# Existence of minimal nodal solutions for the Nonlinear Schroedinger equations with $V(\infty)=0$ 

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#### Abstract

We consider the problem $\Delta u+V(x) u=f^{\prime}(u)$ in $\mathbb{R}^{N}$. Here the nonlinearity has a double power behavior and $V$ is invariant under an orthogonal involution, with $V(\infty)=0$. An existence theorem of one pair of solutions which change sign exactly once is given.


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## 1 Introduction

It is well known that stationary states of Nonlinear Schroedinger equations lead to problems of the type

$$
\left\{\begin{array}{l}
-\Delta u+V(x) u=f^{\prime}(u), \quad x \in \mathbb{R}^{N} ;  \tag{P}\\
E_{V}(u)<\infty
\end{array}\right.
$$

where the energy functional is defined by

$$
\begin{equation*}
E_{V}(u)=\frac{1}{2} \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x+\frac{1}{2} \int_{\mathbb{R}^{N}} V(x) u^{2}(x) d x-\int_{\mathbb{R}^{N}} f(u) d x . \tag{1}
\end{equation*}
$$

We consider a function $f \in C^{2}(\mathbb{R}, \mathbb{R})$ even with $f(0)=f^{\prime}(0)=0$ that satisfies the following requirements:

[^0]( $f_{1}$ ) there exists $\mu>2$ such that
\[

$$
\begin{equation*}
0<\mu f(s) \leq f^{\prime}(s) s<f^{\prime \prime}(s) s^{2} \text { for all } s \neq 0 \tag{2}
\end{equation*}
$$

\]

( $f_{2}$ ) there exist positive numbers $c_{0}, c_{2}, p, q$ with $2<p<2^{*}<q$ such that

$$
\begin{gather*}
\begin{cases}c_{0}|s|^{p} \leq f(s) & \text { for }|s| \geq 1 ; \\
c_{0}|s|^{q}<f(s) & \text { for }|s| \leq 1 ;\end{cases}  \tag{3}\\
\left\{\begin{aligned}
\left|f^{\prime \prime}(s)\right| \leq c_{2}|s|^{p-2} & \text { for }|s| \geq 1 \\
\left|f^{\prime \prime}(s)\right| \leq c_{2}|s|^{q-2} & \text { for }|s| \leq 1
\end{aligned}\right. \tag{4}
\end{gather*}
$$

where $2^{*}=\frac{2 N}{N-2}$.
We assume $V \in L^{N / 2}\left(\mathbb{R}^{N}\right)$ and

$$
\begin{equation*}
\|V\|_{L^{N / 2}}<S:=\inf _{u \in \mathcal{D}^{1,2}} \frac{\int_{\mathbb{R}^{N}}|\nabla u|^{2}}{\left(\int_{\mathbb{R}^{N}}|u|^{2^{*}}\right)^{2 / 2^{*}}} \tag{5}
\end{equation*}
$$

In the case of a single-power nonlinearity some paper has been devoted to the existence of positive solutions when potential $V$ vanishes at infinity. Among others we recall [2, 6] and we quote the references therein.

In pioneering work Berestycki and Lions [11, 12] showed the existence of a positive solution in the case $V \equiv 0$ when $f^{\prime \prime}(0)=0, f$ has a supercritical growth near the origin and subcritical at infinity.

More recently in the papers [3, 8, 8, 10, 18] the double-power growth condition $\left(f_{2}\right)$ has been used to obtain the existence of positive solutions for different problems of the tipe ( $\mathscr{P}$ ). In particular, in [8] the authors proved that if $V \geq 0$ and $V>0$ on a set of positive measure the problem ( $\mathscr{P})$ has no ground state solution, i.e. there is no solution $u$ of ( $\mathscr{P})$ which minimizes the functional $E_{V}$ on the Nehari manifold $\mathcal{N}_{V}$, defined by

$$
\begin{equation*}
\mathcal{N}_{V}=\left\{u \in \mathcal{D}^{1,2}:\left\langle\nabla E_{V}(u), u\right\rangle=0, u \neq 0\right\} \tag{6}
\end{equation*}
$$

On the contrary there exists a ground state solution either if $V \leq 0$ and $V<0$ on a set of positive measure, or $V \equiv 0$.

In this paper we are interested in the existence of sign changing solutions. Besides the difficulty posed by the lack of compactness we have another problem: there is no natural regular constraint for sign changing solution of problem ( $\mathscr{P})$. To overcome this difficulty we consider the problem

$$
\left\{\begin{array}{l}
-\Delta u+V(x) u=f^{\prime}(u), \quad x \in \mathbb{R}^{N} \\
E_{V}(u)<\infty \\
u(\tau x)=-u(x)
\end{array}\right.
$$

where $\tau$ is a non trivial orthogonal involution that is a linear orthogonal transformation on $\mathbb{R}^{N}$ such that $\tau \neq \mathrm{Id}$ and $\tau^{2}=\mathrm{Id}$ (Id being the identity on $\mathbb{R}^{N}$ ).

