

THE QUANTUM PERMUTATION GROUP OF AN INFINITE COUNTABLE SET

DEBASHISH GOSWAMI AND ADAM SKALSKI

ABSTRACT. Two different models for a Hopf–von Neumann algebra of bounded functions on the quantum permutation group on infinitely many elements are proposed, one based on projective limits of enveloping von Neumann algebras related to finite quantum permutation groups, and the second on universal properties with respect to infinite magic unitaries.

Classical groups first entered mathematics as collections of all symmetries of a given object, be it a finite set, a polygon, a metric space or a manifold. Original definitions of quantum groups (also in the topological context, see [Wor] and [KuV]) had rather algebraic character. Recent years however have brought many developments in the theory of quantum symmetry groups, i.e. quantum groups defined as universal objects acting (in the sense of quantum group actions) on a given structure. The first examples of that type were introduced in [Wan], where S. Wang defined the quantum group of permutations of a finite set, \mathbb{S}_n . It turns out that the algebra of ‘continuous functions on a quantum permutation group of n elements’, $C(\mathbb{S}_n)$, is generated by entries of a universal n by n magic unitary, i.e. a unitary matrix whose entries are orthogonal projections. Later the theory was extended to quantum symmetry groups of finite graphs ([Bic]), finite metric spaces ([Ban₂]), and to quantum isometry spaces of compact noncommutative manifolds ([Gos]). In all these cases the structure whose (quantum) symmetries are studied has finite or compact flavour, so that the resulting quantum symmetry group is compact.

In this paper we study possible definitions of the quantum permutation group of an infinite countable set. Even in the classical context there is a natural choice here – we can either consider the group of all permutations of \mathbb{N} , $\text{Perm}(\mathbb{N})$, or the group of all finite permutations of \mathbb{N} , usually denoted by S_∞ . From the analytic point of view the second group arises more naturally, as it is a direct limit of finite permutation groups S_n . Hence this will be the group whose quantum version we want to discuss here. As on the level of groups we have embeddings $S_n \hookrightarrow S_{n+1}$, on the level of the algebras we obtain surjective morphisms $C(S_{n+1}) \twoheadrightarrow C(S_n)$. Hence it is natural to expect that the algebra of continuous functions on the quantum version of S_∞ will arise as the inverse (projective) limit of algebras $C(\mathbb{S}_n)$. As projective limits of C^* -algebras do not behave well (which is easy to understand even in the commutative setting: a direct limit of locally compact spaces need not be locally compact), we work with von Neumann algebras. This allows us to construct in this note the algebra W_∞ , a candidate for (some version of) $L^\infty(\mathbb{S}_\infty)$ as a limit of enveloping von Neumann algebras of $C(\mathbb{S}_n)$ and to study its universal properties. Another possible approach to infinite quantum permutation groups exploits the

2000 *Mathematics Subject Classification.* Primary 81R50, Secondary 37B40.

Key words and phrases. Quantum permutation group, von Neumann algebra, projective limit.

fact that the algebras $C(\mathbb{S}_n)$ are defined in terms of universal magic unitaries, so by analogy one can investigate a universal von Neumann algebra generated by entries of an infinite magic unitary. We show that such an algebra exists and is a proper subalgebra of W_∞ . We do not know if either of these Hopf–von Neumann algebras fits into the theory of locally compact quantum groups developed in [KuV].

The detailed plan of the paper is as follows: in Section 1 we discuss projective limits of von Neumann algebras; although these results are not difficult and may be well-known, we could not locate a suitable reference. We also include several technical lemmas on extending maps to the projective limits. A short Section 2 contains applications of these results to projective limits of Hopf von Neumann algebras. In Section 3 we recall basic facts on Wang’s quantum permutation groups and describe the first of two possible candidates for the algebra $L^\infty(\mathbb{S}_\infty)$, constructed as the projective limit of the enveloping von Neumann algebras of $C(\mathbb{S}_n)$. In Section 4 we propose an alternative approach in terms of a universal ‘infinite magic unitary’ and explain why this leads to a different Hopf–von Neumann algebra.

The spatial tensor of C^* -algebras will be denoted \otimes , and the ultraweak tensor product of von Neumann algebras $\overline{\otimes}$. For a projection a von Neumann algebra \mathbf{M} its lattice of projections will be denoted $\mathcal{P}(\mathbf{M})$ and the central carrier of $p \in \mathcal{P}(\mathbf{M})$ (i.e. the smallest projection in $Z(\mathbf{M})$ dominating p) will be denoted $z(p)$.

1. PROJECTIVE LIMITS OF VON NEUMANN ALGEBRAS

In this section we define, establish existence and prove basic properties of inductive limits of von Neumann algebras. Although most results remain valid for general directed index sets, we consider only projective systems indexed by \mathbb{N} .

Definition 1.1. A sequence $(\mathbf{M}_n)_{n \in \mathbb{N}}$ is a *projective system of von Neumann algebras* if it is a sequence of von Neumann algebras equipped with surjective normal $*$ -homomorphisms $\phi_n : \mathbf{M}_{n+1} \rightarrow \mathbf{M}_n$ (the maps ϕ_n form a part of the definition, but we omit them from the notation). Define the following class of von Neumann algebras: $\mathfrak{M} = \{\mathbf{M} : \forall n \in \mathbb{N} \exists \psi_n : \mathbf{M} \rightarrow \mathbf{M}_n, \text{ a surjective normal } * \text{-homomorphism such that } \psi_n = \phi_n \circ \psi_{n+1}\}$. We say that $\mathbf{M} \in \mathfrak{M}$ is a *final object* for \mathfrak{M} if for each $\mathbf{N} \in \mathfrak{M}$ there exists a surjective normal morphism $\psi : \mathbf{N} \rightarrow \mathbf{M}$ such that $\psi_n^{(\mathbf{M})} \circ \psi = \psi_n^{(\mathbf{N})}$ for all $n \in \mathbb{N}$.

Note that it is not clear at the moment whether even if a final object for \mathfrak{M} exists, it is unique.

Theorem 1.2. *Let $(\mathbf{M}_n)_{n \in \mathbb{N}}$ be a projective system of von Neumann algebras. Then the class \mathfrak{M} admits a (unique) final object.*

Proof. The construction is based on the properties of weak $*$ -closed two-sided ideals in von Neumann algebras. Let $n \geq 2$ and $\mathfrak{l}_n = \text{Ker}(\phi_{n-1})$. Let $r_n \in \mathcal{P}(Z(\mathbf{M}_n))$ be the projection such that $\mathfrak{l}_n = r_n \mathbf{M}_n$ (recall that $r_n := \sup\{p \in \mathcal{P}(\mathbf{M}_{n-1}) : \phi_n(p) = 0\}$). A well-known (and easy to check) fact states that the map $\phi_{n-1}|_{r_n^\perp \mathbf{M}_n} : r_n^\perp \mathbf{M}_n \rightarrow \mathbf{M}_{n-1}$ is an isomorphism. Let $\mathbf{B}_n = r_n \mathbf{M}_n$ and define additionally $\mathbf{B}_1 = \mathbf{M}_1$. Then each \mathbf{M}_n has a natural decomposition of the form $\mathbf{M}_n = \bigoplus_{k=1}^n \mathbf{B}_k$, and additionally this decomposition is ‘well behaved’ with respect to the maps ϕ_n . Not surprisingly, the final object in \mathfrak{M} will be isomorphic to $\prod_{n=1}^\infty \mathbf{B}_n$. Below we give a detailed proof of this fact.

