# Hodograph solutions of the dispersionless coupled KdV hierarchies, critical points and the Euler-Poisson-Darboux equation 

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#### Abstract

It is shown that the hodograph solutions of the dispersionless coupled KdV (dcKdV) hierarchies describe critical and degenerate critical points of a scalar function which obeys the Euler-PoissonDarboux equation. Singular sectors of each dcKdV hierarchy are found to be described by solutions of higher genus dcKdV hierarchies. Concrete solutions exhibiting shock type singularities are presented.


Key words: Integrable systems. Hodograph equations. Euler-PoissonDarboux equation.
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## 1 Introduction

In the present paper we study hierarchies of hydrodynamical systems describing quasiclassical deformations of hyperelliptic curves $[1,2]$

$$
\begin{equation*}
p^{2}=u(\lambda), \quad u(\lambda):=\lambda^{m}-\sum_{i=0}^{m-1} \lambda^{i} u_{i}, \quad m \geq 1 \tag{1}
\end{equation*}
$$

These hierarchies are of interest for several reasons. First, there are hierarchies of important hydrodynamical type systems among them. For $m=1$ one has the Burgers-Hopf hierarchy [3, 4] associated with the dispersionless KdV equation $u_{t}=\frac{3}{2} u u_{x}$. For $m=2$ it is the hierarchy of higher equations for the 1-layer Benney system (classsical long wave equation)

$$
\left\{\begin{array}{l}
u_{t}+u u_{x}+v_{x}=0  \tag{2}\\
v_{t}+(u v)_{x}=0 .
\end{array}\right.
$$

The system (2) and the corresponding hierarchy are quasiclassical limits of the nonlinear Schrödinger (NLS) equation and the NLS hierarchy [5]. For $m \geq 3$ these hierarchies turn to describe the singular sectors of the above $m=1,2$ hierarchies [1].

Second, all these hierarchies are the dispersionless limits of integrable coupled KdV (cKdV) hierarchies [6]-[8] associated to Schrödinger spectral problems

$$
\begin{equation*}
\partial_{x x} \psi=v(\lambda, x) \psi, \tag{3}
\end{equation*}
$$

with potentials which are polynomials in the spectral parameter $\lambda$

$$
v(\lambda, x):=\lambda^{m}-\sum_{i=0}^{m-1} \lambda^{i} v_{i}(x) \quad m \geq 1
$$

The cKdV hierarchies have been studied in [6]-[8], they have bi-Hamiltonian structures and, as a consequence of this property, the dispersionless expansions of their solutions possess interesting features such as the quasi-triviality property [9]-[10]. Moreover, the cKdV hierarchies arise also in the study of the singular sectors of the KdV and AKNS hierarchies [11, 12]. Henceforth we will refer to the hierarchies of hydrodynamical systems associated with the curves (1) for a fixed $m$ as the $m$-th dispersionless coupled $\mathrm{KdV}\left(\mathrm{dcKdV}_{m}\right)$ hierarchies. The Hamiltonian structures of the $\mathrm{dcKdV}_{m}$ hierarchies have been studied in [13]. At last, it should be noticed that the $\mathrm{dcKdV}_{m}$ hierarchies are closely connected with the higher genus Whitham hierarchies introduced in [14].

In our analysis of the hodograph equations for the $\mathrm{dcKdV}_{m}$ hierarchies we use Riemann invariants $\beta_{i}$ (roots of the polynomial $u(\lambda)$ in (1)) which provide a specially convenient system of coordinates. We show that the $\mathrm{dcKdV}_{m}$ hodograph equations have the form

$$
\begin{equation*}
\frac{\partial W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i}}=0, \quad i=1, \ldots, m \tag{4}
\end{equation*}
$$

where $\boldsymbol{t}=\left(t_{1}, t_{2}, \ldots\right)$ are times of the hierarchy and

$$
\begin{equation*}
W_{m}(\boldsymbol{t}, \boldsymbol{\beta}):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{\sum_{n \geq 0} t_{n} \lambda^{n}}{\sqrt{\prod_{i=1}^{m}\left(1-\beta_{i} / \lambda\right)}} . \tag{5}
\end{equation*}
$$

Here $\gamma$ denotes a large positively oriented circle $|\lambda|=r$. Thus, the hodograph solutions of the $\mathrm{dcKdV}_{m}$ hierarchies describe critical points of the functions $W_{m}(\boldsymbol{t}, \boldsymbol{\beta})$. These functions turn to be very special as they satisfy a well-known system of equations in differential geometry: the Euler-Poisson-Darboux (EPD) equations [15]

$$
\begin{equation*}
2\left(\beta_{i}-\beta_{j}\right) \frac{\partial^{2} W_{m}}{\partial \beta_{i} \partial \beta_{j}}=\frac{\partial W_{m}}{\partial \beta_{i}}-\frac{\partial W_{m}}{\partial \beta_{j}} . \tag{6}
\end{equation*}
$$

The system (6) has also appeared in the theory of the Whitham equations arising in the small dispersion limit of the KdV equations [17]-[19], and in the theory of hydrodynamic chains [20].

We also study the singular sectors $\mathcal{M}_{m}^{\text {sing }}$ of the spaces of hodograph solutions for the $\mathrm{dcKdV}_{m}$ hierarchies. They are given by the points $(\boldsymbol{t}, \boldsymbol{\beta})$ such that

$$
\begin{equation*}
\operatorname{rank}\left(\frac{\partial^{2} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i} \partial \beta_{j}}\right)<m \tag{7}
\end{equation*}
$$

The varieties $\mathcal{M}_{m}^{\text {sing }}$ provide us with special classes of degenerate critical points of the function $W_{m}$ within the general theory of critical points developed by V. I. Arnold and others about fourty years ago [23,24]. The use of equations (4)-(6) simplify drastically the analysis of the structure of these singular sectors. In particular, we prove that there is a nested sequence of subvarieties

$$
\begin{equation*}
\mathcal{M}_{m}^{\operatorname{sing}} \supset \mathcal{M}_{m, 1}^{\operatorname{sing}} \supset \mathcal{M}_{m, 2}^{\operatorname{sing}} \supset \cdots \mathcal{M}_{m, q}^{\operatorname{sing}} \supset \cdots \tag{8}
\end{equation*}
$$

which represents subsets of the singular sector $\mathcal{M}_{m}^{\operatorname{sing}}$ of the $\mathrm{dcKdV}{ }_{m}$ hierarchy with increasing singular degree $q$, such that each $\mathcal{M}_{m, q}^{\operatorname{sing}}$ is determined by a class of hodograph solutions of the dcKdV ${ }_{m+2 q}$ hierarchy.

The paper is organized as follows. The dcKdV ${ }_{m}$ hierarchies are described in Section 2. Equations (4)-(6) are derived in Section 3. Section 4 deals with the analysis of the singular sectors of the $\mathrm{dcKaV}_{m}$ hierarchies in terms of their associated hodograph equations. The relation between singular points of the $d c K d V_{m}$ hodograph equations and solutions of higher dcKdV ${ }_{m+2 q}$ hodograph equations is stated in Section 4. Some concrete examples involving shock singularities of the Burgers-Hopf equation and the 1-layer Benney system are presented in Section 5.

