{A}{2})$.

Theorem 7. *Let*

$$Tf(\mathbf{x}) = \frac{\partial}{\partial x_i} \int_{\mathbb{R}^3} e_j(\mathbf{x} - \mathbf{y}) f(\mathbf{y}) d\mathbf{y}, \quad i, j = 1, 2, 3, \quad (42)$$

$f \in C_0^\infty(\mathbb{R}^3)$, $p \in (1, \infty)$ and let g stands for one of weights η_B^A , ν_B^A , $\mu_B^{A,\omega}$. Let A, B be such that g is an A_p weight in \mathbb{R}^3 . Then T maps $C_0^\infty(\mathbb{R}^3)$ into $L^p(\mathbb{R}^3, g)$ and

$$\|Tf\|_{p,(g),\mathbb{R}^3} \leq C \|f\|_{p,(g),\mathbb{R}^3}. \quad (43)$$

Moreover T can be continuously extended onto $L^p(\mathbb{R}^3, g)$.

Corollary 3. *The same holds also for the case of an exterior domain Ω .*

Theorem 8.

- Let $B \in (-\frac{1}{p}, \frac{p-1}{p})$ and $A + B \in (-\frac{3}{p}, \frac{3(p-1)}{p})$. Then the weight η_B^A is an A_p weight in \mathbb{R}^3 for $p \in (1, \infty)$.
- Let moreover $A \in (-\frac{3}{p}, \frac{3(p-1)}{p})$ and $\omega \in [0, A]$. Then the weights ν_B^A and $\mu_B^{A,\omega}$ are A_p weights in \mathbb{R}^3 for $p \in (1, \infty)$.

2.3. Transport equation. Next we consider the steady transport equation

$$z + \mathbf{w} \cdot \nabla z = f \quad \text{in } \Omega. \quad (44)$$

This equation is scalar, nevertheless all theorems below hold also for the vector case. The following theorems are proved in [9] even for more complicated cases.

Theorem 9. (i) Let $\Omega \in C^{0,1}$ be an exterior domain, $\mathbf{w} \in C^{k-1}(\Omega)$, $\mathbf{w} \cdot \mathbf{n} = 0$ on $\partial\Omega$, $\nabla^k \mathbf{w} \in L^3(\Omega)$, $f \in W^{k,q}$ for $q \in (1, 3)$, $kq > 3$. Then there exists $\alpha > 0$ such that if

$$\|\nabla \mathbf{w}\|_{C^{k-2}} + \|\nabla^k \mathbf{w}\|_3 < \alpha, \quad (45)$$

then there exists unique solution $z \in W^{k,q}(\Omega)$ to (44) satisfying the estimate

$$\|z\|_{k,q} \leq C(\alpha) \|f\|_{k,q}. \quad (46)$$

(ii) Let $\Omega \in C^{0,1}$ be an exterior domain, $\mathbf{w} \in C^k(\Omega)$, $\mathbf{w} \cdot \mathbf{n} = 0$ on $\partial\Omega$, $f \in W^{k,q}$ for $kq > 3$. Then there exists $\alpha > 0$ such that if

$$\|\nabla \mathbf{w}\|_{C^{k-1}} < \alpha, \quad (47)$$

then there exists unique solution $z \in W^{k,q}(\Omega)$ to (44) satisfying the estimate

$$\|z\|_{k,q} \leq C(\alpha) \|f\|_{k,q}. \quad (48)$$

Theorem 10. Let Ω, k, q, \mathbf{w} and f satisfy the assumptions of Theorem 9 (ii). Moreover let $g \in C^k(\Omega)$ be a positive weight such that $W^{k,q}(\Omega, g) \subset W^{k,q}(\Omega)$ and let

$$\|\mathbf{w} \cdot \nabla \ln g\|_{C^{k-1}} + |\mathbf{w} \cdot \nabla \ln g|_{k,q} \quad (49)$$

be sufficiently small. Let $f \in W^{k,q}(\Omega, g)$. Then z , the solution to (44), belongs to $W^{k,q}(\Omega, g)$ and

$$\|z\|_{k,q,(g)} \leq C \|f\|_{k,q,(g)}. \quad (50)$$

3. THE PROOF OF THEOREM 1

3.1. Existence of solution. Here we briefly sketch the method of constructing the solution to the system (7)-(10). It is based on the following version of the Banach fixed point theorem.

Theorem 11. Let X, Y be Banach spaces such that X is reflexive and $X \hookrightarrow Y$. Let H be nonempty, closed, convex and bounded subset of X and let $\mathcal{M} : H \mapsto H$ be a mapping such that

$$\|\mathcal{M}(u) - \mathcal{M}(v)\|_Y \leq \delta \|u - v\|_Y \quad \forall u, v \in H, \quad (51)$$

$\delta \in [0, 1)$. Then \mathcal{M} has a unique fixed point in H .

