

UNIQUENESS OF THE FOLIATION OF CONSTANT MEAN CURVATURE SPHERES IN ASYMPTOTICALLY FLAT 3-MANIFOLDS

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Abstract

In this paper I study the constant mean curvature surface in asymptotically flat 3-manifolds with general asymptotics. Under some weak condition, I prove that outside some compact set in the asymptotically flat 3-manifold with positive mass, the foliation of stable spheres of constant mean curvature is unique.

1 Introduction

A three-manifold M with a Riemannian metric g and a two-tensor K is called an initial data set (M, g, K) if g and K satisfy the constraint equations

$$\begin{aligned} R_g - |K|_g^2 + (tr_g(K))^2 &= 16\pi\rho \\ div_g(K) - d(tr_g(K)) &= 8\pi J \end{aligned} \quad (1.1)$$

where R_g is the scalar curvature of the metric g , $tr_g(K)$ denotes $g^{ij}K_{ij}$, ρ is the observed energy density, and J is the observed momentum density.

Definition 1.1. Let $q \in (\frac{1}{2}, 1]$. We say (M, g, K) is asymptotically flat (AF) if it is a initial data set, and there is a compact subset $\tilde{K} \subset M$ such that $M \setminus \tilde{K}$ is diffeomorphic to $R^3 \setminus B_1(0)$ and there exists coordinate $\{x^i\}$ such that

$$g_{ij}(x) = \delta_{ij} + h_{ij}(x) \quad (1.2)$$

$$h_{ij}(x) = O_5(|x|^{-q}) \quad K_{ij}(x) = O_1(|x|^{-1-q}) \quad (1.3)$$

Also, ρ and J satisfy

$$\rho(x) = O(|x|^{-2-2q}) \quad J(x) = O(|x|^{-2-2q}) \quad (1.4)$$

Here, $f = O_k(|x|^{-q})$ means $\partial^l f = O(|x|^{-l-q})$ for $l = 0, \dots, k$. $M \setminus \tilde{K}$ is called an end of this asymptotically flat manifold.

We can define mass for the asymptotically flat manifolds as follows:

$$m = \lim_{r \rightarrow \infty} \frac{1}{16\pi} \int_{|x|=r} (h_{ij,j} - h_{jj,i}) v_g^i d\mu_g \quad (1.5)$$

where v_g and $d\mu_g$ are the normal vector and volume form with respect to the metric g . From [1], we know the mass is well defined when $q > 1/2$.

Definition 1.2. We say (M, g, K) is asymptotically flat satisfying the Regge-Teitelboim condition (AF-RT) if it is AF, and g, K satisfy these asymptotically even/odd conditions

$$h_{ij}^{odd}(x) = O_2(|x|^{-1-q}) \quad K_{ij}^{even}(x) = O_1(|x|^{-2-q}) \quad (1.6)$$

Also, ρ and J satisfy

$$\rho^{odd}(x) = O(|x|^{-3-2q}) \quad J^{odd}(x) = O(|x|^{-3-2q}) \quad (1.7)$$

where $f^{odd}(x) = f(x) - f(-x)$ and $f^{even}(x) = f(x) + f(-x)$.

For (AF-RT) manifolds, the center of mass C is defined by

$$C^\alpha = \frac{1}{16\pi m} \lim_{r \rightarrow \infty} \left(\int_{|x|=r} x^\alpha (h_{ij,i} - h_{ii,j}) v_g^j d\mu_g - \int_{|x|=r} (h_{i\alpha} v_g^i - h_{ii} v_g^\alpha) d\mu_g \right). \quad (1.8)$$

From [3], we know it is well defined.

The constant mean curvature surface is stable means the second variation operator has non-negative eigenvalues when restricted to the functions with 0 mean value, i.e.

$$\int_{\Sigma} (|A|^2 + Ric(v_g, v_g)) f^2 d\mu \leq \int_{\Sigma} |\nabla f|^2 d\mu \quad (1.9)$$

for function f with $\int_{\Sigma} f d\mu = 0$, where A is the second fundamental form, and $Ric(v_g, v_g)$ is the Ricci curvature in the normal direction with respect to the metric g .

We discuss the existence and uniqueness of constant mean curvature spheres that separate the origin from the infinity in the AF-RT manifolds. The following two theorems are due to Lan-Hsuan Huang [2]:

Theorem 1.3. (Existence) If (M, g, K) is the AF-RT with $q \in (\frac{1}{2}, 1]$, there exists a foliation by spheres $\{\Sigma_R\}$ with constant mean curvature $H(\Sigma_R) = \frac{2}{R} + O(R^{-1-q})$ in the exterior region of M . Each leaf Σ_R is a $c_0 R^{1-q}$ -graph over $S_R(C)$ and is strictly stable.

Set $r(x) = (\Sigma(x_i^2))^{1/2}$. For the constant mean curvature sphere Σ which separates infinity from K , we define

$$\begin{aligned} r_0(\Sigma) &= \inf\{r(x) | x \in \Sigma\} \\ r_1(\Sigma) &= \sup\{r(x) | x \in \Sigma\} \end{aligned} \quad (1.10)$$

Theorem 1.4. (Uniqueness) Assume that (M, g, K) is AF-RT with $q \in (\frac{1}{2}, 1]$ and $m > 0$. There exists σ_1 and C_1 so that if Σ has the following properties:

- Σ is topologically a sphere
- Σ has constant mean curvature $H = H(\Sigma_R)$ for some $R \geq \sigma_1$
- Σ is stable
- $r_1 \leq C_1 r_0^{\frac{1}{a}}$ for some a satisfying $\frac{5-q}{2(2+q)} < a \leq 1$

then $\Sigma = \Sigma_R$.

Our main uniqueness result is

Theorem 1.5. Suppose (M, g, K) is AF-RT 3-manifold with positive mass, and g can be expressed on the end $M \setminus \tilde{K}$ as follows:

$$g_{ij} = \delta_{ij} + h_{ij}^1(\theta)/r + Q \quad (1.11)$$

where $\theta = (\theta_1, \theta_2)$ is the coordinate on $S^2 \subset R^3$. If g satisfies the following properties:

- $h_{ij}^1(\theta) \in C^5(S^2)$
- $Q = O_5(|x|^{-2})$

Then for any $k > 2$, there exists some $\varepsilon > 0$ depending on k such that if

$$\|h_{ij}(\theta) - \delta_{ij}(\theta)\|_{W^{k,2}(S^2)} \leq \varepsilon, \quad (1.12)$$

there is a compact domain \tilde{K} such that if a foliation $\{\Sigma\}$ of stable constant mean curvature spheres which separates infinity from \tilde{K} have

$$\lim_{r_0 \rightarrow \infty} \frac{\log(r_1(\Sigma))}{r_0(\Sigma)^{1/4}} = 0 \quad (1.13)$$

then this foliation is the same one as in Theorem 1.3.

Remark 1.6. If we replace $\|h_{ij}(\theta) - \delta_{ij}(\theta)\|_{W^{k,2}} \leq \varepsilon$ by $\|h_{ij}(\theta) - C\delta_{ij}(\theta)\|_{W^{k,2}} \leq \varepsilon$ for any constant $C > 0$, we can also get this theorem, but ε will depend on k and C .

Remark 1.7. RT condition is needed to apply the theorems of Huang and if we assume the scalar curvature satisfies $R = O(r^{-3-\varepsilon})$ for some $\varepsilon > 0$, then we do not need the constraint equation.

Remark 1.8. Here I can only deal with the case when $q = 1$. When $q \in (1/2, 1)$ it seems that $\|h_{ij}(\theta) - \delta_{ij}(\theta)\|_{W^{k,2}(S^2)} \leq \varepsilon$ is not a proper condition.

The above theorem is about the uniqueness of the foliation. For the uniqueness of a single CMC sphere we have:

Corollary 1.9. We assume the same condition on the metric as the above Theorem. Then for any constants $C > 0$ and $\beta > 0$, there exist some compact set $K(C, \beta) \subset M$, such that any stable sphere Σ that separates $K(C, \beta)$ from the infinity with

$$\frac{(\log(r_1(\Sigma)))^{1+\beta}}{r_0(\Sigma)^{1/4}} \leq C \quad (1.14)$$

belongs to the foliation in Theorem 1.3.

The paper is organized much like [9]: In Section 2 we do apriori estimate on the stable constant mean curvature sphere based on the Simon's identity. In Section 3, we introduce blow-down analysis in three different scales. In Section 4 we recall the asymptotic analysis from [10] and prove a technical lemma. In Section 5 we introduce the asymptotically harmonic coordinate. In Section 6 we introduce a sense of the center of mass and prove the theorem.

2 Curvature estimates

From now on let Σ be a constant mean curvature sphere in the asymptotically flat end (M, g) which separates the origin from the infinity. First we have the following estimate as Lemma 5.2 in [5].

Lemma 2.1. Let $X = x^i \frac{\partial}{\partial x^i}$ be the Euclidean coordinate vectorfield and $r = (\Sigma(x^i)^2)^{1/2}$ and with respect to the metric g , v is the outward normal vector field, $d\mu$ is the volume form of Σ . Then we have the estimate:

$$\int_{\Sigma} \langle X, v \rangle r^{-4} d\mu \leq H^2 |\Sigma| \quad (2.1)$$

Moreover for each $a \geq a_0 > 2$ and r_0 sufficiently large, we have:

$$\int_{\Sigma} r^{-a} d\mu \leq C(a_0) r_0^{2-a} H^2 |\Sigma| \quad (2.2)$$

Proof. Because the mean curvature H is constant, then for some smooth vector field Y on Σ , we have the divergence formula:

$$\int_{\Sigma} \operatorname{div}_{\Sigma} Y d\mu = H \int_{\Sigma} \langle Y, v \rangle d\mu. \quad (2.3)$$

We choose $Y = Xr^{-a}$, $a \geq 2$ and e_{α} is the orthonormal basis on Σ , $\alpha = 1, 2$. Suppose $e_{\alpha} = a_{\alpha}^i \frac{\partial}{\partial x^i}$, it is obvious that a_{α}^i is bounded because the manifold is asymptotically flat. Then we have:

$$\begin{aligned} \operatorname{div}_{\Sigma} Y &= \operatorname{div}_{\Sigma}(Xr^{-a}) = \langle \nabla_{e_{\alpha}}(Xr^{-a}), e_{\alpha} \rangle \\ &= r^{-a} \operatorname{div}_{\Sigma} X - ar^{-a-2} a_{\alpha}^i a_{\alpha}^j x^i x^j + O(r^{-a-q}) \\ &= r^{-a} \operatorname{div}_{\Sigma} X - \alpha r^{-a-2} |X^{\tau}|^2 + O(r^{-a-q}) \end{aligned} \quad (2.4)$$

where X^τ is the tangent projection of X .

$$|\operatorname{div}_\Sigma X - 2| = O(r^{-q}) \quad (2.5)$$

Note that $|X^\tau|^2 = r^2 - \langle X, v \rangle^2 + O(r^{2-q})$, then combine all of these we have:

$$\begin{aligned} & |(2-a) \int_\Sigma r^{-a} d\mu + a \int_\Sigma \langle X, v \rangle^2 r^{-a-2} d\mu - H \int_\Sigma \langle X, v \rangle r^{-a} d\mu| \\ & \leq C \int_\Sigma r^{-a-q} d\mu \end{aligned} \quad (2.6)$$

Choosing $a = 2$, from Hölder inequality, we have:

$$\int_\Sigma \langle X, v \rangle^2 r^{-4} d\mu \leq \frac{1}{4} H^2 |\Sigma| + C \int_\Sigma r^{-2-q} d\mu \quad (2.7)$$

then choose $a = 2 + q$,

$$\int_\Sigma r^{-2-q} d\mu \leq 4r_0^{-q} \left(\int_\Sigma \langle X, v \rangle^2 r^{-4} d\mu + H^2 |\Sigma| + C \int_\Sigma r^{-2-q} d\mu \right) \quad (2.8)$$

then combine this with (2.7), we have:

$$\int_\Sigma \langle X, v \rangle^2 r^{-4} d\mu \leq H^2 |\Sigma| \quad (2.9)$$

then again from (2.6), we have for $a \geq a_0 > 2$, we derive:

$$\int_\Sigma r^{-a} \leq C(a_0 - 2)^{-1} r_0^{2-a} H^2 |\Sigma| \quad (2.10)$$

Then we can derive the integral estimate for $|\mathring{A}|$ from the stability of the surface as in [5] Proposition 5.3, i.e. we have