## 2 The dcKdV ${ }_{m}$ hierarchies

Given a positive integer $m \geq 1$ we consider the set $M_{m}$ of algebraic curves (1). For $m=2 g+1$ (odd case) and $m=2 g+2$ (even case) these curves are, generically, hyperelliptic Riemann surfaces of genus $g$. We will denote by $\boldsymbol{q}=\left(q_{1}, \ldots, q_{m}\right)$ any of the two sets of parameters $\boldsymbol{u}:=\left(u_{0}, \ldots, u_{m-1}\right)$ or $\boldsymbol{\beta}:=\left(\beta_{1}, \ldots, \beta_{m}\right)$ which determine the curves (1)

$$
\begin{equation*}
u(\lambda)=\lambda^{m}-\sum_{i=0}^{m-1} \lambda^{i} u_{i}=\prod_{i=1}^{m}\left(\lambda-\beta_{i}\right) \tag{9}
\end{equation*}
$$

Obviously, for any fixed $\boldsymbol{\beta}$ all the permutations $\sigma(\boldsymbol{\beta}):=\left(\beta_{\sigma(1)}, \ldots, \beta_{\sigma(m)}\right)$ represent the same element of $M_{m}$. Note also that

$$
\begin{equation*}
u_{i}=(-1)^{m-i-1} \mathrm{~s}_{m-i}(\boldsymbol{\beta}), \tag{10}
\end{equation*}
$$

where $s_{k}$ are the elementary symmetric polynomials

$$
\mathrm{s}_{k}=\sum_{1 \leq i_{1}<\ldots<i_{k} \leq m} \beta_{i_{1}} \cdots \beta_{i_{k}}
$$

We next introduce the $\mathrm{dcKdV}_{m}$ hierarchy as a particular systems of commuting flows

$$
\boldsymbol{q}(\boldsymbol{t}), \quad \boldsymbol{t}:=\left(x:=t_{0}, t_{1}, t_{2}, \ldots\right)
$$

on $M_{m}$. In order to define these flows we use the set $\mathcal{L}$ of formal power series

$$
f(z)=\sum_{n=-\infty}^{+\infty} c_{n} z^{n}
$$

where

$$
z:=\lambda^{1 / 2} \text { for } m=2 g+1 ; \quad z:=\lambda \text { for } m=2 g+2
$$

For any given $m \geq 1$ a distinguished element of $\mathcal{L}$ is provided by the branch of $p=\sqrt{u(\lambda)}$ such that as $z \rightarrow \infty$ has an expansion of the form

$$
\begin{cases}p(z, \boldsymbol{q})=z^{2 g+1}\left(1+\sum_{n \geq 1} \frac{b_{n}(\boldsymbol{q})}{z^{2 n}}\right), & m=2 g+1  \tag{11}\\ p(z, \boldsymbol{q})=z^{g+1}\left(1+\sum_{n \geq 1} \frac{b_{n}(\boldsymbol{q})}{z^{n}}\right), & m=2 g+2\end{cases}
$$

We define the following splittings $\mathcal{L}=\mathcal{L}_{(+, \boldsymbol{q})} \bigoplus \mathcal{L}_{(-, \boldsymbol{q})}$

$$
\begin{equation*}
f_{(+, \boldsymbol{q})}(z):=\left(\frac{f(z)}{p(z, \boldsymbol{q})}\right)_{\oplus} p(z, \boldsymbol{q}), \quad f_{(-, \boldsymbol{q})}(z):=\left(\frac{f(z)}{p(z, \boldsymbol{q})}\right)_{\ominus} p(z, \boldsymbol{q}) \tag{12}
\end{equation*}
$$

where $f_{\oplus}$ and $f_{\ominus}$ stand for the standard projections on positive and strictly negative powers of $z$, respectively

$$
f_{\oplus}(z):=\sum_{n=0}^{N} c_{n} z^{n}, \quad f_{\ominus}(z):=\sum_{n=-\infty}^{-1} c_{n} z^{n}
$$

The dcKdV $V_{m}$ flows $\boldsymbol{q}(\boldsymbol{t})$ are characterized by the following condition: There exists a family of functions $S(z, \boldsymbol{t}, \boldsymbol{q}(\boldsymbol{t}))$ in $\mathcal{L}$ satisfying

$$
\begin{equation*}
\partial_{t_{n}} S(z, \boldsymbol{t}, \boldsymbol{q}(\boldsymbol{t}))=\Omega_{n}(z, \boldsymbol{q}(\boldsymbol{t})), \quad n \geq 0 \tag{13}
\end{equation*}
$$

where

$$
\Omega_{n}(z, \boldsymbol{q}):=\left(\lambda(z)^{n+m / 2}\right)_{(+, \boldsymbol{q})}=\left\{\begin{array}{ll}
\left(z^{2 n+2 g+1}\right)_{(+, \boldsymbol{q})}, \quad m=2 g+1  \tag{14}\\
\left(z^{n+g+1}\right)_{(+, \boldsymbol{q})}, & m=2 g+2,
\end{array} \quad n \geq 0\right.
$$

We notice that

$$
\begin{equation*}
\Omega_{n}(z, \boldsymbol{q})=\left(\lambda^{n} R(\lambda(z), \boldsymbol{q})\right)_{\oplus} p \tag{15}
\end{equation*}
$$

where $R$ is the generating function

$$
\begin{equation*}
R(\lambda, \boldsymbol{q}):=\sqrt{\frac{\lambda^{m}}{u(\lambda)}}=\sum_{n \geq 0} \frac{R_{n}(\boldsymbol{q})}{\lambda^{n}}, \quad \lambda \rightarrow \infty . \tag{16}
\end{equation*}
$$

The coefficients $R_{n}(\boldsymbol{q})$ are polynomials in the coordinates $\boldsymbol{q}$, for example

$$
R_{0}=1, \quad R_{1}=\frac{1}{2} u_{m-1}, \quad R_{2}=\frac{1}{2} u_{m-2}+\frac{3}{8} u_{m-1}^{2}, \quad \ldots
$$

Functions $S$ which satisfy (13) will be referred to as action functions of the $\mathrm{dcKdV}_{m}$ hierarchy. This kind of generating functions $S$ has been already used in the theory of dispersionless integrable systems (see e.g. [14]). It can be proved [1] that (13) is a compatible system of equations for $S$. In fact its general solution will be determined in the next section. We notice that for $n=0$ the equation (13) reads

$$
\begin{equation*}
\partial_{x} S(z, \boldsymbol{t}, \boldsymbol{q}(\boldsymbol{t}))=p(z, \boldsymbol{q}(\boldsymbol{t})), \tag{17}
\end{equation*}
$$

so that (13) is equivalent to the system

$$
\begin{equation*}
\partial_{t_{n}} p(z, \boldsymbol{q}(t))=\partial_{x} \Omega_{n}(z, \boldsymbol{q}(\boldsymbol{t})), \quad n \geq 0 . \tag{18}
\end{equation*}
$$

We will henceforth refer to the $\mathrm{dcKdV}_{m}$ hierarchy for $m=2 g+1$ and $m=2 g+2$ as the BurgersHopf $\left(\mathrm{BH}_{g}\right)$ and the dispersionless Jaulent-Miodek $\left(\mathrm{dJM}_{g}\right)$ hierarchies, respectively. Observe that both hierarchies, $\mathrm{BH}_{g}$ and $\mathrm{dJM}_{g}$ determine deformations of hyperelliptic Riemann surfaces of genus $g$. In our work we will always consider an arbitrary but finite number of these flows.

Since $u=u(\lambda(z), \boldsymbol{q})=p(z, \boldsymbol{q})^{2}$, the operator $J=J(\lambda, u)$ defined by

$$
\begin{aligned}
& J:=2 p \cdot \partial_{x} \cdot p=2 u \partial_{x}+u_{x}, \\
& J=\sum_{i=0}^{m} \lambda^{i} J_{i}, \quad J_{m}=2 \partial_{x}, \quad J_{i}=-\left(2 u_{i} \partial_{x}+u_{i, x}\right), \quad u_{m}:=-1,
\end{aligned}
$$

satisfies $J R=0$. Then from (18) it follows that

$$
\begin{equation*}
\partial_{n} u=J\left(\lambda^{n} R(\lambda, \boldsymbol{u})\right)_{\oplus}=-J\left(\lambda^{n} R(\lambda, \boldsymbol{u})\right)_{\ominus} \tag{19}
\end{equation*}
$$

which constitutes the $\mathrm{dcKdV}_{m}$ hierarchy in terms of the coordinates $u_{i}$

$$
\begin{equation*}
\partial_{n} u_{i}=\sum_{l-k=i, k \geq 1} J_{l} R_{n+k}(\boldsymbol{u}), \quad i=0, \ldots, m-1 . \tag{20}
\end{equation*}
$$

From (18) it also follows that

$$
\partial_{t_{n}} \log p(z, \boldsymbol{q})=\frac{\partial_{x}\left[\left(\lambda(z)^{n} R(\lambda(z), \boldsymbol{q})\right)_{\oplus} p\right]}{p(z, \boldsymbol{q})}
$$

and then, identifying the residues of both sides at $\lambda=\beta_{i}$, we get

$$
\begin{equation*}
\partial_{n} \beta_{i}=\omega_{n, i}(\boldsymbol{\beta}) \partial_{x} \beta_{i}, \quad i=1, \ldots, m \tag{21}
\end{equation*}
$$

where

$$
\begin{equation*}
\omega_{n, i}(\boldsymbol{\beta}):=\left.\left(\lambda^{n} R(\lambda, \boldsymbol{\beta})\right)_{\oplus}\right|_{\lambda=\beta_{i}} . \tag{22}
\end{equation*}
$$

The systems (21) are the equations of the $\mathrm{dcKdV}_{m}$ hierarchy in terms of the coordinates $\beta_{i}$. Observe that we have two $\mathrm{dcKdV}{ }_{m}$ hierarchies, $\mathrm{BH}_{g}$ and $\mathrm{dJM}_{g}$, which determine deformations of hyperelliptic Riemann surfaces of genus $g$. It can be shown $[2,13]$ that the $\mathrm{dcKdV}_{m}$ flows are bi-Hamiltonian systems.