We assume $V(\tau(x))=V(x)$.
By the nontrivial orthogonal involution $\tau$ on $\mathbb{R}^{N}$ we can define a self adjoint linear isometry on $\mathcal{D}^{1,2}$ which we also denote $\tau$. We define

$$
\begin{align*}
& \tau: \mathcal{D}^{1,2} \rightarrow \mathcal{D}^{1,2} \\
& (\tau u)(x):=-u(\tau(x)) \tag{7}
\end{align*}
$$

If $u(\tau x)=-u(x)$, it will be called $\tau$-antisymmetric. Note that non trivial antisymmetric solutions are changing sign or nodal solutions. Nodal solutions which change sign exactly once will be called minimal nodal solutions.

We define

$$
\mathcal{N}_{V}^{\tau}=\mathcal{N}_{V} \cap \mathcal{D}_{\tau}^{1,2} \quad \text { where } \quad \mathcal{D}_{\tau}^{1,2}=\left\{u \in \mathcal{D}^{1,2}: \tau u=u\right\}
$$

The non trivial antisymmetric solutions of $\left(\mathscr{P}_{\tau}\right)$ are the critical points of $E_{V}$ on $\mathcal{N}_{V}^{\tau}$

We set now

$$
\begin{array}{ll}
\mu_{V}=\mu_{V}\left(\mathbb{R}^{N}\right):=\inf _{\mathcal{N}_{V}} E_{V} ; & \mu_{0}=\mu_{0}\left(\mathbb{R}^{N}\right):=\inf _{\mathcal{N}_{0}} E_{0} \\
\mu_{V}^{\tau}=\mu_{V}^{\tau}\left(\mathbb{R}^{N}\right):=\inf _{\mathcal{N}_{V}^{\tau}} E_{V} ; & \mu_{0}^{\tau}=\mu_{0}^{\tau}\left(\mathbb{R}^{N}\right):=\inf _{\mathcal{N}_{0}^{\tau}} E_{0} \tag{9}
\end{array}
$$

We shall prove the following results.
Theorem 1. If $V(x)>0$ for almost every $x$, then

$$
\mu_{V}^{\tau}=\mu_{0}^{\tau}=2 \mu_{0}
$$

and $\mu_{V}^{\tau}$ is not achieved. Then the problem (師) has no solution of minimal energy.

We consider the following class of potentials.

$$
V_{y}(x)=\left\{\begin{array}{cl}
a|x-y|-1 & |x-y|<1  \tag{10}\\
a|x-\tau y|-1 & |x-\tau y|<1 \\
0 & \text { elsewhere }
\end{array}\right.
$$

where $a \in \mathbb{R}$ is chosen such that $\|V\|_{L^{N / 2}}<S, S$ as in (5). We can prove the following existence result.

Theorem 2. For the potential $V_{y}$ such that $|y-\tau y|$ is sufficiently large we have that $\mu_{V}^{\tau}<\mu_{0}^{\tau}$ and it is achieved. Then the problem ( $\mathscr{P}_{\tau}$ ) has at least one pair of antisymmetric solutions which change sign exactly once, and the energy of these solutions is minimal.

We want to mention some recent work about sign changing solutions. The existence of a sequence of nodal solutions and some properties for the number of their nodal domains has been obtained in 5 considering the problem in a bounded smooth domain $\Omega$ with $V \equiv 0$, and in [4] in $\mathbb{R}^{N}$ with essinf $V>0$.

In [15] there is a theorem of multiplicity of solutions for the problem $-\Delta u+V u=q(x)|u|^{p-2} u$ where $V(x)$ and $q(x)$ tend to some positive number $V_{\infty}$ and $q_{\infty}$ respectively as $|x| \rightarrow \infty$. However no precise information is given whether there are sign changing solutions or not. If $V \equiv 1$ and $q(x)$ suitable chosen, with $\|q-1\|_{\infty}$ small, Hirano [16] prove the existence of at least two pairs of sign changing solutions.

In [14] the equation $-\Delta u+\lambda u=|u|^{p^{*}-2} u, \lambda>-\lambda_{1}$ on a symmetric domain is considered and the effect of the domain topology on the number of minimal nodal solutions in studied.

The plan of the paper is the following.
In section 2 we recall some technical result concerning the appropriate function spaces required by the growth properties of $f$; the proof of these results are contained in [7, 8, 10]. In section 3 we prove a splitting lemma which is a variant of a well known result of [19]; this lemma is the ingredient to handle the problem with lack of compactness. In section 4 we prove our results.