Observe first that the class \mathfrak{M} is non-empty. Indeed, define $M_\infty = \{(m_n)_{n=1}^\infty \in \prod_{n=1}^\infty M_n : \phi_n(m_{n+1}) = m_n\}$. Then M_∞ is a weak*-closed subalgebra of $\prod_{n=1}^\infty M_n$, hence a von Neumann algebra. It is clear that the projections on the individual coordinates are normal *-homomorphisms; they satisfy the intertwining relation with ϕ_n by construction. Surjectivity follows from the existence of isometric lifts for selfadjoint elements in C^* -algebras (hence bounded lifts for arbitrary elements of M_n to elements in M_∞). In fact M_∞ will be (isomorphic to) the final object for \mathfrak{M} .

Let $N \in \mathfrak{M}$ and denote by J_n the kernel of the corresponding map $\psi_n : N \rightarrow M_n$. Let $w_n \in \mathcal{P}(Z(N))$ be the projection such that $J_n = w_n^\perp N$. As in the first part of the proof, $\psi_n|_{w_n^\perp N} : w_n^\perp N \rightarrow M_n$ is an isomorphism. Write $z_n := w_n^\perp$. As $J_{n+1} \subset J_n$, the sequence $(z_n)_{n=1}^\infty$ is increasing. Define additionally $z_\infty = \lim_{n \rightarrow \infty} z_n$, $p_1 = z_1$ and $p_n = z_n - z_{n-1}$ for $n \geq 2$, so that $z_\infty = \sum_{n=1}^\infty p_n$. As all projections p_n are central, we obtain a natural increasing sequence of von Neumann algebras $\oplus_{k=1}^n p_k N$ whose union is weak*-dense in $z_\infty N$. It is easy to see that this yields a natural isomorphism $z_\infty N \approx \prod_{n=1}^\infty p_n N$.

Note that $z_\infty N \in \mathfrak{M}$ - indeed, the only thing to check is that the maps $\psi_n|_{z_\infty N} : M_n$ are surjections, and this follows from the stated above surjectivity of $\psi_n|_{z_n N}$. Our claim is that $z_\infty N$ is the final object of \mathfrak{M} . Indeed, it suffices to show that if W is another von Neumann algebra in \mathfrak{M} , then $z_\infty^{(W)} W$ is isomorphic to $z_\infty N$ and the isomorphism intertwines the corresponding maps into M_n . For the first statement it suffices to describe the algebras $p_n N$ in terms of the projective sequence with which we started. Let $n \geq 2$. Consider the diagram

$$\begin{array}{ccc}
& z_{n-1} N \oplus p_n N = z_n N & \\
& \swarrow \psi_{n-1}|_{z_{n-1} N} & \downarrow \psi_n|_{z_n N} \\
M_{n-1} & & \\
& \swarrow \phi_{n-1}|_{r_n^\perp M_n} & \\
& r_n^\perp M_n \oplus B_n = M_n &
\end{array}$$

in which all arrows are isomorphisms. It immediately implies that $p_n N$ is isomorphic to B_n (note that for $n = 1$ this also holds). Moreover looking at the diagram above we see that if we denote the corresponding isomorphism between $p_n N$ and B_n by γ_n , we can check inductively that $\gamma_1 \oplus \dots \oplus \gamma_n : z_n N \rightarrow M_n$ coincides with ψ_n , which assures that the natural isomorphism between $z_\infty^{(W)} W$ and $z_\infty N$ intertwines the respective ψ_n and $\psi_n^{(W)}$ maps.

We can check that for $N := M_\infty$ we have $z_\infty = 1_{M_\infty}$. Indeed, if $(m_n)_{n=1}^\infty \in w_\infty M_\infty$ then $(m_n)_{n=1}^\infty \in \text{Ker}(\psi_n)$ for each $n \in \mathbb{N}$, so $(m_n)_{n=1}^\infty = 0$.

It remains to prove uniqueness. Suppose then that N is a final object in \mathfrak{M} and let W be a final object in \mathfrak{M} constructed above. Note that if $\psi_n^{(W)} : W \rightarrow M_n$ denote the usual surjections, the construction above implies that $\bigcap_{n=1}^\infty \text{Ker}(\psi_n^{(W)}) = \{0\}$. There is a surjective map $\psi : W \rightarrow N$ such that $\psi_n = \psi_n^{(W)} \circ \psi$ for all $n \in \mathbb{N}$. Thus we must have $\bigcap_{n=1}^\infty \text{Ker}(\psi_n) = \{0\}$, or equivalently $z_\infty = 1_N$, where z_∞ is constructed for N as above. Then $N = Nz_\infty$ and the arguments above show that $N \approx W$. \square

Definition 1.3. Let $(M_n)_{n \in \mathbb{N}}$ be a projective system of von Neumann algebras. The final object in the class \mathfrak{M} will be called the projective limit of $(M_n)_{n \in \mathbb{N}}$ and denoted M_∞ .

In the next section we will show that if $(M_n)_{n \in \mathbb{N}}$ is a projective system of Hopf-von Neumann algebras, with the normal surjections ϕ_n intertwining the respective coproducts, then M_∞ has a natural Hopf-von Neumann algebra structure. To this end we present here several lemmas related to constructing maps acting on/to/between projective limits.

Lemma 1.4. Let $(M_n)_{n \in \mathbb{N}}$ be as in Theorem 1.2 and let us adopt the notations in the proof that theorem. Define additionally for each $n \in \mathbb{N}$ the map $\iota_n : M_n \rightarrow M_\infty$ to be the inverse of $\psi_n|_{z_n M_\infty}$ (or more precisely the composition of that inverse with the embedding of $z_n M_\infty$ into M_∞). Then we have the following: for each $n \in \mathbb{N}$, $x \in M_{n+1}$

$$\iota_n(\phi_n(x)) = z_n \iota_{n+1}(x)$$

Proof. It is a direct consequence of the diagram above, this time interpreted as follows:

$$\begin{array}{ccc}
 & z_n M_\infty & \xleftarrow{z_n} M_\infty \\
 & \nearrow \iota_n & \uparrow \iota_{n+1} \\
 M_n & & M_{n+1} \\
 & \nwarrow \phi_n &
 \end{array}$$

- note that now the maps are not necessarily isomorphisms. □

Lemma 1.5. Suppose that $(N_n)_{n=1}^\infty$, W are von Neumann algebras and that $N = \prod_{n \in \mathbb{N}} N_n$. For each $n \in \mathbb{N}$ denote the central projection in N corresponding to N_n by p_n . Let (for each $n \in \mathbb{N}$) $\kappa_n : W \rightarrow \prod_{k=1}^n N_k$ be a normal contractive map and suppose that (for each $w \in W, n \in \mathbb{N}$)

$$(1.1) \quad \kappa_n(w) = \sum_{k=1}^n p_k \kappa_{n+1}(w).$$

Then there exists a unique normal contraction $\kappa : W \rightarrow N$ such that

$$\kappa_n(w) = \sum_{k=1}^n p_k \kappa(w).$$

If each κ_n is a $*$ -homomorphism (respectively, a $*$ -antihomomorphism), κ is also $*$ -homomorphic (respectively, $*$ -antihomomorphic).