We next present some examples of interesting flows in the $\mathrm{dcKdV}_{m}$ hierarchies. The dcKdV ${ }_{1}$ hierarchy is associated to the curve

$$
p^{2}-u(\lambda)=0, \quad u(\lambda)=\lambda-v, \quad v:=u_{0}=\beta_{1} .
$$

The corresponding flows are given by

$$
\partial_{t_{n}} v=c_{n} v^{n} v_{x}, \quad c_{n}:=\frac{(2 n+1)!!}{2^{n} n!}, \quad n \geq 1
$$

and constitute the Burgers-Hopf hierarchy $\mathrm{BH}_{0}$. In particular the $t_{1}$-flow is the Burgers-Hopf equation

$$
\partial_{t} v=\frac{3}{2} v v_{x}
$$

which is in turn the dispersionless limit of the KdV equation.
The $\mathrm{dcKdV}_{2}\left(\mathrm{dJM}_{0}\right)$ hierarchy is associated to the curve

$$
\begin{gathered}
p^{2}-u(\lambda)=0, \quad u(\lambda)=\lambda^{2}-\lambda u_{1}-u_{0}=\left(\lambda-\beta_{1}\right)\left(\lambda-\beta_{2}\right), \\
u_{1}=\beta_{1}+\beta_{2}, \quad u_{0}=-\beta_{1} \beta_{2} .
\end{gathered}
$$

The $t_{1}$-flow of this hierarchy is given by the disperssionless Jaulent-Miodek system

$$
\left\{\begin{array}{l}
\partial_{t_{1}} u_{0}=u_{0} u_{1 x}+\frac{1}{2} u_{1} u_{0 x},  \tag{23}\\
\partial_{t_{1}} u_{1}=u_{0 x}+\frac{3}{2} u_{1} u_{1 x}
\end{array}\right.
$$

which under the changes of dependent variables

$$
u=-u_{1}, \quad v=u_{0}+\frac{u_{1}^{2}}{4}
$$

becomes the 1-layer Benney system (2). In terms of the Riemann invariants $\beta_{1}$ and $\beta_{2}$

$$
u=-\left(\beta_{1}+\beta_{2}\right), \quad v=\left(\beta_{1}-\beta_{2}\right)^{2} / 4
$$

the system (2) takes the well-known form

$$
\left\{\begin{array}{l}
\partial_{t_{1}} \beta_{1}=\frac{1}{2}\left(3 \beta_{1}+\beta_{2}\right) \beta_{1 x}  \tag{24}\\
\partial_{t_{1}} \beta_{2}=\frac{1}{2}\left(3 \beta_{2}+\beta_{1}\right) \beta_{2 x}
\end{array}\right.
$$

For $v>0$ the 1-layer Benney system is hyperbolic while for $v<0$ it is elliptic.
Finally, we consider the $\mathrm{BH}_{1}$ hierarchy. Its associated curve is given by

$$
\begin{aligned}
& p^{2}-u(\lambda)=0, \quad u(\lambda)=\lambda^{3}-\lambda^{2} u_{2}-\lambda u_{1}-u_{0}=\left(\lambda-\beta_{1}\right)\left(\lambda-\beta_{2}\right)\left(\lambda-\beta_{3}\right), \\
& u_{1}=\beta_{1}+\beta_{2}+\beta_{3}, \quad u_{2}=-\left(\beta_{1} \beta_{2}+\beta_{1} \beta_{3}+\beta_{2} \beta_{3}\right), \quad u_{3}=\beta_{1} \beta_{2} \beta_{3} .
\end{aligned}
$$

The first flow takes the forms

$$
\left\{\begin{array} { l } 
{ \partial _ { t _ { 1 } } u _ { 0 } = \frac { 1 } { 2 } u _ { 2 } u _ { 0 x } + u _ { 0 } u _ { 2 x } , }  \tag{25}\\
{ \partial _ { t _ { 1 } } u _ { 1 } = u _ { 0 x } + \frac { 1 } { 2 } u _ { 2 } u _ { 1 x } + u _ { 1 } u _ { 2 x } , } \\
{ \partial _ { t _ { 1 } } u _ { 2 } = u _ { 1 x } + \frac { 3 } { 2 } u _ { 2 } u _ { 2 x } }
\end{array} \Longleftrightarrow \left\{\begin{array}{l}
\partial_{t_{1}} \beta_{1}=\frac{1}{2}\left(3 \beta_{1}+\beta_{2}+\beta_{3}\right) \beta_{1 x} \\
\partial_{t_{1}} \beta_{2}=\frac{1}{2}\left(\beta_{1}+3 \beta_{2}+\beta_{3}\right) \beta_{2 x} \\
\partial_{t_{1}} \beta_{3}=\frac{1}{2}\left(\beta_{1}+\beta_{2}+3 \beta_{3}\right) \beta_{3 x}
\end{array}\right.\right.
$$

## 3 Hodograph equations for $\mathrm{dcKdV}_{m}$ hierarchies and the Euler-Poisson-Darboux equation

Let us introduce the function

$$
\begin{equation*}
W_{m}(\boldsymbol{t}, \boldsymbol{q}):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} U(\lambda, \boldsymbol{t}) R(\lambda, \boldsymbol{q})=\sum_{n \geq 0} t_{n} R_{n+1}(\boldsymbol{q}) \tag{26}
\end{equation*}
$$

where $\gamma$ denotes a large positively oriented circle $|\lambda|=r, U(\lambda, \boldsymbol{t}):=\sum_{n \geq 0} t_{n} \lambda^{n}$ and $R(\lambda, \boldsymbol{q})$ is the function defined in (16).

Theorem 1. If the functions $\boldsymbol{q}(\boldsymbol{t})=\left(q_{1}\left(\boldsymbol{t}, \ldots, q_{m}(\boldsymbol{t})\right)\right.$ satisfy the system of hodograph equations

$$
\begin{equation*}
\frac{\partial W_{m}(\boldsymbol{t}, \boldsymbol{q})}{\partial q_{i}}=0, \quad i=1, \ldots, m \tag{27}
\end{equation*}
$$

then $\boldsymbol{q}(\boldsymbol{t})$ is a solution of the dcKd $V_{m}$ hierarchy.
Proof. We are going to prove that the function

$$
\begin{equation*}
S(z, \boldsymbol{t}, \boldsymbol{q}(\boldsymbol{t}))=\sum_{n \geq 0} t_{n} \Omega_{n}(z, \boldsymbol{q}(\boldsymbol{t}))=(U(\lambda(z), \boldsymbol{t}) R(\lambda(z), \boldsymbol{q}(\boldsymbol{t})))_{\oplus} p(z, \boldsymbol{q}(\boldsymbol{t})), \tag{28}
\end{equation*}
$$

is an action function for the $d c \mathrm{KdV} V_{m}$ hierarchy. By differentiating (28) with respect to $t_{n}$ we have that

$$
\begin{equation*}
\partial_{n} S=\Omega_{n}+\left(U \partial_{n} R\right)_{\oplus} p+(U R)_{\oplus} \partial_{n} p, \tag{29}
\end{equation*}
$$

We now use the coordinates $\boldsymbol{\beta}=\left(\beta_{1}, \ldots, \beta_{m}\right)$ so that we may take advantage of the identities

$$
\begin{equation*}
\partial_{\beta_{i}} p=-\frac{1}{2} \frac{p}{\lambda-\beta_{i}}, \quad \partial_{\beta_{i}} R=\frac{1}{2} \frac{R}{\lambda-\beta_{i}} . \tag{30}
\end{equation*}
$$

