# Symplectic homology product via Legendrian surgery

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#### Abstract

This research announcement continues the study of the symplectic homology of Weinstein manifolds undertaken in [2] where the symplectic homology, as a vector space, was expressed in terms of the Legendrian homology algebra of the attaching spheres of critical handles. Here we express the product and BV-operator of symplectic homology in that context.

# 1 Introduction

Weinstein manifolds are symplectic counterparts of affine (Stein) complex manifolds. They are exact symplectic manifolds which admit symplectic handle decompositions with isotropic core disks. In particular the index of any handle is at most half of the dimension of the manifold. We refer the reader to [2, 15, 10] for precise definitions. Middle dimensional handles are called *critical* and handles below middle dimension *subcritical*. Given

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a Weinstein manifold X, its symplectic homology  $S\mathbb{H}(X)$  is a symplectic invariant which was first defined by A. Floer and H. Hofer in [11]. It vanishes if X is subcritical, see [3]. A general Weinstein manifold X is obtained by attaching critical handles to a subcritical manifold  $X_0$  along a collection  $\Lambda$  of Legendrian spheres in the contact manifold  $Y_0$  which is the ideal boundary of  $X_0$ . In [2],  $S\mathbb{H}(X)$  was expressed in terms of the Legendrian homology algebra of  $\Lambda \subset Y_0$ .

Symplectic homology comes with TQFT-like operations which on the homology level are all expressed in terms of a (pair-of-pants) product and a BV-operator, see [14, 12]. In this paper we express the product and the BV-operator in terms of the Legendrian homology algebra of  $\Lambda$  and some additional operations on it. The BV-operator is determined by the algebra itself, whereas the expression of the product involves a duality operation which can be identified with the first term in the extension of Legendrian homology algebra to Legendrian rational symplectic field theory (SFT).

The paper is organized as follows. In Section 2 we introduce algebraic constructions which we later apply to Legendrian homology algebras. In Section 3 we express the Legendrian homology algebra results from [2] in the terminology of Section 2 and also include a discussion on the construction of a version of rational SFT in the Legendrian setting (completion of this construction is work in progress). In Section 4 we present the symplectic homology product and the BV-operator in terms of Legendrian homology algebra and a duality operator and in Section 5 we use the result to compute the symplectic homology, with product and BV-operator, of cotangent bundles of spheres. While proofs are mostly omitted, we note that the proofs of the algebraic results from Section 2 are fairly straightforward.

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# 2 Algebraic preliminaries

In this section we describe a number of constructions in a purely algebraic setting. The constructions will be applied in geometrically relevant situations in Sections 3 and 4.

#### 2.1 A tensor algebra and associated objects

We associate to a differential graded algebra A over a field K a pair of dual A-modules and a corresponding tensor algebra. We also discuss cyclic versions of these objects, which are vector spaces over K.

#### **2.1.A** The algebra A

Let R be an algebra over a field K of characteristic 0 generated by idempotents  $1_1, \ldots, 1_m$ , where  $1_k 1_l = \delta_{kl} 1_k$ , where  $\delta_{kl}$  is the Kronecker delta. We define a graded module V over R, which is generated by a set  $\mathcal{C}$  decomposed as a disjoint union  $\mathcal{C} = \bigcup_{1 \leq i,j \leq m} \mathcal{C}_{ij}$ . The degree of an element  $a \in V$ will be denoted by |a|. The decomposition of  $\mathcal{C}$  induces a decomposition  $V = \bigoplus_{1 \leq i,j \leq m} V^{(ij)}$ , where  $V^{(ij)}$  is generated by  $\mathcal{C}_{ij}$ . When applied on the left, the element  $1 \leq i \leq 1$  are acts as the identity of  $V^{(ij)}$  and as the 0 more

the element  $1_j$ , j = 1, ..., m, acts as the identity on  $V^{(ij)}$  and as the 0-map on  $V^{(il)}$ ,  $l \neq j$ . Similarly, when applied on the right, the element  $1_j$  acts as the identity on  $V^{(ji)}$  and as 0 on  $V^{(li)}$ ,  $l \neq j$ .

We define an extension  $\hat{V}$  of the module V as follows. The module  $\hat{V}$  has a set of additional generators  $\mathcal{X} = \{x_1, \ldots, x_m\}$  where  $|x_j| = 0, j = 1, \ldots, m$ , and we define

$$\widehat{V} := K \langle \mathcal{X} \rangle \oplus V[1],$$

where  $K\langle \mathcal{X} \rangle$  denotes the *m*-dimensional module freely generated by the set  $\mathcal{X}$  and where V[1] denotes V with grading shifted up by 1. The module  $\widehat{V}$  inherits a decomposition  $\widehat{V} = \bigoplus_{1 \leq i,j \leq m} \widehat{V}^{(ij)}$  from V, where

$$\widehat{V}^{(ij)} = \begin{cases} V^{(ij)}[1] & \text{if } i \neq j, \\ Kx_j \oplus V^{(jj)}[1] & \text{if } i = j. \end{cases}$$

Here the idempotents act as before on  $V \approx V[1]$  as a subset of  $\hat{V}$ , and as follows on the new generators:  $1_j x_i = \delta_{ji} x_j$  and  $x_i 1_j = \delta_{ij} x_j$ ,  $1 \le i, j \le m$ .

We will use the following notation. If  $a \in V$  then  $\hat{a}$  denotes the element in  $V[1] \subset \hat{V}$  that corresponds to a. Furthermore, if F and G are either Vor  $\hat{V}$  and if (u, v) is an ordered pair of elements in  $F^{(ij)} \oplus G^{(kl)}$  then (u, v) is composable if j = k.

Let

$$A := R \oplus V \oplus V_{R}^{\otimes^{2}} \oplus \dots$$
(2.1)

be the tensor algebra over R generated by V with its natural grading. Then A is generated additively by composable monomials  $w = c_1 \dots c_k, c_j \in V, j = 1, \dots, k$ , where we say that w is composable provided  $(c_j, c_{j+1})$  is composable for  $1 \leq j \leq k - 1$ . (Here, as often in the following, we will suppress the tensor sign in the notation for tensor products). The decomposition of V induces a corresponding decomposition of  $A, A = \bigoplus_{1 \leq i,j \leq m} A^{(ij)}$ , where  $A^{(ij)}$  is generated by composable monomials  $c_1 \dots c_k$  with  $c_1 \in V^{(is)}$  for some  $s \in \{1, \dots, m\}$  and  $c_k \in V^{(tj)}$  for some  $t \in \{1, \dots, m\}$ .

We will assume that the algebra A has a differential  $d: A \to A$  of degree -1 which leaves the subspaces  $A^{(ij)}, i, j \in \{1, \ldots, m\}$  invariant (i.e.,  $d^2 = 0$ , d satisfies the Leibniz rule, and  $d(A^{(ij)}) \subset A^{(ij)}$ ). We will usually write H(A) for the homology algebra  $H_*(A, d)$ .

#### **2.1.B** The A-module M and the K-complex $M^{\text{cyc}}$

Define

$$M = M(A) := A \bigotimes_{R} \widehat{V} \bigotimes_{R} A.$$
(2.2)

Then M is a graded left-right module over the differential graded algebra Awhich is additively generated by composable monomials  $w = c_1 \dots c_k u f_1 \dots f_l$ , where  $c_i \in \mathcal{C}$ ,  $i = 1, \dots, k$ ,  $u \in \widehat{\mathcal{C}} \cup \mathcal{X}$ , and  $f_j \in \mathcal{C}$ ,  $j = 1, \dots, l$ . We say that a composable monomial w is cyclically composable if the pair  $(f_l, c_1)$  is composable. Let  $M^{\text{diag}} \subset M$  be the submodule generated (over K) by cyclically composable monomials and define  $M^{\text{cyc}}$  as the quotient space of  $M^{\text{diag}}$  in which monomials which differ by graded cyclic permutation are identified. If  $w \in M^{\text{diag}}$  then we write  $\llbracket w \rrbracket$  for its equivalence class in  $M^{\text{cyc}}$ , i.e. if  $\pi \colon M^{\text{diag}} \to M^{\text{cyc}}$  denotes the natural projection then  $\llbracket w \rrbracket := \pi(w) \in M^{\text{cyc}}$ .

The differential  $d: A \to A$  induces a differential  $d_M: M \to M$  as follows. If  $w_1 u w_2 \in M$ ,  $u \in \hat{V}$ , and  $w_1, w_2 \in A$ , then

$$d_M(w_1u\,w_2) := (dw_1)u\,w_2 + (-1)^{|w_1|}w_1(d_Mu)w_2 + (-1)^{|w_1u|}w_1u(dw_2).$$
(2.3)

Here

$$d_M u = \begin{cases} ax_i - x_j a + S(da) & \text{if } u = \hat{a}, \ a \in V^{(ij)}, \\ 0 & \text{if } u = x_j, \ j \in \{1, \dots, m\}, \end{cases}$$
(2.4)

where the *R*-module homomorphism  $S: A \to M$  is defined by

$$S(a_1 \dots a_k) = \widehat{a}_1 a_2 a_3 \dots a_k + (-1)^{|a_1|} a_1 \widehat{a}_2 a_3 \dots a_k + \dots + (-1)^{|a_1 \dots a_{k-1}|} a_1 a_2 \dots a_{k-1} \widehat{a}_k,$$
(2.5)

and  $S(1_j) = 0$ , for j = 1, ..., m.

**Lemma 2.1.** The map  $d_M: M \to M$  is a differential (i.e.  $d_M^2 = 0$ ) which leaves  $M^{\text{diag}}$  invariant and descends to a differential  $M^{\text{cyc}} \to M^{\text{cyc}}$ , still denoted  $d_M$ , in such a way that the natural projection  $\pi: (M^{\text{diag}}, d_M) \to (M^{\text{cyc}}, d_M)$  is a chain map.

We write H(M) and  $H(M^{\text{cyc}})$ , respectively, for the homologies  $H_*(M, d_M)$ and  $H_*(M^{\text{cyc}}, d_M)$ . Note that H(M) is a left-right module over H(A), while  $H(M^{\text{cyc}})$  is just a K-vector space.

Remark 2.2. We can write

$$M^{\text{cyc}} = \bigoplus_{j=1}^{m} Ax_j \oplus A \underset{R}{\otimes} V[1] = 'M^{\text{cyc}} \oplus ''M^{\text{cyc}}.$$

Let  $p: M^{\text{cyc}} \to {}^{\prime\prime}M^{\text{cyc}}$  be the projection and let  ${}^{\prime\prime}d_{M^{\text{cyc}}} := p \circ (d_{M^{\text{cyc}}})|_{{}^{\prime\prime}M^{\text{cyc}}}$ . The second summand  ${}^{\prime\prime}M^{\text{cyc}}$  can be viewed as a non-commutative version of the "Lie algebra of vector fields" on the space V. Indeed, we can identify the cyclic monomial  $[\![w\hat{u}]\!] \in M^{\text{cyc}}, w \in A, u \in V$ , with the vector field  $w \frac{\partial}{\partial u}$ . The differential  ${}^{\prime\prime}d_{M^{\text{cyc}}}$  can be naturally viewed as a vector field on V, and hence an element of  ${}^{\prime\prime}M^{\text{cyc}}$ , given by  $\sum_{c \in \mathcal{C}} dc \frac{\partial}{\partial c}$ . For any  $X \in {}^{\prime\prime}M^{\text{cyc}}$  its differential  ${}^{\prime\prime}d_M X$  is just the Lie bracket  $[{}^{\prime\prime}d_M, X]$  of vector fields, or equivalently the Lie derivative  $L_{{}^{\prime\prime}d_M} X$ .