Lemma 2.2. Suppose Σ is a stable constant mean curvature sphere in the asymptotically flat manifold. We have for r_0 sufficiently large

$$\int_\Sigma |\mathring{A}|^2 d\mu \leq C r_0^{-q} \quad (2.11)$$

$$H^2 |\Sigma| \leq C \quad (2.12)$$

$$\int_\Sigma H^2 d\mu = 16\pi + O(r_0^{-q}) \quad (2.13)$$

Proof. Since Σ is stable, we have

$$\int_\Sigma |\nabla f|^2 d\mu \geq \int_\Sigma (|A|^2 + \operatorname{Ric}(v, v)) f^2 d\mu \quad (2.14)$$

for any function f , with $\int_{\Sigma} f d\mu = 0$, where A is the second fundamental form of Σ and Ric is the Ricci curvature of M

Choose ψ to be a conformal map of degree 1 from Σ to the standard S^2 in R^3 . Each component ψ_i of ψ can be chosen such that $\int \psi_i d\mu = 0$, see [8]. We have for each ψ_i

$$\int_{\Sigma} |\nabla \psi_i|^2 d\mu = \frac{8\pi}{3} \quad (2.15)$$

since $\sum \psi_i^2 \equiv 1$ we conclude that

$$\int_{\Sigma} |A|^2 + Ric(v, v) d\mu \leq 8\pi \quad (2.16)$$

From Gauss equation

$$\frac{1}{2}|A|^2 + Ric(v, v) - \frac{1}{2}R + K = \frac{1}{2}H^2 \quad (2.17)$$

we have:

$$|A|^2 + Ric(v, v) = \frac{1}{2}|\mathring{A}|^2 + \frac{3}{4}H^2 + \frac{1}{2}R - K \quad (2.18)$$

where K is the Gauss curvature of Σ and \mathring{A} is defined as $\mathring{A}_{ij} = A_{ij} - \frac{H}{2}g_{ij}$

Then we have:

$$\int_{\Sigma} \frac{1}{2}|\mathring{A}|^2 + \frac{3}{4}H^2 |\Sigma| \leq 12\pi + r_0^{-q} H^2 |\Sigma| \quad (2.19)$$

because $R = O(r^{-2-2q})$.

So we have $H^2 |\Sigma| \leq 16\pi$.

Using the Gauss equation in a different way, we have

$$\begin{aligned} \int_{\Sigma} |\mathring{A}|^2 d\mu &= \int_{\Sigma} |A|^2 - \frac{H^2}{2} d\mu \\ &= \frac{1}{2} \int_{\Sigma} |A|^2 + Ric(v, v) d\mu + \frac{1}{2} \int_{\Sigma} R - 3Ric(v, v) - 2K d\mu \\ &\leq \int_{\Sigma} r^{-2-q} d\mu \\ &= O(r_0^{-q}). \end{aligned} \quad (2.20)$$

Then from Gauss equation (2.17) again, we have:

$$\int_{\Sigma} H^2 d\mu = 4 \int_{\Sigma} K d\mu + O(r_0^{-q}) = 16\pi + O(r_0^{-q}) \quad (2.21)$$

Lemma 2.3. Suppose that M is a constant mean curvature surface in an asymptotically flat end $(R^3 \setminus B_1(0), g)$. Then

$$\int_{\Sigma} H_e^2 d\mu_e = 16\pi + O(r_0^{-q}) \quad (2.22)$$

Proof. We follow the calculation of Huisken and Ilmanen [4],

$$g_{ij} = \delta_{ij} + h_{ij} \quad (2.23)$$

Suppose

$$g_{ij}|_{\Sigma} = f_{ij}, \delta_{ij}|_{\Sigma} = \varepsilon_{ij} \quad (2.24)$$

f^{ij} and ε^{ij} are the corresponding inverse matrices. $v, \omega, A, H, d\mu$ represents the normal vector, the dual form of v , the second fundamental form, the mean curvature and the volume form of Σ in the metric g . And $v_e, \omega_e, A_e, H_e, \mu_e$ represents the corresponding ones in Euclidean metric. Through easy calculation, we have

$$f^{ij} - \varepsilon^{ij} = -f^{ik} h_{kl} f^{lj} \pm C|h|^2 \quad (2.25)$$

$$g^{ij} - \delta^{ij} = -g^{ik} h_{kl} g^{lj} \pm C|h|^2 \quad (2.26)$$

$$\omega = \frac{\omega_e}{|\omega_e|} \quad v^i = g^{ij} \omega_j \quad (2.27)$$

$$(\omega_e)_i = \omega_i \pm C|P| \quad v_e^i = v^i + C|h| \quad 1 - |\omega_e| = \frac{1}{2} h_{ij} v^i v^j \quad (2.28)$$

$$\Gamma_{ij}^k = \frac{1}{2} g^{kl} (\bar{\nabla}_i h_{jl} + \bar{\nabla}_j h_{il} - \bar{\nabla}_l h_{ij}) \pm C|h| \pm C|\bar{\nabla}h| \quad (2.29)$$

and Γ_{ij}^k is the Christoffel symbol for $\bar{\nabla} - \bar{\nabla}_e$, where we denote the gradient for the metric g and δ by $\bar{\nabla}$ and $\bar{\nabla}_e$.

We have the formula:

$$|\omega_e|_g A_{ij} = (A_e)_{ij} - (\omega_e)_k \Gamma_{ij}^k \quad (2.30)$$

So we have

$$\begin{aligned} H - H_e &= f^{ij} A_{ij} - \varepsilon^{ij} (A_e)_{ij} \\ &= (f^{ij} - \varepsilon^{ij}) A_{ij} + \varepsilon^{ij} A_{ij} (1 - |\omega_e|_g) + \varepsilon^{ij} (|\omega_e|_g A_{ij} - (A_e)_{ij}) \end{aligned} \quad (2.31)$$

from (2.25)(2.26)(2.28), we have

$$\varepsilon^{ij} A_{ij} (1 - |\omega_e|_g) = \frac{1}{2} H v^i v^j h_{ij} \pm C|h|^2 |A| \quad (2.32)$$

and using (2.25)(2.26)(2.28)(2.29)(2.30) we have:

$$\begin{aligned}
& \varepsilon^{ij}(|\omega_e|A_{ij} - (A_e)_{ij}) \\
&= -\varepsilon^{ij}(\omega_e)_k \Gamma_{ij}^k \\
&= -\frac{1}{2}f^{ij}\omega_k g^{kl}(\bar{\nabla}_i h_{jl} + \bar{\nabla}_j h_{il} - \bar{\nabla}_l h_{ij}) \pm C|h||\bar{\nabla}h| \\
&= -f^{ij}v^l \bar{\nabla}_i h_{jl} + \frac{1}{2}f^{ij}v^l \bar{\nabla}_l h_{ij} \pm C|h||\bar{\nabla}h|
\end{aligned} \tag{2.33}$$

At last , we have

$$\begin{aligned}
H - H_e &= -f^{ik}h_{kl}f^{lj}A_{ij} + \frac{1}{2}Hv^i v^j h_{ij} - f^{ij}v^l \bar{\nabla}_i h_{jl} \\
&+ \frac{1}{2}f^{ij}v^l \bar{\nabla}_l h_{ij} \pm C|h||\bar{\nabla}h| \pm C|h|^2|A|
\end{aligned} \tag{2.34}$$

$$\begin{aligned}
\int_{\Sigma} H_e^2 d\mu_e &= (1 + O(r_0^{-q})) \int_{\Sigma} H_e^2 d\mu \\
&\leq (1 + O(r_0^{-q})) \left(\int_{\Sigma} H^2 d\mu + \int_{\Sigma} (H_e - H)^2 + 2|H(H_e - H)| d\mu \right) \\
&\leq (1 + O(r_0^{-q})) (16\pi + O(r_0^{-q})) + \int_{\Sigma} (H_e - H)^2 \\
&+ \left(\int_{\Sigma} H^2 d\mu \right)^{\frac{1}{2}} \left(\int_{\Sigma} (H_e - H)^2 d\mu \right)^{\frac{1}{2}}
\end{aligned} \tag{2.35}$$

$$\begin{aligned}
\int (H_e - H)^2 d\mu &\leq \int O(|x|^{-2q})|A|^2 + H^2 O(|x|^{-2q}) + O(|x|^{-2-2q}) d\mu \\
&\leq \int O(|x|^{-2q})H^2 + O(|x|^{-2q})|\dot{A}|^2 + O(|x|^{-2-2q}) d\mu \\
&= O(r_0^{-2q})
\end{aligned} \tag{2.36}$$

so we have

$$\int_{\Sigma} H_e^2 d\mu_e \leq 16\pi + O(r_0^{-q}) \tag{2.37}$$

On the other hand, by Euler formula,

$$K_e = \frac{1}{4}H_e^2 - \frac{1}{2}|\dot{A}_e|^2. \tag{2.38}$$

So we have

$$\int H_e^2 d\mu_e \geq 16\pi \tag{2.39}$$

which implies:

$$\int_{\Sigma} H_e^2 d\mu_e = 16\pi + O(r_0^{-q}) \tag{2.40}$$

Based on Michael and Simon, we have the following Sobolev inequality.

Lemma 2.4. Suppose that Σ is a constant mean curvature surface in an asymptotically flat end $(R^3 \setminus B_1(0), g)$ with $r_0(\Sigma)$ sufficiently large, and that $\int_{\Sigma} H^2 \leq C$. Then

$$\left(\int_{\Sigma} f^2 d\mu\right)^{\frac{1}{2}} \leq C\left(\int_{\Sigma} |\nabla f| d\mu + \int_{\Sigma} H|f| d\mu\right). \quad (2.41)$$

Proof. Note that it is valid for the surface in Euclidean Space. So by the uniform equivalence of the metric g and δ , we have:

$$\left(\int |f|^2 d\mu\right)^{\frac{1}{2}} \leq C\left(\int |f|^2 d\mu_e\right)^{\frac{1}{2}} \leq C\left(\int |\nabla f| + H|f| + |H - H_e||f| d\mu\right) \quad (2.42)$$

To bound the last term on the right, we have:

$$\begin{aligned} \int |H - H_e||f| d\mu &\leq \int O(|x|^{-q})|A||f| + O(|x|^{-q})H|f| \\ &\quad + O(|x|^{-1-q})|f| d\mu \\ &\leq O(r_0^{-q}) \int H|f| + \left(\int |\mathring{A}|^2 d\mu\right)^{\frac{1}{2}} O(r_0^{-q}) \|f\|_{L^2} \\ &\quad + O(r_0^{-q}) \|f\|_{L^2} \end{aligned} \quad (2.43)$$

So we can choose r_0 sufficiently large and get the desired result.

Lemma 2.5. Suppose that Σ is a constant mean curvature surfaces in an asymptotically flat end $(R^3 \setminus B_1(0), g)$ with $r_0(\Sigma)$ sufficiently large, then:

$$C_1 H^{-1} \leq \text{diam}(\Sigma) \leq C_2 H^{-1} \quad (2.44)$$

In particular, if the surface Σ separates the infinity from the compact part, then:

$$C_1 H^{-1} \leq r_1(\Sigma) \leq C_2 H^{-1} \quad (2.45)$$

Proof. We already know that:

$$\int_{\Sigma} H_e^2 d\mu_e = 16\pi + O(r_0^{-q}) \quad (2.46)$$

Then from [7] Lemma 1.1, we know that

$$\sqrt{\frac{2|\Sigma|_e}{F(\Sigma)}} \leq \text{diam}(\Sigma) \leq C\sqrt{|\Sigma|_e F(\Sigma)} \quad (2.47)$$

where $F(\Sigma) = \frac{1}{2} \int_{\Sigma} H_e^2$ is the Willmore functional and $|\Sigma|_e$ is the volume of Σ with respect to the Euclidean metric. But the Euclidean metric is uniformly equivalent to g , so we get the result.

Now to get the pointwise estimate for \mathring{A} , we use the Simons identity and the Moser's iteration argument.