Thus we deduce that

$$
\begin{equation*}
\left(U \partial_{n} R\right)_{\oplus} p+(U R)_{\oplus} \partial_{n} p=\frac{1}{2} \sum_{i=1}^{m}\left[\left(\frac{U R}{\lambda-\beta_{i}}\right)_{\oplus}-\frac{(U R)_{\oplus}}{\lambda-\beta_{i}}\right] p \partial_{n} \beta_{i} . \tag{31}
\end{equation*}
$$

On the other hand

$$
\begin{equation*}
\frac{\partial W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i}}=\frac{1}{2} \oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U(\lambda, \boldsymbol{t}) R(\lambda, \boldsymbol{\beta})}{\lambda-\beta_{i}}=\frac{1}{2} \oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{(U(\lambda, \boldsymbol{t}) R(\lambda, \boldsymbol{\beta}))_{\oplus}}{\lambda-\beta_{i}} . \tag{32}
\end{equation*}
$$

Hence the hodograph equations (27) can be written as

$$
\begin{equation*}
\left.(U(\lambda, \boldsymbol{t}) R(\lambda, \boldsymbol{\beta}(\boldsymbol{t})))_{\oplus}\right|_{\lambda=\beta_{i}}=0, \quad i=1, \ldots, m . \tag{33}
\end{equation*}
$$

Thus we have that $\left(U(\lambda, \boldsymbol{t}) R(\lambda, \boldsymbol{\beta}(\boldsymbol{t}))_{\oplus}\right.$ is a polynomial in $\lambda$ which vanish at $\lambda=\beta_{i}(\boldsymbol{t})$ for all $i$. As a consequence

$$
\frac{(U R)_{\oplus}}{\lambda-\beta_{i}}=\left(\frac{(U R)_{\oplus}}{\lambda-\beta_{i}}\right)_{\oplus}=\left(\frac{U R}{\lambda-\beta_{i}}\right)_{\oplus} .
$$

Then from (29) and (31) we deduce that $\partial_{n} S=\Omega_{n}$ and therefore the statement follows.

Using (26) we obtain that the hodograph equations (27) can be expressed as

$$
\begin{equation*}
\sum_{n \geq 0} t_{n} \frac{\partial R_{n+1}(\boldsymbol{q})}{\partial q_{i}}=0, \quad i=1, \ldots, m \tag{34}
\end{equation*}
$$

Furthermore, from (21), (22) and (33) the hodograph equations (27) can be also written as [1]

$$
\begin{equation*}
\sum_{n \geq 0} t_{n} \omega_{n, i}(\boldsymbol{\beta})=0, \quad i=1, \ldots, m, \tag{35}
\end{equation*}
$$

which represent the hodograph transform for the $\mathrm{dcKdV}_{m}$ hierarchy of flows in hydrodynamic form.

Notice also that we may shift the time parameters $t_{n} \rightarrow t_{n}-c_{n}$ in (34) to get solutions depending on an arbitrary number of constants.

It is easy to see that the generating function

$$
R(\lambda, \boldsymbol{\beta}):=\sqrt{\frac{\lambda^{m}}{u(\lambda)}}=\sqrt{\frac{\lambda^{m}}{\prod_{i=1}^{m}\left(\lambda-\beta_{i}\right)}},
$$

is a symmetric solution of the EPD equation

$$
\begin{equation*}
2\left(\beta_{i}-\beta_{j}\right) \frac{\partial^{2} R}{\partial \beta_{i} \partial \beta_{j}}=\frac{\partial R}{\partial \beta_{i}}-\frac{\partial R}{\partial \beta_{j}} . \tag{36}
\end{equation*}
$$

Consequently, the same property is satisfied by $W(\boldsymbol{t}, \boldsymbol{\beta})$ for all $\boldsymbol{t}$. Thus, we have proved

Theorem 2. The solutions $(\boldsymbol{t}, \boldsymbol{\beta})$ of the hodograph equations

$$
\begin{equation*}
\frac{\partial W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i}}=0, \quad i=1, \ldots, m \tag{37}
\end{equation*}
$$

are the critical points of the solution

$$
W_{m}(\boldsymbol{t}, \boldsymbol{\beta}):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U(\lambda, \boldsymbol{t})}{\sqrt{\prod_{i=1}^{m}\left(1-\beta_{i} / \lambda\right)}}
$$

of the EPD equation

$$
\begin{equation*}
2\left(\beta_{i}-\beta_{j}\right) \frac{\partial^{2} W_{m}}{\partial \beta_{i} \partial \beta_{j}}=\frac{\partial W_{m}}{\partial \beta_{i}}-\frac{\partial W_{m}}{\partial \beta_{j}} \tag{38}
\end{equation*}
$$

Let us denote by $\mathcal{M}_{m}$ the variety of points $(\boldsymbol{t}, \boldsymbol{\beta}) \in \mathbb{C}^{\infty} \times \mathbb{C}^{m}$ which satisfy the hodograph equations (37). From (32) it is clear that for any permutation $\sigma$ of $\{1, \ldots, m\}$ the functions

$$
\begin{equation*}
F_{i}(\boldsymbol{t}, \boldsymbol{\beta}):=\frac{\partial W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i}} \tag{39}
\end{equation*}
$$

satisfy

$$
\begin{equation*}
F_{i}(\boldsymbol{t}, \sigma(\boldsymbol{\beta}))=F_{\sigma(i)}(\boldsymbol{t}, \boldsymbol{\beta}) \tag{40}
\end{equation*}
$$

Then, it is clear that $\mathcal{M}_{m}$ is invariant under the action of the group of permutations

$$
(\boldsymbol{t}, \boldsymbol{\beta}) \in \mathcal{M}_{m} \Longrightarrow(\boldsymbol{t}, \sigma(\boldsymbol{\beta})) \in \mathcal{M}_{m}
$$

If $(\boldsymbol{t}, \boldsymbol{\beta})$ is a solution of (37) such that $\beta_{i} \neq \beta_{j}$ for all $i \neq j$ then it will be called an unreduced solution of (37). In this case the EPD equation (38) implies that

$$
\begin{equation*}
\frac{\partial^{2} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i} \partial \beta_{j}}=0, \quad \forall i \neq j \tag{41}
\end{equation*}
$$

Given $2 \leq r \leq m$, a solution $(\boldsymbol{t}, \boldsymbol{\beta})$ of (37) such that exactly $r$ of its components are equal will be called a $r$-reduced solution of (37).

The formulation (27) of the hodograph equations for the $\mathrm{dcK}_{\mathrm{C}} \mathrm{V}_{m}$ hierarchies allows us to apply the theory of critical points of functions to analyze the solutions of these hierarchies, while (38) indicates that the functions $W_{m}$ are of a very special class.

The EPD equation (38) arose in the study of cyclids [15], where solutions $W$ of the above form have been found too. Much later it appeared in the theory of Whitham equations describing the small dispersion limit of the KdV equation $[17,19]$.

We note that hodograph equations of a form close to (27) have been presented in [20] and [22]. Furthermore, linear equations of the EPD type and their connection with hydrodynamic chains have been studied in [21] too.

Finally, we emphasize that the functions $W_{m}$ depend on the parameters $t_{1}, t_{2}, \ldots$ (times of the hierarchy). Since "degenerate critical points appear naturally in cases when the functions depend on parameters " $[23,24]$, one should expect the existence of families of degenerate critical points for the functions $W_{m}$. Their connection with the singular sectors in the spaces of solutions for dcKdV ${ }_{m}$ will be considered in the next section.