The proof of existence of solutions in Sobolev spaces is based on the method described for example in [11] and [10]. The solution is obtained as a limit of successive approximations

$$(\mathbf{u}_{n+1}, q_{n+1}) = \mathcal{M}(\mathbf{u}_n, q_n), \quad n \geq 0, \quad (52)$$

where the mapping \mathcal{M} was introduced in (15).

Theorem 12. Let $\Omega \in C^{k+1}$ be an exterior domain in \mathbb{R}^3 and let $\mathbf{f} = \operatorname{div} \mathbf{H}$, $\mathbf{H} \in W^{k,2}(\Omega)$, $k \geq 3$. Let $\mathcal{R}_0, \mathcal{W}_0$ be sufficiently small. Then for any $\mathcal{R} \in (0, \mathcal{R}_0)$, $\mathcal{W} \in (0, \mathcal{W}_0)$ there exists (\mathbf{u}, q) a solution to the system (7)-(10) such that $\mathbf{u} \in L^4(\Omega)$ and $\nabla \mathbf{u}, q \in W^{k,2}(\Omega)$.

Proof. We use Theorem 11 for the following choice of spaces: $X = V_k$, $Y = V_{k-1}$, where

$$V_k = \{(\mathbf{u}, q) : \mathbf{u} \in L^4(\Omega), \nabla \mathbf{u}, q \in W^{k,2}(\Omega)\} \quad (53)$$

with the norm

$$\|(\mathbf{u}, q)\|_{V_k} = \mathcal{R}^{\frac{1}{4}} \|\mathbf{u}\|_4 + \|\nabla \mathbf{u}\|_{k,2} + \|q\|_{k,2} \quad (54)$$

We use Theorems 2 and 9 for $q = 2$ and the observation that the right-hand side of the transport equation (9) (we denote it by $\mathbf{B}(\mathbf{f}, \mathbf{w}, s)$) can be written in the divergence form as

$$\begin{aligned} \mathbf{B}(\mathbf{f}, \mathbf{w}, s) = \operatorname{div} \left[\mathcal{R}\mathbf{H} - \mathcal{R}\mathbf{w} \otimes \mathbf{w} + \mathcal{W}(\nabla \mathbf{w})^T (\nabla \mathbf{w} + (\nabla \mathbf{w})^T) - \right. \\ \left. - \mathcal{W}(\nabla \mathbf{w})^T s + \mathcal{R}\mathcal{W} \frac{\partial \mathbf{w}}{\partial x_1} \otimes \mathbf{w} + \mathcal{R}\mathcal{W} \frac{\partial \mathbf{w}}{\partial x_1} \otimes \mathbf{e}_1 \right] =: \operatorname{div} \mathbf{C}(\mathbf{H}, \mathbf{w}, s) \end{aligned} \quad (55)$$

assuming $\mathbf{f} = \operatorname{div} \mathbf{H}$. As a part of the proof we obtain the following estimates

$$\begin{aligned} \mathcal{R}^{\frac{1}{4}} \|\mathbf{u}_n\|_4 + \|\nabla \mathbf{u}_n\|_{k,2} + \|q_n\|_{k,2} &\leq K \\ \|\mathbf{C}(\mathbf{H}, \mathbf{u}_n, q_n)\|_{k,2} &\leq K \end{aligned} \quad (56)$$

for all $n \geq 0$, for some constant $K > 0$ depending only on $\mathcal{R}_0, \mathcal{W}_0$ and \mathbf{H} . \square

3.2. Weighted estimates. In this section we study weighted estimates which are crucial to obtain asymptotic behaviour of the solution. As we have used Theorem 11 to prove existence of solutions in Sobolev spaces, it is now sufficient to prove that \mathcal{M} maps sufficiently large balls in proper weighted spaces into themselves. Then choosing (\mathbf{u}_1, q_1) from this ball the solution, as the limit of the sequence (\mathbf{u}_n, q_n) , belongs to the same ball. We estimate the sequence in the following space

$$\begin{aligned} V = \left\{ (\mathbf{u}, q) : \mathbf{u} \in L^p(\Omega, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R})), \nabla \mathbf{u}, \nabla^2 \mathbf{u} \in L^p(\Omega, \mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R})), \right. \\ \left. q, \nabla q \in L^p(\Omega, \mu_{\frac{1}{2}-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R})) \right\} \end{aligned} \quad (57)$$

with the norm

$$\begin{aligned} \|(\mathbf{u}, q)\|_V = \|\mathbf{u}\|_{p, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}), \Omega} + \\ + \|\nabla \mathbf{u}, \nabla^2 \mathbf{u}\|_{p, \mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}), \Omega} + \|q, \nabla q\|_{p, \mu_{\frac{1}{2}-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}), \Omega}, \end{aligned} \quad (58)$$

where p is sufficiently large and $\omega < \frac{1}{2} - \frac{3}{2p}$.