Lemma 2.6. (Simons identity [11]) Suppose N is a hypersurface in a Riemannian manifold (M, g) , then the second fundamental form satisfies the following identity:

$$\begin{aligned} \Delta A_{ij} &= \nabla_i \nabla_j H + H A_{ik} A_{jk} - |A|^2 A_{ij} + H R_{3i3j} - A_{ij} R_{3k3k} + A_{jk} R_{klil} \\ &+ A_{ik} R_{kljl} - 2A_{lk} R_{iljk} + \bar{\nabla}_j R_{3kik} + \bar{\nabla}_k R_{3ijk} \end{aligned} \quad (2.48)$$

where R_{ijkl} and $\bar{\nabla}$ are the curvature and gradient operator of (M, g) , then from this we easily deduce for constant mean curvature surface we have the next inequality for \dot{A} :

$$\begin{aligned} -|\dot{A}|\Delta|\dot{A}| &\leq |\dot{A}|^4 + CH|\dot{A}|^3 + CH^2|\dot{A}|^2 + C|\dot{A}|^2|x|^{-2-q} \\ &+ CH|\dot{A}|^2|x|^{-2-q} + C|\dot{A}|^3|x|^{-3-q} \end{aligned} \quad (2.49)$$

We also need an inequality for $\nabla \dot{A}$ because we also want to estimate the higher derivative:

$$\begin{aligned} -|\nabla \dot{A}|\Delta|\nabla \dot{A}| &\leq C|\nabla \dot{A}|^2(|\dot{A}|^2 + H|\dot{A}| + H^2 + O(|x|^{-2-q})) \\ &+ |\nabla \dot{A}|((|\dot{A}|^2 + H|\dot{A}| + H^2)O(|x|^{-2-q}) + (|\dot{A}| + H)O(|x|^{-3-q}) + O(|x|^{-4-q})) \end{aligned} \quad (2.50)$$

Lemma 2.7.

$$\|\dot{A}^2\|_{L^2} + \|\nabla|\dot{A}|\|_{L^2} + \|\nabla \dot{A}\|_{L^2} + \|H|\dot{A}|\|_{L^2} \leq Cr_0^{-1-q} \quad (2.51)$$

Proof. See [2] Lemma 4.5

Then we can get the pointwise estimates for \dot{A} and $\nabla \dot{A}$.

Theorem 2.8. [9] Suppose that $(R^3 \setminus B_1(0), g)$ is an asymptotically flat end. Then there exist positive numbers σ_0, δ_0 such that for any constant mean curvature surface in the end, which separates the infinity from the compact part, we have:

$$|\dot{A}|^2(x) \leq C|x|^{-2} \int_{B_{\delta_0|x|}(x)} |\dot{A}|^2 d\mu + C|x|^{-2-2q} \leq C|x|^{-2} r_0^{-q} \quad (2.52)$$

$$|\nabla \dot{A}|^2(x) \leq C|x|^{-2} \int_{B_{\delta_0|x|}(x)} |\nabla \dot{A}|^2 d\mu + C|x|^{-4-2q} \leq C|x|^{-2} r_0^{-2-2q} \quad (2.53)$$

provided that $r_0 \geq \sigma_0$.

Proof. In the Sobolev inequality (2.41) we take $f = u^2$, then we get:

$$\begin{aligned} \left(\int_{\Sigma} u^4 d\mu \right)^{\frac{1}{2}} &\leq C \left(2 \int_{\Sigma} |u| |\nabla u| d\mu + \int_{\Sigma} H u^2 d\mu \right) \\ &\leq C \left(\int_{\Sigma} u^2 \right)^{\frac{1}{2}} \left(\int_{\Sigma} |\nabla u|^2 d\mu \right)^{\frac{1}{2}} + C \left(\int_{\text{supp}(u)} H^2 d\mu \right)^{\frac{1}{2}} \left(\int_{\Sigma} u^4 d\mu \right)^{\frac{1}{2}} \end{aligned} \quad (2.54)$$

Lemma 2.9. For any $\varepsilon > 0$, we can find a uniform δ_0 sufficiently small such that if for any $x \in \Sigma$, we have that:

$$\int_{B_{\delta_0|x|}(x)} H^2 \leq \varepsilon \quad (2.55)$$

Proof. In fact we need only to prove that there exist C

$$|B_{\delta_0|x|}(x)| \leq C\delta_0^2|x|^2 \quad (2.56)$$

because then,

$$H^2|B_{\delta_0|x|}(x)| \leq C\delta_0^2|x|^2 H^2 \leq C\delta_0^2 \quad (2.57)$$

From [7] the proof of lemma 1.1, we know that, for any $x \in \Sigma$, $B_\sigma(x)$ denotes the Euclidean ball of radius σ with center x in R^3 , $\Sigma_\sigma = \Sigma \cap B_\sigma(x)$, then there exists C such that for $0 < \sigma \leq \rho < \infty$

$$\sigma^{-2}|\Sigma_\sigma| \leq C(\rho^{-2}|\Sigma_\rho| + F(\Sigma_\rho)) \quad (2.58)$$

where $F(\Sigma_\rho)$ is the Willmore functional. C doesn't depend on Σ, σ, ρ .

Let $\rho \rightarrow \infty$, $\rho^{-2}|\Sigma_\rho| \rightarrow 0$, so we have:

$$\sigma^{-2}|\Sigma_\sigma| \leq CF(\Sigma) \leq C \quad (2.59)$$

so we prove the lemma.

So if $\text{supp}(u) \subset B_{\delta_0|x|}(x)$, we have the following scaling invariant Sobolev inequality:

$$\left(\int_\Sigma u^4 d\mu\right)^{\frac{1}{2}} \leq C\left(\int_\Sigma u^2\right)^{\frac{1}{2}}\left(\int_\Sigma |\nabla u|^2 d\mu\right)^{\frac{1}{2}} \quad (2.60)$$

Lemma 2.10. [9] Suppose that a nonnegative function $v \in L^2$ solves

$$-\Delta v \leq fv + h \quad (2.61)$$

on $B_{2R}(x_0)$, where

$$\int_{B_{2R}(x_0)} f^2 d\mu \leq CR^{-2} \quad (2.62)$$

and $h \in L^2(B_{2R}(x_0))$. And suppose that

$$\left(\int_\Sigma u^4 d\mu\right)^{\frac{1}{2}} \leq C\left(\int_\Sigma u^2\right)^{\frac{1}{2}}\left(\int_\Sigma |\nabla u|^2 d\mu\right)^{\frac{1}{2}} \quad (2.63)$$

holds for all u with support inside $B_{2R}(x_0)$. Then

$$\sup_{B_R(x_0)} v \leq CR^{-1}\|v\|_{L^2(B_{2R}(x_0))} + CR\|h\|_{L^2(B_{2R}(x_0))} \quad (2.64)$$

See [9] Lemma 2.6 for the proof of this lemma.

Then we find that:

$$\begin{aligned} -\Delta|\mathring{A}| &\leq (|\mathring{A}|^2 + H^2 + H|\mathring{A}| + C|x|^{-2-q})|\mathring{A}| + CH|x|^{-2-q} + C|x|^{-3-q} \\ &= f_1|\mathring{A}| + h_1 \end{aligned} \quad (2.65)$$

$$\begin{aligned} -\Delta|\nabla\mathring{A}| &\leq C|\nabla\mathring{A}|(|\mathring{A}|^2 + H|\mathring{A}| + H^2 + O(|x|^{-3})) \\ &\quad + ((|\mathring{A}|^2 + H|\mathring{A}| + H^2)O(|x|^{-3}) + (|\mathring{A}| + H)O(|x|^{-4}) + O(|x|^{-5})) \\ &= f_2|\nabla\mathring{A}| + h_2. \end{aligned} \quad (2.66)$$

We need to prove that $\|f_1\|_{L^2(B_{2\delta_0|x|}(x))}^2, \|f_2\|_{L^2(B_{2\delta_0|x|}(x))}^2 \leq C|x|^{-2}$, see [9] Theorem 2.5 for the proof. and it is easy to show that $\|h_1\|_{L^2(B_{2\delta_0|x|}(x))}^2 = O(|x|^{-4-2q})$ and $\|h_2\|_{L^2(B_{2\delta_0|x|}(x))}^2 = O(|x|^{-6-2q})$.

Remark 2.11. We can also do the same kind of estimate for $\nabla^2\mathring{A}$, where we need the third derivative of curvature. It is needed by the $C^{2,\alpha}$ convergence of the surface in the next section. This is the reason why we require the metric g to be smooth up to 5th order.

3 Blow down analysis

Now like [9], we blow down the surface in three different scales. First we consider

$$\tilde{N} = \frac{1}{2}HN = \left\{\frac{1}{2}Hx : x \in N\right\} \quad (3.1)$$

Suppose that there is a sequence of constant mean curvature surfaces $\{N_i\}$ such that

$$\lim_{i \rightarrow \infty} r_0(N_i) = \infty \quad (3.2)$$

we have known that

$$\lim_{i \rightarrow \infty} \int_{N_i} H_e^2 d\sigma = 16\pi \quad (3.3)$$

Hence, by the curvature estimates established in the previous section combining the proof of Theorem 3.1 in [7], we have

Lemma 3.1. Suppose that $\{N_i\}$ is a sequence of constant mean curvature surfaces in a given asymptotically flat end $(R^3 \setminus B_1(0), g)$ and that

$$\lim_{i \rightarrow \infty} r_0(N_i) = \infty. \quad (3.4)$$

And suppose that N_i separates the infinity from the compact part. Then, there is a subsequence of $\{\tilde{N}_i\}$ which converges in Gromov-Hausdorff distance to a

round sphere $S_1^2(a)$ of radius 1 and centered at $a \in R^3$. Moreover, the convergence is in $C^{2,\alpha}$ sense away from the origin.

Then, we use a smaller scale r_0 to blow down the surface

$$\widehat{N} = r_0(N)^{-1}N = \{r_0^{-1}x : x \in N\}. \quad (3.5)$$

Lemma 3.2. Suppose that $\{N_i\}$ is a sequence of constant mean curvature surfaces in a given asymptotically flat end $(R^3 \setminus B_1(0), g)$ and that

$$\lim_{i \rightarrow \infty} r_0(N_i) = \infty. \quad (3.6)$$

And suppose that

$$\lim_{i \rightarrow \infty} r_0(N_i)H(N_i) = 0. \quad (3.7)$$

Then there is a subsequence of $\{\widehat{N}_i\}$ converges to a 2-plane at distance 1 from the origin. Moreover the convergence is in $C^{2,\alpha}$ in any compact set of R^3 .

We must understand the behavior of the surfaces N_i in the scales between $r_0(N_i)$ and $H^{-1}(N_i)$. We consider the scale r_i such that

$$\lim_{i \rightarrow \infty} \frac{r_0(N_i)}{r_i} = 0 \quad \lim_{i \rightarrow \infty} r_i H(N_i) = 0 \quad (3.8)$$

and blow down the surfaces

$$\overline{N}_i = r_i^{-1}N = \{r_i^{-1}x : x \in N\}. \quad (3.9)$$

Lemma 3.3. Suppose that $\{N_i\}$ is a sequence of constant mean curvature surfaces in a given asymptotically flat end $(R^3 \setminus B_1(0), g)$ and that

$$\lim_{i \rightarrow \infty} r_0(N_i) = \infty \quad (3.10)$$

And suppose that r_i are such that

$$\lim_{i \rightarrow \infty} \frac{r_0(N_i)}{r_i} = 0 \quad \lim_{i \rightarrow \infty} r_i H(N_i) = 0 \quad (3.11)$$

Then there is a subsequence of $\{\overline{N}_i\}$ converges to a 2-plane at the origin in Gromov-Hausdorff distance. Moreover the convergence is $C^{2,\alpha}$ in any compact subset away from the origin.