To illustrate the statements given above we next present some simple examples. For the dcKdV ${ }_{2}$ hierarchy we have

$$
\begin{aligned}
W_{2}(\boldsymbol{t}, \boldsymbol{\beta}) & =\frac{x}{2}\left(\beta_{1}+\beta_{2}\right)+\frac{t_{1}}{8}\left(3 \beta_{1}^{2}+2 \beta_{1} \beta_{2}+3 \beta_{2}^{2}\right)+\frac{t_{2}}{16}\left(5 \beta_{1}^{3}+3 \beta_{1}^{2} \beta_{2}+3 \beta_{1} \beta_{2}^{2}+5 \beta_{2}^{3}\right) \\
& +\frac{t_{3}}{128}\left(35 \beta_{1}^{4}+20 \beta_{1}^{3} \beta_{2}+18 \beta_{1}^{2} \beta_{2}^{2}+20 \beta_{1} \beta_{2}^{3}+35 \beta_{2}^{4}\right)+\cdots
\end{aligned}
$$

The hodograph equations with $t_{n}=0$ for $n \geq 4$, take the form

$$
\left\{\begin{array}{l}
8 x+4 t_{1}\left(3 \beta_{1}+\beta_{2}\right)+3 t_{2}\left(5 \beta_{1}^{2}+2 \beta_{1} \beta_{2}+\beta_{2}^{2}\right)+\frac{t_{3}}{8}\left(140 \beta_{1}^{3}+60 \beta_{1}^{2} \beta_{2}+36 \beta_{1} \beta_{2}^{2}+20 \beta_{2}^{3}\right)=0  \tag{42}\\
8 x+4 t_{1}\left(\beta_{1}+3 \beta_{2}\right)+3 t_{2}\left(\beta_{1}^{2}+2 \beta_{1} \beta_{2}+5 \beta_{2}^{2}\right)+\frac{t_{3}}{8}\left(140 \beta_{2}^{3}+60 \beta_{2}^{2} \beta_{1}+36 \beta_{2} \beta_{1}^{2}+20 \beta_{1}^{3}\right)=0
\end{array}\right.
$$

For the $\mathrm{dcKdV}_{3}$ hierarchy we have

$$
\begin{aligned}
W_{3}(\boldsymbol{t}, \boldsymbol{\beta})= & \frac{x}{2}\left(\beta_{1}+\beta_{2}+\beta_{3}\right)+\frac{t_{1}}{8}\left(3 \beta_{1}^{2}+3 \beta_{2}^{2}+3 \beta_{3}^{2}+2 \beta_{1} \beta_{2}+2 \beta_{1} \beta_{3}+2 \beta_{2} \beta_{3}\right) \\
+ & \frac{t_{2}}{16}\left(5 \beta_{1}^{3}++5 \beta_{2}^{3}+5 \beta_{3}^{3}+3 \beta_{1}^{2} \beta_{2}+3 \beta_{1}^{2} \beta_{3}+3 \beta_{1} \beta_{2}^{2}+3 \beta_{2}^{2} \beta_{3}+3 \beta_{1} \beta_{3}^{2}\right. \\
& \left.+3 \beta_{2} \beta_{3}^{2}+2 \beta_{1} \beta_{2} \beta_{2}\right)+\cdots
\end{aligned}
$$

The hodograph equations with $t_{n}=0$ for $n \geq 3$ are

$$
\left\{\begin{array}{l}
8 x+4 t_{1}\left(3 \beta_{1}+\beta_{2}+\beta_{3}\right)+t_{2}\left(15 \beta_{1}^{2}+3 \beta_{2}^{2}+3 \beta_{3}^{2}+6 \beta_{1} \beta_{2}+6 \beta_{1} \beta_{3}+2 \beta_{2} \beta_{3}\right)=0  \tag{43}\\
8 x+4 t_{1}\left(\beta_{1}+3 \beta_{2}+\beta_{3}\right)+t_{2}\left(3 \beta_{1}^{2}+15 \beta_{2}^{2}+3 \beta_{3}^{2}+6 \beta_{1} \beta_{2}+2 \beta_{1} \beta_{3}+6 \beta_{2} \beta_{3}\right)=0 \\
8 x+4 t_{1}\left(\beta_{1}+\beta_{2}+3 \beta_{3}\right)+t_{2}\left(3 \beta_{1}^{2}+3 \beta_{2}^{2}+15 \beta_{3}^{2}+2 \beta_{1} \beta_{2}+6 \beta_{1} \beta_{3}+6 \beta_{2} \beta_{3}\right)=0
\end{array}\right.
$$

## 4 Singular sectors of $\mathrm{dcKdV}_{m}$ hierarchies

We say that $(\boldsymbol{t}, \boldsymbol{\beta}) \in \mathcal{M}_{m}$ is a regular point if it is a nondegenerate critical point of the function $W_{m}$. That it is to say, if it satisfies [23, 24]

$$
\begin{equation*}
\operatorname{det}\left(\frac{\partial^{2} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i} \partial \beta_{j}}\right) \neq 0 \tag{44}
\end{equation*}
$$

The set of regular points of $\mathcal{M}_{m}$ will be denoted by $\mathcal{M}_{m}^{\text {reg }}$ and the points of its complementary set $\mathcal{M}_{m}^{\operatorname{sing}}:=\mathcal{M}_{m}-\mathcal{M}_{m}^{\text {reg }}$, where the second differential of $W_{m}$ is a degenerate quadratic form, will be called singular points. We will also refer to $\mathcal{M}_{m}^{\text {reg }}$ and $\mathcal{M}_{m}^{\text {sing }}$ as the regular and singular sectors of the $\mathrm{dcKdV}_{m}$ hierarchy. So $\mathcal{M}_{m}^{\text {sing }}$ describes families of degenerate critical points of the
function $W_{m}$. Near a regular point the variety $\mathcal{M}_{m}^{\text {reg }}$ can be uniquely described as $(\boldsymbol{t}, \boldsymbol{\beta}(\boldsymbol{t})$ ) where $\boldsymbol{\beta}(\boldsymbol{t})$ is a solution of the $\mathrm{dcKdV}_{m}$ hierarchy.

The aim of this section is to analyze the structure of $\mathcal{M}_{m}^{\operatorname{sing}}$ by taking advantage of the special properties of the set of coordinates $\boldsymbol{\beta}$.

In general, the singular sectors of $\mathrm{dcKdV}_{m}$ hierarchies with $m \geq 2$ contain both reduced and unreduced points. For example, the hodograph equations (42) of the $\mathrm{dcKdV}_{2}$ hierarchy have reduced singular points given by ( $x, t_{1}, t_{2}, t_{3}, \beta_{1}=\beta_{2}$ ) where

$$
72 x t_{3}^{2}=-9 t_{2}^{2}+36 t_{1} t_{2} t_{3}+\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{9 t_{2}^{2}-24 t_{1} t_{3}},
$$

and

$$
\beta_{1}=\beta_{2}=-\frac{3 t_{2}+\sqrt{9 t_{2}^{2}-24 t_{1} t_{3}}}{12 t_{3}} .
$$

Furthermore, there are also unreduced singular points $\left(x, t_{1}, t_{2}, t_{3}, \beta_{1}, \beta_{2}\right)$ determined by the constraint

$$
360 x t_{3}^{3}=-45 t_{3} t_{2}^{3}+180 t_{1} t_{3}^{2} t_{2}+\sqrt{15}\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}
$$

and

$$
\beta_{1}=\frac{-3 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{12 t_{3}^{2}}, \quad \beta_{2}=-\frac{5 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{20 t_{3}^{2}}
$$

From (41) it follows at once that
Theorem 3. Let $(\boldsymbol{t}, \boldsymbol{\beta})$ be an unreduced solution of the hodograph equations (37), then $(\boldsymbol{t}, \boldsymbol{\beta})$ is a singular point if and only if at least one of the derivatives

$$
\frac{\partial^{2} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i}^{2}}, \quad i=1, \ldots, m
$$

vanishes.
Notice that since the function $W_{m}$ satisfies the EPD equation (38), its partial derivatives at unreduced points $(\boldsymbol{t}, \boldsymbol{\beta})$

$$
\frac{\partial^{q} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{1}^{q_{1}} \cdots \partial \beta_{m}^{q_{m}}}, \quad q:=q_{1}+\cdots+q_{m}
$$

can always be expressed as a linear combination of diagonal derivatives $\partial_{\beta_{i}}^{k_{i}} W_{m}$ with $k_{i} \leq q_{i}$. Thus, for each vector $\boldsymbol{q}=\left(q_{1}, \ldots, q_{m}\right) \in \mathbb{N}^{m}$ with at least one $q_{i} \geq 1$ it is natural to introduce an associated subvariety $\mathcal{M}_{m, \boldsymbol{q}}^{\operatorname{sing}}$ of $\mathcal{M}_{m}^{\operatorname{sing}}$ defined as the set of unreduced solutions $(\boldsymbol{t}, \boldsymbol{\beta})$ of the hodograph equations (37) such that

$$
\begin{equation*}
\frac{\partial^{k_{i}} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{i}^{k_{i}}}=0, \quad \forall k_{i} \leq q_{i}+1 . \tag{45}
\end{equation*}
$$

In particular, for $\boldsymbol{q}=(0, \ldots, 0, q)$ with $q \geq 1$ we denote by $\mathcal{M}_{m, q}^{\operatorname{sing}}$ the subvariety associated to $\boldsymbol{q}=(0, \ldots, 0, q)$. That is to say, $\mathcal{M}_{m, q}^{\operatorname{sing}}$ is the set of solutions $(\boldsymbol{t}, \boldsymbol{\beta})$ of the hodograph equations (37) such that

$$
\begin{equation*}
\frac{\partial^{2} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{m}^{2}}=\frac{\partial^{3} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{m}^{3}}=\ldots=\frac{\partial^{q+1} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})}{\partial \beta_{m}^{q+1}}=0 . \tag{46}
\end{equation*}
$$

These subvarieties define a nested sequence

$$
\begin{equation*}
\mathcal{M}_{m}^{\operatorname{sing}} \supset \mathcal{M}_{m, 1}^{\operatorname{sing}} \supset \mathcal{M}_{m, 2}^{\operatorname{sing}} \supset \cdots \mathcal{M}_{m, q}^{\operatorname{sing}} \supset \cdots \tag{47}
\end{equation*}
$$

and represent sets of points whose singular degree increases with $q$. Moreover, due to the covariance of the functions $F_{i}=\partial_{\beta_{i}} W_{m}$ under permutations there is no need of introducing alternative sequences of the form (46) based on systems of equations corresponding to the remaining coordinates $\beta_{j}$ for $j \neq m$.