This interpretation corresponds to the "satellite philosophy" from [9]. Another interpretation was pointed out to the authors by M. Abouzaid:  $H(M^{\text{cyc}})$  is isomorphic to the Hochschild homology  $HH_*(A, A)$ .

#### 2.1.C The extended tensor algebra U and the space $U^{\text{bal}}$

Recall that the module V is generated over R by a set  $\mathcal{C}$  and that the module M is generated over A as a left-right module by  $\widehat{\mathcal{C}} \cup \mathcal{X}$ , where  $\widehat{\mathcal{C}} = \{\widehat{c}: c \in \mathcal{C}\}$  and  $\mathcal{X} = \{x_1, \ldots, x_m\}$ . Let  $\widehat{V}^*$  denote the K-vector space dual to  $\widehat{V}$ . Then  $\widehat{V}^*$  is the *direct product* of the 1-dimensional K-vector spaces which are generated by elements of  $\widehat{\mathcal{C}}^* \cup \mathcal{X}^*$ , where  $\widehat{\mathcal{C}}^* = \{\widehat{c}^*: c \in \mathcal{C}\}$  and  $\mathcal{X}^* = \{x_j^*: x_j \in \mathcal{X}\}$ . Furthermore, we have

$$\widehat{V}^* = \bigoplus_{1 \le i,j \le m} \widehat{V}^{*(ij)}, \quad \text{where} \quad \widehat{V}^{*(ij)} := (\widehat{V}^{(ij)})^*.$$

In  $\widehat{V}^*$ , the idempotent  $1_j$ ,  $j = 1, \ldots, m$ , acts from the left as the identity on  $\widehat{V}^{*(ji)}$  and as 0 on  $\widehat{V}^{*(li)}$ ,  $l \neq j$ , and acts from the right as the identity on  $\widehat{V}^{*(ij)}$  and as 0 on  $\widehat{V}^{*(il)}$ ,  $l \neq j$ . Fix an integer  $n \geq 2$  (it will correspond to half the dimension of the Liouville domain  $X_0$  in Section 3). We endow  $\widehat{V}^*$  with a grading defined as follows:

$$|\widehat{c}^*| = n - 3 - |\widehat{c}|, \text{ for } \widehat{c} \in \widehat{\mathcal{C}},$$
(2.6)

$$|x_j^*| = n - 3$$
, for  $j = 1, \dots, m$ . (2.7)

Define

$$M^* := A \underset{R}{\otimes} \widehat{V}^* \underset{R}{\otimes} A. \tag{2.8}$$

We define the tensor algebra

$$U(0,0) := \bigoplus_{k \ge 0} \left( M \oplus M^* \right)_A^{\otimes^k}.$$
(2.9)

as a certain completion of the free associative algebra generated by  $\mathcal{C}\cup\widehat{\mathcal{C}}\cup\widehat{\mathcal{C}}^*\cup$  $\mathcal{X}\cup\mathcal{X}^*$ . Its elements are infinite series in the  $(\widehat{\mathcal{C}}^*\cup\mathcal{X}^*)$ -variables with polynomial coefficients in the  $(\widehat{\mathcal{C}}\cup\mathcal{X})$ -variables (compare to the above definitions of M as an infinite sum and  $M^*$  as an infinite product). For simplicity of notation and since all completions we consider below are of the same type as that of (2.9), we will, as in that equation, suppress completions from notation and use the direct sum symbol in decompositions. Consider the extension U of U(0,0) obtained by adding multiplicative generators  $\hbar, \hbar^{-1}$  and  $\sigma, \sigma^{-1}$  of degrees

$$|\sigma| = 1, \quad |\sigma^{-1}| = -1, \quad |\hbar| = n - 3, \quad |\hbar^{-1}| = -(n - 3)$$

which satisfy the following relations:

$$\sigma\sigma^{-1} = \sigma^{-1}\sigma = 1, \quad \hbar\hbar^{-1} = \hbar^{-1}\hbar = 1, \quad \sigma\hbar = (-1)^{n-3}\hbar\sigma,$$

and

$$\sigma u = \begin{cases} (-1)^{|u|} u\sigma, & u \in \mathcal{C}; \\ (-1)^{|u|+1} u\sigma, & u \in \widehat{\mathcal{C}} \cup \mathcal{X} \cup \widehat{\mathcal{C}}^* \cup \mathcal{X}^*, \end{cases}$$
$$\hbar u = \begin{cases} (-1)^{|u|(n-3)} u\hbar, & u \in \mathcal{C} \cup \mathcal{X} \cup \widehat{\mathcal{C}}; \\ (-1)^{(|u|+n-3)(n-3)} u\hbar, & u \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*. \end{cases}$$

The algebra U can be decomposed according to tensor type:  $U = \bigoplus_{p,q \ge 0} U_q^p$ , where  $U_q^p \subset U$  is the K-subspace generated by monomials with q factors from M and p from  $M^*$ . We decompose further:  $U_q^p = \bigoplus_{s,h \in \mathbb{Z}} U_q^p(s,h)$ , where  $U_q^p(s,h) \subset U_q^p$  is the subspace spanned by all monomials of total  $\sigma$ degree s and total  $\hbar$ -degree h. In particular,  $U_0^0(0,0) = A$ ,  $U_1^0(0,0) = M$ , and  $U_0^1(0,0) = M^*$ . We also write  $U^p := \bigoplus_{q \ge 0} U_q^p$ ,  $U^+ := \bigoplus_{p > 0} U^p$ ,  $U_q := \bigoplus_{p \ge 0} U_q^p$ ,  $U_q(s,h) := \bigoplus_{p \ge 0} U_q^p(s,h)$ , and  $U^p(s,h) := \bigoplus_{q \ge 0} U_q^p(s,h)$ . For  $X \in U_q^p(s,h)$ , let

$$\operatorname{st}(X) := p + h$$
 and  $\operatorname{jp}(X) := p + q + s$ .

In analogy with the definition of  $M^{\text{diag}}$  in Section 2.1.B, we define  $U^{\text{diag}}$ . Further, define  $U^{\text{cyc}}$  as the quotient of  $U^{\text{diag}}$  obtained by dividing by the following graded cyclic permutation rule: identify  $a_1 \ldots a_k \in U^{\text{diag}}$  with  $\varepsilon_1 \varepsilon_2 \varepsilon_3 a_2 \ldots a_k a_1 \in U^{\text{diag}}$ , where

$$\begin{split} \varepsilon_1 &= (-1)^{|a_1||a_2\dots a_k|}, \\ \varepsilon_2 &= \begin{cases} 1 & \text{if } a_1 \in \mathcal{C} \cup \{\hbar, \hbar^{-1}\}, \\ (-1)^{jp(a_2\dots a_k)} & \text{if } a_1 \in \widehat{\mathcal{C}} \cup \mathcal{X} \cup \widehat{\mathcal{C}}^* \cup \mathcal{X}^* \cup \{\sigma, \sigma^{-1}\}, \\ \varepsilon_3 &= \begin{cases} 1 & \text{if } a_1 \in \mathcal{C} \cup \widehat{\mathcal{C}} \cup \mathcal{X} \cup \{\sigma, \sigma^{-1}\}, \\ (-1)^{\text{st}(a_2\dots a_k)} & \text{if } a_1 \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^* \cup \{\hbar, \hbar^{-1}\}. \end{cases} \end{split}$$

The vector space  $U^{\text{cyc}}$  inherits the decomposition by tensor type:

$$U^{\text{cyc}} = \bigoplus_{p,q \ge 0} (U^p_q)^{\text{cyc}} = \bigoplus_{p,q \ge 0} \bigoplus_{s,h \in \mathbb{Z}} (U^p_q(s,h))^{\text{cyc}}$$

If the monomials X and X' differ by a graded cyclic permutation, i.e. if  $[\![X]\!] = [\![X']\!]$ , then, for  $z \in \{\sigma, \sigma^{-1}, \hbar, \hbar^{-1}\}$ ,  $[\![Xz]\!] = [\![X'z]\!]$  and  $[\![zX]\!] = [\![zX']\!]$ . We thus define  $z[\![X]\!] := [\![zX]\!]$  and  $[\![X]\!]z := [\![Xz]\!]$ , and note that this definition does not depend on choice of representative X.

A monomial  $X \in (U_q^p(s,h))^{\text{cyc}}$  is balanced if h = -1 and p - q - s = 0. The subspace of  $U^{\text{cyc}}$  generated by balanced monomials will be denoted by  $U^{\text{bal}}$ . It contains balanced versions  $M^{\text{bal}}$  and  $(M^*)^{\text{bal}}$  of  $M^{\text{cyc}}$  and  $(M^*)^{\text{cyc}}$ , respectively. More precisely,

$$M^{\text{bal}} := \left( U_1^0(-1, -1) \right)^{\text{cyc}} = U_1^0 \cap U^{\text{bal}},$$
$$(M^*)^{\text{bal}} := \left( U_0^1(1, -1) \right)^{\text{cyc}} = U_0^1 \cap U^{\text{bal}}.$$

#### **2.1.D** The operator $S: U^{cyc} \to U^{cyc}$

Consider a monomial  $X \in (U_q^p(s,h))^{\text{cyc}}$  and fix a factor  $c \in \mathcal{C}$  in X. Then  $X = [\![X'_c c X''_c]\!]$  for monomials  $X'_c, X''_c \in U$ . Define

$$\mathcal{S}(X) = \sum_{c} \llbracket X'_{c} \sigma^{-1} \widehat{c} X''_{c} \rrbracket \in \left( U^{p}_{q+1}(s-1,h) \right)^{\text{cyc}},$$

where the sum is taken over all factors  $c \in C$  in X. It is straightforward to check that this gives a well defined operation and by definition  $\mathcal{S}(U^{\text{bal}}) \subset U^{\text{bal}}$ . In particular, if  $X = [\![c_1 \dots c_m \widehat{a}^*]\!] \in (U_0^1(0,0))^{\text{cyc}}$  then  $\mathcal{S}(X) \in (U_1^0(-1,0))^{\text{cyc}}$  satisfies

$$S(c_1 \dots c_m \widehat{a}^*) = \sigma^{-1} \sum_{j=1}^m (-1)^{|c_1 \dots c_{j-1}|} c_1 \dots c_{j-1} \widehat{c}_j c_{j+1} \dots c_m \widehat{a}^*$$
  
=  $\sigma^{-1} S(c_1 \dots c_m) \widehat{a}^*,$ 

where S is the operator  $A \to M$  defined in (2.5) above. Furthermore,

$$\mathcal{S}^{k}(c_{1} \dots c_{m} \widehat{a}^{*}) = k! \sigma^{-k} \sum_{j_{1} < \dots < j_{k}}^{m} (-1)^{\sum_{l=1}^{k} |c_{1} \dots c_{j_{l}-1}|} c_{1} \dots \widehat{c}_{j_{1}} \dots \widehat{c}_{j_{k}} \dots c_{m} \widehat{a}^{*}.$$

#### 2.1.E Symmetry breaking and contraction operations

We next effectively break the cyclic symmetry of  $U^{\text{bal}}$ . An *excited* monomial is a cyclic monomial  $X \in U^{\text{bal}}$  with one of its factors from  $\widehat{\mathcal{C}} \cup \widehat{\mathcal{C}}^* \cup \mathcal{X} \cup \mathcal{X}^*$ distinguished. We call the distinguished factor of an excited monomial X the *excited* generator of X. The K-vector space generated by excited monomials from  $U^{\text{bal}}$  will be denoted by  $\mathbf{U}^{\text{bal}}$ . Note that  $\mathbf{U}^{\text{bal}}$  inherits decompositions from  $U^{\text{bal}}$ ,  $\mathbf{U}^{\text{bal}} = \bigoplus_{p,q} (\mathbf{U}_q^p)^{\text{bal}}$ , etc. We set  $(\mathbf{U}_0^0)^{\text{bal}} = 0$ . Define the *excitation operator*  $\mathcal{E} : U^{\text{bal}} \to \mathbf{U}^{\text{bal}}$  to be the linear map which maps a monomial Xthe sum of all its excitations.