Remark 6. *There exist a constant C depending only on Ω, A, B, ω such that for $p > 3$, $A, B \geq 0$, $\omega \in [0, A]$ and $\mathcal{R} \leq 1$*

$$\|g\|_{L^\infty(\Omega, \mu_B^{A, \omega}(\cdot, \mathcal{R}))} \leq C(\|g\|_{L^p(\Omega, \mu_B^{A, \omega}(\cdot, \mathcal{R}))} + \|\nabla g\|_{L^p(\Omega, \mu_B^{A, \omega}(\cdot, \mathcal{R}))}). \quad (59)$$

This is an easy consequence of the Sobolev imbedding theorem and the fact that there is a constant C independent of \mathcal{R} such that

$$\|g\|_{L^p(\Omega, \nabla \mu_B^{A, \omega}(\cdot, \mathcal{R}))} \leq C \|g\|_{L^p(\Omega, \mu_B^{A, \omega}(\cdot, \mathcal{R}))}. \quad (60)$$

Let us mention that using this Remark we get $\mathbf{u} \in L^\infty(\Omega, \mu_{1-\frac{3}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))$ with p arbitrarily large and therefore almost the same asymptotic structure as \mathcal{O} .

Let us assume

$$\|(\mathbf{w}, s)\|_V \leq C_0, \quad (61)$$

where C_0 is sufficiently large constant which will be determined later. It is important to mention that this constant is determined by the estimates (56) and is independent of \mathcal{R} and \mathcal{W} .

Our aim is to prove that also

$$\|\mathcal{M}(\mathbf{w}, s)\|_V = \|(\mathbf{u}, q)\| \leq C_0. \quad (62)$$

We recall that we also assume that (\mathbf{w}, s) satisfy (56).

Throughout the rest of this paper we will use the following notation to simplify things

$$X^\omega = L^p(\Omega, \mu_{1-\frac{3}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R})) \quad (63)$$

We will estimate both $\mathbf{B}(\mathbf{f}, \mathbf{w}, s)$ and $\mathbf{C}(\mathbf{H}, \mathbf{w}, s)$ in $X^{2\omega}$. Due to the presence of the Reynolds and Weissenberg numbers in front of each term on the right hand side it is sufficient to show the presence of all terms in $X^{2\omega}$, smallness of these terms is achieved by assuming \mathcal{R}, \mathcal{W} sufficiently small. We will proceed term by term and denote the terms on the right hand side of (16) by B_1, \dots, B_6 and the corresponding terms of \mathbf{C} by C_1, \dots, C_6 . First we use assumption

$$\mathbf{f}, \mathbf{H} \in X^{2\omega}, \quad (64)$$

which allows us to estimate B_1, C_1 in $X^{2\omega}$. We estimate B_2 in the following way

$$\|\mathbf{w} \nabla \mathbf{w}\|_{X^{2\omega}}^p \leq \|\mathbf{w}\|_{L^\infty(\Omega, \mu_{1-\frac{3}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \|\nabla \mathbf{w}\|_{X^\omega}^p \left\| \eta_{2-p}^{3-p}(\cdot, \mathcal{R}) \right\|_{L^\infty(\Omega)}. \quad (65)$$

The last term is finite for $p > 3$ and therefore using Remark 6

$$\|B_2\|_{X^{2\omega}} \leq C \mathcal{R} C_0^2. \quad (66)$$

We proceed in the similar way also in the divergence form case. Here

$$\|\mathbf{w} \otimes \mathbf{w}\|_{X^{2\omega}}^p \leq \|\mathbf{w}\|_{L^\infty(\Omega, \mu_{1-\frac{3}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \|\mathbf{w}\|_{L^p(\Omega, \mu_{1-\frac{3}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \left\| \eta_{2-p}^{3-\frac{3p}{2}}(\cdot, \mathcal{R}) \right\|_{L^\infty(\Omega)} \quad (67)$$

and therefore for $p > 6$ we get

$$\|C_2\|_{X^{2\omega}} \leq C \mathcal{R} C_0^2. \quad (68)$$

Similar procedure works also for terms B_3, B_4, B_5 and C_3, C_4, C_5 , we only show the estimates for B_3 and B_4 which are most restrictive.