4 Asymptotically analysis

First we revise Proposition 2.1 in [10]. We prove a different version. Let us denote:

$$\|u\|_{1,i}^2 = \int_{[(i-1)L, iL] \times S^1} |u|^2 + |\nabla u|^2 dt d\theta \quad (4.1)$$

Lemma 4.1. Suppose $u \in W^{1,2}(\Sigma, R^k)$ satisfies

$$\Delta u + A \cdot \nabla u + B \cdot u = h \quad (4.2)$$

in Σ , where $\Sigma = [0, 3L] \times S^1$. And suppose that L is given and large. Then there exists a positive number δ_0 such that if

$$|h|_{L^2(\Sigma)} \leq \delta_0 \max_{1 \leq i \leq 3} |u|_{1,i} \quad (4.3)$$

and

$$|A|_{L^\infty(\Sigma)} \leq \delta_0 \quad |B|_{L^\infty(\Sigma)} \leq \delta_0 \quad (4.4)$$

then,

- (a) $\|u\|_{1,3} \leq e^{-\frac{1}{2}L} \|u\|_{1,2}$ implies $\|u\|_{1,2} < e^{-\frac{1}{2}L} \|u\|_{1,1}$
- (b) $\|u\|_{1,1} \leq e^{-\frac{1}{2}L} \|u\|_{1,2}$ implies $\|u\|_{1,2} < e^{-\frac{1}{2}L} \|u\|_{1,3}$
- (c) If both $\int_{L \times S^1} u d\theta$ and $\int_{2L \times S^1} u d\theta \leq \delta_0 \max_{1 \leq i \leq 3} \|u\|_{1,i}$, then either $\|u\|_{1,2} < e^{-\frac{1}{2}L} \|u\|_{1,1}$ or $\|u\|_{1,2} < e^{-\frac{1}{2}L} \|u\|_{1,3}$

Proof. Suppose that $u \in W^{1,2}(\Sigma)$ and u is harmonic, we can deduce that u satisfies (a)(b)(c) with

- (c') If both $\int_{L \times S^1} u d\theta$ and $\int_{2L \times S^1} u d\theta = 0$, then either $\|u\|_{1,2} < e^{-\frac{1}{2}L} \|u\|_{1,1}$ or $\|u\|_{1,2} < e^{-\frac{1}{2}L} \|u\|_{1,3}$

A harmonic function u can be written as:

$$u = a_0 + b_0 t + \sum_{n=1}^{\infty} \{e^{nt} (a_n \cos n\theta + b_n \sin n\theta) + e^{-nt} (a_{-n} \cos n\theta + b_{-n} \sin n\theta)\} \quad (4.5)$$

Then it follows that:

$$\begin{aligned} \|u\|_{1,i}^2 &= 2\pi((a_0^2 + b_0^2)L + a_0 b_0 L^2(2i-1) + \frac{1}{3}b_0^2 L^3(3i^2 - 3i + 1)) \\ &+ \frac{\pi}{2} \sum_{n=1}^{\infty} \left\{ \frac{e^{2nL-1}}{n} (e^{2(i-1)nL} (a_n^2 + b_n^2) + e^{-2niL} (a_{-n}^2 + b_{-n}^2)) + 4L(a_n a_{-n} + b_n b_{-n}) \right\} \\ &+ \pi \sum_{n=1}^{\infty} \left\{ \frac{e^{2nL-1}}{n} (e^{2(i-1)nL} (n^2 a_n^2 + n^2 b_n^2) + e^{-2niL} (n^2 a_{-n}^2 + n^2 b_{-n}^2)) \right. \\ &\left. + 4L(n^2 a_n a_{-n} + n^2 b_n b_{-n}) \right\} \end{aligned} \quad (4.6)$$

$i = 1, 2, 3$

If L is fixed and sufficiently large, then we have

$$\|u\|_{1,2}^2 < \frac{1}{2} (e^L \|u\|_{1,3}^2 + e^{-L} \|u\|_{1,1}^2) \quad (4.7)$$

which implies (a). We get (b) in the same way. For (c'), we have $a_0 = b_0 = 0$ then we have

$$\|u\|_{1,2}^2 < \frac{1}{2}e^{-L}(\|u\|_{1,3}^2 + \|u\|_{1,1}^2) \quad (4.8)$$

which implies (c')

The second step is to pass limits. If the proposition were false, then one would have a sequence of $\delta_k \rightarrow 0$ and a sequence of solution u_k with $\|h_k\|_{L^2} \leq \delta_k$, $|A_k| \leq \delta_k$ and $|B_k| \leq \delta_k$ solves:

$$\Delta u_k + A_k \cdot \nabla u_k + B_k \cdot u_k = h_k \quad (4.9)$$

We may assume $\max_{1 \leq i \leq 3} \|u_k\|_{1,i} = 1$ otherwise we can normalize them. Then we know that there is a subsequence that converges to some $u \in W^{1,2}(\Sigma)$ weakly. And u is a harmonic function. From the interior $W^{2,p}$ estimate we know the convergence is strongly $W^{1,2}$ in I_2 , which implies that u is not trivially zero. Because, with the assumption of the proof by contradiction, the middle one is the largest.

And because $u_i \rightharpoonup u$ weakly in $W^{1,2}(\Sigma)$ sense. So $u_i \rightharpoonup u$ in $W^{1,2}(I_1)$ and $W^{1,2}(I_3)$ sense, then we have:

$$\liminf_{i \rightarrow \infty} \|u_i\|_{1,1} \geq \|u\|_{1,1}, \liminf_{i \rightarrow \infty} \|u_i\|_{1,3} \geq \|u\|_{1,3} \quad (4.10)$$

and

$$\lim_{i \rightarrow \infty} \|u_i\|_{1,2} = \|u\|_{1,2} \quad (4.11)$$

then u_i converges to some non-trivial harmonic function u which violates one of (a)(b) or (c), which proves the lemma.

From now on we assume $q = 1$.

Given a surface N in R^3 , recall from, for example, (8.5) in [6], that

$$\Delta_e v + |\nabla_e v|^2 v = \nabla_e H_e \quad (4.12)$$

where v is the Gauss map from $N \rightarrow S^2$. For the constant mean curvature surfaces in the asymptotically flat end $(R^3 \setminus B_1(0), g)$, we have

Lemma 4.2.

$$|\nabla_e H_e|(x) \leq C|x|^{-2}r_0^{-1} \quad (4.13)$$

Proof. Because the metric g and the Euclidean metric are uniformly equivalent. So we just prove that

$$|\nabla H_e|(x) \leq C|x|^{-2}r_0^{-1} \quad (4.14)$$

From (2.34), we know that:

$$\begin{aligned}
|\nabla H_e| &\leq |\bar{\nabla} h_{ij}| |A| + |h_{ij}| |A|^2 + |h_{ij}| |\nabla \dot{A}_{ij}| + H|A| |h_{ij}| + H|\bar{\nabla} h_{ij}| \\
&+ |A| |\bar{\nabla} h_{ij}| + |\bar{\nabla}^2 h| \\
&\leq |x|^{-2} r_0^{-1}
\end{aligned} \tag{4.15}$$

Suppose Σ is a constant mean curvature surface in the asymptotically flat end. Set

$$A_{r_1, r_2} = \{x \in \Sigma : r_1 \leq |x| \leq r_2\} \tag{4.16}$$

and A_{r_1, r_2}^0 stands for the standard annulus in R^2 . We are concerned with the behavior of v on $A_{Kr_0(\Sigma), sH^{-1}(\Sigma)}$ of Σ where K will be fixed large and s will be fixed small. The lemma below gives us a good coordinate on the surface.

Lemma 4.3. Suppose Σ is a constant mean curvature surface in a given asymptotically flat end $(R^3 \setminus B_1(0), g)$. Then, for any $\varepsilon > 0$ and L fixed and large, there are M, s and K such that, if $r_0 \geq M$ and $Kr_0(\Sigma) < r < sH^{-1}(\Sigma)$, then $(r^{-1}A_{r, eLr}, r^{-2}g_e)$ may be represented as $(A_{1, eL}^0, \bar{g})$ and

$$\|\bar{g} - |dx|^2\|_{C^1(A_{1, eL}^0)} \leq \varepsilon. \tag{4.17}$$

In other words, in the cylindrical coordinates $(S^1 \times [\log r, L + \log r, \bar{g}_c])$

$$\|\bar{g}_c - (dt^2 + d\theta^2)\|_{C^1(S^1 \times [\log r, L + \log r])} \leq \varepsilon \tag{4.18}$$

Proof. Suppose this is not true. Then we can assume that such K (or such s) cannot be found. Then by Lemma 3.2. for some $\varepsilon_0 > 0$, there is a sequence Σ_n with $r_0(\Sigma_n) \rightarrow \infty$, and $\tilde{l}_n \rightarrow \infty$, such that:

$$((Kr_0 e^{\tilde{l}_n L})^{-1} A_{Kr_0 e^{\tilde{l}_n L}, Kr_0 e^{(\tilde{l}_n + 1)L}}, (Kr_0 e^{\tilde{l}_n L})^{-2} g_e) \tag{4.19}$$

is not ε_0 close to $(A_{1, eL}^0, \bar{g})$.

By Lemma 3.1. We know that

$$\frac{Kr_0 e^{\tilde{l}_n L}}{sH^{-1}(\Sigma_n)} \rightarrow 0 \tag{4.20}$$

must hold because we have choose s sufficiently small.

So if we assume $r_n = Kr_0 e^{\tilde{l}_n L}$, we have:

$$\lim_{n \rightarrow \infty} \frac{r_n}{Kr_0} = \infty, \lim_{n \rightarrow \infty} \frac{r_n}{sH^{-1}} = 0 \tag{4.21}$$

We blow down the surface using r_n , and have a contradiction with Lemma 3.3. This proves the lemma.

Now consider the cylindrical coordinates (t, θ) on $(S^1 \times [\log Kr_0, \log sH^{-1}])$, then the tension field

$$|\tau(v)| = r^2 |\nabla_e H_e| \leq Cr_0^{-1} \quad (4.22)$$

for $t \in [\log Kr_0, \log sH^{-1}]$. Thus,

$$\int_{S^1 \times [t, t+L]} |\tau(v)|^2 dt d\theta \leq Cr_0^{-2} \quad (4.23)$$

Let I_i stand for $S^1 \times [\log Kr_0 + (i-1)L, \log Kr_0 + iL]$, and N_i stand for $I_{i-1} \cup I_i \cup I_{i+1}$. On Σ_n we assume $\log(sH^{-1}) - \log(Kr_0) = l_n L$. And like [10], first we prove that,

Lemma 4.4. For each $i \in [3, l_n - 2]$, there exists a geodesic γ such that

$$\int_{I_i} |\tilde{\nabla}(v - \gamma)|^2 dt d\theta \leq C(e^{-iL} + e^{-(l_n-i)L})s^2 + Cr_0^{-1} \quad (4.24)$$

where $\tilde{\nabla}$ is the gradient on $S^1 \times [\log(Kr_0), \log(sH^{-1})]$

Proof. By Theorem 2.8, we have

$$[v]_{C^\alpha(I_i)} \leq \|\tilde{\nabla}v\|_{L^\infty} \leq C(r_0^{-\frac{1}{2}} + s) \quad (4.25)$$

then if r_0 sufficiently large and s sufficiently small, we have $[v]_{C^\alpha(N_i)}$ is very small.

To apply the Lemma 4.1 to prove this lemma we choose to points P and Q on S^2 (the image of Gauss map) satisfying

$$\begin{aligned} |P - \frac{1}{2\pi} \int_{(i-1)L \times S^1} v d\theta| &\leq C \max_{(i-1)L \times S^1} |v - P|^2 \\ |Q - \frac{1}{2\pi} \int_{iL \times S^1} v d\theta| &\leq C \max_{iL \times S^1} |v - Q|^2 \end{aligned} \quad (4.26)$$

Note that S^2 is compact and smooth, so by (4.25) we can always find such P and Q and P, Q are very close. So there is a unique geodesic γ_i connecting P and Q whose velocity is sufficiently small.

So if we write down the equation satisfied by $v - \gamma_i$ on $S^1 \times [\log(Kr_0), \log(sH^{-1})]$

$$\tilde{\Delta}u + A \cdot \tilde{\nabla}u + B \cdot u = \tau \quad (4.27)$$

where $u = v - \gamma_i$, we have:

$$\begin{aligned} |A| &\leq C(|\tilde{\nabla}v| + |\tilde{\nabla}\gamma_i|) \leq \delta_0 \\ |B| &\leq C \min\{|\tilde{\nabla}v|^2, |\tilde{\nabla}\gamma_i|^2\} \leq \delta_0 \end{aligned} \quad (4.28)$$

If Lemma 4.1 (C') cannot be used, the only reason is that

$$\|v - \gamma_i\|_{1,i} \leq C\|\tau\|_{L^2(N_i)} \quad (4.29)$$

which implies

$$\int_{I_i} |\tilde{\nabla}(v - \gamma_i)|^2 dt d\theta \leq Cr_0^{-2} \quad (4.30)$$

which implied (4.24).

