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A Note on Proper Affine Symmetry in Bianchi

Types VI_0 and VII_0 Space-Times

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Abstract

A study proper affine symmetry in the most general form of the Bianchi types VI_0 and VII_0 space-times is given by using holonomy and decomposability, the rank of the 6×6 Riemann matrix and direct integration techniques. It is shown that the very special classes of the above space-times admit proper affine symmetry.

Keywords: Affine vector fields, holonomy and decomposability, direct integration technique.

1 INTRODUCTION

In this paper we will explore all the possibilities when the Bianchi types VI_0 and VII_0 space-times admit proper affine symmetry. We use holonomy and decomposability, the rank of the 6×6 Rieman matrix and direct integration techinques to study proper affine symmetry in the above space-times. Throughout

M represents a four dimensional, connected, Hausdorff space-time manifold with Lorentz metric g of signature (-, +, +, +). The curvature tensor associated with g_{ab} , through the Levi-Civita connection, is denoted in component form by R^{a}_{bcd} . The usual covariant, partial and Lie derivatives are denoted by a semicolon, a comma and the symbol *L*, respectively. Round and square brackets denote the usual symmetrization and skew-symmetrization, respectively. Here, *M* is assumed non-flat in the sense that the curvature tensor does not vanish over any non-empty open subset of *M*.

A vector field X on M is called an affine vector field if it satisfies

$$X_{a;bc} = R_{abcd} X^{d}, \qquad (1)$$

where $R_{abcd} = g_{af} R^{f}_{bcd} = g_{af} (\Gamma^{f}_{bd,c} - \Gamma^{f}_{bc,d} + \Gamma^{f}_{ce} \Gamma^{e}_{bd} - \Gamma^{f}_{ed} \Gamma^{e}_{bc})$. If one decomposes $X_{a;b}$ on M into its symmetric and skew-symmetric parts

$$X_{a;b} = \frac{1}{2}H_{ab} + G_{ab}, \ (H_{ab} (\equiv X_{a;b} + X_{b;a}) = H_{ba}, \ G_{ab} = -G_{ba})$$
(2)

then equation (1) is equivalent to

(*i*)
$$H_{ab;c} = 0$$
 (*ii*) $G_{ab;c} = R_{abcd} X^{d}$ (*iii*) $G_{ab;c} X^{c} = 0.$ (3)

The proof of the above equation (1) implies (3) or equations (3) implies (1) can be found in [2, 3]. If $H_{ab} = 2cg_{ab}, c \in R$, then the vector field X is called homothetic (and *Killing* if c = 0). The vector field X is said to be proper affine if it is not homothetic vector field and also X is said to be proper homothetic vector field if it is not Killing vector field on M [2]. Define the subspace S_p of the tangent space T_pM to M at p as those $k \in T_pM$ satisfying

$$R_{abcd}k^d = 0. (4)$$

2 Affine Vector Fields

Suppose that M is a simple connected space-time. Then the holonomy group of M is a connected Lie subgroup of the idenity component of the Lorentz group and is thus characterized by its subalgebra in the Lorentz algebra. These have been labeled into fifteen types $R_1 - R_{15}$ [1]. It follows from [2] that the only such space-times which could admit proper affine vector fields are those which admit nowhere zero covariantly constant second order symmetric tensor field H_{ab} . This forces the holonomy type to be either R_2 , R_3 , R_4 , R_6 , R_7 , R_8 , R_{10} , R_{11} or R_{13} [2]. A study of the affine vector fields for the above holonomy type can be found in [2]. It follows from [4] that the rank of the 6×6 Riemann matrix of the above space-times which have holonomy type R_2 , R_3 , R_4 , R_6 , R_7 , R_7 ,

 R_8 , R_{10} , R_{11} or R_{13} is at most three. Hence for studying affine vector fields we are interested in those cases when the rank of the 6×6 Riemann matrix is less than or equal to three.

3 Main Results

Consider Bianchi types VI₀ and VII₀ space-times in usual coordinate system (t, x, y, z) (labeled by (x^0, x^1, x^2, x^3) , respectively) with line element [5] $ds^2 = -dt^2 + (A f^2(z) + B h^2(z))dx^2 + (A h^2(z) + B f^2(z))dy^2 + 2(A + B)f(z)h(z)dx dy + C(t)dz^2$, (5)

where A, B and C are nowhere zero functions of t only. For $f(z) = \cosh z$, $h(z) = \sinh z$ or $f(z) = \cos z$, $h(z) = \sin z$ the above space-time (5) becomes Bianchi type VI₀ or VII₀, respectively. The above space-time admits three linearly independent Killing vector fields which are

$$\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, -y \frac{\partial}{\partial x} - x \frac{\partial}{\partial y} + \frac{\partial}{\partial z}.$$
 (6)