Proof. Let $w \in W$. Define

$$\kappa(w) = \sum_{n=1}^\infty p_n \kappa_n(w) = \lim_{n \rightarrow \infty} \kappa_n(w).$$

The equality of both expressions follows from the formula (1.1) and the properties of weak $*$ topology in N (recall that we have a natural Banach space isomorphism $N_* \approx \bigoplus_{n=1}^\infty (N_n)_*$, where the last sum is of the l^1 -type). Similarly, normality of κ

follows from the explicit description of the predual of \mathbb{N} and normality of each κ_n . The statement on algebraic properties of κ is easy to check, and the uniqueness is clear. \square

The last two results have a following consequence.

Proposition 1.6. *Suppose that $(M_n)_{n \in \mathbb{N}}$ and $(N_n)_{n \in \mathbb{N}}$ are projective systems of von Neumann algebras, with connecting maps respectively denoted by $(\phi_n^{(M)})_{n \in \mathbb{N}}$ and $(\phi_n^{(N)})_{n \in \mathbb{N}}$ and the maps from the final objects M_∞ and N_∞ respectively denoted by $(\psi_n^{(M)})_{n \in \mathbb{N}}$ and $(\psi_n^{(N)})_{n \in \mathbb{N}}$. Let $\lambda_n : M_n \rightarrow N_n$ ($n \in \mathbb{N}$) be normal contractive maps such that*

$$\lambda_n \circ \phi_n^{(M)} = \phi_n^{(N)} \circ \lambda_{n+1}, \quad n \in \mathbb{N}.$$

Then there exists a unique map $\lambda_\infty : M_\infty \rightarrow N_\infty$ such that

$$\lambda_n \circ \psi_n^{(M)} = \psi_n^{(N)} \circ \lambda_\infty, \quad n \in \mathbb{N}.$$

If each λ_n is a $*$ -homomorphism (respectively, a $*$ -antihomomorphism, a unital map), λ is also $*$ -homomorphic (respectively, $*$ -antihomomorphic, unital).

Proof. Use the notation of Theorem 1.2 and Lemma 1.4, adorning respective maps with (M) and (N) . Define $\tilde{\lambda}_n : M_\infty \rightarrow z_n^{(N)} N_\infty$ ($n \in \mathbb{N}$) as $\tilde{\lambda}_n = \iota_n^{(M)} \circ \lambda_n \circ \psi_n^{(N)}$. Then

$$\begin{aligned} z_n^{(N)} \tilde{\lambda}_{n+1}(\cdot) &= z_n^{(N)} (\iota_{n+1}^{(N)} \circ \lambda_{n+1} \circ \psi_{n+1}^{(M)})(\cdot) = \iota_n^{(N)} \circ \phi_n^{(N)} \circ \lambda_{n+1} \circ \psi_{n+1}^{(M)} \\ &= \iota_n^{(N)} \circ \lambda_n \circ \phi_n^{(M)} \circ \psi_{n+1}^{(M)} = \iota_n^{(N)} \circ \lambda_n \circ \psi_n^{(M)} = \tilde{\lambda}_n, \end{aligned}$$

where in the second equality we used Lemma 1.4. Apply now Lemma 1.5 for $\kappa_n := \lambda_n$, $W := M_\infty$ and $N_n := p_n^{(N)} N_\infty$. This yields a map $\lambda_\infty : M_\infty \rightarrow N_\infty$ such that

$$\tilde{\lambda}_n = z_n^{(N)} \lambda_\infty(\cdot).$$

Straightforward identifications using the commuting diagrams presented earlier end the proof of the main statement. As before, uniqueness and algebraic properties of λ_∞ follow easily. \square

The above lemma provides a simple corollary describing a construction of maps acting from M_∞ into some other von Neumann algebra.

Corollary 1.7. *Let $(M_n)_{n \in \mathbb{N}}$ be a projective system of von Neumann algebras; adopt the notations of Theorem 1.2. Let (for each $n \in \mathbb{N}$) $\mu_n : M_n \rightarrow W$ be a normal $*$ -homomorphism and suppose that (for each $n \in \mathbb{N}$)*

$$\mu_n \circ \phi_n = \mu_{n+1}.$$

Then there exists a unique normal $*$ -homomorphism $\mu : M_\infty \rightarrow W$ such that

$$\mu = \mu_n \circ \psi_n.$$

Proof. It suffices to apply Proposition 1.6 to the projective systems $(M_n)_{n \in \mathbb{N}}$ and $(N_n)_{n \in \mathbb{N}}$, where $N_n := W$ and $\phi_n = \text{id}_W$ for all $n \in \mathbb{N}$. \square

2. PROJECTIVE LIMITS OF HOPF-VON NEUMANN ALGEBRAS

Here we apply the results of Section 1 to construct the projective limit of a projective sequence of Hopf-von Neumann algebras.

Definition 2.1. A Hopf-von Neumann algebra is a von Neumann algebra equipped with a coproduct, i.e. a unital normal $*$ -homomorphism $\Delta : M \rightarrow M \overline{\otimes} M$ which is coassociative:

$$(\text{id}_M \otimes \Delta)\Delta = (\Delta \otimes \text{id}_M)\Delta.$$

Definition 2.2. A sequence $(M_n)_{n \in \mathbb{N}}$ is called a projective system of Hopf-von Neumann algebras if it is a projective system of von Neumann algebras, each M_n is a Hopf-von Neumann algebra (with the coproduct $\Delta_n : M_n \rightarrow M_n \otimes M_n$) and the surjective normal homomorphisms $\phi_n : M_{n+1} \rightarrow M_n$ satisfy the conditions

$$(\phi_n \otimes \phi_n)\Delta_{n+1} = \Delta_n \phi_n.$$

Theorem 2.3. *Let $(M_n)_{n \in \mathbb{N}}$ be a projective system of Hopf-von Neumann algebras. Then M_∞ is also a Hopf-von Neumann algebra: there exists a unique coproduct $\Delta : M_\infty \rightarrow M_\infty \overline{\otimes} M_\infty$ such that*

$$(2.1) \quad \Delta_n \psi_n = (\psi_n \otimes \psi_n)\Delta, \quad n \in \mathbb{N}.$$

In addition if each Δ_n is injective, so is Δ .