If Lemma 4.1 (C') can be used, then applying it for $u = v - \gamma_i$ over N_i , we have either

$$\|u\|_{1,i} < e^{-\frac{1}{2}L} \|u\|_{1,i-1} \quad (4.31)$$

or

$$\|u\|_{1,i} < e^{-\frac{1}{2}L} \|u\|_{1,i+1}. \quad (4.32)$$

Suppose the first one happens (without loss of generality). Then we may push this relation to the left because (4.28) hold regardless of t 's position. If the theorem can be used on N_{j+1} but not on N_j for some $j \geq 2$, then we have

$$\|u\|_{1,i} < e^{-\frac{1}{2}(i-j)L} \|u\|_{1,j} \leq Ce^{-\frac{1}{2}(i-j)L} r_0^{-1} \leq Cr_0^{-1}. \quad (4.33)$$

If the theorem can be used until I_2 , then we have

$$\begin{aligned} e^{\frac{L}{2}} \|u\|_{1,2} &\leq \|u\|_{1,1} = \left(\int_{I_1} u^2 dt d\theta \right)^{\frac{1}{2}} + \left(\int_{I_1} |\tilde{\nabla}u|^2 dt d\theta \right)^{\frac{1}{2}} \\ &\leq \left(\int_{I_2} u^2 dt d\theta \right)^{\frac{1}{2}} + \left(\int_{I_1} (u(t, \theta) - u(t+L, \theta))^2 dt d\theta \right)^{\frac{1}{2}} + \left(\int_{I_1} |\tilde{\nabla}u|^2 dt d\theta \right)^{\frac{1}{2}} \end{aligned} \quad (4.34)$$

So we have

$$\begin{aligned} (e^{\frac{L}{2}} - 1) \|u\|_{1,2} &\leq \left(\int_{I_1} \left(\int_0^L \left| \frac{\partial u}{\partial t}(t+s, \theta) \right|^2 ds \right) dt d\theta \right)^{\frac{1}{2}} + \left(\int_{I_1} |\tilde{\nabla}u|^2 dt d\theta \right)^{\frac{1}{2}} \\ &\leq \int_0^L \left(\int_{I_1} \left| \frac{\partial u}{\partial t}(t+s, \theta) \right|^2 dt d\theta \right)^{\frac{1}{2}} ds + \left(\int_{I_1} |\tilde{\nabla}u|^2 dt d\theta \right)^{\frac{1}{2}} \\ &\leq C \left(\int_{I_1 \cup I_2} |\tilde{\nabla}u|^2 dt d\theta \right)^{\frac{1}{2}} \\ &\leq C \left(\int_{I_1 \cup I_2} |\tilde{\nabla}v|^2 dt d\theta \right)^{\frac{1}{2}} + C \left(\int_{I_1 \cup I_2} |\tilde{\nabla}\gamma_i|^2 dt d\theta \right)^{\frac{1}{2}} \\ &\leq C(r_0^{-\frac{1}{2}} + s) \end{aligned} \quad (4.35)$$

So we have the estimate

$$\|u\|_{1,i} \leq Ce^{-\frac{i-2}{2}L} \|u\|_{1,2} \leq Ce^{-\frac{i}{2}L} (r_0^{-\frac{1}{2}} + s) \quad (4.36)$$

If $\|u\|_{1,i} < e^{-\frac{1}{2}L}\|u\|_{1,i+1}$ happens, we will have similarly

$$\|u\|_{1,i} \leq C e^{-\frac{l_n-i}{2}L} (r_0^{-\frac{1}{2}} + s) \quad (4.37)$$

Finally we get

$$\|u\|_{1,i} \leq C (e^{-\frac{i}{2}L} + e^{-\frac{l_n-i}{2}L})s + C r_0^{-\frac{1}{2}} \quad (4.38)$$

which implies (4.24).

Then to get the energy decay, we use the Hopf differential

$$\Phi = |\partial_t v|^2 - |\partial_\theta v|^2 - 2\sqrt{-1}\partial_t v \cdot \partial_\theta v \quad (4.39)$$

We know that the L^1 norm of Φ is invariant under conformal change of the coordinates. (t, θ) is the coordinate of $A_{Kr_0 e^{(i-2)L}, Kr_0 e^{(i+1)L}}$, we find another coordinate for it: set $r_i = Kr_0 e^{iL}$, then $(r_i^{-1} A_{Kr_0 e^{(i-2)L}, Kr_0 e^{(i+1)L}}, r_i^{-2} g_e)$ can be represented as $(A_{e^{-2L}, e^L}^0, \bar{g})$, where $\|\bar{g} - |dx|^2\|_{C^1(A_{e^{-2L}, e^L}^0)} \leq \varepsilon$. Assume this Euclidean coordinate is (x, y) , so:

$$\int_{S^1 \times [\log Kr_0 + (i-1)L, \log Kr_0 + iL]} |\Phi| dt d\theta = \int_{A_{e^{-L}, 1}^0} |\Phi| dx dy \quad (4.40)$$

To estimate the right hand side, we use the Cauchy integral formula on $\Omega = A_{e^{-2L}, e^L}^0$, and set $\Omega' = A_{e^{-L}, 1}^0$, for any $z \in \Omega'$

$$\Phi(v)(z) = \frac{1}{2\pi\sqrt{-1}} \int_{\partial\Omega} \frac{\Phi(w)}{w-z} dw + \frac{1}{2\pi\sqrt{-1}} \int_{\Omega} \frac{\partial\Phi(w)}{\partial\bar{w}} \frac{dw \wedge d\bar{w}}{w-z} \quad (4.41)$$

We know

$$|\partial_x v|, |\partial_y v| \leq CKr_0 e^{iL} |A| \leq CKr_0 e^{iL} (|x|^{-1} r_0^{-\frac{1}{2}} + r_1^{-1}) \leq C(r_0^{-\frac{1}{2}} + s e^{-(l_n-i)L}) \quad (4.42)$$

so we have:

$$\left| \frac{1}{2\pi\sqrt{-1}} \int_{\partial\Omega} \frac{\Phi(w)}{w-z} dw \right| \leq C(r_0^{-1} + s^2 e^{-2(l_n-i)L}) \quad (4.43)$$

For the second term, notice that by easy calculation

$$\frac{\partial\Phi(w)}{\partial\bar{w}} = \partial v \cdot \bar{\tau}(v) \quad (4.44)$$

where $\bar{\tau}(v)$ is the tension field under this coordinate. And

$$|\bar{\tau}(v)| \leq (Kr_0 e^{iL})^2 |\nabla_e H_e| \leq C r_0^{-1} \quad (4.45)$$

so we have:

$$\frac{1}{2\pi\sqrt{-1}} \int_{\Omega} \frac{\partial\Phi(w)}{\partial\bar{w}} \frac{dw \wedge d\bar{w}}{w-z} \leq Cr_0^{-1} \quad (4.46)$$

Then we get:

$$\int_{\Omega'} |\Phi| \leq C(r_0^{-1} + s^2 e^{-2(l_n-i)L}) \quad (4.47)$$

By direct calculation

$$\begin{aligned} & \int_{S^1 \times [Kr_0 e^{(i-1)L}, Kr_0 e^{iL}]} |\partial_t v|^2 dt d\theta \\ & \leq \int_{S^1 \times [Kr_0 e^{(i-1)L}, Kr_0 e^{iL}]} |\Phi| dt d\theta + \int_{S^1 \times [Kr_0 e^{(i-1)L}, Kr_0 e^{iL}]} |\partial_\theta v|^2 dt d\theta \end{aligned} \quad (4.48)$$

and we can get the estimate of $\int_{S^1 \times [Kr_0 e^{(i-1)L}, Kr_0 e^{iL}]} |\partial_\theta v|^2 dt d\theta$ directly by (4.24). So we get the estimate:

$$\int_{S^1 \times [Kr_0 e^{(i-1)L}, Kr_0 e^{iL}]} |\tilde{\nabla} v|^2 dt d\theta \leq C(e^{-iL} + e^{-(l_n-i)L})s^2 + Cr_0^{-1} \quad (4.49)$$

Proposition 4.5. Suppose that $\{\Sigma_n\}$ is a sequence of constant mean curvature surfaces in a given asymptotically flat end $(R^3 \setminus B_1(0), g)$ and that

$$\lim_{i \rightarrow \infty} r_0(\Sigma_n) = \infty \quad (4.50)$$

And suppose that

$$\lim_{n \rightarrow \infty} r_0(\Sigma_n)H(\Sigma_n) = 0 \quad (4.51)$$

Then there exist a large number K , a small number s and n_0 such that, when $n \geq n_0$,

$$\max_{I_i} |\tilde{\nabla} v| \leq C(e^{-\frac{i}{2}L} + e^{-\frac{(l_n-i)}{2}L})s + Cr_0^{-\frac{1}{2}} \quad (4.52)$$

where

$$I_i = S^1 \times [\log(Kr_0(\Sigma_n)) + (i-1)L, \log(Kr_0(\Sigma_n)) + iL] \quad (4.53)$$

and

$$i \in [0, l_n] \quad \log(Kr_0(\Sigma_n)) + l_n L = \log(sH^{-1}(\Sigma_n)) \quad (4.54)$$

Proof. We just use the interior estimate of the elliptic equation

$$\tilde{\Delta}v + |\tilde{\nabla}v|^2v = \tau \quad (4.55)$$

We know $\|\tilde{\nabla}v\|_\infty \leq C(r_0^{-\frac{1}{2}} + s)$, and $\|\tau\|_\infty \leq Cr_0^{-1}$. Assume that :

$$I_i \subset\subset \tilde{I}_i \subset\subset N_i \quad (4.56)$$

then for some $p > 2$

$$\begin{aligned} \sup_{I_i} |\tilde{\nabla}v| &\leq C\|\tilde{\nabla}v\|_{W^{1,p}(I_i)} \leq C(\|v\|_{L^p(\tilde{I}_i)} + r_0^{-1}) \leq C(\|v\|_{L^2(N_i)} + r_0^{-1}) \\ &\leq C(e^{-\frac{i}{2}L} + e^{-\frac{(l_n-i)}{2}L})s + Cr_0^{-\frac{1}{2}} \end{aligned} \quad (4.57)$$

This analysis improves our understanding of the blowdowns that we discussed in the previous section. Namely,

Corollary 4.6. Assume the same condition as the above proposition and in addition $\lim_{r_0 \rightarrow \infty} \frac{\log(r_1)}{r_0^{1/4}} = 0$. Then the limit plane in Lemma3.2 and Lemma3.3 are all orthogonal to the same vector a . In fact, we may choose s small and i large enough so that,

$$|v(x) + a| \leq \varepsilon \quad (4.58)$$

for all $x \in \Sigma_n$ and $|x| \leq sH^{-1}(\Sigma_n)$

Proof. We want to prove that

$$Osc_{B_{sH^{-1}} \cap \Sigma_n} v \quad (4.59)$$

is sufficiently small if $r_0(\Sigma_n)$ large and s small. We already know that

$$Osc_{B_{Kr_0} \cap \Sigma_n} v \quad (4.60)$$

is very small from Lemma 3.2, so we need only to prove that

$$Osc_{(B_{sH^{-1}} \setminus B_{Kr_0}) \cap \Sigma_n} v \quad (4.61)$$

is small.

From the proposition above we find that

$$\begin{aligned} Osc_{(B_{sH^{-1}} \setminus B_{Kr_0}) \cap \Sigma_n} v &\leq \sum_{i=1}^{l_n} Osc_{I_i} v \leq C \sum_{i=1}^{l_n} \sup_{I_i} |\tilde{\nabla}v| \\ &\leq C \sum_{i=1}^{l_n} ((e^{-\frac{i}{2}L} + e^{-\frac{(l_n-i)}{2}L})s + r_0^{-\frac{1}{2}}) \leq Cs + l_n r_0^{-\frac{1}{2}} \end{aligned} \quad (4.62)$$

From $C^{-1}r_1 \leq H^{-1} \leq Cr_1$ and the condition $\lim_{r_0 \rightarrow \infty} \frac{\log(r_1)}{r_0^{1/4}} = 0$, we have

$$l_n r_0^{-\frac{1}{2}} = L^{-1}(\log(sH^{-1}) - \log(Kr_0))r_0^{-\frac{1}{2}} \leq C \frac{\log r_1}{r_0^{\frac{1}{2}}} \rightarrow 0 \quad (4.63)$$

as $r_0 \rightarrow \infty$, so we prove the lemma.

Corollary 4.7. Assume the same condition as Proposition 4.5. Let $v_n = v(p_n)$ for some $p_n \in I_{\frac{l_n}{2}}$. Then

$$\sup_{I_i} |v - v_n| \leq C(e^{-\frac{1}{2}iL} + e^{-\frac{1}{4}l_n L})s + l_n r_0^{-\frac{1}{2}} \quad (4.64)$$

for $i \in [0, \frac{1}{2}l_n]$

$$\sup_{I_i} |v - v_n| \leq C(e^{-\frac{1}{4}l_n L} + e^{-\frac{1}{2}(l_n - i)L})s + l_n r_0^{-\frac{1}{2}} \quad (4.65)$$

for $i \in [\frac{1}{2}l_n, l_n]$

5 Harmonic Coordinates

We assume that the metric g can be expanded in the coordinate $\{x_i\}$ as

$$g_{ij} = \delta_{ij} + h_{ij} = \delta_{ij} + h_{ij}^1(\theta)/r + Q$$

where θ is the coordinate on the unit sphere S^2 , and $h_{ij}^1(\theta)$ is a function extended constantly along the radius direction. And Q satisfies

$$\sup r^{2+k} |\partial^k Q| \leq C \quad (5.1)$$

for $k = 0, 1, \dots, 5$

First, note that:

$$\begin{aligned} \Delta_g x_k &= \frac{1}{\sqrt{g}} \frac{\partial}{\partial x_i} (\sqrt{g} g^{ij} \frac{\partial}{\partial x_j} x_k) \\ &= \frac{\partial}{\partial x_i} g^{ik} + \frac{1}{2} g^{ik} g^{mn} g_{mn,i} \\ &= -g^{mn} \Gamma_{mn}^k = O(|x|^{-2}) \end{aligned} \quad (5.2)$$

Now our aim is to find asymptotically harmonic coordinate, i.e. some coordinate y^i such that $\Delta_g y^k = O(|x|^{-3})$

$$\begin{aligned} \Delta_g x^k &= -g^{jl} g^{ik} \frac{1}{2} \left(\frac{\partial}{\partial x_j} h_{li} + \frac{\partial}{\partial x_l} h_{ji} - \frac{\partial}{\partial x_i} h_{jl} \right) \\ &= -g^{jl} g^{ik} \frac{1}{2} \left(r^{-2} \left((h_{li,j}^1(\theta) - h_{li}^1(\theta) \frac{x_j}{r}) \right. \right. \\ &\quad \left. \left. + (h_{ji,l}^1(\theta) - h_{ji}^1(\theta) \frac{x_l}{r}) - (h_{jl,i}^1(\theta) - h_{jl}^1(\theta) \frac{x_i}{r}) \right) \right) + \partial Q \\ &= -g^{jl} g^{ik} \frac{1}{2} r^{-2} f_{li,j}^1(\theta) + O(|x|^{-3}) \end{aligned} \quad (5.3)$$

We also know that $g^{ij} = \delta^{ij} - h_{ij}^1(\theta)/r + O(r^{-2})$

Then :

$$\Delta_g x^k = -\frac{1}{2}r^{-2}f_{jkj}^1(\theta) + O(r^{-3}) \quad (5.4)$$

Suppose $0 = \xi_0 > \xi_1 \geq \xi_2 \geq \dots$ are the eigenvalues of $\Delta|_{S^2}$, and $A_n(\theta)$ are the corresponding orthonormal eigenvectors.