The non-zero components of the Riemann tensor are

$$\begin{split} R_{0101} &= -\frac{1}{4} \left[\left(\frac{2A\ddot{A} - \dot{A}^2}{A} \right) f^2(z) + \left(\frac{2B\ddot{B} - \dot{B}^2}{B} \right) h^2(z) \right] = \alpha_1, \\ R_{0102} &= -\frac{1}{4} \left[\left(\frac{2A\ddot{A} - \dot{A}^2}{A} \right) + \left(\frac{2B\ddot{B} - \dot{B}^2}{B} \right) \right] f(z)h(z) = \alpha_1, \\ R_{0113} &= R_{0223} &= -\frac{1}{4ABC} \left[A^2\dot{B}C + \dot{A}B^2C + 2AB\dot{C}(A + B) - 3ABC\left(\dot{A} + \dot{B}\right) \right] f(z)h(z) = \alpha_8, \\ R_{0123} &= -\frac{1}{4} \left[\left\{ \left(A + B\left(\frac{\dot{A}}{A} + \frac{\dot{C}}{C} \right) - 2\left(\dot{A} + \dot{B} \right) \right\} h^2(z) + \left\{ \left(A + B\left(\frac{\dot{B}}{B} + \frac{\dot{C}}{C} \right) - 2\left(\dot{A} + \dot{B} \right) \right\} f^2(z) \right] \right] \\ &= \alpha_9, \ R_{0202} &= -\frac{1}{4} \left[\left\{ \left(A + B\left(\frac{\dot{A}}{A} + \frac{\dot{C}}{C} \right) - 2\left(\dot{A} + \dot{B} \right) \right\} h^2(z) + \left(\frac{2B\ddot{B} - \dot{B}^2}{B} \right) f^2(z) \right] = \alpha_2, \\ R_{0213} &= -\frac{1}{4} \left[\left\{ \left(A + B\left(\frac{\dot{A}}{A} + \frac{\dot{C}}{C} \right) - 2\left(\dot{A} + \dot{B} \right) \right\} f^2(z) + \left\{ \left(A + B\left(\frac{\dot{B}}{B} + \frac{\dot{C}}{C} \right) - 2\left(\dot{A} + \dot{B} \right) \right\} h^2(z) \right] \right] \\ &= \alpha_{10}, \ R_{0303} &= -\frac{1}{4} \left(\frac{2C\ddot{C} - \dot{C}^2}{C} \right) = \alpha_3, \ R_{0312} &= -\frac{1}{4} \left(A + B\left(\frac{\dot{A}B - A\dot{B}}{AB} \right) \right) = \alpha_{11}, \\ R_{1212} &= \frac{1}{4ABC} \left[\left(A + B\right)^2 + \dot{A}\dot{B}C \right] \left[Af^2(z) + Bh^2(z) \right] \left[Ah^2(z) + Bf^2(z) \right] = \alpha_4, \\ R_{1313} &= \frac{1}{4AB} \left[\left\{ A^3 - AB\left(2A + 3B - \dot{A}\dot{C} \right) \right\} f^2(z) - \left\{ AB\left(3A + 2B - \dot{B}\dot{C} \right) - B^3 \right\} h^2(z) \right] = \alpha_5, \\ R_{1323} &= \frac{1}{4AB} \left[\left\{ A^3 - AB\left(2A + 3B - \dot{A}\dot{C} \right) \right\} h^2(z) - \left\{ AB\left(3A + 2B - \dot{B}\dot{C} \right) - B^3 \right\} h^2(z) \right] = \alpha_6. \end{split}$$

Writing the curvature tensor with components R_{abcd} at p as a 6×6 symmetric matrix [6]

$$R_{abcd} = \begin{pmatrix} \alpha_{1} & \alpha_{7} & 0 & 0 & \alpha_{8} & \alpha_{9} \\ \alpha_{7} & \alpha_{2} & 0 & 0 & \alpha_{10} & \alpha_{8} \\ 0 & 0 & \alpha_{3} & \alpha_{11} & 0 & 0 \\ 0 & 0 & \alpha_{11} & \alpha_{4} & 0 & 0 \\ \alpha_{8} & \alpha_{10} & 0 & 0 & \alpha_{5} & \alpha_{12} \\ \alpha_{9} & \alpha_{8} & 0 & 0 & \alpha_{12} & \alpha_{6} \end{pmatrix}.$$
(7)

As mentioned in section 2, the space-times which can admit proper affine vector fields have holonomy type R_2 , R_3 , R_4 , R_6 , R_7 , R_8 , R_{10} , R_{11} or R_{13} and the rank of the 6×6 Riemann matrix is at most three. Therefore we are only interested in those cases when the rank of the 6×6 Riemann matrix is less than or equal to three. Hence there exist following two possibilities when the rank of the 6×6 symmetric matrix is less or equal to three which are:

(A) Rank=3, $\alpha_1 = \alpha_2 = \alpha_3 = \alpha_7 = \alpha_8 = \alpha_9 = \alpha_{10} = \alpha_{11} = 0$, $\alpha_4 \neq 0, \alpha_5 \neq 0, \alpha_6 \neq 0$ and $\alpha_{12} \neq 0$. (B) Rank=1, $\alpha_1 \neq 0$, and

(B) Rank=1,
$$\alpha_4 \neq 0$$
, and

$$\alpha_1 = \alpha_2 = \alpha_3 = \alpha_5 = \alpha_6 = \alpha_7 = \alpha_8 = \alpha_9 = \alpha_{10} = \alpha_{11} = \alpha_{12} = 0$$

We will consider each case in turn.

Case (A):

In this case we have $\alpha_4 \neq 0, \alpha_5 \neq 0, \alpha_6 \neq 0, \alpha_{12} \neq 0$,

 $\alpha_1 = \alpha_2 = \alpha_3 = \alpha_7 = \alpha_8 = \alpha_9 = \alpha_{10} = \alpha_{11} = 0$, the rank of the 6×6 Riemann matrix is three and there exists a unique (up to a multiple) no where zero time like vector field $t_a = t_{,a}$ solution of equation (4) and $t_{a;b} \neq 0$. From the above constraints we have B(t) = dA(t), C(t) = eA(t) and $A(t) = (at+b)^2$, where $a,b,d,e \in R(d,e>0)$. The line element in this case takes the form

$$ds^{2} = -dt^{2} + (a t + b)^{2} \left[\frac{(f^{2}(z) + e h^{2}(z))dx^{2} + (d h^{2}(z) + f^{2}(z))dy^{2} + 2(1 + d)f(z)h(z)dx dy + e dz^{2}}{2(1 + d)f(z)h(z)dx dy + e dz^{2}} \right]. (8)$$

Substituting the above information into the affine equations one find that

$$X^{0} = c_{4}t + c_{5}, X^{1} = -c_{1}y + c_{2}, X^{2} = -c_{1}x + c_{3}, X^{3} = c_{1},$$
(9)

where $c_1, c_2, c_3, c_4, c_5 \in R$. One can write the above equation (9) subtracting Killing vector fields as

$$X = (c_4 t + c_5, 0, 0, 0). \tag{10}$$

Clearly in this case one can easily see that the above space-times (8) admit proper affine symmetry. It is important to note that the constants a and b can not be zero simultaneously.

Case (B):

In this case we have

 $\alpha_1 = \alpha_2 = \alpha_3 = \alpha_5 = \alpha_6 = \alpha_7 = \alpha_8 = \alpha_9 = \alpha_{10} = \alpha_{11} = \alpha_{12} = 0$, $\alpha_4 \neq 0$ and the rank of the 6×6 Riemann matrix is one. Here, there exist two linear independent solutions $t_a = t_{,a}$ and $z_a = z_{,a}$ of equation (4). The vector field t_a is covariantly constant whereas z_a is not covariantly constant. From the above constraints we have $A(t) = B(t) = C(t) = (t+q)^2$, where $q \in R$. The line element takes the form

$$ds^{2} = -dt^{2} + (t+q)^{2} \begin{bmatrix} (f^{2}(z) + h^{2}(z))dx^{2} + (h^{2}(z) + f^{2}(z))dy^{2} + \\ 4f(z)h(z)dx dy + dz^{2} \end{bmatrix}.$$
 (11)

Affine vector fields in this case

$$X^{0} = c_{4}t + c_{5}z + c_{6}, X^{1} = -c_{1}y + c_{2}, X^{2} = -c_{1}x + c_{3}, X^{3} = c_{1},$$
(12)

where $c_1, c_2, c_3, c_4, c_5, c_6 \in R$. One can write the above equation (12) subtracting Killing vector fields as

$$X = (c_4 t + c_5 z + c_6, 0, 0, 0).$$
⁽¹³⁾

In this case clearly the above space-times (11) admit proper affine symmetry.

SUMMARY

In this paper an attempt is made to explore all the possibilities when the Bianchi types VI_0 and VII_0 space-times admit proper affine vector fields. A different approach is adopted to study proper affine vector fields of the above space-times by using holonomy and decomposability, the rank of the 6×6 Riemann matrix and direct integration techniques. From the above study we obtain the following results:

(i) We obtain the space-time (8) that admits proper affine vector fields when the rank of the 6×6 Riemann matrix is three and there exists a unique nowhere zero independent timelike vector field, which is the solution of equation (4) and is not covariantly constant (for details see case A).

(ii) The space-time (11) is obtained, which admits proper affine vector fields (see case B) when the rank of the 6×6 Riemann matrix is one and there exist two independent solutions of equation (4) but only one independent covariantly constant vector field.

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