The next result states that the varieties $\mathcal{M}_{m, q}^{\operatorname{sing}}$ of the $d c K d V_{m}$ hierarchy are closely related to the $(2 q+1)$-reduced solutions of the $\mathrm{dcKdV}_{m+2} q$ hierarchy.

Notice that given $2 \leq r \leq m$, the hodograph equations for $r$-reduced solutions

$$
\beta_{m-r+1}=\beta_{m-r+2}=\ldots=\beta_{m}
$$

of the $\mathrm{dcKdV}_{m}$ hierarchy reduce to the system

$$
F_{i}(\boldsymbol{t}, \boldsymbol{\beta})=0, \quad i=1, \ldots, m-r+1
$$

of $m-r+1$ equations for the $m-r+1$ unknowns $\left(\beta_{1}, \ldots, \beta_{m-r+1}\right)$. Now we prove
Theorem 4. If $(\boldsymbol{t}, \boldsymbol{\beta}) \in \mathcal{M}_{m, q}^{\operatorname{sing}}$ where $\boldsymbol{t}=\left(t_{0}, t_{1}, \ldots\right)$ and $\boldsymbol{\beta}=\left(\beta_{1}, \ldots, \beta_{m}\right)$, then if we define

$$
\boldsymbol{t}^{(m+2 q)}:=\left(t_{q}, t_{q+1}, \ldots\right), \quad \boldsymbol{\beta}^{(m+2 q)}:=(\beta_{1}, \ldots, \beta_{m}, \overbrace{\beta_{m}, \ldots, \beta_{m}}^{2 q}),
$$

it follows that $\left(\boldsymbol{t}^{(m+2 q)}, \boldsymbol{\beta}^{(m+2 q)}\right)$ is a $(2 q+1)$-reduced solution of the hodograph equations for the $d c K d V_{m+2 q}$ hierarchy.

Proof. To proof this statement we will use superscripts $(m)$ and $(m+2 q)$ to distinguish objects corresponding to different hierarchies. By assumption we have that

$$
\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right) \in \mathcal{M}_{m, q}^{\operatorname{sing}}
$$

Thus, taking (30) into account, we have that (46) can be rewritten as

$$
\begin{cases}F_{i}^{(m)}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\lambda-\beta_{i}^{(m)}}=0, & i=1, \ldots, m  \tag{48}\\ F_{m, j}^{(m)}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\left(\lambda-\beta_{m}^{(m)}\right)^{j}}=0, & j=2, \ldots, q+1\end{cases}
$$

Now a $(2 q+1)$-reduced solution of the hodograph equations for the $\operatorname{dcKdV}_{m+2}$ is characterized by

$$
\begin{equation*}
F_{i}^{(m+2 q)}\left(\boldsymbol{t}^{(m+2 q)}, \boldsymbol{\beta}^{(m+2 q)}\right):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U^{(m+2 q)}\left(\lambda, \boldsymbol{t}^{(m+2 q)}\right) R^{(m+2 q)}\left(\lambda, \boldsymbol{\beta}^{(m+2 q)}\right)}{\lambda-\beta_{i}^{(m+2 q)}}=0 \tag{49}
\end{equation*}
$$

where $i=1, \ldots, m$. But it is clear that

$$
\begin{equation*}
R^{(m+2 q)}\left(\lambda, \boldsymbol{\beta}^{(m+2 q)}\right)=\frac{\lambda^{q}}{\left(\lambda-\beta_{m}^{(m)}\right)^{q}} R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right) \tag{50}
\end{equation*}
$$

Hence if we set

$$
t_{i}^{(m+2 q)}:=t_{i+q}^{(m)}, \quad i \geq 0,
$$

we have

$$
\begin{equation*}
U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right)=x^{(m)}+\lambda t_{1}^{(m)}+\cdots+\lambda^{q-1} t_{q-1}^{(m)}+\lambda^{q} U^{(m+2 q)}\left(\lambda, \boldsymbol{t}^{(m+2 q)}\right) . \tag{51}
\end{equation*}
$$

Then it follows that

$$
\begin{equation*}
F_{i}^{(m+2 q)}\left(\boldsymbol{t}^{(m+2 q)}, \boldsymbol{\beta}^{(m+2 q)}\right)=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\left(\lambda-\beta_{i}^{(m)}\right)\left(\lambda-\beta_{m}^{(m)}\right)^{q}}, \quad i=1, \ldots, m . \tag{52}
\end{equation*}
$$

Furthermore, for any given $i=1, \ldots, m$ we have

$$
\begin{aligned}
& F_{i}^{(m)}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\lambda-\beta_{i}^{(m)}} \\
& =\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{\left(\lambda-\beta_{m}^{(m)}\right)^{q} U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\left(\lambda-\beta_{i}^{(m)}\right)\left(\lambda-\beta_{m}^{(m)}\right)^{q}} \\
& =\sum_{k=0}^{q} c_{1, k}\left(\boldsymbol{\beta}^{(m)}\right) I_{i, k}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right),
\end{aligned}
$$

and

$$
\begin{aligned}
& F_{m, j}^{(m)}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\left(\lambda-\beta_{m}^{(m)}\right)^{j}} \\
& =\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{\left(\lambda-\beta_{i}^{(m)}\right)\left(\lambda-\beta_{m}^{(m)}\right)^{q-j} U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\left(\lambda-\beta_{i}^{(m)}\right)\left(\lambda-\beta_{m}^{(m)}\right)^{q}} \\
& =\sum_{k=0}^{q-j+1} c_{j, k}\left(\boldsymbol{\beta}^{(m)}\right) I_{i, k}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right), \quad j=2, \ldots, q+1 .
\end{aligned}
$$

where the functions $c_{j, k}\left(\boldsymbol{\beta}^{(m)}\right)$ are the coefficients of the polynomials

$$
\left\{\begin{array}{l}
\left(\lambda-\beta_{m}^{(m)}\right)^{q}=\sum_{k=0}^{q} c_{1 k}\left(\boldsymbol{\beta}^{(m)}\right) \lambda^{k} ;  \tag{53}\\
\left(\lambda-\beta_{i}^{(m)}\right)\left(\lambda-\beta_{m}^{(m)}\right)^{q-j}=\sum_{k=0}^{q-j+1} c_{j k}\left(\boldsymbol{\beta}^{(m)}\right) \lambda^{k}, \quad j=2, \ldots, q+1 .
\end{array}\right.
$$

and

$$
\begin{equation*}
I_{i, k}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right):=\oint_{\gamma} \frac{\mathrm{d} \lambda}{2 i \pi} \frac{\lambda^{k} U^{(m)}\left(\lambda, \boldsymbol{t}^{(m)}\right) R^{(m)}\left(\lambda, \boldsymbol{\beta}^{(m)}\right)}{\left(\lambda-\beta_{i}^{(m)}\right)\left(\lambda-\beta_{m}^{(m)}\right)^{q}} \tag{54}
\end{equation*}
$$

Now, for any given $i=1, \ldots, m$ the system (46) implies

$$
\left\{\begin{array}{l}
F_{i}^{(m)}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)=0, \\
F_{m, j}^{(m)}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)=0, \quad j=2, \ldots, q+1,
\end{array}\right.
$$

and, as a consequence, we deduce the following system of $q$ homogeneous linear equations for the $q$ functions $I_{i, k}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)$

$$
\sum_{k=0}^{q-j+1} c_{j, k}\left(\boldsymbol{\beta}^{(m)}\right) I_{i, k}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)=0, \quad j=1, \ldots, q+1
$$

Because of the linear independence of the polynomials (53) these equations are linearly independent and, therefore, all the functions $I_{i, k}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)$ vanish. Finally, from (52) we conclude that $I_{i, 0}\left(\boldsymbol{t}^{(m)}, \boldsymbol{\beta}^{(m)}\right)=0$ is equivalent to $F_{i}^{(m+2 q)}\left(\boldsymbol{t}^{(m+2 q)}, \boldsymbol{\beta}^{(m+2 q)}\right)=0$ and the statement follows.