Proof. Observe that the sequence $(M_n \overline{\otimes} M_n)_{n \in \mathbb{N}}$, together with surjective connecting maps $\phi_n \otimes \phi_n : M_{n+1} \overline{\otimes} M_{n+1} \rightarrow M_n \overline{\otimes} M_n$ forms a projective limit of von Neumann algebras; moreover a projective limit of this sequence can be easily identified with $M_\infty \overline{\otimes} M_\infty$. Hence an application of Proposition 1.6 yields the existence and uniqueness of a unital normal $*$ -homomorphism $\Delta : M_\infty \rightarrow M_\infty \overline{\otimes} M_\infty$ satisfying (2.1).

Coassociativity of Δ can be proved in an analogous way, exploiting the uniqueness part of Proposition 1.6.

If each Δ_n is injective, $x \in M_\infty$ and $\Delta(x) = 0$, then by (2.1) we have (for each $n \in \mathbb{N}$) $\psi_n(x) = 0$. Via identifications in Theorem 1.2 we see that $z_n x = 0$ for all $n \in \mathbb{N}$, which implies that $x = 0$. \square

We could also consider Hopf-von Neumann algebras with a *counit*, i.e. a normal character $\epsilon : M \rightarrow \mathbb{C}$ such that

$$(\epsilon \otimes \text{id}_M)\Delta = (\text{id}_M \otimes \epsilon)\Delta = \text{id}_M$$

Then for $(M_n)_{n \in \mathbb{N}}$ to be a projective system of Hopf-von Neumann algebras we additionally require that

$$\epsilon_n \circ \phi_n = \epsilon_{n+1}, \quad n \in \mathbb{N}.$$

Lemma 1.7 and a simple calculation imply that if the above conditions are satisfied, then M_∞ admits a natural counit.

We finish this section with a short discussion of projective limits of actions of Hopf-von Neumann algebras.

Definition 2.4. Let W be a von Neumann algebra and (M, Δ) be a Hopf-von Neumann algebra. We say that $\alpha : W \rightarrow W \overline{\otimes} M$ is a (Hopf-von Neumann algebraic) action of M on W if it is a normal unital injective $*$ -homomorphism such that

$$(\text{id}_W \otimes \Delta)\alpha = (\alpha \otimes \text{id}_M)\alpha.$$

A combination of Theorem 2.3 and Lemma 1.5 yields the following result, which says that the Hopf-von Neumann algebraic actions behave well under passing to projective limits.

Theorem 2.5. *Let \mathcal{W} be a von Neumann algebra and let $(M_n)_{n \in \mathbb{N}}$ be a projective system of Hopf-von Neumann algebras. Denote by M_∞ the Hopf-von Neumann algebra arising as the projective limit in the sense of Theorem 2.3. Let $(\alpha_n)_{n \in \mathbb{N}}$ be a sequence of actions of M_n on \mathcal{W} such that for each $n \in \mathbb{N}$*

$$(id_{\mathcal{W}} \otimes \phi_n)\alpha_{n+1} = \alpha_n,$$

where ϕ_n are connecting maps defining the system $(M_n)_{n \in \mathbb{N}}$. Then there exists a unique action α of M_∞ on \mathcal{W} such that for each $n \in \mathbb{N}$

$$(id_{\mathcal{W}} \otimes \psi_n)\alpha = \alpha_n.$$

Proof. Similar to that of Theorem 2.3, using the fact that the von Neumann algebra $\mathcal{W} \overline{\otimes} M_\infty$ is the projective limit of the system $(\mathcal{W} \overline{\otimes} M_n)_{n \in \mathbb{N}}$, with the connecting maps $id_{\mathcal{W}} \otimes \phi_n$, and then applying Proposition 1.6. \square

3. THE HOPF-VON NEUMANN ALGEBRA OF ‘ALL FINITE QUANTUM PERMUTATIONS OF AN INFINITE SET’ AS A PROJECTIVE LIMIT

Let $C(\mathbb{S}_n)$ denote the algebra of continuous functions on the quantum permutation group of the n -point set. Recall ([Wan]) that it is the universal C^* -algebra generated by the collection of orthogonal projections $\{p_{ij} : i, j = 1, \dots, n\}$ such that for each $i = 1, \dots, n$ there is $\sum_{j=1}^n p_{ij} = \sum_{j=1}^n p_{ji} = 1$. The coproduct, counit and (bounded, *-antihomomorphic) antipode are defined on $C(\mathbb{S}_n)$ by the formulas ($i, j = 1, \dots, n$)

$$\Delta_n(p_{ij}) = \sum_{k=1}^n p_{ik} \otimes p_{kj},$$

$$\epsilon_n(p_{ij}) = \delta_{ij}, \quad \kappa_n(p_{ij}) = p_{ji}.$$

Denote the enveloping von Neumann algebra of $C(\mathbb{S}_n)$ by W_n . Standard arguments show that maps Δ_n , ϵ_n and κ_n have unique normal extensions to W_n , which will be denoted by the same symbols – so that for example $\Delta_n : W_n \rightarrow W_n \overline{\otimes} W_n$.

For each $n \in \mathbb{N}$ we denote by ω_n the natural surjection (and a compact quantum group morphism) from $C(\mathbb{S}_{n+1})$ to $C(\mathbb{S}_n)$, which corresponds to mapping $\begin{bmatrix} P & 0 \\ 0 & 1 \end{bmatrix} \rightarrow P$ and whose existence follows from the universal properties. This induces in a standard way the surjection on the level of universal enveloping von Neumann algebras (it is enough to define $\phi_n = \omega_n^{**} : C(\mathbb{S}_{n+1})^{**} \rightarrow C(\mathbb{S}_n)^{**}$ - the fact that ϕ_n is multiplicative is the standard consequence of the definition of the Arens multiplication, surjectivity follows from the fact that images of normal representations of von Neumann algebras are ultraweakly closed). Hence the sequence of algebras $(W_n)_{n=1}^\infty$ forms a projective system of von Neumann algebras. As ω_n intertwines the respective coproducts on the level of C^* -algebras, so does ϕ_n on the level of von Neumann algebras; similarly $\epsilon_{n+1} \circ \phi_n = \epsilon_n$ for all $n \in \mathbb{N}$. Hence Theorem 2.3 implies that the projective limit of $(W_n)_{n \in \mathbb{N}}$ is a Hopf-von Neumann algebra, denoted further by W_∞ . We formulate it as a theorem:

Theorem 3.1. *The sequence $(W_n := C(\mathbb{S}_n)^{**})_{n=1}^\infty$ is a projective system of Hopf-von Neumann algebras with counits. Hence its projective limit denoted by W_∞ is also a Hopf-von Neumann algebra with a counit.*

Proof. A direct consequence of Theorem 2.3 and the discussion before the theorem. \square

In general we cannot expect Hopf-von Neumann algebras to possess antipodes. Here we have however the following fact.

Theorem 3.2. *The Hopf-von Neumann algebra W_∞ admits a unique *-antihomomorphic involutive map $\kappa : W_\infty \rightarrow W_\infty$ such that*

$$\kappa_n \circ \psi_n = \psi_n \circ \kappa, \quad n \in \mathbb{N},$$

where $\psi_n : W_\infty \rightarrow W_n$ are the canonical surjections.