Note that  $\mathbf{U}^{\text{bal}}$  contain excited versions  $\mathbf{M}^{\text{bal}} = (\mathbf{U}_1^0)^{\text{bal}}$  and  $(\mathbf{M}^*)^{\text{bal}} = (\mathbf{U}_0^1)^{\text{bal}}$  of  $M^{\text{bal}}$  and  $(M^*)^{\text{bal}}$ , respectively, which are isomorphic to their nonexcited counterparts. (In the context of  $\mathbf{U}^{\text{bal}}$  we think of  $M^{\text{bal}}$  and  $(M^*)^{\text{bal}}$ as monomials in  $(U_1^0)^{\text{bal}}$  and  $(U_0^1)^{\text{bal}}$  without excited generator.)

We will use the following two notations to identify excited generators in monomials. First, in concrete calculations we underline the excited generator in a monomial in  $U^{\text{bal}}$  and write e.g.  $\underline{x}, \underline{x}^* \widehat{\underline{c}}$ , and  $\widehat{\underline{c}}^*$ . Second, if X is a monomial in  $\mathbf{U}^{\text{bal}}$  and if u is a generator which is a factor in X then we let

$$Q_X(u) = \begin{cases} 1 & \text{if } u \text{ is excited,} \\ 0 & \text{otherwise.} \end{cases}$$

We next define an associative product on  $\mathbf{U}^{\text{bal}}$ . Consider two excited cyclic monomials  $X, Y \in \mathbf{U}^{\text{bal}}$ . Fix a factor  $u \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*$  from X and a factor  $v \in \widehat{\mathcal{C}} \cup \mathcal{X}$  from Y. Write  $X = \varepsilon_1 [\![X'_u \hbar^{-1} u]\!]$  and  $Y = \varepsilon_2 [\![vY'_v]\!]$ , where  $\varepsilon_1$  is the sign which arises from cyclically rotating X so that u is its last factor and then inserting a factor  $\hbar^{-1}$  in front of u and where  $\varepsilon_2$  is the sign which arises from cyclically rotating Y so that v is its first factor. Define the *contraction* operation  $\star: \mathbf{U} \otimes \mathbf{U} \to \mathbf{U}$ ,

$$X \star Y := \sum_{(u,v)} \varepsilon_1 \varepsilon_2 u(v) \llbracket X'_u Y'_v \rrbracket, \qquad (2.10)$$

where the sum is taken over all pairs (u, v) of factors  $u \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*$  from X and  $v \in \widehat{\mathcal{C}} \cup \mathcal{X}$  from Y, and where

$$u(v) = \begin{cases} \hbar & \text{if } u = v^*, v \notin \mathcal{X} \text{ and } Q_X(u) \neq Q_Y(v), \\ (-1)^{n-1}\hbar & \text{if } u = \underline{x}_j^* \text{ and } v = x_j, \text{ for some } j = 1, \dots, m, \\ \hbar & \text{if } u = x_j^* \text{ and } v = \underline{x}_j, \text{ for some } j = 1, \dots, m, \\ 0 & \text{otherwise.} \end{cases}$$

If  $Y \in (\mathbf{U}_q^1)^{\text{bal}}$  then we define, for  $j = 1, \ldots, q$ , the partial contraction  $X \underset{i}{\star} Y$  as follows:

$$X \star_{j} Y := \sum_{u} \varepsilon_{1} \varepsilon_{2} u(v_{j}) \llbracket X'_{u} Y'_{v_{j}} \rrbracket,$$

where the sum is taken over all  $u \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*$  from X and where  $v_j$  is the  $j^{\text{th}}$  factor from  $\widehat{\mathcal{C}} \cup \mathcal{X}$  in Y counting *counter-clockwise* from the unique factor from  $\widehat{\mathcal{C}}^* \cup \mathcal{X}^*$ . Similarly, if  $X \in (\mathbf{U}_1^p)^{\text{bal}}$  then we define, for  $i = 1, \ldots, p$  the *partial contraction*  $X \stackrel{i}{\star} Y$ :

$$X \stackrel{i}{\star} Y := \sum_{v} \varepsilon_1 \varepsilon_2 u_i(v) \llbracket X'_{u_i} Y'_v \rrbracket,$$

where the sum is taken over all  $v \in \widehat{\mathcal{C}} \cup \mathcal{X}$  from Y and where  $u_i$  is the *i*-th factor from  $\widehat{\mathcal{C}}^* \cup \mathcal{X}^*$  in X counting *clockwise* from the unique factor from  $\widehat{\mathcal{C}} \cup \mathcal{X}$ .

**Remark 2.3.** Let  $X, Y \in \mathbf{U}^{\text{bal}}$  then  $X \star Y \in \mathbf{U}^{\text{bal}}$ . Further if  $X \in (\mathbf{U}_q^1)^{\text{bal}}$  then  $X \star Y \in \mathbf{U}^{\text{bal}}$ , and if  $Y \in (\mathbf{U}_1^p)^{\text{bal}}$  then  $X \star Y \in \mathbf{U}^{\text{bal}}$ .

Define the *bracket* or *commutator*  $[,]: \mathbf{U} \otimes \mathbf{U} \to \mathbf{U}$  as

$$[X, Y] = X \star Y - (-1)^{|X||Y|} Y \star X.$$
(2.11)

Lemma 2.4. The bracket satisfies the graded Jacobi identity:

$$(-1)^{|X||Z|}[X, [Y, Z]] + (-1)^{|Z||Y|}[Z, [X, Y]] + (-1)^{|Y||X|}[Y, [Z, X]] = 0.$$

Let

$$\tau \mathbf{M}^{\text{bal}} = \mathbf{M}^{\text{bal}} \oplus M^{\text{bal}}$$
 and  $(\tau \mathbf{M}^*)^{\text{bal}} = (\mathbf{M}^*)^{\text{bal}} \oplus (M^*)^{\text{bal}}$ 

We associate with  $X \in (\mathbf{U}_1^p)^{\text{bal}}$  a multi-linear operation  $X_{\downarrow}$ :  $(\tau \mathbf{M}^{\text{bal}})^{\otimes^p} \to \tau \mathbf{M}^{\text{bal}}$  defined by

$$X_{\downarrow}(A_1 \otimes \cdots \otimes A_p) = \left( \dots \left( \left( X \stackrel{1}{\star} A_1 \right) \stackrel{1}{\star} A_2 \right) \stackrel{1}{\star} \dots \right) \stackrel{1}{\star} A_p.$$
(2.12)

Note that if  $A_1, \ldots, A_p \in \mathbf{M}^{\text{bal}}$  then  $X_{\downarrow}(A_1 \otimes \cdots \otimes A_p) \in \mathbf{M}^{\text{bal}}$ , that if exactly one of the arguments lies in  $M^{\text{bal}}$  then the image lies in  $M^{\text{bal}}$ , and that the operation vanishes on *p*-tuples with more than component in  $M^{\text{bal}}$ .

Similarly,  $Y \in (\mathbf{U}_q^1)^{\text{bal}}$  defines a multi-linear operation  $Y^{\uparrow}$ :  $((\tau \mathbf{M}^*)^{\text{bal}})^{\otimes^q} \to (\tau \mathbf{M}^*)^{\text{bal}}$ ,

$$Y^{\uparrow}(B_1 \otimes \dots \otimes B_q) = B_q \mathop{\star}_1 \left( \dots \mathop{\star}_1 \left( B_2 \mathop{\star}_1 \left( B_1 \mathop{\star}_1 Y \right) \right) \dots \right), \qquad (2.13)$$

if  $B_1, \ldots, B_q \in (\mathbf{M}^*)^{\text{bal}}$  then  $Y^{\uparrow}(B_1 \otimes \cdots \otimes B_q) \in (\mathbf{M}^*)^{\text{bal}}$ , if exactly one of the arguments lies in  $(M^*)^{\text{bal}}$  then the image lies in  $(M^*)^{\text{bal}}$ , and the operation vanishes on q-tuples with more than one component in  $(M^*)^{\text{bal}}$ .

In the special case q = 2, we associate with  $Y \in (\mathbf{U}_2^1)^{\text{bal}}$  two more operations,

$$Y^{\uparrow\downarrow}\colon (\tau \mathbf{M}^*)^{\mathrm{bal}} \otimes \tau \mathbf{M}^{\mathrm{bal}} \to \tau \mathbf{M}^{\mathrm{bal}}$$
(2.14)

and

$$Y^{\downarrow\uparrow} \colon \tau \mathbf{M}^{\mathrm{bal}} \otimes (\tau \mathbf{M}^*)^{\mathrm{bal}} \oplus \to \tau \mathbf{M}^{\mathrm{bal}}, \qquad (2.15)$$

where

$$Y^{\uparrow\downarrow}(X,Z) = \left(X \underset{1}{\star} Y\right) \star Z, \quad X \in (\tau \mathbf{M}^*)^{\mathrm{bal}}, \ Z \in \tau \mathbf{M}^{\mathrm{bal}},$$
$$Y^{\downarrow\uparrow}(Z,X) = \left(X \underset{2}{\star} Y\right) \star Z, \quad Z \in \tau \mathbf{M}^{\mathrm{bal}}, \ X \in (\tau \mathbf{M}^*)^{\mathrm{bal}}.$$

Similarly,  $X \in (\mathbf{U}_0^2)^{\text{bal}}$  induces an operation  $X^{\circlearrowright} \colon \tau \mathbf{M}^{\text{bal}} \to (\tau \mathbf{M}^*)^{\text{bal}}$ ,

$$X^{\circlearrowright}(Y) = X \star Y. \tag{2.16}$$

We observe that  $X^{\circlearrowright}(\mathbf{M}^{\mathrm{bal}}) \subset (\mathbf{M}^*)^{\mathrm{bal}}$  and  $X^{\circlearrowright}(M^{\mathrm{bal}}) \subset (M^*)^{\mathrm{bal}}$ .