## 5 Examples

## dcKdV ${ }_{1}$ hierarchy

The hodograph equation for the $\mathrm{dcKdV}_{1}$ hierarchy with $t_{n}=0$ for all $n \geq 3$ reduce to

$$
\begin{equation*}
8 x+12 t_{1} \beta_{1}+15 t_{2} \beta_{1}^{2}=0 . \tag{55}
\end{equation*}
$$

The singular variety $\mathcal{M}_{1,1}^{\operatorname{sing}}$ for (55) is determined by adding to (55) the equation

$$
\begin{equation*}
2 t_{1}+5 t_{2} \beta_{1}=0 \tag{56}
\end{equation*}
$$

so that for $t_{2} \neq 0$ we have $\beta_{1}=-\frac{2 t_{1}}{5 t_{2}}$. Substituting this result in (55) we find a constraint for the flow parameters

$$
x=\frac{3}{10} \frac{t_{1}^{2}}{t_{2}},
$$

which is the shock region for the solution of (55) given by

$$
\begin{equation*}
\beta_{1}=\frac{2}{15 t_{2}}\left(-3 t_{1}+\sqrt{3\left(3 t_{1}^{2}-10 t_{2} x\right)}\right) . \tag{57}
\end{equation*}
$$

There are two sectors $\mathcal{M}_{1,1, k}^{\operatorname{sing}}(k=1,2)$ in $\mathcal{M}_{1,1}^{\operatorname{sing}}$

$$
\begin{array}{ll}
\mathcal{M}_{1,1,1}^{\mathrm{sing}}: & x=t_{1}=t_{2}=0, \quad \beta_{1} \text { arbitrary } \\
\mathcal{M}_{1,1,2}^{\mathrm{sing}}: & \left(x, t_{1}, t_{2}, \beta_{1}\right) \text { such that } t_{2} \neq 0, x=\frac{3}{10} \frac{t_{1}^{2}}{t_{2}} \text { and } \beta_{1}=-\frac{2}{5} \frac{t_{1}}{t_{2}} \tag{58}
\end{array}
$$

To see the relationship with the $\mathrm{dcKdV}_{3}$ hierarchy we notice that

$$
x^{(3)}=t_{1}, \quad t_{1}^{(3)}=t_{2}
$$

and

$$
\boldsymbol{\beta}^{(3}=\left(\beta_{1}, \beta_{1}, \beta_{1}\right)=-\frac{2}{5} \frac{x^{(3)}}{t_{1}^{(3)}}(1,1,1)
$$

which is a 3-reduced solution of the first flow (25) of the $\mathrm{dcKdV}_{3}$ hierarchy.
The dcKdV $V_{1}$ hodograph equation with $t_{n}=0$ for all $n \geq 6$ is

$$
693 t_{5} \beta_{1}^{5}+630 t_{4} \beta_{1}^{4}+560 t_{3} \beta_{1}^{3}+480 t_{2} \beta_{1}^{2}+384 t_{1} \beta_{1}+256 x=0
$$

Let us first consider the singular variety $\mathcal{M}_{1,1}^{\operatorname{sing}}$ with $t_{n}=0$ for all $n \geq 4$. It is is determined by the equations

$$
\begin{aligned}
& 560 t_{3} \beta_{1}^{3}+480 t_{2} \beta_{1}^{2}+384 t_{1} \beta_{1}+256 x=0 \\
& 1680 t_{3} \beta_{1}^{2}+960 t_{2} \beta_{1}+384 t_{1}=0
\end{aligned}
$$

Thus an open subset of $\mathcal{M}_{1,1}^{\operatorname{sing}}$ can be parametrized by the equations

$$
\begin{aligned}
& x=\frac{-25 t_{2}^{3}+105 t_{1} t_{2} t_{3}+\sqrt{5} \sqrt{125 t_{2}^{6}-1050 t_{1} t_{3} t_{2}^{4}+2940 t_{1}^{2} t_{2}^{2} t_{3}^{2}-2744 t_{1}^{3} t_{3}^{3}}}{245 t_{3}^{2}} \\
& \beta_{1}=-\frac{2\left(-25 t_{2}^{3}+70 t_{1} t_{2} t_{3}+\sqrt{5} \sqrt{\left(5 t_{2}^{2}-14 t_{1} t_{3}\right)^{3}}\right)}{35 t_{3}\left(14 t_{1} t_{3}-5 t_{2}^{2}\right)}
\end{aligned}
$$

It determines the following 3-reduced solution of the two first flows of the dcKdV hierachy $\left(x^{(3)}=\right.$ $\left.t_{1}, t_{1}^{(3)}=t_{2}, t_{2}^{(3)}=t_{3}\right)$

$$
\beta_{1}^{(3)}=\beta_{2}^{(3)}=\beta_{3}^{(3)}=-\frac{2\left(-25\left(t_{1}^{(3)}\right)^{3}+70 x^{(3)} t_{1}^{(3)} t_{2}^{(3)}+\sqrt{5} \sqrt{\left(5\left(t_{1}^{(3)}\right)^{2}-14 x^{(3)} t_{2}^{(3)}\right)^{3}}\right)}{35 t_{2}^{(3)}\left(14 x^{(3)} t_{2}^{(3)}-5\left(t_{1}^{(3)}\right)^{2}\right)}
$$

Next, for the sector $\mathcal{M}_{1,2}^{\operatorname{sing}}$ if we set $t_{n}=0$ for all $n \geq 5$, we obtain the equations

$$
\begin{aligned}
& 630 t_{4} \beta_{1}^{4}+560 t_{3} \beta_{1}^{3}+480 t_{2} \beta_{1}^{2}+384 t_{1} \beta_{1}+256 x=0 \\
& 2520 t_{4} \beta_{1}^{3}+1680 t_{3} \beta_{1}^{2}+960 t_{2} \beta_{1}+384 t_{1}=0 \\
& 7560 t_{4} \beta_{1}^{2}+3360 t_{3} \beta_{1}+960 t_{2}=0
\end{aligned}
$$

From these equations we find

$$
\begin{aligned}
& t_{1}=\frac{5\left(-49 t_{3}^{3}+189 t_{2} t_{3} t_{4}+\sqrt{7} \sqrt{343 t_{3}^{6}-2646 t_{2} t_{4} t_{3}^{4}+6804 t_{2}^{2} t_{3}^{2} t_{4}^{2}-5832 t_{2}^{3} t_{4}^{3}}\right)}{1701 t_{4}^{2}}, \\
& x=\frac{5\left(-98 t_{3}^{4}+378 t_{2} t_{4} t_{3}^{2}+2 \sqrt{7} \sqrt{\left(7 t_{3}^{2}-18 t_{2} t_{4}\right)^{3}} t_{3}-243 t_{2}^{2} t_{4}^{2}\right)}{10206 t_{4}^{3}}, \\
& \beta_{1}=-\frac{2\left(-49 t_{3}^{3}+126 t_{2} t_{3} t_{4}+\sqrt{7} \sqrt{\left(7 t_{3}^{2}-18 t_{2} t_{4}\right)^{3}}\right)}{63 t_{4}\left(18 t_{2} t_{4}-7 t_{3}^{2}\right)} .
\end{aligned}
$$