Proof. As $\omega_n \circ \kappa_{n+1}|_{C(\mathbb{S}_{n+1})} = \kappa_n \circ \omega_n|_{C(\mathbb{S}_{n+1})}$, we also have a similar relation on the level of maps between the enveloping von Neumann algebras, with ω_n replaced by ϕ_n . Hence Proposition 1.6 implies the existence of the map κ as above; the fact it is involutive is a consequence of the analogous property of all κ_n . \square

It would be of course more natural to use for the projective limit construction instead of $C(\mathbb{S}_n)^{**}$ the algebras $L^\infty(\mathbb{S}_n)$, the von Neumann completions of $C(\mathbb{S}_n)$ in the GNS representation with respect to the respective Haar states. The problem lies in the fact that the maps ω_n cannot extend to ‘reduced’ versions of the algebras of $C(\mathbb{S}_n)$, so also not to continuous maps $L^\infty(\mathbb{S}_{n+1}) \rightarrow L^\infty(\mathbb{S}_n)$. The first statement is a consequence of the fact that $C(\mathbb{S}_n)$ is not *coamenable* for $n \geq 5$, as follows from the quantum version of the Kesten criterion for amenability ([Ban₁]).

The fact that we can only construct the projective limit using the universal completions is related to the problem described in the next remark.

Remark 3.3. Recently C. Köstler and R. Speicher introduced a notion of *quantum exchangeability* or *invariance under quantum permutations* for a family of quantum random variables (see Definition 2.4 in [KSp]). This notion was later studied by S. Curran in [Cur] and extended to finite sequences; the basic idea is that a sequence of random variables is quantum exchangeable if its distribution (understood as a state on a von Neumann algebra generated by the variables in question) is invariant under natural actions of all Wang’s quantum permutation groups \mathbb{S}_n . Classically exchangeability can be defined as the invariance of the distribution under the action of the infinite permutation group; it would be natural to expect a similar result in the quantum context. It is not clear whether our definition would allow such a formulation; although Theorem 2.5 offers a way of constructing actions of the projective limit, the natural actions of quantum permutation groups considered in [KSp] are defined only on the Hopf algebraic level. As shown in Theorem 3.3 of [Cur] (see also Section 5.6 of that paper), in the presence of quantum exchangeability the actions can be extended to the reduced von Neumann algebraic completions $L^\infty(\mathbb{S}_n)$, but to apply Theorem 2.5 to obtain the action of W_∞ on the von Neumann algebra in question we would need to be able to extend the original actions to $C(\mathbb{S}_n)^{**}$.

4. UNIVERSAL VON NEUMANN ALGEBRA GENERATED BY AN INFINITE MAGIC UNITARY

In this section we shall define a quantum analogue of the algebra of functions on the permutation group of a countably infinite set as the universal von Neumann algebra generated by the entries of an ‘infinite magic unitary’.

We begin with a C^* -algebraic construction.

Definition 4.1. Let \mathfrak{C} denote the category with objects $(\mathbf{C}, \{q_{ij}, i, j = 1, \dots, \infty\})$, where \mathbf{C} is a (possibly nonunital) C^* -algebra generated by a family of orthogonal projections $\{q_{ij} : i, j \in \mathbb{N}\}$ and such that there exists a faithful (and nondegenerate) representation (π, \mathbf{H}) of \mathbf{C} such that for each $i \in \mathbb{N}$

$$(4.1) \quad \sum_{j=1}^{\infty} \pi(q_{ij}) = \sum_{j=1}^{\infty} \pi(q_{ji}) = 1_{B(\mathbf{H})},$$

with the convergence understood in the strong operator topology. A morphism from $(\mathbf{C}, \{q_{ij}\})$ to $(\mathbf{C}', \{q'_{ij}\})$ is given by a (necessarily nondegenerate) C^* -homomorphism from \mathbf{C} to \mathbf{C}' which maps q_{ij} to q'_{ij} for all $i, j \in \mathbb{N}$.

Theorem 4.2. *The category \mathfrak{C} has a universal (initial) object.*

Proof. Consider the (formal) $*$ -algebra \mathcal{B} generated by symbols $\{b_{ij} : i, j \in \mathbb{N}\}$ which are selfadjoint idempotents

$$(4.2) \quad b_{ij} = b_{ij}^* = b_{ij}^2,$$

and satisfy the orthogonality relations

$$(4.3) \quad b_{ij}b_{ik} = 0, \quad b_{ji}b_{ki} = 0 \text{ for } k \in \mathbb{N} \text{ such that } j \neq k.$$

It is easy to see that this $*$ -algebra admits many nontrivial representations on Hilbert spaces. For example, for any $n \in \mathbb{N}$, we can denote the canonical generators of $C(\mathbb{S}_n)$ by $\{q_{ij}^{(n)} : i, j = 1, \dots, n\}$ and put $b_{ij}^{(n)} = q_{ij}^{(n)}$ for $i, j \leq n$, $b_{ij}^{(n)} = 0$ otherwise. Clearly, $b_{ij}^{(n)}$ satisfy the required relations, so that we get a $*$ -homomorphism $\rho_n : \mathcal{B} \rightarrow C(\mathbb{S}_n)$ sending b_{ij} to $b_{ij}^{(n)}$ and we can compose it with any faithful representation of $C(\mathbb{S}_n)$. Since each b_{ij} is a self-adjoint projection, the norm of its image under any representation on a Hilbert space must be less than or equal to 1. This implies that the universal norm defined by $\|b\| := \sup_{\pi} \|\pi(b)\|$, where π varies over all representations of \mathcal{B} on a Hilbert space, is finite. The completion of \mathcal{B} under this norm will be denoted by \mathbf{B} . It is the universal C^* -algebra generated by $\{b_{ij} : i, j \in \mathbb{N}\}$ satisfying relations (4.2)-(4.3). We shall denote the universal enveloping von Neumann algebra of \mathbf{B} by \mathbf{B}^{**} and identify \mathbf{B} as a C^* -subalgebra of \mathbf{B}^{**} .

Observe that for fixed $i \in \mathbb{N}$, $p_i^{(n)} := \sum_{j=1}^n b_{ij}$ is an increasing family of projections in $\mathbf{B} \subset \mathbf{B}^{**}$, so it will converge in the ultraweak topology of \mathbf{B}^{**} to some projection, say, p_i . Similarly, for fixed $j \in \mathbb{N}$, we write $p^j := \lim_{n \rightarrow \infty} \sum_{i=1}^n b_{ij}$ in \mathbf{B}^{**} . Let w be the smallest central projection in \mathbf{B}^{**} which dominates $1 - p_i, 1 - p^j$ for all $i, j \in \mathbb{N}$ and let $z = 1 - w$. Consider the C^* -algebra $\mathbf{A} := z\mathbf{B} \subset \mathbf{B}^{**}$. Clearly, \mathbf{A} is generated as a C^* -algebra by projections $\{q_{ij} := zb_{ij} : i, j \in \mathbb{N}\}$. We claim that $(\mathbf{A}, \{q_{ij} : i, j \in \mathbb{N}\})$ is in \mathfrak{C} and is indeed the universal C^* -algebra in this category.