$$\|\nabla^2 \mathbf{w} \nabla \mathbf{w}\|_{X^{2\omega}}^p \leq \|\nabla \mathbf{w}\|_{L^\infty(\Omega, \mu_{1-\frac{3}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \|\nabla^2 \mathbf{w}\|_{L^p(\Omega, \mu_{1-\frac{3}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \left\| \eta_{2-p}^{3-\frac{3p}{2}}(\cdot, \mathcal{R}) \right\|_{L^\infty(\Omega)} \quad (69)$$

$$\|\nabla \mathbf{w} \nabla s\|_{X^{2\omega}}^p \leq \|\nabla \mathbf{w}\|_{L^\infty(\Omega, \mu_{1-\frac{3}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \|\nabla s\|_{L^p(\Omega, \mu_{\frac{1}{2}-\frac{3}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \left\| \eta_{2-\frac{3p}{2}}^{3-\frac{3p}{2}}(\cdot, \mathcal{R}) \right\|_{L^\infty(\Omega)} \quad (70)$$

Linear terms B_6, C_6 are trivial. We get

$$\begin{aligned} \|B_3, B_4, C_3, C_4\|_{X^{2\omega}} &\leq C\mathcal{W}C_0^2, \\ \|B_5, C_5\|_{X^{2\omega}} &\leq C\mathcal{R}\mathcal{W}C_0^2, \\ \|B_6, C_6\|_{X^{2\omega}} &\leq \mathcal{R}^{1-\omega}\mathcal{W}C_0 \end{aligned} \quad (71)$$

and putting all calculations together we end up with

$$\begin{aligned} \|\mathbf{B}(\mathbf{f}, \mathbf{w}, s)\|_{X^{2\omega}} &\leq C((\mathcal{R} + \mathcal{W})C_0^2 + \mathcal{R}^{1-\omega}\mathcal{W}C_0), \\ \|\mathbf{C}(\mathbf{H}, \mathbf{w}, s)\|_{X^{2\omega}} &\leq C((\mathcal{R} + \mathcal{W})C_0^2 + \mathcal{R}^{1-\omega}\mathcal{W}C_0), \end{aligned} \quad (72)$$

for $p > 6$. Now we can use Theorem 10 on the equation (16) to get

$$\|\mathbf{z}\|_{X^{2\omega}} \leq C\|\mathbf{B}(\mathbf{f}, \mathbf{w}, s)\|_{X^{2\omega}} \leq C((\mathcal{R} + \mathcal{W})C_0^2 + \mathcal{R}^{1-\omega}\mathcal{W}C_0). \quad (73)$$

Moreover we can write the equation (16) in the following form

$$\mathbf{z} = \operatorname{div} [\mathbf{C}(\mathbf{H}, \mathbf{w}, s) - \mathcal{W}\mathbf{z} \otimes (\mathbf{w} + \mathbf{e}_1)] \quad (74)$$

since $\operatorname{div}(\mathbf{w} + \mathbf{e}_1) = 0$. Hence $\mathbf{z} = \operatorname{div} \mathbf{Z}$ for some tensor \mathbf{Z} and

$$\|\mathbf{Z}\|_{X^{2\omega}} \leq \|\mathbf{C}(\mathbf{H}, \mathbf{w}, s)\|_{X^{2\omega}} + \mathcal{W}\|\mathbf{z}\|_{X^{2\omega}} \|(\mathbf{w} + \mathbf{e}_1)\|_\infty \leq C((\mathcal{R} + \mathcal{W})C_0^2 + \mathcal{R}^{1-\omega}\mathcal{W}C_0). \quad (75)$$

Now we can proceed with the Oseen equation (17) with the right hand side $\mathbf{z} = \operatorname{div} \mathbf{Z}$. We use the integral representation formulas (27) - (29) and (30) - (31) to estimate (\mathbf{u}, q) in V . We can split \mathbf{u} into $\mathbf{u} = \mathbf{u}^V + \mathbf{u}^S$, where \mathbf{u}^V denotes the volume integral and \mathbf{u}^S denotes the surface integrals. Similarly we split $\nabla \mathbf{u}$, $\nabla^2 \mathbf{u}$, q and ∇q .

We start with the estimates of the volume parts. For the estimate of $\mathbf{u}^V \in L^p(\Omega, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))$ we use Theorem 6 and its Corollary and get

$$\|\mathbf{u}^V\|_{L^p(\mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq C\mathcal{R}^{-1+\omega}\|\mathbf{Z}\|_{X^{2\omega}} \leq C((\mathcal{R}^\omega + \mathcal{R}^{-1+\omega}\mathcal{W})C_0^2 + \mathcal{W}C_0), \quad (76)$$

which can be made sufficiently small by choosing \mathcal{R}, \mathcal{W} small.