#### 2.1.F The Hamiltonian

We define special elements  $H_q^1 \in \mathbf{U}^{\text{bal}}$  which derive from the differential on A and which, in the spirit of SFT, see [9], are called *Hamiltonians*.

#### Definition 2.5.

$$\begin{split} (h_1^1)' &:= \sum_{c \in \mathcal{C}} \llbracket S(dc) \hbar^{-1} \widehat{c}^* \rrbracket \\ (h_1^1)'' &:= \sum_{i,j=1}^m \sum_{c \in \mathcal{C}_{ij}} \left( \llbracket c x_j \hbar^{-1} \widehat{c}^* \rrbracket - \llbracket x_i c \hbar^{-1} \widehat{c}^* \rrbracket \right) \\ h_1^1 &:= (h_1^1)' + (h_1^1)'' \\ h_2^1 &:= \frac{1}{2!} \mathcal{S}(h_1^1)' + \mathcal{S}(h_1^1)'' + (-1)^{n-1} \sigma^{-1} \sum_{j=1}^m \llbracket x_j x_j \hbar^{-1} x_j^* \rrbracket \\ h_q^1 &:= \frac{1}{q!} \mathcal{S}^{q-1} h_1^1, \quad \text{for } q > 2, \\ H_q^1 &:= \mathcal{E}(h_q^1), \quad H^1 = \sum_{q \ge 1} H_q^1. \end{split}$$

We use the Hamiltonian  $H_1^1$  to define a differential on  $\tau \mathbf{U}^{\text{bal}} := \mathbf{U}^{\text{bal}} \oplus M^{\text{bal}} \oplus (M^*)^{\text{bal}}$ . First define the *exterior differential*  $d: \tau \mathbf{U}^{\text{bal}} \to \tau \mathbf{U}^{\text{bal}}$  to be the map which acts as the algebra differential  $d: A \to A$  on generators of  $\mathbf{U}$  from A, which maps all other generators to 0, and which satisfies the graded

Leibniz rule. Then  $d^2 = 0$  and we observe that  $H^1$  satisfies the following master equation

$$dH^1 + H^1 \star H^1 = 0. (2.17)$$

Since  $|H^1| = -1$ ,  $H^1 \star H^1 = \frac{1}{2}[H^1, H^1]$  and (2.17) can be rewritten in Maurer-Cartan form:

$$dH^1 + \frac{1}{2}[H^1, H^1] = 0. (2.18)$$

Decomposing  $H^1 = \sum_{q \ge 1} H_q^1$ , (2.17) is equivalent to the following sequence of identities for the  $(\mathbf{U}_q^1)^{\text{bal}}$ -components of  $H^1 \star H^1$ :

$$dH_q^1 + \sum_{k=1}^q \sum_{j=1}^k H_{q-k+1}^1 \star H_k^1 = 0, \quad q = 1, 2, 3, \dots$$
 (2.19)

The first identity in this sequence is

$$lH_1^1 + H_1^1 \star H_1^1 = 0, (2.20)$$

**Definition 2.6.** Define the differential  $d_U: \tau \mathbf{U}^{\text{bal}} \to \tau \mathbf{U}^{\text{bal}}$ ,

$$d_U X := dX + [H_1^1, X].$$

(Note that (2.20) is equivalent to  $d_U^2 = 0$ .)

The differential interacts with the bracket in the following way.

Lemma 2.7. For  $X, Y \in \tau \mathbf{U}^{\text{bal}}$ ,

$$d_U([X,Y]) = [d_U X, Y] + (-1)^{|X|} [X, d_U Y].$$

In particular, the operation [,] descends to the homology  $H(\tau \mathbf{U}^{\text{bal}}, d_U)$ .

We observe that  $d_U$  leaves  $M^{\text{bal}}$ ,  $\mathbf{M}^{\text{bal}}$ ,  $(M^*)^{\text{bal}}$ , and  $(\mathbf{M}^*)^{\text{bal}}$  invariant. For simple notation we write  $d_{\tau M}$  for the restriction of  $d_U$  to any one of these subspaces. Consider the isomorphisms  $\beta \colon M^{\text{cyc}} \to M^{\text{bal}}$  and  $\underline{\beta} \colon M^{\text{cyc}} \to \mathbf{M}^{\text{bal}}$ ,

$$\beta(X) = X\hbar^{-1}\sigma^{-1}$$
 and  $\underline{\beta} = \mathcal{E} \circ \beta$ , respectively. (2.21)

**Lemma 2.8.** The maps  $\beta \colon (M^{\text{cyc}}, d_M) \to (M^{\text{bal}}, d_{\tau}M) \text{ and } \underline{\beta} \colon (M^{\text{cyc}}, d_M) \to (\mathbf{M}^{\text{bal}}, d_{\tau M}) \text{ are chain isomorphisms.}$ 

The following result is a consequence of Lemma 2.7.

**Corollary 2.9.** If  $X \in \mathbf{U}^{\text{bal}}$  satisfies  $d_U X = 0$  and has the correct tensorial type to define one of the operations (2.12) - (2.16), see Section 2.1.E, then this operation is a chain map. If furthermore  $X = d_U Y$  for some  $Y \in \mathbf{U}^{\text{bal}}$  then the operation is chain homotopic to 0.

#### 2.2 Products

We introduce product operations on  $\mathbf{M}^{\text{bal}}$  and  $(\mathbf{M}^*)^{\text{bal}}$ .

#### **2.2.A** The product on $H((\mathbf{M}^*)^{\text{bal}})$

(

Recall the operation  $(H_2^1)^{\uparrow}$ :  $(\tau \mathbf{M}^*)^{\text{bal}} \otimes (\tau \mathbf{M}^*)^{\text{bal}} \to (\tau \mathbf{M}^*)^{\text{bal}}$  defined by the formula

$$\left(H_{2}^{1}\right)^{\uparrow}\left(X,Y\right) = Y \star \left(X \star H_{2}^{1}\right),$$

see (2.13). Define  $X \diamond Y := (-1)^{|X|} (H_2^1)^{\uparrow} (X, Y)$ . Then  $(\mathbf{M}^*)^{\text{bal}} \diamond (\mathbf{M}^*)^{\text{bal}} \subset (\mathbf{M}^*)^{\text{bal}}$ .

**Proposition 2.10.** The operation  $\diamond$ :  $(\mathbf{M}^*)^{\text{bal}} \otimes (\mathbf{M}^*)^{\text{bal}} \to (\mathbf{M}^*)^{\text{bal}}$  descends to an associative product on the homology  $H((\mathbf{M}^*)^{\text{bal}})$  which satisfies the following graded commutativity relation for the grading shifted by 1:

$$X \diamond Y = (-1)^{(|X|+1)(|Y|+1)} Y \diamond X.$$

The homology class of the cycle  $E := (-1)^{n-1} \sum_{i=1}^{k} \hbar^{-1} \underline{x}_{i}^{*} \sigma \in (\mathbf{M}^{*})^{\text{bal}}$  is the unit for this product on  $H((\mathbf{M}^{*})^{\text{bal}})$ . In fact, E is a unit for  $\diamond$  already on the chain level, before passing to homology.

Consider the second and third identities satisfied by the Hamiltonian  $H^1$ :

$$dH_2^1 + [H_1^1, H_2^1] = 0, (2.22)$$

$$dH_3^1 + [H_1^1, H_3^1] + H_2^1 \star H_2^1 + H_2^1 \star H_2^1 = 0.$$
(2.23)

Here (2.22) is the chain map equation for  $(H_2^1)^{\uparrow}$ :  $(\mathbf{M}^*)^{\text{bal}} \otimes (\mathbf{M}^*)^{\text{bal}} \rightarrow (\mathbf{M}^*)^{\text{bal}}$ , and (2.23) shows that this product is associative on the homology level. Indeed,

$$\begin{split} (X \diamond Y) \diamond Z &= (-1)^{|Y|+1} Z \star ((Y \star (X \star H_2^1)) \star H_1^2) \\ &= (-1)^{|Y|} Z \star (Y \star (X \star (H_2^1 \star H_1^2))) \\ &= (-1)^{|Y|+|X|} (Z \star (Y \star H_2^1)) \star (X \star H_1^2) \\ &= X \diamond (Y \diamond Z). \end{split}$$

The unit property can be checked directly if one observes that the sum G of all terms in  $H_2^1$  with excited generator from  $\widehat{\mathcal{C}}^* \cup \mathcal{X}^*$  and which contain at least one  $x_j$ -factor,  $j = 1, \ldots, m$  is

$$G = \sigma^{-1} \left( (-1)^{n-1} \sum_{j} \llbracket x_j x_j \hbar^{-1} \underline{x}_j^* \rrbracket + \sum_{i,j} \sum_{c \in \mathcal{C}_{ij}} \llbracket x_i \widehat{c} \hbar^{-1} \underline{\widehat{c}}^* \rrbracket + \llbracket \widehat{c} x_j \hbar^{-1} \underline{\widehat{c}}^* \rrbracket \right).$$

Hence, taking into account that |E| = 1 we have for any  $Y \in (\mathbf{M}^*)^{\text{bal}}$ 

$$\begin{split} E \diamond Y &= Y \star \left( \sum_{i=1}^{m} (-1)^{n-1} \hbar^{-1} \underline{x}_{i}^{*} \sigma \star G \right) \\ &= Y \star \left( \sum_{j} (-1)^{n-1} \llbracket x_{j} \hbar^{-1} \underline{x}_{j}^{*} \rrbracket + \sum_{c \in \mathcal{C}} \llbracket \widehat{c} \hbar^{-1} \widehat{\underline{c}}^{*} \rrbracket \right) = Y \end{split}$$

To verify the commutativity relation, let X be a  $d_M$ -cycle in  $\mathbf{M}^{\text{bal}}$ . Applying the  $\mathcal{S}$ -operator to the equation  $dX + [H_1^1, X] = 0$  gives  $d_U(\mathcal{S}X) = X \star H_2^1$ and thus, on the level of homology,  $X \star H_2^1 + X \star H_2^1 = 0$  and we calculate

$$\begin{aligned} X \diamond Y &= (-1)^{|X|} Y \star (X \star H_2^1) = (-1)^{|X|+1} Y \star (X \star H_2^1) \\ &= (-1)^{|X|+1+|X||Y|} X \star (Y \star H_2^1) = (-1)^{(|X|+1)(|Y|+1)} Y \diamond X. \end{aligned}$$

#### **2.2.B** The product on $H(\mathbf{M}^{\text{bal}})$

If there exists an element

$$H' = \sum_{p \ge 2} H^p \in \bigoplus_{p \ge 2} (\mathbf{U}^p)^{\text{bal}}$$

such that  $H = H^1 + H' = \sum_{p \ge 1} H^p$  satisfies the master equation

$$dH + H \star H = dH + \frac{1}{2}[H, H] = 0,$$
 (2.24)

then we say that H' is *compatible* with  $H^1$  and we call  $H = H^1 + H'$  the *full (rational) Hamiltonian*. As we shall see later, when A is a Legendrian

homology algebra then we can always find H' compatible with  $H^1$  using a relative SFT formalism associated with parallel copies of the Legendrian submanifold in the spirit of [5].