First of all, it follows from the definition of z that for each $i \in \mathbb{N}$ we have $\sum_{j=1}^{\infty} q_{ij} = 1 = \sum_{j=1}^{\infty} q_{ji}$ in the ultraweak topology inherited from the inclusion

$z\mathbf{B}^{**} \subseteq \mathbf{B}^{**}$, i.e. the ultraweak topology of $\mathbf{B}(z\mathbf{H}_u)$ where \mathbf{H}_u denotes the universal Hilbert space on which \mathbf{B}^{**} acts. We complete the proof of the lemma by showing the universality of \mathbf{A} . To this end, let \mathbf{D} be a C^* -algebra generated by elements $\{t_{ij} : i, j \in \mathbb{N}\}$ satisfying the relations (4.1), where the infinite series in (4.1) converge in the ultraweak topology of the von Neumann algebra $\pi(\mathbf{D})''$ for a fixed faithful representation (π, \mathbf{H}) of \mathbf{D} . By the definition of \mathbf{B} , we get a $*$ -homomorphism from \mathbf{B} onto \mathbf{D} which sends b_{ij} to t_{ij} (for each $i, j \in \mathbb{N}$). This composed with π extends to a unital, normal $*$ -homomorphism, say ρ , from \mathbf{B}^{**} onto $\pi(\mathbf{D})''$. In particular, $\rho(p_i) = \sum_{j=1}^{\infty} t_{ij} = 1$, and $\rho(p^i) = \sum_{j=1}^{\infty} t_{ji} = 1$ for all $i \in \mathbb{N}$, so $1 - p_i, 1 - p^i$ belong to the ultraweakly closed two-sided ideal $\mathbf{l} := \text{Ker } \rho$ of \mathbf{B}^{**} . Thus, if we denote by w_0 the central projection in \mathbf{B}^{**} such that $\mathbf{l} = w_0\mathbf{B}^{**}$, then w_0 dominates $1 - p_i$ and $1 - p^i$ for all $i \in \mathbb{N}$, and hence by the definition of w , we have $w_0 \geq w$. It follows that $w \in \mathbf{l}$, i.e. $\rho(w) = 0$, or in other words, $\rho(z) = 1$. This implies $\rho(b) = \rho(zb)$ for all $b \in \mathbf{B}$, so that we get a $*$ -homomorphism $\rho_1 := \rho|_{\mathbf{A}}$ from \mathbf{A} to \mathbf{D} which satisfies $\rho_1(q_{ij}) = t_{ij}$ for all $i, j \in \mathbb{N}$. This completes the proof of the universality of \mathbf{A} . \square

Denote the von Neumann algebra $z\mathbf{B}^{**}$ by \mathbf{A}_{∞} , and note that it should not be confused with the universal enveloping von Neumann algebra of \mathbf{A} , which may be bigger. Note that the proof of the above theorem indeed provides also universal property of the von Neumann algebra \mathbf{A}_{∞} , as stated in the next corollary.

Corollary 4.3. *The von Neumann algebra \mathbf{A}_{∞} is the (unique up to an isomorphism of von Neumann algebras) universal object in the category of all von Neumann algebras \mathbf{N} which are generated (in the ultraweak topology) by projections $\{n_{ij} : i, j \in \mathbb{N}\}$ satisfying $\sum_{j=1}^{\infty} n_{ij} = \sum_{j=1}^{\infty} n_{ji} = 1_{\mathbf{N}}$ (convergence in the ultraweak topology).*

Using the von Neumann algebraic universality we have the following result.

Proposition 4.4. *The von Neumann algebra \mathbf{A}_{∞} admits a natural coproduct $\Delta_{\mathbf{A}} : \mathbf{A}_{\infty} \rightarrow \mathbf{A}_{\infty} \overline{\otimes} \mathbf{A}_{\infty}$ and a counit $\epsilon_{\mathbf{A}} : \mathbf{A}_{\infty} \rightarrow \mathbb{C}$.*

Proof. Consider for each $i, j \in \mathbb{N}$

$$x_{ij} := \sum_{k=1}^{\infty} q_{ik} \otimes q_{kj}$$

as an element of $\mathbf{A}_{\infty} \overline{\otimes} \mathbf{A}_{\infty}$. We note that the series converges in the ultraweak topology of the von Neumann algebra $\mathbf{A}_{\infty} \overline{\otimes} \mathbf{A}_{\infty}$, the summands being mutually orthogonal projections. It is easy to check using the defining properties of q_{ij} that for each $i, j \in \mathbb{N}$ there is $x_{ij}^2 = x_{ij} = x_{ij}^*$, and $\sum_{k=1}^{\infty} x_{ik} = \sum_{k=1}^{\infty} x_{ki} = 1_{\mathbf{A}_{\infty} \overline{\otimes} \mathbf{A}_{\infty}}$. By the universality of the von Neumann algebra stated in Corollary 4.3, we obtain a normal unital $*$ -homomorphism $\Delta_{\mathbf{A}} : \mathbf{A}_{\infty} \rightarrow \mathbf{A}_{\infty} \overline{\otimes} \mathbf{A}_{\infty}$ given by $\Delta_{\mathbf{A}}(q_{ij}) = x_{ij}$, $i, j \in \mathbb{N}$, which is easily seen to be coassociative. Similarly, we have a normal $*$ -homomorphism $\epsilon_{\mathbf{A}} : \mathbf{A}_{\infty} \rightarrow \mathbb{C}$ given on generators by $\epsilon_{\mathbf{A}}(q_{ij}) = \delta_{ij}$. Note that the existence of the counit implies in particular that $\Delta_{\mathbf{A}}$ is injective. \square

The algebra \mathbf{A}_{∞} is also equipped with a kind of an antipode.

Proposition 4.5. *The prescription*

$$\kappa_{\mathbf{A}}(q_{ij}) = q_{ji}, \quad i, j \in \mathbb{N}$$

extends to a normal involutive $$ -antihomomorphism of \mathbf{A}_{∞} .*

Proof. View generators q_{ij} as the elements of the opposite von Neumann algebra A_∞^{op} and denote them by $\{q_{ij}^o : i, j \in \mathbb{N}\}$. Once again using the universality as in Corollary 4.3, it is easy to see that the map $q_{ij} \mapsto q_{ji}^o$ canonically induces a normal unital $*$ -homomorphism from A_∞ to A_∞^{op} , which can be viewed as a $*$ -antihomomorphism on A_∞ . \square

Let us now compare the construction above with that from the previous section. Recall the projective system $(W_n)_{n=1}^\infty$ of Hopf - von Neumann algebras introduced in Section 3. Let \mathfrak{W} denote the corresponding category of von Neumann algebras (as in Definition 1.1).