For the estimates of $(\nabla \mathbf{u})^V$ and $(\nabla^2 \mathbf{u})^V$ we use Theorem 5 and its Corollary and get

$$\|(\nabla \mathbf{u})^V\|_{L^p(\mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq C\mathcal{R}^\omega\|\mathbf{Z}\|_{X^{2\omega}} \leq C((\mathcal{R}^{1+\omega} + \mathcal{R}^\omega\mathcal{W})C_0^2 + \mathcal{R}\mathcal{W}C_0) \quad (77)$$

$$\|(\nabla^2 \mathbf{u})^V\|_{L^p(\mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq C\mathcal{R}^\omega\|\mathbf{z}\|_{X^{2\omega}} \leq C((\mathcal{R}^{1+\omega} + \mathcal{R}^\omega\mathcal{W})C_0^2 + \mathcal{R}\mathcal{W}C_0) \quad (78)$$

Again, terms on the right-hand sides can be made sufficiently small by choosing \mathcal{R}, \mathcal{W} small. For the estimates of q^V and $(\nabla q)^V$ we use Theorem 7 and its Corollary and we obtain

$$\begin{aligned} \|q^V\|_{L^p(\mu_{\frac{1}{2}-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} &\leq C\|\mathbf{Z}\|_{L^p(\mu_{\frac{1}{2}-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq \\ &\leq C\|\mathbf{Z}\|_{X^{2\omega}} \leq C((\mathcal{R} + \mathcal{W})C_0^2 + \mathcal{R}^{1-\omega}\mathcal{W}C_0) \end{aligned} \quad (79)$$

$$\begin{aligned} \|(\nabla q)^V\|_{L^p(\mu_{\frac{1}{2}-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} &\leq C\|\mathbf{z}\|_{L^p(\mu_{\frac{1}{2}-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq \\ &\leq C\|\mathbf{z}\|_{X^{2\omega}} \leq C((\mathcal{R} + \mathcal{W})C_0^2 + \mathcal{R}^{1-\omega}\mathcal{W}C_0) \end{aligned} \quad (80)$$

Again the right-hand sides can be made small same way as before.

Next we proceed with the surface integrals. Here we distinguish three cases

$$\begin{aligned}\Omega_1 &= \{\mathbf{x} \in \Omega, |\mathbf{x}| \leq 1\} \\ \Omega_2 &= \left\{ \mathbf{x} \in \Omega, 1 \leq |\mathbf{x}| \leq \frac{1}{\mathcal{R}} \right\} \\ \Omega_3 &= \left\{ \mathbf{x} \in \Omega, |\mathbf{x}| \geq \frac{1}{\mathcal{R}} \right\}.\end{aligned}\tag{81}$$

In the case Ω_1 all our weights ~ 1 and we do not use the integral representation formulas. We rather use the following estimate

$$\|\mathbf{u}\|_{L^p(\Omega_1)} \leq C(1 + \|\nabla \mathbf{u}\|_{W^{1,2}(\Omega_1)}) \leq C(1 + \|\nabla \mathbf{u}\|_{W^{1,2}(\Omega)}) \leq C(1 + K)\tag{82}$$

which is due to Friedrichs inequality and (56). Arising term $C(1 + K)$ can be made small in comparison with C_0 by choosing C_0 large enough. Together with (76) we get

$$\|\mathbf{u}^S\|_{L^p(\Omega_1, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq C(1 + K) + C((\mathcal{R}^\omega + \mathcal{R}^{-1+\omega}\mathcal{W})C_0^2 + \mathcal{W}C_0).\tag{83}$$

We use analogous procedure also for $\nabla \mathbf{u}, \nabla^2 \mathbf{u}$ and get

$$\|\nabla \mathbf{u}, \nabla^2 \mathbf{u}\|_{L^p(\Omega_1)} \leq C(1 + \|\nabla \mathbf{u}\|_{W^{3,2}(\Omega_1)}) \leq C(1 + \|\nabla \mathbf{u}\|_{W^{3,2}(\Omega)}) \leq C(1 + K)\tag{84}$$

and therefore

$$\|(\nabla \mathbf{u})^S, (\nabla^2 \mathbf{u})^S\|_{L^p(\Omega_1, \mu_{1-\frac{2}{p}}^{\frac{3}{2}-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))} \leq C(1 + K) + C((\mathcal{R}^{1+\omega} + \mathcal{R}^\omega \mathcal{W})C_0^2 + \mathcal{R}\mathcal{W}C_0).\tag{85}$$

Analogously for the pressure.