Assume now that H, and in particular  $H_0^2$ , which satisfies (2.24) exists and is given. Then we define the operation  $\boxdot: \tau \mathbf{M}^{\text{bal}} \otimes \tau \mathbf{M}^{\text{bal}} \to \tau \mathbf{M}^{\text{bal}}$ ,

$$X \boxdot Y = \left(H_0^2 \star H_2^1\right)^{\downarrow} (X, Y) = \left(\left(H_0^2 \star H_2^1\right)^{\frac{1}{4}} X\right) \star Y.$$

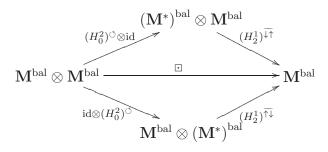
Note that if  $X, Y \in \mathbf{M}^{\text{bal}} \subset \tau \mathbf{M}^{\text{bal}}$  then  $X \boxdot Y \in \mathbf{M}^{\text{bal}}$ .

#### Theorem 2.11.

1. The operation  $\Box$ :  $\mathbf{M}^{\text{bal}} \otimes \mathbf{M}^{\text{bal}} \to \mathbf{M}^{\text{bal}}$  descends to an operation on the homology  $H(\mathbf{M}^{\text{bal}})$  where it can be expressed by the following three equivalent expressions:

$$\begin{aligned} X &\boxdot Y = -\left(H_0^2 \star H_2^1\right)^{\downarrow}(X, Y) = (-1)^{|X|} \left(H_2^1\right)^{\uparrow\downarrow} \left(\left(H_0^2\right)^{\circlearrowright}(X), Y\right) \\ &= (-1)^{|X||Y|+|X|+1} \left(H_2^1\right)^{\downarrow\uparrow} \left(X, \left(H_0^2\right)^{\circlearrowright}(Y)\right). \end{aligned}$$

Equivalently, the following diagram commutes on the homology level,



where

$$\begin{split} (H_2^1)^{\uparrow\downarrow}(X,Y) &:= (-1)^{|X|} (H_2^1)^{\uparrow\downarrow}(X,Y), \\ (H_2^1)^{\widecheck{\downarrow\uparrow}}(Y,X) &:= (-1)^{|Y||X|+|Y|+1} (H_2^1)^{\downarrow\uparrow}(Y,X), \end{split}$$

for  $X \in (\mathbf{M}^*)^{\text{bal}}$  and  $Y \in \mathbf{M}^{\text{bal}}$ .

2. The operation  $\boxdot$  is associative and graded commutative on homology, i.e. if  $X, Y \in H(\mathbf{M}^{\text{bal}})$  then  $X \boxdot Y = (-1)^{|X||Y|} Y \boxdot X$ .

$$(I) \swarrow + \checkmark + \checkmark + \checkmark + \checkmark + \checkmark = 0$$
$$(II) \swarrow + \checkmark + \checkmark + \checkmark + \checkmark + \checkmark = 0$$
$$(II) \swarrow + \checkmark + \checkmark + \checkmark + \checkmark + \checkmark + \land = 0$$

Figure 1: A graphical proof of associativity. A tree with p upwards and q downwards 1-valent vertices represents  $H_q^p$ . The dot represents the  $\star$ -operator. Equations (I) and (II) are consequences of  $dH + H \star H = 0$  (terms from dH are not shown in the figures).

The above result is a consequence of the master equation (2.24) which in particular implies the identities

$$dH_0^2 + [H_1^1, H_0^2] = dH_0^2 + H_0^2 \star H_1^1 = 0, \qquad (2.25)$$

$$dH_1^2 + [H_1^1, H_1^2] + H_0^2 \star H_2^1 + H_0^2 \star H_2^1 = 0.$$
(2.26)

A straightforward calculation shows that  $d_U(H_0^2 \star H_2^1) = (d_U H_0^2) \star H_2^1 - H_0^2 \star (d_U H_2^1)$ . By (2.25),  $d_U H_0^2 = 0$ , by (2.22),  $d_U H_2^1 = 0$ , and therefore  $H_0^2 \star H_2^1$  is a cycle in  $(\mathbf{U}_1^2)^{\text{bal}}$ . Similarly,  $-H_0^2 \star H_2^1$  is a cycle and (2.26) implies that their homology classes coincide. This shows that the product descends to homology and implies the first equation of part 1 of the theorem. The alternative expressions for  $\Box$  in part 1 are immediate from the definitions of the operations  $(H_0^2)^{\bigcirc}$ ,  $(H_2^1)^{\downarrow\uparrow}$ , and  $(H_2^1)^{\uparrow\downarrow}$ .

To verify the commutativity we observe that given a cycle  $X \in \mathbf{M}^{\text{bal}}$ ,

using the homological identity  $(H_0^2 \star X) \star H_2^1 = -(H_0^2 \star X) \star H_2^1$ , we calculate

$$\begin{split} X &\boxdot Y = \left( (H_0^2 \star H_2^1)^{\frac{1}{\star}} X \right) \star Y = (-1)^{|X|} ((H_0^2 \star X) \star H_2^1) \star Y \\ &= (-1)^{|X|+1} ((H_0^2 \star X) \star H_2^1) \star Y = (-1)^{|X||Y|+1} ((H_0^2 \star H_2^1)^{\frac{1}{\star}} Y) \star X \\ &= (-1)^{|X||Y|} ((H_0^2 \star H_2^1)^{\frac{1}{\star}} Y) \star X = (-1)^{|X||Y|} Y \boxdot X. \end{split}$$

Finally, the proof of associativity is pictorially illustrated on Figure 1.

The products  $\boxdot$  on  $\mathbf{M}^{\text{bal}}$  and  $\diamond$  on  $(\mathbf{M}^*)^{\text{bal}}$  are related via the chain map  $(H_0^2)^{\circlearrowright}: \mathbf{M}^{\text{bal}} \to (\mathbf{M}^*)^{\text{bal}}$  in the following way.

**Proposition 2.12.** The map induced by  $(H_0^2)^{\bigcirc}$  on homology is a homomorphism of rings  $\Phi: (H(\mathbf{M}^{\mathrm{bal}}), \boxdot) \to (H((\mathbf{M}^*)^{\mathrm{bal}}), \diamondsuit)$ . Moreover, if there is a unit e for  $\boxdot$  on  $H(\mathbf{M}^{\mathrm{bal}})$ , i.e.  $e \boxdot X = X \boxdot e$ , for all  $X \in H(\mathbf{M}^{\mathrm{bal}})$ , then  $\Phi$  is injective.

Indeed, if  $X, Y \in \mathbf{M}^{\text{bal}}$  then, taking into account that  $|\Phi(X)| = |X| - 1$ , we calculate

$$\begin{split} \Phi(X) \diamond \Phi(Y) &= (-1)^{|X|+1} (H_0^2 \star Y) \star ((H_0^2 \star X) \underset{1}{\star} H_2^1) \\ &= H_0^2 \star (((Y \star H_0^2) \underset{1}{\star} H_2^1) \star X) \\ &= (-1)^{|X||Y|} H_0^2 \star (((H_0^2 \underset{1}{\star} H_2^1) \underset{1}{\star} Y) \star X) \\ &= (-1)^{|X||Y|} \Phi(Y \boxdot X) = \Phi(X \boxdot Y). \end{split}$$

To verify the injectivity of  $\Phi$  we use one of equivalent definitions of  $\boxdot$  from Theorem 2.11, part 1:

$$X = X \boxdot e = (-1)^{|X|+1} \left(H_2^1\right)^{\downarrow\uparrow} \left(\Phi(X), e\right),$$

and hence,  $\Phi(X) = 0$  implies X = 0.

Using the isomorphism  $\beta : M^{\text{cyc}} \to \mathbf{M}^{\text{bal}}, \ \underline{\beta}(X) = \mathcal{E}(X\sigma^{-1}\hbar^{-1})$  defined above we transport the product to  $M^{\text{cyc}}$ . Given  $X, Y \in M^{\text{cyc}}$ , we define

$$X \boxtimes Y := \underline{\beta}^{-1} \left( \underline{\beta}(X) \boxdot \underline{\beta}(Y) \right)$$
  
=  $\left( \left( (H_0^2 \star H_2^1)^{\frac{1}{\star}} \mathcal{E}(X\sigma^{-1}\hbar^{-1}) \right) \star \mathcal{E}(Y\sigma^{-1}\hbar^{-1}) \right) \hbar\sigma$   
=  $(-1)^{(n-2)|Y|+1} \left( (H_0^2 \star H_2^1)^{\frac{1}{\star}} \mathcal{E}(X) \right) \star \mathcal{E}(Y)\sigma^{-1}\hbar^{-1}.$  (2.27)

Note that the homomorphism  $\beta$  shifts the grading by 2 - n, and hence the commutation equation for  $\Box$  translates to the equation

$$X \boxtimes Y = (-1)^{(|X|-n)(|X|-n)} Y \boxtimes X.$$

Let us also point out that the degree of the product  $\boxtimes$  is equal to -n.

## 3 Legendrian algebra constructions

In this section we first reformulate the main result relating Legendrian homology algebra and symplectic homology [2, Corollary 5.7] in the terminology of Section 2, and second give a brief sketch of the geometry that enters the construction of the full rational Hamiltonian in the framework of Legendrian SFT with focus on the parts of it that enter the expression of the product.

### **3.1** The complex $LH^{Ho}(\Lambda)$ expressed as $M^{cyc}$

Let  $(\overline{X}_0, \lambda)$  be a Liouville domain with contact boundary  $(Y_0, \ker(\lambda))$  and let  $\Lambda = \Lambda_1 \cup \ldots \Lambda_m \subset Y_0$  be a union of Legendrian spheres. Let  $\mathcal{C}$  be the set of Reeb chords of  $\Lambda$  and consider the decomposition  $\mathcal{C} = \bigcup_{1 \leq i,j \leq m} \mathcal{C}_{ij}$ , where a chord in  $\mathcal{C}_{ij}$  starts on  $\Lambda_j$  and ends on  $\Lambda_i$ . Let V be the graded module  $K\langle \mathcal{C} \rangle$  generated by the set  $\mathcal{C}$  with decomposition  $V = \bigoplus_{1 \leq i,j \leq m} V^{(ij)}$ , where  $V^{(ij)}$  is generated by  $\mathcal{C}_{ij}$ . Then the Legendrian homology algebra  $LHA(\Lambda)$  is the tensor algebra A associated to V as in (2.1) and the differential  $d_{LHA}$  gives a differential graded algebra  $(A, d) := (LHA(\Lambda), d_{LHA})$ , see [2, Section 4.1]. Let M = M(A) be the chain complex defined in (2.2). The following result is a consequence of [2, Remark 7.4].

**Lemma 3.1.** The chain complex  $M^{\text{cyc}}$  is canonically isomorphic to the chain complex  $LH^{\text{Ho}}(\Lambda)$  defined in [2, Section 4.5].