Set:

$$y^k = x^k + \sum_{n=0}^{\infty} f_n^k(r)A_n(\theta) \quad (5.5)$$

We have:

$$\Delta_g y^k = \Delta_g x^k + \sum_{n=0}^{\infty} \Delta_{R^3}(f_n^k(r)A_n(\theta)) + \sum_{n=0}^{\infty} (\Delta_g - \Delta_{R^3})(f_n^k(r)A_n(\theta)) \quad (5.6)$$

Solve the equation:

$$\Delta_g x^k + \sum_{n=0}^{\infty} \Delta_{R^3}(f_n^k(r)A_n(\theta)) = O(|x|^{-3}) \quad (5.7)$$

Assume

$$\frac{1}{2}f_{jkj}^1(\theta) = \sum_{n=0}^{\infty} \lambda_n^k A_n(\theta) \quad (5.8)$$

so we have:

$$\sum_{n=0}^{\infty} \Delta_{R^3}(f_n^k(r)A_n(\theta)) = r^{-2} \sum_{n=0}^{\infty} \lambda_n^k A_n(\theta) \quad (5.9)$$

$$\frac{1}{r^2}(2r f_n^{k'} + r^2 f_n^{k''} + f_n^k(r)\xi_n) = \lambda_n^k, n = 0, \dots, \infty \quad (5.10)$$

$$n = 0, \quad f_0^k = \lambda_0^k \log(r) \quad (5.11)$$

$$n > 0, \quad f_n^k = \frac{\lambda_n^k}{\xi_n} \quad (5.12)$$

and this solution satisfies that:

$$\sum_{n=0}^{\infty} (\Delta_g - \Delta_{R^3})(f_n^k(r)A_n(\theta)) = O(|x|^{-3}) \quad (5.13)$$

so if

$$y^k = x^k + \frac{1}{2\sqrt{\pi}}\lambda_0^k \log r + \sum_{n=1}^{\infty} \frac{\lambda_n^k}{\xi_n} A_n(\theta) \quad (5.14)$$

then we must have:

$$\Delta y^k = O(|x|^{-3}) \quad (5.15)$$

Note that

$$\Delta|_{S^2} \sum_{n=1}^{\infty} \frac{\lambda_n^k}{\xi_n} A_n(\theta) = \sum_{n=1}^{\infty} \lambda_n^k A_n(\theta) = \frac{1}{2} f_{jkj}^1(\theta) - \frac{1}{2} \overline{f_{jkj}^1(\theta)} \quad (5.16)$$

where $\overline{f_{jkj}^1(\theta)}$ is its mean value on the unit sphere.

Set

$$g_k^1(\theta) = \sum_{n=1}^{\infty} \frac{\lambda_n^k}{\xi_n} A_n(\theta) = \Delta^{-1} \left(\frac{1}{2} f_{jkj}^1(\theta) - \frac{1}{2} \overline{f_{jkj}^1(\theta)} \right) \quad (5.17)$$

$$\frac{\partial y^k}{\partial x^i} = \delta_{ik} + \frac{\lambda_0^k}{2\sqrt{\pi}} \frac{1}{r} \frac{x^i}{r} + g_k^1(\theta)_i \frac{1}{r} \quad (5.18)$$

$$\frac{\partial x^i}{\partial y^k} = \delta_{ik} + O(|x|^{-1}) \quad (5.19)$$

So we get:

$$\tilde{g}_{ij} = g\left(\frac{\partial}{\partial y^i}, \frac{\partial}{\partial y^j}\right) = \delta_{ij} + O(|x|^{-1}) \quad (5.20)$$

Suppose

$$\tilde{g}_{ij} = \delta_{ij} + \tilde{h}_{ij} \quad (5.21)$$

Now I want to discuss the ellipticity of \tilde{h}_{ij}

$$\tilde{h}_{ij} = h_{ij} - \frac{1}{2r\sqrt{\pi}} \left(\lambda_0^i \frac{x^j}{r} + \lambda_0^j \frac{x^i}{r} \right) - \frac{(g_{i,j}^1(\theta) + g_{j,i}^1(\theta))}{r} \quad (5.22)$$

Where $g_{i,j}^1(\theta)$ denotes the constant extension along the radius direction of function $\frac{\partial g_i^1(\theta)}{\partial x_j} |_{S^2}$

Example 5.1. : For the metric $g_{ij} = \delta_{ij} + \frac{\delta_{ij}}{r}$, we have:

$$\Delta_g x^k = -\frac{1}{2} \frac{x^k}{r^3} + O(|x|^{-3}) \quad (5.23)$$

We know that on S^2 , we have $\Delta|_{S^2} x^k = -2x^k$. So if we let:

$$y^k = x^k - \frac{1}{4} \frac{x^k}{r} \quad (5.24)$$

We have $\Delta_g y^k = O(|x|^{-3})$, then:

$$\frac{\partial y^k}{\partial x^i} = \delta_{ki} - \frac{1}{4} \left(\frac{\delta_{ki}}{r} - \frac{x^k x^i}{r^3} \right) \quad (5.25)$$

$$\tilde{h}_{ij} = \frac{3\delta_{ij}}{2r} - \frac{x^i x^j}{2r^3} + O(r^{-2}) \quad (5.26)$$

Lemma 5.2. Suppose in some coordinate $\{x^i\}$, $g_{ij} = \delta_{ij} + h_{ij}^1(\theta)/r + Q$, then for any $m > 2$ there exists $\varepsilon > 0$, if $\|h_{ij}^1(\theta) - \delta_{ij}(\theta)\|_{W^{m,2}(S^2)} \leq \varepsilon$ then in the asymptotically harmonic coordinate $\{y^i\}$ we get above, we have

$$\tilde{g}_{ij} = \delta_{ij} + \tilde{h}_{ij} \quad (5.27)$$

where $\tilde{h}_{ij} = O(|y|^{-1})$, and $|y|\tilde{h}_{ij}$ is uniformly elliptic.

Proof: We know easily from (5.18) that $\tilde{h}_{ij} = O(|x|^{-1})$ and that $\lim_{|x| \rightarrow \infty} \frac{|y|}{|x|} = 1$, then $\tilde{h}_{ij} = O(|y|^{-1})$. So we need only to prove that $|y|\tilde{h}_{ij}$ is uniformly elliptic.

First we know from $\|h_{ij}^1(\theta) - \delta_{ij}(\theta)\|_{W^{m,2}(S^2)} \leq \varepsilon$ that

$$\left\| \frac{1}{2} f_{jkj}^1(\theta) - \frac{1}{2} \frac{x^k}{r} \right\|_{W^{m-1,2}(S^2)} \leq C\varepsilon \quad (5.28)$$

Note that $\frac{1}{2} f_{jkj}^1(\theta) = \sum_{n=0}^{\infty} \lambda_n^k A_n(\theta)$ and x^k is an eigenvector of Δ_{S^2} , so we can assume that $A_1(\theta) = C_k x^k|_{S^2}$ without loss of generality.

$$\left\| \lambda_0^k A_0(\theta) + (\lambda_1^k C_k - \frac{1}{2}) x^k + \sum_{n=2}^{\infty} \lambda_n^k A_n(\theta) \right\|_{W^{m-1,2}(S^2)} \leq \varepsilon \quad (5.29)$$

so we get

$$|\lambda_0^k| \leq \varepsilon, (\lambda_1^k C_k - \frac{1}{2}) \leq \varepsilon, \sum_{n=2}^{\infty} (|\xi_n|^{\frac{m-1}{2}} \lambda_n^k)^2 \leq \varepsilon \quad (5.30)$$

Note that from (5.14)

$$\frac{\partial y^k}{\partial x^i} = \delta_{ik} + \frac{\lambda_0^k}{2\sqrt{\pi}} \frac{1}{r} \frac{x^i}{r} - \frac{1}{2} (\frac{1}{2} \pm \varepsilon) \left(\frac{\delta_{ik}}{r} - \frac{x_i x_k}{r^3} \right) + \sum_{n=2}^{\infty} \frac{\lambda_n^k}{\xi_n} \frac{\partial A_n(\theta)}{\partial x_i} \quad (5.31)$$

where the last term on the right can be estimated, for some $p > 0$

$$\begin{aligned} \left| \sum_{n=2}^{\infty} \frac{\lambda_n^k}{\xi_n} \frac{\partial A_n(\theta)}{\partial x_i} \right| &\leq \sum_{n=2}^{\infty} \frac{|\lambda_n^k|}{|\xi_n|} \frac{|\nabla_{S^2} A_n(\theta)|}{r} \\ &\leq \sum_{n=2}^{\infty} \frac{|\lambda_n^k|}{|\xi_n|} \frac{\|A_n(\theta)\|_{W^{2+p,2}}}{r} \\ &\leq \sum_{n=2}^{\infty} \frac{|\lambda_n^k|}{|\xi_n|} \frac{|\xi_n|^{1+\frac{p}{2}} \|A_n(\theta)\|_{L^2}}{r} \\ &\leq \frac{1}{r} \sum_{n=2}^{\infty} |\lambda_n^k| |\xi_n|^{\frac{m-1}{2}} |\xi_n|^{\frac{p-m+1}{2}} \\ &\leq \frac{1}{r} \left(\sum_{n=2}^{\infty} (|\lambda_n^k| |\xi_n|^{\frac{m-1}{2}})^2 \right)^{\frac{1}{2}} \left(\sum_{n=2}^{\infty} |\xi_n|^{p-m+1} \right)^{\frac{1}{2}} \end{aligned} \quad (5.32)$$

let $p = \frac{m-2}{2}$, then from $\xi_n = O(n)$ we have

$$\sum_{n=2}^{\infty} |\xi_n|^{p-m+1} \leq C \quad (5.33)$$

so

$$\left| \sum_{n=2}^{\infty} \frac{\lambda_n^k}{\xi_n} \frac{\partial A_n(\theta)}{\partial x_i} \right| \leq \frac{C\varepsilon}{r} \quad (5.34)$$

then we have:

$$\frac{\partial y^k}{\partial x^i} = \delta_{ik} - \frac{1}{4} \left(\frac{\delta_{ik}}{r} - \frac{x_i x_k}{r^3} \right) + \frac{C\varepsilon}{r} \quad (5.35)$$

so we can deduce that:

$$\tilde{h}_{ij} = h_{ij} + \frac{\delta_{ij}}{2r} - \frac{x^i x^j}{2r^3} + \frac{C\varepsilon}{r} \quad (5.36)$$

because $|h_{ij}^1(\theta) - \delta_{ij}(\theta)|_{W^{m,2}(S^2)} \leq \varepsilon$, we have rh_{ij} is uniformly elliptic. And the eigenvalues of $\frac{x^i x^j}{r^2}$ are between 0 and 1, so $|y|\tilde{h}_{ij}$ is uniformly elliptic from $\lim_{r \rightarrow \infty} \frac{|y|}{r} = 1$ for ε sufficiently small.

So all the analysis in Section 2,3,4 can be done in the asymptotically harmonic coordinate $\{y_i\}$.