Next we proceed with the case Ω_2 . We start with \mathbf{u}^S and denote four terms in the surface integral (27) by $\mathbf{u}^{S,1}, \dots, \mathbf{u}^{S,4}$. For $\mathbf{u}^{S,1}$ we have

$$\begin{aligned}|\mathbf{u}^{S,1}(\mathbf{x})|^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} &\leq \\ &\leq \mathcal{R}^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left| \int_{\partial\Omega} \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) d\mathbf{y} \right|^p \leq \\ &\leq C\mathcal{R}^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left| \mathcal{O}_{ij}(\mathbf{x}, \mathcal{R}) + \nabla \mathcal{O}_{ij}\left(\frac{\mathbf{x}}{2}, \mathcal{R}\right) \right|^p \leq \\ &\leq C\mathcal{R}^p (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left(\frac{1}{|\mathbf{x}|^{p-p\omega}} + \frac{1}{|\mathbf{x}|^{2p-p\omega}} \right),\end{aligned}\tag{86}$$

where the crucial estimate is the following

$$|\nabla^k \mathcal{O}(\mathbf{x}, \mathcal{R})| \leq C \frac{\mathcal{R}^{\frac{k}{2}}}{|\mathbf{x}|^{1+\frac{k}{2}}}\tag{87}$$

for $k \geq 0$. We use this estimate throughout the rest of the procedure in the case Ω_2 . Therefore

$$\|\mathbf{u}^{S,1}\|_{L^p(\Omega_2, \mu_{1-\frac{2}{p}}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p \leq C\mathcal{R}^p \int_{\Omega_2} \left(\frac{1}{|\mathbf{x}|^{p-p\omega}} + \frac{1}{|\mathbf{x}|^{2p-p\omega}} \right) d\mathbf{x}\tag{88}$$

Arising functions are integrable and $\int_{\Omega_2} |\mathbf{x}|^{p\omega-p} d\mathbf{x} \leq \frac{4\pi}{p-p\omega-3} \leq 4\pi$ for $p > 6$ and $\omega < \frac{1}{2} - \frac{3}{2p}$, i.e. integrals of such functions over Ω_2 are bounded independently of \mathcal{R} by universal constant 4π . Therefore this term can be estimated choosing C_0 large and at this point we do not require \mathcal{R} to be small, even if we have \mathcal{R}^p at our disposal. This fact will play a role in estimating $\mathbf{u}^{S,2}$, $\mathbf{u}^{S,3}$ and $\mathbf{u}^{S,4}$, where there is no power of \mathcal{R} available.

Next for $\mathbf{u}^{S,2}$ we proceed similarly

$$\begin{aligned} & |\mathbf{u}^{S,2}(\mathbf{x})|^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \leq \\ & \leq C |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \times \\ & \quad \times \left| \nabla \mathcal{O}_{ij}(\mathbf{x}, \mathcal{R}) + \nabla^2 \mathcal{O}_{ij}\left(\frac{\mathbf{x}}{2}, \mathcal{R}\right) + e_i(\mathbf{x}) + \nabla e_i\left(\frac{\mathbf{x}}{2}\right) \right|^p \leq \\ & \leq C (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left(\frac{1}{|\mathbf{x}|^{2p-p\omega}} + \frac{1}{|\mathbf{x}|^{3p-p\omega}} \right) \end{aligned} \quad (89)$$

and therefore

$$\|\mathbf{u}^{S,2}\|_{L^p(\Omega_2, \mu_{1-\frac{3}{2p}, \omega}^{\cdot, \mathcal{R}})}^p \leq C \int_{\Omega_2} \left(\frac{1}{|\mathbf{x}|^{2p-p\omega}} + \frac{1}{|\mathbf{x}|^{3p-p\omega}} \right) d\mathbf{x} \quad (90)$$

and we are in similar situation as in the case $\mathbf{u}^{S,1}$.