In [2, Corollary 5.7] it was shown that if  $S\mathbb{H}(X_0) = 0$  (e.g. if  $X_0$  is a subcritical Weinstein manifold) and if X is obtained from  $X_0$  by attaching Lagrangian *n*-handles along  $\Lambda$ , then there is a quasi-isomorphism

$$\Phi \colon SH(X) \to LH^{\mathrm{Ho}}(\Lambda) = M^{\mathrm{cyc}} \stackrel{\beta}{=} \mathbf{M}^{\mathrm{bal}}.$$
(3.1)

#### 3.2 Legendrian rational SFT

The differential in the Legendrian algebra  $A = LHA(\Lambda)$  is defined by counting elements of moduli spaces of rigid holomorphic disks in  $\mathbb{R} \times Y_0$ , with boundary on  $\mathbb{R} \times \Lambda$ , and with exactly one positive and several negative punctures, see e.g. [2, Section 4.1]. The formalism of Legendrian rational SFT comes from algebraic structures associated with holomorphic disks with an arbitrary number of positive and negative punctures. As is well known, in order to construct such a theory one needs to handle boundary cusp-degenerations of holomorphic disks. The problem of boundary cuspdegeneration can be solved by incorporating string topology operations into the SFT formalism. The corresponding theory has not yet been constructed in sufficient generality for it to be applicable in the setting of this paper. (See, however [13] for the case of 1-dimensional Legendrian knots in  $\mathbb{R}^3$ ).

However, the algebraic formalism described in Section 2 requires a different version of the full rational Hamiltonian H, in the spirit of [5]. It arises naturally when one considers moduli spaces of holomorphic disks with boundaries on multiple parallel copies of  $\mathbb{R} \times \Lambda$ .

Let us first consider the term  $H_0^2$ . This is the only term besides  $H_2^1$  which is needed for defining the product on  $M^{\text{cyc}}$ , see Figure 3. We recall here that  $H_k^1$  for any  $k \ge 1$  is determined by the differential on the Legendrian algebra A.

For each i = 1, ..., m fix a point  $y_i \in \Lambda_i$ . Let  $Z = \llbracket \underline{u}_1 w_1 u_2 w_2 \rrbracket \in (\mathbf{U}_0^2)^{\text{bal}}$ , where  $w_j = c_j^1 \dots c_j^{r_j} \in A$  and  $u_j \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*$ , j = 1, 2, be an excited cyclic monomial. We associate to it a moduli space  $\mathcal{M}(Z) = \mathcal{M}'(Z) \cup \mathcal{M}''(Z)$ . Here the moduli space  $\mathcal{M}'(Z)$  consists of holomorphic disks anchored in  $X_0$ 

$$f: (D, \partial D \setminus \mathbf{z}) \to (\mathbb{R} \times Y_0, \mathbb{R} \times \Lambda),$$

where

$$\mathbf{z} = \left\{ z_{u_1}, \, z_{c_1^1}, \dots, z_{c_1^{r_1}}, \, z_{u_2}, \, z_{c_2^1}, \dots, z_{c_2^{r_2}} \right\} \subset \partial D$$

is a set of cyclically ordered boundary punctures on  $\partial D$ , with the following properties:

- (ch<sup>-</sup>) at the puncture  $z_{c_i^i}$ , f is asymptotic to the chord  $c_j^i$  at  $-\infty$ ;
- (ch<sup>+</sup>) if  $u_j = \hat{c}^* \in \widehat{\mathcal{C}}^*$  then at the puncture  $z_{u_j} f$  is asymptotic to the chord c at  $+\infty$ ;

(pt<sup>+</sup>) if  $u_j = x_i^* \in \mathcal{X}^*$  then, f extends continuously over  $z_{u_j}$ , and  $f(z_{u_j}) \in \mathbb{R} \times \{y_i\}$ .

The moduli space  $\mathcal{M}''(Z)$  consists of holomorphic disks connected by Morse flow lines. To define it, choose auxiliary Morse functions  $\phi_i \colon \Lambda_i \to \mathbb{R}$ ,  $i = 1, \ldots, m$ , each with exactly two critical points. For  $(s, t) \in \{0, 1, \ldots, r_1\} \times \{0, 1, \ldots, r_2\}$  define the auxiliary moduli space  $\mathcal{M}_{(s,t)}(Z)$  of pairs  $(f_1, f_2)$  of holomorphic disks anchored in  $X_0$ ,

$$f_{\sigma} \colon (D, \partial D \setminus \mathbf{z}_{\sigma}) \to (\mathbb{R} \times Y_0, \mathbb{R} \times \Lambda), \ \sigma = 1, 2,$$

where

$$\mathbf{z}_{1} = \left\{ z_{u_{1}}, \, z_{c_{1}^{1}}, \dots, z_{c_{1}^{s}}, \, \zeta_{1}, z_{c_{2}^{t+1}}, \dots, z_{c_{2}^{r_{2}}} \right\} \subset \partial D, \\ \mathbf{z}_{2} = \left\{ z_{u_{2}}, \, z_{c_{2}^{1}}, \dots, z_{c_{2}^{t}}, \, \zeta_{2}, z_{c_{1}^{s+1}}, \dots, z_{c_{2}^{r_{2}}} \right\} \subset \partial D$$

are sets of cyclically ordered boundary punctures on  $\partial D$ , with the following properties. At  $z_{u_{\sigma}}$ ,  $f_{\sigma}$  satisfies condition (ch<sup>+</sup>) if  $u_{\sigma} \in \widehat{\mathcal{C}}^*$  or condition (pt<sup>+</sup>) if  $u_{\sigma} \in \mathcal{X}^*$ . At  $z_{c_i^j}$ ,  $f_{\sigma}$  satisfies condition (ch<sup>-</sup>). At the puncture  $\zeta_{\sigma}$ ,  $\sigma = 1, 2,$  $f_{\sigma}$  extends continuously and the following condition holds:

(mo) there exists a trajectory of  $-\nabla \phi_i$  which starts at  $f_1(\zeta_1)$  and ends at  $f_2(\zeta_2)$ , where  $\phi_i \colon \Lambda_i \to \mathbb{R}$  is the auxiliary Morse function of the component in which  $f_{\sigma}(\zeta_{\sigma})$  lies. Here the gradient is taken with respect to the metric induced by the symplectic form and the almost complex structure.

Finally, define

$$\mathcal{M}''(Z) := \bigcup_{(s,t) \in \{0,...,r_1\} \times \{0,...,r_2\}} \mathcal{M}_{(s,t)}(Z).$$

Given  $u \in \widehat{\mathcal{C}} \cup \widehat{\mathcal{C}}^*$ , let  $\widetilde{u} = a$  if  $u = \widehat{a}$  or  $u = \widehat{a}^*$ . Define  $|\widetilde{x}_j| = |\widetilde{x}_j^*| = -1$ ,  $j = 1, \ldots, m$ . Then we have the dimension formula

$$\dim(\mathcal{M}(Z)) = |\tilde{u}_1| + |\tilde{u}_2| - |w_1| - |w_2| - (n-3).$$

In particular in terms of the grading on U, for the balanced monomial  $\sigma^2 \hbar^{-1} Z$ we have

$$\dim(\mathcal{M}(Z)) = -|\sigma^2 \hbar^{-1} Z|.$$

Define  $H_0^2 \in (\mathbf{U}_0^2)^{\text{bal}}$  as  $\mathcal{E}(h_0^2)$ , where

$$h_0^2 = \sigma^2 \hbar^{-1} \sum_{|\sigma^2 \hbar^{-1} Z| = -1} \# \mathcal{M}(Z) Z, \qquad (3.2)$$

where the sum is taken over all balanced monomials  $\sigma^2 \hbar^{-1} Z$  which additively generate  $(\mathbf{U}_0^2)^{\text{bal}}$  and for which  $|\sigma^2 \hbar^{-1} Z| = -1$ . Here  $\# \mathcal{M}(Z)$  is the algebraic number of 1-dimensional components of the moduli space  $\mathcal{M}(Z)$ . As will be briefly discussed below, the above definition of  $H_0^2$  can generalized to more than two positive punctures and as a consequence we get the following result.

**Proposition 3.2.** There exists an element  $H'' \in \mathbf{U}^{\text{bal}} \ominus \left( (\mathbf{U}^1)^{\text{bal}} \oplus (\mathbf{U}^2_0)^{\text{bal}} \right)$ such that  $H' = H_0^2 + H''$  complements  $H^1$  to the full rational SFT Hamiltonian (i.e. if  $H = H^1 + H_0^2 + H''$  then  $dH + H \star H = 0$ ).

Here the component  $H_q^2 \in (\mathbf{U}_q^2)^{\text{bal}}$ , q > 0 of the Hamiltonian is determined by  $H_0^2$  defined above  $H_q^2 := \mathcal{E}(\frac{1}{q!}\mathcal{S}^q(h_0^2))$ , where  $h_0^2 = \frac{1}{2}\pi(H_0^2)$ , where  $\pi: (\mathbf{U}_0^2)^{\mathrm{bal}} \to (U_0^2)^{\mathrm{bal}}$  is the natural projection forgetting excitation. In general, the components  $H_q^p \in (\mathbf{U}_q^p)^{\text{bal}}$ , p > 2 satisfies a similar relation  $H_q^p = \frac{1}{q!} \mathcal{S}^q(h_0^p)$ ,  $h_0^p = \frac{1}{p} \pi(H_0^p)$ , and are defined through a construction which generalizes that used in the definition of  $H_0^2$  above. These higher order components of the Hamiltonian are not needed for the definition of the product (but some of them are needed for proving its properties) and will not be defined in this paper. We give however a brief description:  $H_q^p$  is defined using moduli spaces  $\mathcal{M}(Z)$  of so-called generalized holomorphic disks. Generalized holomorphic disks are combined objects consisting of holomorphic disks with boundary on one copy of  $\Lambda$  connected by gradient flow trees associated to auxiliary Morse functions on  $\Lambda$ , an example is depicted in Figure 2, see [7] for similar considerations. The elements of the moduli space  $\mathcal{M}''(Z)$  defined above are examples of generalized holomorphic disks where the gradient flow tree consists of only one gradient trajectory. A more general configuration is shown in Figure 2. More precisely, if  $Z = \underline{u}_1 w_1 u_2 w_2 \dots u_k w_k$  is a monomial in  $\mathbf{U}_a^p(0,0)$ , where  $u_i \in \widehat{\mathcal{C}} \cup \mathcal{X} \cup \widehat{\mathcal{C}}^* \cup \mathcal{X}^*$  and where  $w_i \in A, i = 1, \ldots, k$ , then we consider the moduli space  $\mathcal{M}(Z)$  of generalized holomorphic disks with asymptotic data according to Z, i.e. positive punctures at  $\widetilde{u}_i$  if  $u_i \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*$ , negative punctures at  $\widetilde{u}_j$  if  $u_j \in \widehat{\mathcal{C}} \cup \mathcal{X}$ , auxiliary negative punctures at the Reeb chords in the monomials  $w_i$ , and these punctures appearing along the

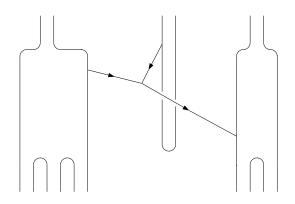


Figure 2: A generalized holomorphic disk in  $Y_0 \times \mathbb{R}$  with boundary on  $\mathbb{R} \times \Lambda$ .

boundary of the disk in the cyclic order given by Z. The dimension of  $\mathcal{M}(Z)$  is given by

$$\dim(\mathcal{M}(Z)) = (n-3) + \sum_{u_i \in \widehat{\mathcal{C}}^* \cup \mathcal{X}^*} \left( |\widetilde{u}_i| - (n-3) \right) - \sum_{u_j \in \widehat{\mathcal{C}} \cup \mathcal{X}} |\widetilde{u}_j| - \sum_{j=1}^m |w_j|.$$

In particular for  $\sigma^{p-q}\hbar^{-1}Z \in \left(\mathbf{U}_q^p\right)^{\text{bal}}$  we have

$$\dim(\mathcal{M}(Z)) = -|\sigma^{p-q}\hbar^{-1}Z|$$

and  $H_q^p$  is defined as

$$H_q^p = \sigma^{p-q} \hbar^{-1} \sum_{|\sigma^{p-q} \hbar^{-1}Z| = -1} \# \mathcal{M}(Z) Z,$$

where the sum ranges over all  $\sigma^{p-q}\hbar^{-1}Z \in (\mathbf{U}_q^p)^{\text{bal}}$  with  $|\sigma^{p-q}\hbar^{-1}Z| = -1$ .