Proposition 4.6. *The algebra A_∞ of Corollary 4.3 is an element of \mathfrak{W} . Therefore W_∞ is a direct summand of A_∞ .*

Proof. Recall that $A_\infty \approx zB^{**}$ in the notation of Theorem 4.2. The universal property of B implies that for each $n \in \mathbb{N}$ there is a surjection $\gamma_n : B \rightarrow C(\mathbb{S}_n)$ defined by the formula

$$\gamma_n(b_{ij}) = \begin{cases} q_{ij}^{(n)} & i, j \leq n \\ 0 & \text{otherwise} \end{cases}.$$

Let $\psi_n = \gamma_n^{**}$ - it again becomes a surjection, this time onto $W_n = C(\mathbb{S}_n)^{**}$, and it is easy to check that $\psi_n = \phi_n \circ \psi_{n+1}$ for all $n \in \mathbb{N}$. Hence B^{**} is in the class \mathfrak{W} associated with the sequence $(W_n)_{n \in \mathbb{N}}$ according to Definition 1.1.

Define w_n to be the smallest central projection in B^{**} dominating all projections $(p_j^{(n)})^\perp$ and $(q_j^{(n)})^\perp$, where

$$p_j^{(n)} = \sum_{i=1}^n b_{ij}, \quad q_j^{(n)} = \sum_{i=1}^n b_{ji}.$$

Note that we can describe w_n in terms of the central supports of $(p_j^{(n)})^\perp$ and $(q_j^{(n)})^\perp$:

$$(4.4) \quad w_n = \bigvee_{j \in \mathbb{N}} z((p_j^{(n)})^\perp) \vee \bigvee_{j \in \mathbb{N}} z((q_j^{(n)})^\perp).$$

For that it suffices to note that a central projection dominates another, not necessarily central, projection if and only if it dominates its central carrier.

The argument similar to that of the proof of Theorem 4.2, exploiting the fact that $C(\mathbb{S}_n)^{**}$ can be described as the universal von Neumann algebra generated by an n by n magic unitary implies that $\psi_n : w_n^\perp B^{**} \rightarrow C(\mathbb{S}_n)^{**}$ is an isomorphism. Indeed, it is easy to see that for each $j \in \mathbb{N}$ there is $\psi_n(p_j^{(n)}) = \psi_n(q_j^{(n)}) = 1_{C(\mathbb{S}_n)^{**}}$, so that the projection determining the kernel of ψ_n dominates $z_n := w_n^\perp$ and $\psi_n(x) = \psi_n(z_n x)$ for all $x \in B^{**}$. Thus we obtain a surjective map $\psi_n|_{z_n B^{**}} \rightarrow C(\mathbb{S}_n)^{**}$ which preserves the natural magic unitaries in both algebras (observe that $\sum_{i=1}^n z_n b_{ij} = \sum_{i=1}^n z_n b_{ji} = z_n$). The afore-mentioned universality of $C(\mathbb{S}_n)^{**}$ implies that it is an isomorphism.

Hence $\text{Ker}(\psi_n)$ is equal to $w_n B^{**}$ and the intersection $\bigcap_{n \in \mathbb{N}} \text{Ker}(\psi_n)$ is equal to $w_\infty B^{**}$, where $w_\infty = \lim_{n \in \mathbb{N}} w_n$.

Recall that the central projection $w = z^\perp \in B^{**}$ was defined in the proof of Theorem 4.2 as the smallest central projection in B^{**} dominating all projections p_j^\perp and q_j^\perp , where $p_j = \lim_{n \in \mathbb{N}} p_j^{(n)}$ and $q_j = \lim_{n \in \mathbb{N}} q_j^{(n)}$. Hence it is easy to check that $z \geq z_\infty := w_\infty^\perp$ and in particular we can view zB^{**} as an element of \mathfrak{W} . \square

The inclusion $W_\infty \subset A_\infty$ is close to being an inclusion of Hopf - von Neumann algebras. This is formulated in the next proposition.

Proposition 4.7. *View W_∞ as a subalgebra of A_∞ , so that $W_\infty = z_\infty A_\infty$. The normal $*$ -homomorphism $\hat{\Delta} : W_\infty \rightarrow W_\infty \overline{\otimes} W_\infty$ defined by: $\hat{\Delta}(x) = (z_\infty \otimes z_\infty)(\Delta_A(x))$ ($x \in W_\infty$) is unital and coassociative. It in fact coincides with the coproduct on W_∞ constructed as a projective limit in Theorem 2.3.*

Proof. We use the notation of the last proposition. As $W_\infty = z_\infty A_\infty$, it is enough to show that $\Delta_A(z_\infty) \geq z_\infty \otimes z_\infty$, so that $\hat{\Delta} : W_\infty \rightarrow W_\infty \overline{\otimes} W_\infty$ satisfies the required conditions.

Note first that as $\text{Ker}(\psi_n) = w_n A_\infty$, we can check that

$$\text{Ker}(\psi_n \otimes \psi_n) = (z_n \otimes z_n)^\perp (A_\infty \overline{\otimes} A_\infty).$$

The construction of the coproduct on A_∞ implies that the maps $\psi_n : A_\infty \rightarrow C(\mathbb{S}_n)^{**}$ intertwine the respective coproducts (recall that $C(\mathbb{S}_n)^{**}$ has a canonical Hopf-von Neumann algebra structure induced from $C(\mathbb{S}_n)$). As we have $(\psi_n \otimes \psi_n)(\Delta_A(w_n)) = \Delta_n(\psi_n(w_n)) = 0$, the formula displayed above implies that the projection $\Delta_A(w_n)$ is dominated by $(z_n \otimes z_n)^\perp$. Passing to the limit (exploiting normality of the coproduct) we obtain that

$$\Delta_A(w_\infty) \leq (z_\infty \otimes z_\infty)^\perp,$$

so the proof of the first statement of the lemma is finished.

To show the second part, by the uniqueness in Theorem 2.3 it suffices to show that for each $n \in \mathbb{N}$ we have

$$\Delta_n \psi_n|_{W_\infty} = (\psi_n|_{W_\infty} \otimes \psi_n|_{W_\infty}) \hat{\Delta}.$$

In fact we can even show that

$$(4.5) \quad \Delta_n \psi_n = (\psi_n \otimes \psi_n) \Delta_A.$$

Indeed, as maps on both sides of the last equation are normal, it suffices to check they take the same values on each $z b_{ij}$ (where z is now a central projection in \mathbb{B}^{**} defined in Theorem 4.2). Fix then $i, j \in \mathbb{N}$:

$$\begin{aligned} (\psi_n \otimes \psi_n)(\Delta_A(z b_{ij})) &= (\psi_n \otimes \psi_n) \left(\lim_{k \rightarrow \infty} \sum_{l=1}^k z b_{il} \otimes z b_{lj} \right) \\ &= \lim_{k \rightarrow \infty} (\psi_n \otimes \psi_n) \left(\sum_{l=1}^k z b_{il} \otimes z b_{lj} \right) = \sum_{l=1}^n \psi_n(z b_{il}) \otimes \psi_n(z b_{lj}). \end{aligned}$$

Now it is easy to check that $\Delta_n(\psi_n(z b_{ij})) = (\psi_n \otimes \psi_n)(\Delta(z b_{ij}))$, considering separately two cases: first $i, j \leq n$ and then $\max\{i, j\} > n$. Thus (4.5) is proved. \square

Proposition 4.6 does not exclude the possibility of A_∞ actually coinciding with W_∞ , i.e. $z = z_\infty$. Below we show that this is not the case.