We treat $\mathbf{u}^{S,3}$ and $\mathbf{u}^{S,4}$ together

$$\begin{aligned} & |\mathbf{u}^{S,3} + \mathbf{u}^{S,4}(\mathbf{x})|^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \leq \\ & \leq C |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left| \mathcal{O}_{ij}(\mathbf{x}, \mathcal{R}) + \nabla \mathcal{O}_{ij}\left(\frac{\mathbf{x}}{2}, \mathcal{R}\right) \right|^p \times \\ & \quad \times \left(\|\nabla \mathbf{u}\|_{W^{1,2}(\Omega)} + \|q\|_{W^{1,2}(\Omega)} + \|\mathbf{Z}\|_{W^{1,2}(\Omega)} \right) \leq \\ & \leq CK (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left(\frac{1}{|\mathbf{x}|^{p-p\omega}} + \frac{1}{|\mathbf{x}|^{2p-p\omega}} \right) \end{aligned} \quad (91)$$

and therefore

$$\|\mathbf{u}^{S,3} + \mathbf{u}^{S,4}\|_{L^p(\Omega_2, \mu_{1-\frac{3}{2p}, \omega}^{\cdot, \mathcal{R}})}^p \leq CK \int_{\Omega_2} \left(\frac{1}{|\mathbf{x}|^{p-p\omega}} + \frac{1}{|\mathbf{x}|^{2p-p\omega}} \right) d\mathbf{x}. \quad (92)$$

Here we have used also (56).

For higher gradients of \mathbf{u} and pressure and its gradient we use similar procedure, in this case higher gradients are even easier to estimate.

We finish with the case Ω_3 . Here the situation is a little different. We have

$$\begin{aligned}
 & \|\mathbf{u}^{S,1}(\mathbf{x})\|^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \leq \\
 & \leq \mathcal{R}^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left| \int_{\partial\Omega} \mathcal{O}_{ij}(\mathbf{x} - \mathbf{y}, \mathcal{R}) d\mathbf{y} \right|^p \leq \\
 & \leq C\mathcal{R}^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \left| \mathcal{O}_{ij}(\mathbf{x}, \mathcal{R}) + \nabla \mathcal{O}_{ij}\left(\frac{\mathbf{x}}{2}, \mathcal{R}\right) \right|^p \leq \\
 & \leq C\mathcal{R}^p |\mathbf{x}|^{p\omega} (1 + \mathcal{R}|\mathbf{x}|)^{p-3-p\omega} (1 + s(\mathcal{R}\mathbf{x}))^{p-2} \times \\
 & \quad \times \left(\frac{\mathcal{R}^p}{|\mathcal{R}\mathbf{x}|^p (1 + s(\mathcal{R}\mathbf{x}))^p} + \frac{\mathcal{R}^{2p}}{|\mathcal{R}\mathbf{x}|^{\frac{3p}{2}} (1 + s(\mathcal{R}\mathbf{x}))^{\frac{3p}{2}}} \right). \quad (93)
 \end{aligned}$$

Here we have used that

$$\left| \nabla^k \mathcal{O}(\mathbf{x}, \mathcal{R}) \right| \leq C \frac{\mathcal{R}^{\frac{k}{2}}}{|\mathbf{x}|^{1+\frac{k}{2}} (1 + s(\mathcal{R}\mathbf{x}))^{1+\frac{k}{2}}} \quad (94)$$

for $k \geq 0$. Therefore

$$\begin{aligned}
 \|\mathbf{u}^{S,1}\|_{L^p(\Omega_3, \mu_{1-\frac{3}{p}, \omega}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p & \leq C\mathcal{R}^{2p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^3} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^2} \right) d\mathbf{x} + \\
 & + C\mathcal{R}^{3p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^{3+\frac{p}{2}}} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^{2+\frac{p}{2}}} \right) d\mathbf{x} \leq \\
 & \leq C\mathcal{R}^{2p-p\omega-3} \int_{\mathbb{R}^3} \eta_{-2}^{-3}(\mathbf{y}) d\mathbf{y} + C\mathcal{R}^{3p-p\omega-3} \int_{\mathbb{R}^3} \eta_{-2-\frac{p}{2}}^{-3-\frac{p}{2}}(\mathbf{y}) d\mathbf{y} \quad (95)
 \end{aligned}$$

Arising integrals are finite due to (34).

For $\mathbf{u}^{S,2}$ we obtain in the similar way

$$\begin{aligned}
 \|\mathbf{u}^{S,2}\|_{L^p(\Omega_3, \mu_{1-\frac{3}{p}, \omega}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p & \leq C\mathcal{R}^{2p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^{3+\frac{p}{2}}} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^{2+\frac{p}{2}}} \right) d\mathbf{x} + \\
 & + C\mathcal{R}^{2p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^{3+p}} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^{2-p}} \right) d\mathbf{x} + \\
 & + C\mathcal{R}^{3p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^{3+p}} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^{2+p}} \right) d\mathbf{x} + \\
 & + C\mathcal{R}^{3p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^{3+2p}} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^{2-p}} \right) d\mathbf{x}. \quad (96)
 \end{aligned}$$