Lemma 5.3. In the asymptotically harmonic coordinate $\{y^i\}$, we have that

$$-\frac{1}{2} \Delta_g \log |\tilde{g}| = R(g) + O(|y|^{-4}) \quad (5.37)$$

Proof. From direct calculation we have

$$R(g) = \tilde{g}^{jk} \tilde{g}^{il} \tilde{g}_{ml} \left(\frac{\partial \tilde{\Gamma}_{jk}^m}{\partial y^i} - \frac{\partial \tilde{\Gamma}_{ik}^m}{\partial y^j} \right) + O(|y|^{-4}) \quad (5.38)$$

$$\begin{aligned} \tilde{g}^{jk} \tilde{g}^{il} \tilde{g}_{ml} \frac{\partial \tilde{\Gamma}_{jk}^m}{\partial y^i} &= \tilde{g}^{il} \tilde{g}_{ml} \frac{\partial (\tilde{g}^{jk} \tilde{\Gamma}_{jk}^m)}{\partial y^i} + O(|y|^{-4}) \\ &= -\tilde{g}^{il} \tilde{g}_{ml} \frac{\partial \Delta_g y^m}{\partial y^i} + O(|y|^{-4}) = O(|y|^{-4}) \end{aligned} \quad (5.39)$$

$$\begin{aligned} -\tilde{g}^{jk} \tilde{g}^{il} \tilde{g}_{ml} \frac{\partial \tilde{\Gamma}_{ik}^m}{\partial y^j} &= -\frac{1}{2} \tilde{g}^{jk} \tilde{g}^{ip} \frac{\partial^2 \tilde{g}_{ip}}{\partial y^j \partial y^k} + O(|y|^{-4}) \\ &= -\frac{1}{2} \Delta_g \log |\tilde{g}| + O(|y|^{-4}) \end{aligned} \quad (5.40)$$

so we prove the lemma.

Corollary 5.4. If in addition $R = O(|x|^{-3-\tau})$ for some $\tau > 0$, then in the asymptotically harmonic coordinate $\{y^i\}$, we have

$$\sum_{i=1}^3 \tilde{h}_{ii} = 8m(g)/|y| + o(|y|^{-1-\frac{\tau}{2}}) \quad (5.41)$$

Proof: First we know that

$$\lim_{|x| \rightarrow \infty} \frac{|y|}{|x|} = 1, \quad (5.42)$$

then from the lemma above that in the coordinate $\{y^i\}$, we have

$$\Delta_g \log |\tilde{g}| = O(|y|^{-3-\tau}) \quad (5.43)$$

We know that

$$\log |\tilde{g}| = O(|y|^{-1}) \quad (5.44)$$

From the theory of harmonic functions in R^n , we have there exist some constant C such that:

$$\log |\tilde{g}| = \frac{C}{|y|} + o(|y|^{-1-\frac{\tau}{2}}) \quad (5.45)$$

From Bartnik's result, we know the mass is invariant under the change of coordinates because $R(g) \in L^1$.

$$m(g) = \lim_{R \rightarrow \infty} \frac{1}{16\pi} \int_{s_R} (\tilde{h}_{ij,j} - \tilde{h}_{jj,i}) v_g^i d\mu \quad (5.46)$$

Now we have

$$\begin{aligned} \tilde{g}_{ik,k} - \frac{1}{2} \tilde{g}_{kk,i} &= \tilde{g}^{ij} \tilde{g}^{kl} (\tilde{g}_{jk,l} - \frac{1}{2} \tilde{g}_{kl,j}) + O(|y|^{-3}) \\ &= -\Delta_g y^i + O(|y|^{-3}) = O(|y|^{-3}) \end{aligned} \quad (5.47)$$

So we have:

$$\begin{aligned} m(g) &= \lim_{R \rightarrow \infty} \frac{1}{16\pi} \int_{s_R} (-\frac{1}{2} \tilde{h}_{jj,i}) v_g^i d\mu \\ &= -\lim_{R \rightarrow \infty} \frac{1}{32\pi} \int_{s_R} \frac{\partial \log |\tilde{g}|}{\partial y^i} v_g^i d\mu \\ &= \lim_{R \rightarrow \infty} \frac{1}{32\pi} \int_{s_R} \frac{C y^i}{|y|^3} v_g^i d\mu \\ &= \frac{C}{8} \end{aligned} \quad (5.48)$$

So we get the result by easy calculation.

Remark 5.5. In fact we can replace the constraint equation by the condition

$$R = O(|x|^{-3-\tau}) \quad (5.49)$$

for some $\tau > 0$.

6 Proof of the Theorem

Now let's prove Theorem 1.5.

First recall that, for any surface Σ embedded in \mathbb{R}^3 and any given vector $b \in \mathbb{R}^3$, one has

$$\int_{\Sigma} H_e \langle v_e \cdot b \rangle_e d\mu_e = 0 \quad (6.1)$$

where H_e and v_e denote the mean curvature and normal vector field with respect to the Euclidean metric.

On the other hand, if Σ is a constant mean curvature surface in the asymptotically flat end, then

$$\int_{\Sigma} H \langle v_e \cdot b \rangle_e d\mu_e = 0 \quad (6.2)$$

So we have

$$\int_{\Sigma} (H - H_e) \langle v_e \cdot b \rangle_e d\mu_e = 0 \quad (6.3)$$

From now on, our calculation is in the coordinate $\{x^i\}$, which is assumed to be the asymptotically harmonic coordinate. We have calculated $H - H_e$, so we have

$$\begin{aligned} \int_{\Sigma} (H - H_e) \langle v_e \cdot b \rangle_e d\mu_e &= \int_{\Sigma} (-f^{ik} h_{kl} f^{lj} A_{ij} + \frac{1}{2} H v^i v^j h_{ij} - f^{ij} v^l \bar{\nabla}_i h_{jl} \\ &+ \frac{1}{2} f^{ij} v^l \bar{\nabla}_l h_{ij} \pm C|h| |\bar{\nabla} h| \pm C|h|^2 |A|) \langle v_e \cdot b \rangle_e d\mu_e \end{aligned} \quad (6.4)$$

We assume that there exists a sequence of constant mean curvature surfaces Σ_n with

$$\lim_{n \rightarrow \infty} r_0(\Sigma_n) = \infty \quad \lim_{n \rightarrow \infty} H(\Sigma_n) r_0(\Sigma_n) = 0 \quad (6.5)$$

otherwise we have get the result from the uniqueness theorem of Lan-Hsuan Huang. So we can choose s sufficiently small and K sufficiently large with $sH^{-1} > Kr_0$ for r_0 sufficiently large.

We know that

$$|h| = O(|x|^{-1}), |\bar{\nabla} h| = O(|x|^{-2}), |A| \leq CH + C|\dot{A}| \quad (6.6)$$

from the estimate

$$|\dot{A}| \leq r_0^{-\frac{1}{2}} O(|x|^{-1}) \quad (6.7)$$

we have

$$\begin{aligned} & \left| \int_{\Sigma} (\pm C|h| |\bar{\nabla} h| \pm C|h|^2 |A|) \langle v_e \cdot b \rangle_e d\mu_e \right| \leq C \int_{\Sigma} (H|x|^{-2} + |x|^{-3}) \\ &= O(r_0^{-1}) \end{aligned} \quad (6.8)$$

by the estimates in Section 2.

Now we calculate other terms in (6.4)

$$\begin{aligned}
& \int_{\Sigma_n} -f^{ij}v^l(\bar{\nabla}_i h_{jl})v^m b^m d\mu_e \\
= & \frac{1}{2} \int_{\Sigma_n} (f^{ij}h_{jk}f^{kl}A_{li} - Hv^jv^l h_{jl})v^m b^m d\mu_e + \frac{1}{2} \int_{\Sigma_n} f^{ij}v^l h_{jl}A_{ik}f^{km}b^m d\mu_e \\
& - \frac{1}{2} \int_{\Sigma_n} f^{ij}v^l(\bar{\nabla}_i h_{jl})v^m b^m d\mu_e \tag{6.9}
\end{aligned}$$

because $d\mu_e = (1 + O(r^{-1}))d\mu$, $v_e = (1 + O(r^{-1}))v$ and $\langle v_e \cdot b \rangle_e = \langle v \cdot b \rangle_g + O(r^{-1})$.

So we have

$$\begin{aligned}
\int_{\Sigma_n} (H - H_e) \langle v_e \cdot b \rangle_e d\mu_e = & \int_{\Sigma_n} -\frac{1}{2}f^{ik}h_{kl}f^{lj}A_{ij}v^m b^m + f^{ij}v^l h_{jl}A_{ik}f^{km}b^m \\
& - \frac{1}{2}f^{ij}v^l \bar{\nabla}_i h_{jl}v^m b^m + \frac{1}{2}f^{ij}v^l \bar{\nabla}_l h_{ij}v^m b^m + O(r_0^{-1})d\bar{\mu} \tag{6.10}
\end{aligned}$$

Note that

$$A_{ij} = \mathring{A}_{ij} + \frac{f_{ij}}{2}H, \quad \sup |\mathring{A}| \leq r_0^{-\frac{1}{2}}O(|x|^{-1}) \tag{6.11}$$

So we have

$$\begin{aligned}
\int_{\Sigma_n} (H - H_e) \langle v_e \cdot b \rangle_e d\mu_e = & \int_{\Sigma_n} -\frac{H}{4}f^{kl}h_{kl}v^m b^m + \frac{H}{4}f^{jm}h_{jl}v^l b^m \\
& + \frac{1}{2}f^{ij}(\bar{\nabla}_l h_{ij})v^l v^m b^m - \frac{1}{2}f^{ij}(\bar{\nabla}_i h_{jl})v^l v^m b^m \\
& \pm C \int_{\Sigma_n} |x|^{-2}r_0^{-\frac{1}{2}} + O(r_0^{-1}) \tag{6.12}
\end{aligned}$$

In this case we calculate

$$\int_{\Sigma_n} |x|^{-2}r_0^{-\frac{1}{2}}d\mu_e \tag{6.13}$$

We divide the integral into three parts:

$$\int_{\Sigma_n} |x|^{-2}r_0^{-\frac{1}{2}} = \int_{\Sigma_n \cap B_{s_{H-1}}^c(0)} + \int_{\Sigma_n \cap B_{Kr_0}(0)} + \int_{\Sigma_n \cap (B_{s_{H-1}} \setminus B_{Kr_0})} |x|^{-2}r_0^{-\frac{1}{2}} \tag{6.14}$$

Then by the blowdown results in Section 3 we have

$$\int_{\Sigma_n \cap B_{s_{H-1}}^c(0)} |x|^{-2}r_0^{-\frac{1}{2}}d\mu_e = \int_{\tilde{\Sigma}_n \cap B_s^c(0)} |\tilde{x}|^{-2}r_0^{-\frac{1}{2}}d\tilde{\mu} \leq Cr_0^{-\frac{1}{2}} \tag{6.15}$$

$$\int_{\Sigma_n \cap B_{Kr_0}(0)} |x|^{-2} r_0^{-\frac{1}{2}} d\mu_e = \int_{\widehat{\Sigma}_n \cap B_k(0)} |\widehat{x}|^{-2} r_0^{-\frac{1}{2}} d\widehat{\mu} \leq C r_0^{-\frac{1}{2}} \quad (6.16)$$

$$\begin{aligned} \int_{\Sigma_n \cap (B_{sH^{-1}} \setminus B_{Kr_0})} |x|^{-2} r_0^{-\frac{1}{2}} d\mu_e &= \sum_{i=0}^n \int_{\Sigma_n \cap (B_{Kr_0 e^{4iL}} \setminus B_{Kr_0 e^{4(i-1)L})}} |x|^{-2} r_0^{-\frac{1}{2}} d\mu_e \\ &\leq C \sum_{i=0}^n \int_{B_{e^{4L}} \setminus B_1} |\bar{x}|^{-2} r_0^{-\frac{1}{2}} d\bar{\mu} \leq C r_0^{-\frac{1}{2}} l_n L \end{aligned} \quad (6.17)$$

where $e^{l_n L} K r_0 = sH^{-1}$
so if

$$\lim_{r_0 \rightarrow 0} \frac{|\log H|}{r_0^{\frac{1}{2}}} = 0 \quad (6.18)$$

in other words

$$\lim_{r_0 \rightarrow 0} \frac{|\log r_1|}{r_0^{\frac{1}{2}}} = 0 \quad (6.19)$$

we have

$$\int_{\Sigma} |x|^{-2} r_0^{-\frac{1}{2}} d\bar{\mu} \rightarrow 0 \quad (6.20)$$

as $r_0 \rightarrow \infty$

From the property of the asymptotically harmonic coordinate

$$g^{ij} h_{ij} = \frac{8m(g)}{r} + o(r^{-1-\frac{\varepsilon}{2}}) \quad (6.21)$$

$$g^{kl} (g_{ik,l} - \frac{1}{2} g_{kl,i}) = O(|x|^{-3}) \quad (6.22)$$

$$\begin{aligned} &\int_{\Sigma_n} -\frac{H}{4} f^{kl} h_{kl} v^m b^m + \frac{H}{4} f^{jm} h_{jl} v^l b^m + \frac{1}{2} f^{ij} (\bar{\nabla}_i h_{ij} - \bar{\nabla}_i h_{ji}) v^l v^m b^m \\ &= \int_{\Sigma_n} -\frac{H}{4} g^{kl} h_{kl} v^m b^m + \frac{H}{4} g^{jm} h_{jl} v^l b^m \\ &\quad + \frac{1}{2} g^{ij} (\bar{\nabla}_i h_{ij} - \bar{\nabla}_i h_{ji}) v^l v^m b^m + O(|r_0|^{-1}) \\ &= -2m(g) \int_{\Sigma_n} \left(\frac{H}{r} \langle v_e \cdot b_e \rangle_e + \frac{\langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e}{r^3} \right) \\ &\quad + \int_{\Sigma_n} \frac{H}{4} h_{ml} v^l b^m + o(1). \end{aligned} \quad (6.23)$$