Lemma 4.8. *Let $z, z_\infty \in \mathcal{P}(\mathbb{B}^{**})$ be projections discussed in Proposition 4.6. Then $z \neq z_\infty$.*

Proof. Observe that another application of the argument used in Proposition 4.6 implies that

$$(4.6) \quad z^\perp = \bigvee_{j \in \mathbb{N}} z(p_j^\perp) \vee \bigvee_{j \in \mathbb{N}} z(q_j^\perp),$$

so the comparison of the formulas (4.4) and (4.6) shows that the problem of deciding whether $z = z_\infty$ is related to the fact that for a decreasing sequence of projections in a von Neumann algebra, say $(r_n)_{n=1}^\infty$ we can have $z(\lim_{n \in \mathbb{N}} r_n) \neq \lim_{n \in \mathbb{N}} z(r_n)$.

Suppose for the moment that there exists a non-zero normal representation $\pi : \mathbf{B}^{**} \rightarrow B(\mathfrak{h})$ such that $\pi(z) = 1_{B(\mathfrak{h})}$, $\mathbf{N} := \pi(\mathbf{B}^{**})$ is a factor, and if we write $d_{ij} = \pi(b_{ij})$ ($i, j \in \mathbb{N}$) then we have $r_k := \sum_{j=1}^k d_{1j} \neq 1_{B(\mathfrak{h})}$ for all $k \in \mathbb{N}$. Then $z(r_k^\perp) = 1_{\mathbf{N}} = 1_{B(\mathfrak{h})}$ (central carrier understood in \mathbf{N}). As $\pi : \mathbf{B}^{**} \rightarrow \mathbf{N}$ is onto (so in particular it maps $Z(\mathbf{B}^{**})$ into $Z(\mathbf{N})$), we have for each $p \in \mathcal{P}(\mathbf{B}^{**})$ the inequality $z(\pi(p)) \leq \pi(z(p))$. As $r_k^\perp = \pi((p_1^{(k)})^\perp)$, we have therefore (recall ((4.4))

$$\pi(z_k^\perp) \geq \pi(z((p_1^{(k)})^\perp)) \geq z(r_k^\perp) = 1_{B(\mathfrak{h})}.$$

Hence $\pi(z_k) = 0$ and thus also $\pi(z_\infty) = 0$, so z cannot be equal to z_∞ .

It remains to show that such a representation exists. It suffices to exhibit a concrete magic unitary $(d_{ij})_{i,j=1}^\infty$ built of projections on a Hilbert space \mathfrak{h} such that each row and column sum to $1_{B(\mathfrak{h})}$, $\sum_{j=1}^k d_{1j} < 1_{B(\mathfrak{h})}$ for each $k \in \mathbb{N}$ (in other words the first row is not ‘finitely supported’) and the entries generate $B(\mathfrak{h})$ as a von Neumann algebra. Let then $(d_n)_{n=1}^\infty$ be a sequence of non-zero mutually orthogonal projections summing to $1_{B(\mathfrak{h})}$ and consider the matrix:

$$\begin{bmatrix} d_1 & 0 & d_2 & d_3 & d_4 & \cdots \\ d_1^\perp & d_1 & 0 & 0 & 0 & \cdots \\ 0 & d_2 & d_2^\perp & 0 & 0 & \cdots \\ 0 & d_3 & 0 & d_3^\perp & 0 & \cdots \\ 0 & d_4 & 0 & 0 & d_4^\perp & \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}.$$

It is easy to see it gives a magic unitary with the first row ‘infinitely supported’. The generation condition can be achieved by considering a finite sequence of projections $(t_n)_{n=1}^k$ generating the whole $B(\mathfrak{h})$ and adding to a given magic unitary two by two blocks of the form $\begin{bmatrix} t_n & t_n^\perp \\ t_n^\perp & t_n \end{bmatrix}$ (with respective rows and columns completed by zeros). □

Corollary 4.9. W_∞ is a proper von Neumann subalgebra of A_∞ .

Acknowledgment. Part of the work on this paper was done during the visit of the second author to ISI Kolkata in April 2010 funded by the UKIERI project Quantum Probability, Noncommutative Geometry and Quantum Information.

REFERENCES

- [Ban₁] T. Banica, Representations of compact quantum groups and subfactors, *J. Reine Angew. Math.* **509** (1999), 167–198.
- [Ban₂] T. Banica, Quantum automorphism groups of small metric spaces, *Pacific J. Math.* **219** (2005), no. 1, 27–51.
- [BBC] T. Banica, J. Bichon and B. Collins, Quantum permutation groups: a survey, *Banach Center Publ.* **78** (2007), 13–34.
- [Bic] J. Bichon, Quantum automorphism groups of finite graphs, *Proc. Am. Math. Soc.* **131** (2003), no. 3, 665–673.
- [BSp] T. Banica and R. Speicher, Liberation of orthogonal Lie groups. *Adv. Math.* **222** (2009), no. 4, 1461–1501.

- [Cur] S. Curran, Quantum exchangeable sequences of algebras, *Indiana Univ. Math. J.* **58** (2009), no. 3, 1097–1125.
- [Gos] D. Goswami, Quantum Group of Isometries in Classical and Noncommutative Geometry, *Comm. Math. Phys.* **285** (2009), no. 1, 141–160.
- [KSp] C. Köstler and R. Speicher, A noncommutative de Finetti theorem: invariance under quantum permutations is equivalent to freeness with amalgamation, *Comm. Math. Phys.* **291** (2009), no. 2, 473–490.
- [KuV] J. Kustermans and S. Vaes, Locally compact quantum groups, *Ann. Sci. École Norm. Sup. (4)* **33** (2000) no. 6, 837–934.
- [MVD] A. Maes and A. Van Daele, Notes on compact quantum groups, *Nieuw Arch. Wisk(4)* **16** (1998), no. 1-2, 73–112.
- [Wan] S. Wang, Quantum symmetry groups of finite spaces, *Comm. Math. Phys.* **195** (1998), no. 1, 195–211.
- [Wor] S.L. Woronowicz, Compact quantum groups, in A. Connes, K. Gawedzki, and J. Zinn-Justin, editors, *Symétries Quantiques, Les Houches, Session LXIV, 1995*, 845–884.

STAT-MATH UNIT, INDIAN STATISTICAL INSTITUTE, 203, B. T. ROAD, KOLKATA 700 208, INDIA
E-mail address: `goswamid@isical.ac.in`

DEPARTMENT OF MATHEMATICS AND STATISTICS, LANCASTER UNIVERSITY, LANCASTER, LA1 4YF, UK
 INSTITUTE OF MATHEMATICS OF THE POLISH ACADEMY OF SCIENCES, UL.ŚNIADECKICH 8, 00-956
 WARSZAWA, POLAND
E-mail address: `a.skalski@lancaster.ac.uk`