We will use the following notations

- $u^{+}=\max (0, u) ; u^{-}=\min (0, u)$;
- $\mathcal{D}^{1,2}=\mathcal{D}^{1,2}\left(\mathbb{R}^{N}\right)=$ completion of $C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ with respect to the norm

$$
\|u\|_{\mathcal{D}^{1,2}}=\left(\int_{\mathbb{R}^{n}}|\nabla u|^{2}\right)^{1 / 2}
$$

- $v_{y}(x)=v(x+y)$;
- $G_{V}(u)=\left\langle\nabla E_{V}(u), u\right\rangle=\int|\nabla u|^{2}+V u^{2}-f^{\prime}(u) u$
- $g_{u}(t)=g_{u}^{V}(t):=E_{V}(t u)=\int_{\mathbb{R}^{N}} \frac{t^{2}}{2}\left(|\nabla u|^{2}+V u^{2}\right)-f(t u) d x$
- $B_{R}=\left\{x \in \mathbb{R}^{N}:|x|<R\right\} ;$
- $B_{R}(z)=B(z, R)=\left\{x \in \mathbb{R}^{N}:|x-z|<R\right\} ;$
- $A^{C}=\mathbb{R}^{N} \backslash A$, where $A \subset \mathbb{R}^{N}$.


## 2 Variational Setting

In order to study the functional $E_{V}$, by the growth assumption on $f$, is useful to consider the functional space $L^{p}+L^{q}$, where $2<p<2^{*}<q$. Hereafter we recall some result contained in [7, 8, 10].

Given $p \neq q$, we consider the space $L^{p}+L^{q}$ made up of the functions $v: \mathbb{R}^{N} \rightarrow \mathbb{R}$ such that

$$
\begin{equation*}
v=v_{1}+v_{2} \text { with } v_{1} \in L^{p}, v_{2} \in L^{q} . \tag{11}
\end{equation*}
$$

The space $L^{p}+L^{q}$ is a Banach space equipped with the norm:

$$
\begin{equation*}
\|v\|_{L^{p}+L^{q}}=\inf \left\{\left\|v_{1}\right\|_{L^{p}}+\left\|v_{2}\right\|_{L^{q}}: v_{1} \in L^{p}, v_{2} \in L^{q}, v_{1}+v_{2}=v\right\} . \tag{12}
\end{equation*}
$$

It is well known (see, for example [13]) that $L^{p}+L^{q}$ coincides with the dual of $L^{p^{\prime}} \cap L^{q^{\prime}}$. Then:

$$
\begin{equation*}
L^{p}+L^{q}=\left(L^{p^{\prime}} \cap L^{q^{\prime}}\right)^{\prime} \text { with } p^{\prime}=\frac{p}{p-1}, q^{\prime}=\frac{q}{q-1} . \tag{13}
\end{equation*}
$$

We recall some results useful for this paper.
Remark 3. Set $\Gamma_{u}=\left\{u \in \mathbb{R}^{N}:|u(x)| \geq 1\right\}$. We have

1. if $v \in L^{p}+L^{q}$, the following inequalities hold:

$$
\begin{aligned}
& \max \left[\|v\|_{L^{q}\left(\mathbb{R}^{N} \backslash \Gamma_{v}\right)}-1, \frac{1}{1+\left|\Gamma_{v}\right|^{\frac{1}{\tau}}}\|v\|_{L^{p}\left(\Gamma_{v}\right)}\right] \leq \\
\leq & \|v\|_{L^{p}+L^{q}} \leq \\
\leq & \max \left[\|v\|_{L^{q}\left(\mathbb{R}^{N} \backslash \Gamma_{v}\right)},\|v\|_{L^{p}\left(\Gamma_{v}\right)}\right]
\end{aligned}
$$

when $\tau=\frac{p q}{q-p}$;
2. let $\left\{v_{n}\right\} \subset L^{p}+L^{q}$. Then $\left\{v_{n}\right\}$ is bounded in $L^{p}+L^{q}$ if and only if the sequences $\left\{\left|\Gamma_{v_{n}}\right|\right\}$ and $\left\{\|v\|_{L^{q}\left(\mathbb{R}^{N} \backslash \Gamma_{v_{n}}\right)}+\|v\|_{L^{p}\left(\Gamma_{v_{n}}\right)}\right\}$ are bounded.
3. we have $L^{2^{*}} \subset L^{p}+L^{q}$ when $2<p<2^{*}<q$. Then, by Sobolev inequality, we get the continuous embedding

$$
\mathcal{D}^{1,2}\left(\mathbb{R}^{N}\right) \subset L^{p}+L^{q}
$$

Remark 4. If $f$ satisfies the hypothesis that we have made in the previous section, we have that

1. $f^{\prime}$ is a bounded map from $L^{p}+L^{q}$ into $L^{p / p-1} \cap L^{q / q-1}$;
2. $f^{\prime \prime}$ is a bounded map from $L^{p}+L^{q}$ into $L^{p / p-2} \cap L^{q / q-2}$;
3. $f^{\prime \prime}$ is a continuous map from $L^{p}+L^{q}$ into $L^{p / p-2} \cap L^{q / q-2}$.