**Remark 3.3.** In this geometric context, the excitation operator  $\mathcal{E}$  has the following meaning. Given a generalized holomorphic disk with boundary in the ordered union of (symplectizations of) parallel copies of a Legendrian manifold  $\Lambda$  we always insist that the boundary arcs of the disk between punctures corresponding to generators from  $\widehat{\mathcal{C}} \cup \widehat{\mathcal{C}}^* \cup \mathcal{X} \cup \mathcal{X}^*$  be mapped to different copies of  $\Lambda$ . We also require that when going around the boundary of the disk in the counter-clockwise direction we jump at every such puncture except one to a higher copy of  $\Lambda$  in terms of the natural ordering corresponding to the direction of the Reeb vector field. The unique puncture where we jump

down corresponds to the excited generator of the monomial corresponding to the boundary punctures. Here the moduli spaces corresponding to different assignment of the distinguished "jump-down" punctures are canonically diffeomorphic, and the operator  $\mathcal{E}$  just associates with this universal moduli space the moduli spaces corresponding to all possible assignments of the "jump-down" vertex.

**Remark 3.4.** The geometric motivation for the "gluing sign" in the equations  $\underline{x}_j^*(x_j) = (-1)^{n-1}\hbar$ ,  $j = 1, \ldots, m$  is as follows. The algebraic variables  $x_j$  and  $x_j^*$ ,  $j = 1, \ldots, m$  has two dual geometric meanings as follows: when  $x_j^*$  is glued to  $\underline{x}_j, x_j^*$  corresponds to a positive puncture at a short Reeb chord between copies of  $\Lambda$  at the minimum of the Morse function  $\phi \colon \Lambda \to \mathbb{R}$  and  $\underline{x}_j$  corresponds to a negative puncture at this chord, when  $\underline{x}_j^*$  is glued to  $x_j$ ,  $\underline{x}_j^*$  corresponds to a negative puncture at the maximum of  $\phi$  and  $x_j$  to a positive puncture at this maximum. In a coherent orientation scheme there is a sign difference (a factor  $(-1)^{n-1}$ ) between the gluings at these two types of punctures.

# 4 The product and the BV-operator via Legendrian algebra

Let  $X_0$ ,  $\Lambda \subset Y_0$ , and X be as above. Let  $A = LHA(\Lambda)$  and M = M(A). Then (3.1) gives a quasi-isomorphism  $\Phi: SH(X) \to M^{\text{cyc}} \approx \mathbf{M}^{\text{bal}}$ . In this section we present the operations on  $M^{\text{cyc}}$  which correspond to the product and the BV-operator on SH(X) under  $\Phi$ .

#### 4.1 The product

Let  $H \in \mathbf{U}^{\text{bal}}$  be the full rational Hamiltonian determined by  $\Lambda \subset Y_0$ , see (3.2). Then Proposition 3.2 implies that  $H_0^2$  satisfies (2.25). Hence, according to Theorem 2.11 and (2.27), there exists a product  $\boxtimes : M^{\text{cyc}} \otimes M^{\text{cyc}} \to M^{\text{cyc}}$  of degree -n which is associative and commutative on homology and which is defined by conjugating one of the following homotopically equivalent formulas

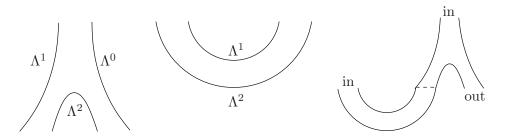


Figure 3: Holomorphic disks used in the expression of the product,  $\Lambda^j$ , j = 0, 1, 2 are parallel copies of  $\Lambda$ . The leftmost disk has three punctures at mixed (i.e. connecting different components) Reeb chords, one positive and two negative, and auxiliary negative punctures at pure (i.e. connecting a component to itself) chords which are not shown in the picture. The middle disk has two mixed positive punctures and auxiliary pure negative punctures, the latter again not shown. The rightmost picture shows how the disks are joined at a common mixed chord between  $\Lambda^2$  and  $\Lambda^1$  to form the tensor, with inputs and output as indicated, which represents the product.

for a product on  $\mathbf{M}^{\text{bal}}$  with  $\beta$ :

$$X \boxdot Y = (H_0^2 \star H_2^1)_{\downarrow}(X, Y) \simeq -(H_0^2 \star H_2^1)_{\downarrow}(X, Y)$$
  
$$\simeq (-1)^{|X|} (H_2^1)^{\downarrow\uparrow} \circ \left( (H_0^2)^{\circlearrowright} (X) \otimes Y \right)$$
  
$$\simeq (-1)^{|X||Y|+|X|+|Y|} (H_2^1)^{\uparrow\downarrow} \circ \left( X \otimes (H_0^2)^{\circlearrowright} (Y) \right), \qquad (4.1)$$

see Figure 3.

**Theorem 4.1.** If  $S\mathbb{H}(X_0) = 0$  (e.g. if  $X_0$  is subcritical) then, under the map on homology induced by the quasi-isomorphism  $\Phi: SH(X) \to M^{\text{cyc}}$ , see (3.1), the pair of pants product on  $S\mathbb{H}(X)$  corresponds to the product  $\boxtimes$  on  $H(M^{\text{cyc}})$ .

#### 4.2 The BV-operator

Let  $A = LHA(\Lambda)$  and consider the linear map  $\Delta \colon M^{\text{cyc}} \to M^{\text{cyc}}$  of degree 1 defined as follows on monomials  $\llbracket uw \rrbracket$ ,  $u \in \widehat{\mathcal{C}} \cup \mathcal{X}$ ,  $w \in A$ :

$$\Delta(\llbracket uw \rrbracket) = \begin{cases} S(w) & \text{if } u \in \mathcal{X}, \\ 0 & \text{if } u \in \widehat{\mathcal{C}}. \end{cases}$$

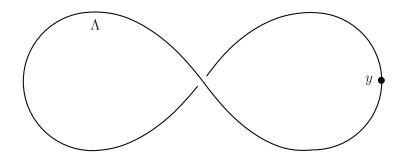


Figure 4: The Legendrian unknot.

**Proposition 4.2.** The map  $\Delta$  is a chain map and if  $S\mathbb{H}(X_0) = 0$  then, under the map on homology induced by the quasi-isomorphism  $\Phi: SH(X) \rightarrow M^{\text{cyc}}$ , see (3.1), the BV-operator on  $S\mathbb{H}(X)$  corresponds to the operator  $\Delta$ on  $H(M^{\text{cyc}})$ .

# **5 Example:** $SH(T^*S^n)$

In this section we compute the symplectic homology with product and BVoperator for cotangent bundles of spheres presented as the result of attaching a Lagrangian *n*-handle to the Legendrian unknot in the boundary of the ball.

Consider the Legendrian unknot in  $\mathbb{R}^{2n-1}$  with coordinates  $(q, p, z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{n-1} \times \mathbb{R}$  and contact form  $dz - p \cdot dq$ . Figure 4 shows the projection of the Legendrian unknot  $\Lambda$  for n = 2 to the qp-plane. To define the Legendrian unknot of dimension n-1 consider the 1-dimensional unknot  $\Lambda$  with Reeb chord over q = p = 0 and symmetric with respect to the rotation by a half turn around the z axis. Include  $\Lambda$  into the  $q_1p_1z$ -subspace of  $\mathbb{R}^{2n-1}$  and take the (n-1)-dimensional unknot as the orbit of  $\Lambda$  under the action by SO(n-1) given by  $\sigma(q, p, z) = (\sigma \cdot q, \sigma \cdot p, z)$  for  $\sigma \in SO(n-1)$ . For simplicity, we let  $\Lambda$  denote the Legendrian unknot in  $\mathbb{R}^{2n-1}$  for any n.

Consider  $\Lambda \subset \mathbb{R}^{2n-1}$  and include it into a Darboux chart in the boundary of the 2*n*-ball  $B^{2n}$ . Then attaching a Lagrangian *n*-handle to  $B^{2n}$  along  $\Lambda$  gives a symplectic manifold which is symplectomorphic to the cotangent bundle  $T^*S^n$  of the *n*-sphere. As explained in [2], in order to determine  $LHA(\Lambda)$  it suffices to study Reeb chords and holomorphic disks that are contained in the Darboux chart. The unknot  $\Lambda$  has only one Reeb chord which we denote *a* and |a| = n - 1. The moduli space of holomorphic disks with positive puncture at a and a marked point on the boundary evaluates with degree  $\pm 1$  to  $\Lambda$ . To see this, note that for n = 2, holomorphic disks that are rigid up to translation correspond to polygons with boundary on the knot diagram and with convex corners, where at positive punctures the boundary orientation goes from the lower to the upper strand and vice versa at negative punctures. There are thus two disks with evaluation maps that give a degree  $\pm 1$  map since we use the null-cobordant spin structure on the circle to orient moduli spaces. For n > 2 we note that we may take the almost complex structure to be invariant with respect to the SO(n-1) action and it follows that the moduli space of holomorphic disks with positive puncture at a is diffeomorphic to  $S^{n-2} \times \mathbb{R}$ and that, when equipped with a marked point on the boundary, it evaluates with degree  $\pm 1$ .

The above description of moduli spaces implies in particular that the differential d on  $A = LHA(\Lambda) = \mathbb{Q}\langle a \rangle$  is trivial, i. e., d = 0. Consequently, the differential on  $\mathbf{M}^{\text{bal}}$  satisfies

$$d_M X = dX + [H_1^1, X] = [H_1^1, X],$$

where

$$H_1^1 = \left( \llbracket \underline{a}\underline{x}\hbar^{-1}\widehat{a}^* \rrbracket - \llbracket \underline{x}a\hbar^{-1}\widehat{a}^* \rrbracket \right).$$
(5.1)

Thus  $H(\mathbf{M}^{\text{bal}})$  and  $H(M^{\text{cyc}})$  (recall the chain isomorphism  $\underline{\beta} \colon M^{\text{cyc}} \to \mathbf{M}^{\text{bal}}$ , see (2.21)) are generated by the cycles listed in Tables 1 and 2 when n-1 is even and odd, respectively.