So we have:

$$\lim_{n \rightarrow \infty} (-2m(g) \int_{\Sigma_n} \left(\frac{H}{r} \langle v_e \cdot b \rangle_e + \frac{\langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e}{r^3} \right) + \int_{\Sigma_n} \frac{H}{4} h_{ml} v^l b^m) = 0 \quad (6.24)$$

Note that:

$$h_{ml}v^l = (h_{ml} - \frac{tr(h)}{2}\delta_{ml})v^l + \frac{tr(h)}{2}v^m \quad (6.25)$$

where $tr(h) = g^{ij}h_{ij}$

Assume that the three eigenvalues of h_{ml} are

$$\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq 0 \quad (6.26)$$

For $p \in \Sigma$ fixed, choose coordinate properly such that

$$h_{ml} - \frac{tr(h)}{2}\delta_{ml} \quad (6.27)$$

can be written as

$$\begin{pmatrix} \lambda_1 - \frac{tr(h)}{2} & 0 & 0 \\ 0 & \lambda_2 - \frac{tr(h)}{2} & 0 \\ 0 & 0 & \lambda_3 - \frac{tr(h)}{2} \end{pmatrix} \quad (6.28)$$

Assume $v = (\tilde{v}^1, \tilde{v}^2, \tilde{v}^3)$, and $(\tilde{v}^1)^2 + (\tilde{v}^2)^2 + (\tilde{v}^3)^2 = 1$. Then we have

$$\sum_{i=1}^3 ((\lambda_i - \frac{tr(h)}{2})\tilde{v}^i)^2 = \frac{(tr(h))^2}{4} - \sum_{i=1}^3 \lambda_i (tr(h) - \lambda_i)(\tilde{v}^i)^2 \quad (6.29)$$

Because of the uniformly ellipticity we have there exists $C > 0$, such that

$$\frac{th(h)}{C} \leq \lambda_3 \leq \lambda_2 \leq \lambda_1 \leq (1 - \frac{1}{C})th(h) \quad (6.30)$$

so

$$\lambda_i(th(h) - \lambda_i) \geq \frac{1}{C}(1 - \frac{1}{C})(tr(h))^2 \quad (6.31)$$

hence

$$\sum_{i=1}^3 ((\lambda_i - \frac{tr(h)}{2})\tilde{v}^i)^2 \leq (\frac{1}{4} - \frac{1}{C}(1 - \frac{1}{C}))(tr(h))^2 \quad (6.32)$$

$$\begin{aligned} \int_{\Sigma_n} \frac{H}{4} h_{ml}v^l b^m &= \int_{\Sigma_n} \frac{H}{4} (\frac{tr(h)}{2} \langle v_e \cdot b \rangle_e + (h_{ml} - \frac{tr(h)}{2}\delta_{ml})v^l b^m) \\ &\leq \int_{\Sigma_n} \frac{Htr(h)}{4} (\frac{1}{2} \langle v_e \cdot b \rangle_e + \sqrt{\frac{1}{4} - \frac{1}{C}(1 - \frac{1}{C})}) \\ &= \int_{\Sigma_n} \frac{Hm(g)}{r} (\langle v_e \cdot b \rangle_e + 1 - \frac{2}{C}) \end{aligned} \quad (6.33)$$

so we have

$$\begin{aligned} \int_{\Sigma_n} (H - H_e) \langle v_e \cdot b \rangle_e &\leq -m \int_{\Sigma_n} \frac{H}{r} \langle v_e \cdot b \rangle_e \\ &+ \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e + \left(1 - \frac{2}{C}\right) m(g) \int_{\Sigma_n} \frac{H}{r} d\mu_e + o(1) \end{aligned} \quad (6.34)$$

as $n \rightarrow \infty$

From Lemma 3.1, we have $\frac{H}{2}\Sigma_n$ subconverges to some sphere $S_1^2(a)$ with $|a| = 1$. Now we choose $b = -a$. Then from the calculation in [9], we have

$$-m(g) \int_{\Sigma_n} \frac{H}{r} \langle v_e \cdot b \rangle_e \rightarrow -\frac{8}{3}\pi m(g) \quad (6.35)$$

$$-m(g) \int_{\Sigma_n} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e \rightarrow -\frac{16}{3}\pi m(g) \quad (6.36)$$

$$\left(1 - \frac{2}{C}\right) m(g) \int_{\Sigma_n} \frac{H}{r} \rightarrow \left(1 - \frac{2}{C}\right) 8\pi m(g) \quad (6.37)$$

as $n \rightarrow \infty$

Because there is a little difference from [9], we prove them again. We notice from Lemma 3.1, we have $\frac{H}{2}\Sigma_n$ subconverges to some sphere $S_1(a)$ with $|a| = 1$, and the first and third integral converges to $-m(g) \int_{S_1(a)} \frac{2}{r} \langle v_e \cdot b \rangle_e = -\frac{8}{3}\pi m(g)$ and $\left(1 - \frac{2}{C}\right) m(g) \int_{S_1(a)} \frac{2}{r} = \left(1 - \frac{2}{C}\right) 8\pi m(g)$ respectively.

To deal with the (6.36), first we notice that

$$\int_{S^2(a)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e = \frac{4}{3}\pi \quad (6.38)$$

then we break up the integral (6.36) into three parts.

$$\begin{aligned} &\int_{\Sigma_n} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \\ &= \int_{\Sigma_n \cap B_{sH^{-1}}^c(0)} + \int_{\Sigma_n \cap B_{K\tau_0}(0)} + \int_{\Sigma_n \cap B_{sH^{-1}} \setminus B_{K\tau_0}} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \end{aligned} \quad (6.39)$$

Then

$$\begin{aligned} &\lim_{n \rightarrow \infty} \int_{\Sigma_n \cap B_{sH^{-1}}^c(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \\ &= \int_{S^2(a) \cap B_s^c} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \end{aligned} \quad (6.40)$$

and

$$\begin{aligned}
& \lim_{n \rightarrow \infty} \int_{\Sigma_n \cap B_{Kr_0}(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \\
&= \int_{P \cap B_K(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e, \tag{6.41}
\end{aligned}$$

where P is the limit plane in Lemma 3.2. From Corollary 4.6, we know the normal vector of P is v_e . Then due to an easy calculation we know

$$\int_P \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e = 4\pi \tag{6.42}$$

From the divergence theorem we have

$$\int_{\Sigma_n} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e = 8\pi \tag{6.43}$$

for any n and

$$\int_{S^2(a)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e = 4\pi \tag{6.44}$$

because the origin is on the sphere $S^2(a)$. Since

$$\lim_{n \rightarrow \infty} \int_{\Sigma_n \cap B_{sH-1}^c(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e = \int_{S^2(a) \cap B_s^c(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e \tag{6.45}$$

$$\lim_{n \rightarrow \infty} \int_{\Sigma_n \cap B_{Kr_0}(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e = \int_{P \cap B_K(0)} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e \tag{6.46}$$

and

$$\int_P \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e = 4\pi \tag{6.47}$$

then we have

$$\lim_{s \rightarrow 0, K \rightarrow \infty} \limsup_{n \rightarrow \infty} \left| \int_{\Sigma_n \cap (B_{sH-1} \setminus B_{Kr_0})} \frac{2}{r^3} \langle x \cdot v_e \rangle_e d\mu_e \right| = 0 \tag{6.48}$$

Now we want to prove that

$$\lim_{s \rightarrow 0, K \rightarrow \infty} \limsup_{n \rightarrow \infty} \left| \int_{\Sigma_n \cap (B_{sH-1} \setminus B_{Kr_0})} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \right| = 0 \tag{6.49}$$

We use Lemma 4.7 to get (6.49) from (6.48), but there is a bit difference from [9].

$$\begin{aligned}
& \int_{\Sigma_n \cap (B_{sH^{-1}} \setminus B_{Kr_0})} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle v_e \cdot b \rangle_e d\mu_e \\
&= \langle v_n \cdot b \rangle_e \int_{\Sigma_n \cap (B_{sH^{-1}} \setminus B_{Kr_0})} \frac{2}{r^3} \langle v_e \cdot b \rangle_e d\mu_e \\
&+ \int_{\Sigma_n \cap (B_{sH^{-1}} \setminus B_{Kr_0})} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle (v_e - v_n) \cdot b \rangle_e d\mu_e \quad (6.50)
\end{aligned}$$

The first term will converge to 0. For the second term, we deal with it in the cylinder coordinate in Section 4:

$$\begin{aligned}
& \left| \int_{\Sigma_n \cap (B_{sH^{-1}} \setminus B_{Kr_0})} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle (v_e - v_n) \cdot b \rangle_e d\mu_e \right| \\
&= \left| \sum_{j=1}^{l_n} \int_{A_{Kr_0 e^{(j-1)L}, Kr_0 e^{jL}}} \frac{2}{r^3} \langle x \cdot v_e \rangle_e \langle (v_e - v_n) \cdot b \rangle_e d\mu_e \right| \\
&\leq C \sum_{j=1}^{l_n} L \max_{I_j} |v_e - v_n| \\
&= C \sum_{j=1}^{l_n/2} L \max_{I_j} |v_e - v_n| + C \sum_{j=l_n/2+1}^{l_n} L \max_{I_j} |v_e - v_n| \quad (6.51)
\end{aligned}$$

From Lemma 4.7

$$\begin{aligned}
& CL \sum_{i=1}^{l_n/2} \sup_{I_i} |v - v_n| + CL \sum_{i=l_n/2+1}^{l_n} \sup_{I_i} |v - v_n| \\
&\leq C(l_n e^{-\frac{1}{4}l_n L} + C)s + l_n^2 r_0^{-\frac{1}{2}} \quad (6.52)
\end{aligned}$$

But from the condition

$$\lim_{n \rightarrow \infty} \frac{\log(r_1(\Sigma_n))}{r_0(\Sigma_n)^{1/4}} = 0 \quad (6.53)$$

we know

$$\lim_{n \rightarrow \infty} l_n^2 r_0^{-\frac{1}{2}} = \lim_{n \rightarrow \infty} \left(\frac{L^{-1}(\log sH^{-1} - \log Kr_0)}{r_0^{\frac{1}{4}}} \right)^2 = 0 \quad (6.54)$$

so (6.49) holds.

Then

$$0 \leq -\frac{8}{3}\pi m(g) - \frac{16}{3}\pi m(g) + \left(1 - \frac{2}{C}\right)8\pi m(g) = -\frac{16}{C}\pi m(g) \quad (6.55)$$

but $m(g) > 0$, this is a contradiction. So for the stable constant mean curvature foliation there exists some constant $C > 0$ such that for any sphere Σ in the foliation,

$$\frac{r_0(\Sigma)}{r_1(\Sigma)} \geq C. \quad (6.56)$$

Then the uniqueness follows from Theorem 1.4.

Proof of the Corollary 1.9. Suppose there is not such $K(C, \beta)$, then we can find a sequence of constant mean curvature spheres Σ_n , with

$$\lim_{n \rightarrow \infty} r_0(\Sigma_n) = \infty \quad \lim_{n \rightarrow \infty} \frac{\log(r_1)}{r_0^{\frac{1}{4}}} = 0 \quad (6.57)$$

and Σ_n do not belong to the foliation. But from the argument above we know this sequence satisfies

$$\frac{r_0(\Sigma_n)}{r_1(\Sigma_n)} \geq C. \quad (6.58)$$

So when n is sufficiently large, Σ_n must belong to the foliation, which ends the proof.

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