Then the associated 5 -reduced solution of the two first flows of the dcKdV 5 hierarchy $\left(x^{(5)}=t_{2}\right.$, $\left.t_{1}^{(5)}=t_{3}, t_{2}^{(5)}=t_{4}\right)$ is given by

$$
\beta_{i}=-\frac{2\left(-49\left(t_{1}^{(5)}\right)^{3}+126 x^{(5)} t_{1}^{(5)} t_{2}^{(5)}+\sqrt{7} \sqrt{\left(7\left(t_{1}^{(5)}\right)^{2}-18 x^{(5)} t_{2}^{(5)}\right)^{3}}\right)}{63 t_{2}^{(5)}\left(18 x^{(5)} t_{2}^{(5)}-7\left(t_{1}^{(5)}\right)^{2}\right)}, \quad i=1, \ldots, 5
$$

## dcKdV ${ }_{2}$ hierarchy

Let us consider the hodograph equations for the $\mathrm{dcKdV}_{2}$ hierarchy with $t_{n}=0$ for all $n \geq 3$. From (42) we have that they take the form

$$
\left\{\begin{array}{l}
8 x+4 t_{1}\left(3 \beta_{1}+\beta_{2}\right)+3 t_{2}\left(5 \beta_{1}^{2}+2 \beta_{1} \beta_{2}+\beta_{2}^{2}\right)=0  \tag{59}\\
8 x+4 t_{1}\left(\beta_{1}+3 \beta_{2}\right)+3 t_{2}\left(\beta_{1}^{2}+2 \beta_{1} \beta_{2}+5 \beta_{2}^{2}\right)=0
\end{array}\right.
$$

The singular variety $\mathcal{M}_{2}^{\operatorname{sing}}$ is determined by (59) together with the additional condition $\left(\operatorname{det}\left(\partial_{\beta_{i} \beta_{j}} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})\right)=0\right)$

$$
\begin{equation*}
-\left(2 t_{1}+3 t_{2}\left(\beta_{1}+\beta_{2}\right)\right)^{2}+9\left(2 t_{1}+t_{2}\left(5 \beta_{1}+\beta_{2}\right)\right)\left(2 t_{1}+t_{2}\left(\beta_{1}+5 \beta_{2}\right)\right)=0 . \tag{60}
\end{equation*}
$$

There elements of $\mathcal{M}_{2}^{\text {sing }}$ are

$$
\begin{align*}
& x=t_{1}=t_{2}=0, \quad\left(\beta_{0}, \beta_{1}\right) \text { arbitrary } \\
& \left(x, t_{1}, t_{2}, \beta_{1}, \beta_{2}\right) \text { such that } t_{2} \neq 0, x=\frac{t_{1}^{2}}{3 t_{2}} \text { and } \beta_{1}=\beta_{2}=-\frac{t_{1}}{3 t_{2}} \tag{61}
\end{align*}
$$

The subvarieties $\mathcal{M}_{2, q}^{\operatorname{sing}}$ are all equal and given by

$$
x=t_{1}=t_{2}=0, \quad\left(\beta_{0}, \beta_{1}\right) \text { arbitrary with } \beta_{0} \neq \beta_{1} .
$$

Notice that the constraint $x=\frac{t_{1}^{2}}{3 t_{2}}$ determines the shock region for the following solution of

$$
\begin{equation*}
\beta_{1}=\frac{-t_{1}+\sqrt{2} \sqrt{t_{1}^{2}-3 t_{2} x}}{3 t_{2}}, \quad \beta_{2}=\frac{-t_{1}-\sqrt{2} \sqrt{t_{1}^{2}-3 t_{2} x}}{3 t_{2}} . \tag{59}
\end{equation*}
$$

Let us now consider the system of hodograph equations (42) for the $\mathrm{dcKdV}_{2}$ hierarchy with $t_{n}=0$ for all $n \geq 4$. The singular variety $\mathcal{M}_{2}^{\text {sing }}$ is now determined by (42) and the condition $\left(\operatorname{det}\left(\partial_{\beta_{i} \beta_{j}} W_{m}(\boldsymbol{t}, \boldsymbol{\beta})\right)=0\right)$

$$
\begin{gathered}
32 t_{1}^{2}+96 t_{2}\left(\beta_{1}+\beta_{2}\right) t_{1}+702 t_{3}^{2} \beta_{1}^{2} \beta_{2}^{2}+72\left(3 t_{2}^{2}+t_{1} t_{3}\right) \beta_{1} \beta_{2}+12\left(3 t_{2}^{2}+13 t_{1} t_{3}\right)\left(\beta_{1}^{2}+\beta_{2}^{2}\right)+ \\
\quad 486 t_{2} t_{3}\left(\beta_{2} \beta_{1}^{2}+\beta_{2}^{2} \beta_{1}\right)+90 t_{2} t_{3}\left(\beta_{1}^{3}+\beta_{2}^{3}\right)+180 t_{3}^{2}\left(\beta_{2} \beta_{1}^{3}+\beta_{2}^{3} \beta_{1}\right)+45 t_{3}^{2}\left(\beta_{1}^{4}+\beta_{2}^{4}\right)=0 .
\end{gathered}
$$

One finds the following six sectors in $\mathcal{M}_{2}^{\text {sing }}$

1. $x=\frac{-9 t_{2}^{3}+36 t_{1} t_{3} t_{2}+\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{9 t_{2}^{2}-24 t_{1} t_{3}}}{72 t_{3}^{2}}, \quad \beta_{1}=\beta_{2}=-\frac{3 t_{2}+\sqrt{9 t_{2}^{2}-24 t_{1} t_{3}}}{12 t_{3}}$,
2. $x=\frac{-9 t_{2}^{3}+36 t_{1} t_{3} t_{2}-\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{9 t_{2}^{2}-24 t_{1} t_{3}}}{72 t_{3}^{2}}, \quad \beta_{1}=\beta_{2}=\frac{-3 t_{2}+\sqrt{9 t_{2}^{2}-24 t_{1} t_{3}}}{12 t_{3}}$,
3. $x=\frac{-45 t_{3} t_{2}^{3}+180 t_{1} t_{3}^{2} t_{2}+\sqrt{15}\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{360 t_{3}^{3}}$,

$$
\beta_{1}=-\frac{5 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{20 t_{3}^{2}}, \quad \beta_{2}=\frac{-3 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{12 t_{3}^{2}},
$$

4. $x=\frac{-45 t_{3} t_{2}^{3}+180 t_{1} t_{3}^{2} t_{2}-\sqrt{15}\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{360 t_{3}^{3}}$,

$$
\beta_{1}=-\frac{3 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{12 t_{3}^{2}}, \quad \beta_{2}=\frac{-5 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{20 t_{3}^{2}},
$$

5. $x=\frac{-45 t_{3} t_{2}^{3}+180 t_{1} t_{3}^{2} t_{2}-\sqrt{15}\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{360 t_{3}^{3}}$,

$$
\begin{aligned}
& \beta_{1}=\frac{-5 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{20 t_{3}^{2}}, \quad \beta_{2}=-\frac{3 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{12 t_{3}^{2}} . \\
& \text { 6. } x=\frac{-45 t_{3} t_{2}^{3}+180 t_{1} t_{3}^{2} t_{2}+\sqrt{15}\left(8 t_{1} t_{3}-3 t_{2}^{2}\right) \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{360 t_{3}^{3}}, \\
& \beta_{1}=\frac{-3 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{12 t_{3}^{2}}, \quad \beta_{2}=-\frac{5 t_{2} t_{3}+\sqrt{15} \sqrt{t_{3}^{2}\left(3 t_{2}^{2}-8 t_{1} t_{3}\right)}}{20 t_{3}^{2}} .
\end{aligned}
$$

It is easy to see that $\mathcal{M}_{2,1}^{\operatorname{sing}}$ is given by the sectors 5 and 6 . To check the connection between these sectors and the $\mathrm{dcKdV}_{4}$ hierarchy it is enough to set

$$
x^{(4)}=t_{1}, \quad t_{1}^{(4)}=t_{2}, \quad t_{2}^{(4)}=t_{3}, \quad \boldsymbol{\beta}^{(4)}=\left(\beta_{1}, \beta_{2}, \beta_{2}, \beta_{2}\right),
$$

and it is immediate to prove that $\left.\boldsymbol{\beta}^{(4)} \boldsymbol{t}^{(4)}\right)$ verifies the equations of the first flow of the $\operatorname{dcKaV}_{4}$ hierarchy

$$
\frac{\partial \beta_{i}}{\partial t_{1}^{(4)}}=\left(\beta_{i}+\frac{1}{2} \sum_{k=1}^{4} \beta_{k}\right) \frac{\partial \beta_{i}}{\partial x^{(4)}}, \quad i=1, \ldots, 4
$$

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