$H(\mathbf{M}^{\mathrm{bal}})$	degree	$H(M^{\rm cyc})$	degree
$[\underline{x}]\hbar^{-1}\sigma^{-1}$	-(n-2)	$\llbracket x \rrbracket$	0
$\llbracket \underline{x} a^{2k+1} \rrbracket \hbar^{-1} \sigma^{-1}$	2k(n-1) + 1	$[\![xa^{2k+1}]\!]$	(2k+1)(n-1)
$\left[ \widehat{\underline{a}} a^{2k} \right] \hbar^{-1} \sigma^{-1}$	2k(n-1) + 2	$[\![\widehat{a}a^{2k}]\!]$	(2k+1)(n-1)+1

Table 1: Homology of  $\mathbf{M}^{\text{bal}}$  and  $M^{\text{cyc}}$  for n-1 odd,  $k = 0, 1, 2, \dots$ 

In order to compute the product we need to determine  $H_2^1$  and  $H_0^2$ . Con-

$H(\mathbf{M}^{\mathrm{bal}})$	degree	$H(M^{\mathrm{cyc}})$	degree
$[\underline{x}]\hbar^{-1}\sigma^{-1}$	-(n-2)	$\llbracket x \rrbracket$	0
$\llbracket \underline{x} a^k \rrbracket \hbar^{-1} \sigma^{-1}$	(k-1)(n-1) + 1	$[\![xa^k]\!]$	k(n-1)
$\left[\!\left[\underline{\widehat{a}}a^{k-1}\right]\!\right]\hbar^{-1}\sigma^{-1}$	(k-1)(n-1)+2	$[\![\widehat{a}a^{k-1}]\!]$	k(n-1) + 1

Table 2: Homology of  $\mathbf{M}^{\text{bal}}$  and  $M^{\text{cyc}}$  for n-1 even,  $k = 1, 2, \ldots$ 

sider first  $H_2^1$  (see Definition 2.5):

$$h_{2}^{1} = \frac{1}{2} Sh_{1}^{1} + \sigma^{-1} \llbracket xx\hbar^{-1}x^{*} \rrbracket$$
  
=  $\sigma^{-1} \left( \llbracket \widehat{a}x\hbar^{-1}\widehat{a}^{*} \rrbracket + \llbracket x\widehat{a}\hbar^{-1}\widehat{a}^{*} \rrbracket + (-1)^{n-1} \llbracket xx\hbar^{-1}x^{*} \rrbracket \right), \qquad (5.2)$ 

see Remark 5.2 for a discussion of the sign in this formula.

**Remark 5.1.** In our description of h in terms of generalized holomorphic disks, the disks which give rise to (5.2) are degenerate. The first two terms comes from the trivial strip of a with an unconstrained Morse puncture lying in one of the two boundary components. The last term comes from the  $\mathbb{R}$ -family of three punctured constant disks lying in  $\{y\} \times \mathbb{R}$  with one puncture constrained to lie in  $\{y\} \times \mathbb{R}$  and remaining two punctures unconstrained.

Next, consider  $h_0^2$ . Since the evaluation map from the moduli space of holomorphic disks with positive puncture at a into  $\Lambda$  has degree  $\pm 1$  there is algebraically  $\pm 1$  disk which satisfies a point constraint at  $y \in \Lambda$  and we conclude from the description of generalized holomorphic disks that

$$h_0^2 = \sigma^2 [\![\hat{a}^* x^*]\!] \hbar^{-1}, \qquad (5.3)$$

see Remark 5.2 for a discussion of the sign in this formula.

**Remark 5.2.** In order to discuss the overall signs in (5.2) and (5.3), recall that orientations of moduli spaces are fixed by orienting capping operators at Reeb chords, see [8] for details. Note that the overall sign in (5.2) is independent of the choice of capping operator for a since variables associated to a appears quadratically in  $H_2^1$ . Similarly, the overall signs in (5.2) depends  $\mathbb{Z}_2$ -linearly on the orientation data at x, and the overall sign in (5.3) depends

 $\mathbb{Z}_2$ -linearly on the orientations of capping operators at both x and a. It follows that we can choose any sign combination in the two equations and in particular the one under consideration.

To compute  $H_0^2 \underset{1}{\star} H_2^1$  acting on  $\mathbf{M}^{\text{bal}} \otimes \mathbf{M}^{\text{bal}}$  we use the components of  $H_2^1$  in which the puncture at the second leg is excited. We will denote the sum of these by  $\dot{H}_2^1$ . The relevant parts of the Hamiltonian are thus (recall gradings:  $|\hat{a}| = n$ ,  $|\hat{a}^*| = n - 3 - n = -3$ , |x| = 0, and  $|x^*| = n - 3$ )

$$H_0^2 = \sigma^2 \left( \left[ \widehat{a}^* \hbar^{-1} \underline{x}^* \right] - \left[ x^* \hbar^{-1} \underline{\widehat{a}}^* \right] \right)$$

and

$$\dot{H}_2^1 = \sigma^{-1} \left( \left[ \widehat{a}\underline{x}\hbar^{-1}\widehat{a}^* \right] + \left[ x\underline{\widehat{a}}\hbar^{-1}\widehat{a}^* \right] + (-1)^{n-1} \left[ x\underline{x}\hbar^{-1}x^* \right] \right)$$

Equation (5.2) then implies

$$\begin{aligned} H_0^2 \star \dot{H}_2^1 &= -H_0^2 \star \left( \left[\!\left[ \widehat{a}\underline{x} \, \hbar^{-1} \widehat{a}^* \right]\!\right] + \left[\!\left[ x \underline{\hat{a}} \hbar^{-1} \widehat{a}^* \right]\!\right] + (-1)^{n-1} \left[\!\left[ x \underline{x} \hbar^{-1} x^* \right]\!\right] \right) \sigma^{-1} \\ &= \sigma^2 \left( \left[\!\left[ x^* \underline{x} \hbar^{-1} \widehat{a}^* \right]\!\right] + (-1)^n \left[\!\left[ \widehat{a}^* \underline{\hat{a}} \hbar^{-1} \widehat{a}^* \right]\!\right] - \left[\!\left[ \widehat{a}^* \underline{x} \hbar^{-1} x^* \right]\!\right] \right) \sigma^{-1} \\ &= \sigma \left( - \left[\!\left[ x^* \underline{x} \widehat{a}^* \hbar^{-1} \right]\!\right] + (-1)^{n-1} \left[\!\left[ \widehat{a}^* \underline{\hat{a}} \widehat{a}^* \hbar^{-1} \right]\!\right] + \left[\!\left[ \widehat{a}^* \underline{x} x^* \hbar^{-1} \right]\!\right] \right) \\ &= \sigma \left( - \left[\!\left[ \underline{x} \widehat{a}^* \hbar^{-1} x^* \right]\!\right] + (-1)^{n-1} \left[\!\left[ \underline{\hat{a}} \widehat{a}^* \hbar^{-1} \widehat{a}^* \right]\!\right] + \left[\!\left[ \underline{x} x^* \hbar^{-1} \widehat{a}^* \right]\!\right] \right) \end{aligned}$$

Using (4.1) for the tensor of the product  $\Box: H(\mathbf{M}^{\text{bal}}) \otimes H(\mathbf{M}^{\text{bal}}) \rightarrow H(\mathbf{M}^{\text{bal}})$  then gives the following expression for the product, with  $\theta = \hbar^{-1}\sigma^{-1}$ : If n-1 is odd then

$$\begin{split} \underbrace{\llbracket \underline{x} \rrbracket \theta \boxdot \llbracket \underline{x} \rrbracket \theta = 0, \quad \llbracket \underline{x} \rrbracket \theta \boxdot \llbracket \underline{x} a^{2k+1} \rrbracket \theta = 0, \quad \llbracket \underline{x} a^{2k+1} \rrbracket \theta \boxdot \llbracket \underline{x} a^{2l+1} \rrbracket \theta = 0, \\ \underbrace{\llbracket \underline{\hat{a}} a^{2k} \rrbracket \theta \boxdot \llbracket \underline{\hat{a}} a^{2l} \rrbracket \theta = \llbracket \underline{\hat{a}} a^{2(l+k)} \rrbracket \theta, \\ \underbrace{\llbracket \underline{\hat{a}} \rrbracket \theta \boxdot \llbracket \underline{x} \rrbracket \theta = \llbracket \underline{x} \rrbracket \theta, \\ \underbrace{\llbracket \underline{\hat{a}} a^{2k} \rrbracket \theta \boxdot \llbracket \underline{x} \rrbracket \theta = \llbracket \underline{x} a^{2k} \rrbracket \theta = \frac{1}{2} d_{\tau M} \llbracket \underline{\hat{a}} a^{2k} \rrbracket \theta \simeq 0 \text{ if } k > 0, \\ \underbrace{\llbracket \underline{\hat{a}} a^{2k} \rrbracket \theta \boxdot \llbracket \underline{x} a^{2l+1} \rrbracket \theta = \llbracket \underline{x} a^{2(k+l)+1} \rrbracket \theta. \end{split}$$

and if n-1 is even then for  $k, l \ge 0$ 

$$\begin{split} \underline{\llbracket x a^{k} \rrbracket \theta \boxdot \llbracket \underline{x} a^{l} \rrbracket \theta &= 0, \\ \underline{\llbracket \widehat{a} a^{k} \rrbracket \theta \boxdot \llbracket \underline{x} a^{l} \rrbracket \theta &= \llbracket \underline{x} a^{k+l} \rrbracket \theta, \\ \underline{\llbracket \widehat{a} a^{k} \rrbracket \theta \boxdot \llbracket \underline{\widehat{a}} a^{l} \rrbracket \theta &= \llbracket \underline{\widehat{a}} a^{k+l} \rrbracket \theta. \end{split}$$

We conclude from this calculation that the commutative product  $\boxtimes : M^{\text{cyc}} \otimes M^{\text{cyc}} \to M^{\text{cyc}}$  which corresponds to the pair of pants product on  $S\mathbb{H}(T^*S^n)$  has unit represented by the class  $[\![\hat{a}]\!] \in M^{\text{cyc}}$  and multiplication according to the above formulas after substituting  $\theta$  with 1 everywhere.

To finish the example we determine also the operator  $\Delta: H(M^{\text{cyc}}) \to H(M^{\text{cyc}})$  which corresponds to the BV-operator on  $S\mathbb{H}(T^*S^n)$ . If n-1 is odd then

$$\Delta(\llbracket x \rrbracket) = 0, \quad \Delta(\llbracket \widehat{a}a^{2k} \rrbracket) = 0, \ k \ge 0, \quad \Delta(\llbracket xa^{2k+1} \rrbracket) = \llbracket \widehat{a}a^{2k} \rrbracket, \ k \ge 0,$$

and if n-1 is even then

$$\Delta(\llbracket x \rrbracket) = 0, \quad \Delta(\llbracket \widehat{a}a^{k-1}\rrbracket) = 0, \ k \ge 1, \quad \Delta(\llbracket xa^k \rrbracket) = \llbracket \widehat{a}a^{k-1}\rrbracket, \ k \ge 1.$$

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