Treating $\mathbf{u}^{S,3}$ and $\mathbf{u}^{S,4}$ together we get

$$\begin{aligned}
 \|\mathbf{u}^{S,3} + \mathbf{u}^{S,4}\|_{L^p(\Omega_3, \mu_{1-\frac{3}{p}, \omega}^{1-\frac{3}{p}, \omega}(\cdot, \mathcal{R}))}^p & \leq C\mathcal{R}^{p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^3} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^2} \right) d\mathbf{x} + \\
 & + C\mathcal{R}^{2p-p\omega} \int_{\Omega_3} \left(\frac{1}{(1 + |\mathcal{R}\mathbf{x}|)^{3+\frac{p}{2}}} \frac{1}{(1 + s(\mathcal{R}\mathbf{x}))^{2+\frac{p}{2}}} \right) d\mathbf{x} \leq \\
 & \leq C\mathcal{R}^{p-p\omega-3} \int_{\mathbb{R}^3} \eta_{-2}^{-3}(\mathbf{y}) d\mathbf{y} + C\mathcal{R}^{2p-p\omega-3} \int_{\mathbb{R}^3} \eta_{-2-\frac{p}{2}}^{-3-\frac{p}{2}}(\mathbf{y}) d\mathbf{y}. \quad (97)
 \end{aligned}$$

Again, we proceed similarly with the estimates of gradients of \mathbf{u} and pressure and its gradient. Putting all calculations together, choosing first C_0 sufficiently

large and then \mathcal{R}, \mathcal{W} sufficiently small we finally end up with

$$\|(\mathbf{u}, q)\|_V < C_0 \quad (98)$$

and the proof of Theorem 1 is finished.

REFERENCES

- [1] Coscia, V., Galdi, G. P.: *Existence, Uniqueness and Stability of Regular Steady Motions of a Second-Grade Fluid*, Int. J. Non-Linear Mechanics, **29**, No. 4, (1994) 493–506.
- [2] Dunn, J. E., Fosdick, R. L.: *Thermodynamics, Stability and Boundedness of Fluids of Complexity 2 and Fluids of Second Grade*, Arch. Rational Mech. Anal., **56** (1974) 191–252.
- [3] Finn, R.: *Estimates at infinity for stationary solution of the Navier-Stokes equations*, Bult. Math. de la Soc. Sci. Math. de la R. P. R., Tome 3, **51**, 4 (1959), 387–418.
- [4] Galdi, G. P.: **An Introduction to the Mathematical Theory of the Navier-Stokes Equations I**, Springer Tracts in Natural Philosophy, Vol. **38**, Springer, New York, (1994).
- [5] Galdi, G. P., Padula, M., Rajagopal, K. R.: *On the conditional stability of the rest state of a fluid of second grade in unbounded domains*, Arch. Rational Mech. Anal., **109**, No. 2, (1990) 173–182.
- [6] Koch, H.: *Partial differential equations and singular integrals*, Dispersive nonlinear problems in mathematical physics, Quad. Mat., **15**, Dept. Math., Seconda Univ. Napoli, Caserta (2004) 59–122.
- [7] Kračmar, S., Novotný, A., Pokorný, M.: *Estimates of three-dimensional Oseen kernels in weighted L^p spaces*, in: da Veiga, B., Sequeira, A., Videman, J. (Eds.): **Applied Functional Analysis**, (1999) 281–316.
- [8] Mogilevskii, I. Sh., Solonnikov V. A.: *Problem on a Stationary Flow of a Second-Grade Fluid in Hölder Classes of Functions*, Zapisky Nauc. Sem. LOMI, **243**, (1997) 154–165.
- [9] Novotný, A.: *About the steady transport equation*, Proceedings of the Fifth Winter School at Paseky, Pitman Research Notes in Mathematics, (1998).
- [10] Novotný, A., Pokorný, M.: *Three-dimensional Steady Flow of Viscoelastic Fluid past an Obstacle*, J. Math. Fluid Mech., **2**, (2000), 294–314.
- [11] Pokorný, M.: **Asymptotic behaviour of solutions to certain PDE’s describing the flow of fluids in unbounded domains**, Ph.D. thesis, Charles University in Prague and University of Toulon and Var, Toulon-La Garde, (1999).
- [12] Smith, D. R.: *Estimates at Infinity for Stationary Solutions of the N. S. Equations in Two Dimensions*, Arch. Rat. Mech. Anal. **20**, (1965), 341–372.
- [13] Truesdell, C., Noll, W.: **The Nonlinear Field Theories of Mechanics**, Handbuch der Physik, III/3, Springer Verlag, Heidelberg (1965).

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