At last we recall some result on Nehari manifolds. For the proofs we refer to [8, 10].
Remark 5. The functional $E_{V}$ is of class $C^{2}$ and it holds

$$
\begin{equation*}
\left\langle\nabla E_{V}(u), v\right\rangle=\int \nabla u \nabla v+V u v-f^{\prime}(u) v d x . \tag{14}
\end{equation*}
$$

Moreover the Nehari manifold defined as

$$
\begin{equation*}
\mathcal{N}_{V}=\left\{u \in \mathcal{D}^{1,2} \backslash 0: \int|\nabla u|^{2}+V u^{2}-f^{\prime}(u) u d x=0\right\} \tag{15}
\end{equation*}
$$

is of class $C^{1}$ and its tangent space at the point $u$ is

$$
\begin{equation*}
T_{u} \mathcal{N}_{V}=\left\{u \in \mathcal{D}^{1,2}: \int 2 \nabla u \nabla v+2 V u v-f^{\prime}(u) v d x-f^{\prime \prime}(u) u v=0\right\} \tag{16}
\end{equation*}
$$

Remark 6. We have

$$
\begin{equation*}
\inf _{u \in \mathcal{N}_{V}}\|u\|_{\mathcal{D}^{1,2}}>0 \tag{17}
\end{equation*}
$$

Proof. At first notice that, by 5

$$
\begin{equation*}
\exists c>0 \text { s.t } \int|\nabla u|^{2}+V u^{2} \geq c\|u\|_{\mathcal{D}^{1,2}}^{2}, \forall u \in \mathcal{D}^{1,2} \tag{18}
\end{equation*}
$$

Now, let $\left\{u_{n}\right\}$ a minimizing sequence in $\mathcal{N}_{V}$. By contradiction, we suppose that $u_{n}$ converges to 0 . We set $t_{n}=\left\|u_{n}\right\|_{\mathcal{D}^{1,2}}$, hence we can write $u_{n}=t_{n} v_{n}$ where $\left\|v_{n}\right\|_{\mathcal{D}^{1,2}}=1$. By claim 3 of Remark 3, the sequence is bounded in
$L^{p}+L^{q}$. Since $u_{n} \in \mathcal{N}_{V}$ and $t_{n}$ converges to 0 , we have

$$
\begin{aligned}
& c t_{n}=\frac{c}{t_{n}}\left\|u_{n}\right\|_{\mathcal{D}^{1,2}}^{2} \leq \frac{1}{t_{n}} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2}+V u_{n}^{2}=\int_{\mathbb{R}^{N}} f^{\prime}\left(t_{n} v_{n}\right) v_{n} \leq \\
& \leq c_{1} t_{n}^{q-1} \int_{\mathbb{R}^{N} \backslash \Gamma_{t_{n} v_{n}}}\left|v_{n}\right|^{q}+c_{1} t_{n}^{p-1} \int_{\Gamma_{t_{n} v_{n}}}\left|v_{n}\right|^{p} \leq \\
& \leq c_{1} t_{n}^{q-1} \int_{\mathbb{R}^{N} \backslash \Gamma_{t_{n} v_{n}}}\left|v_{n}\right|^{q}+c_{1} t_{n}^{p-1} \int_{\Gamma_{v_{n}}}\left|v_{n}\right|^{p} \leq \\
& \leq c_{1} t_{n}^{q-1} \int_{\mathbb{R}^{N} \backslash \Gamma_{v_{n}}}\left|v_{n}\right|^{q}+c_{1} t_{n}^{q-1} \int_{\left(\mathbb{R}^{N} \backslash \Gamma_{t_{n} v_{n}}\right) \cap \Gamma_{v_{n}}} \frac{\left|v_{n}\right|^{p}}{t_{n}^{q-p}}+c_{1} t_{n}^{p-1} \int_{\Gamma_{v_{n}}}\left|v_{n}\right|^{p} \leq \\
& \leq c_{1} t_{n}^{q-1} \int_{\mathbb{R}^{N} \backslash \Gamma_{v_{n}}}\left|v_{n}\right|^{q}+2 c_{1} t_{n}^{p-1} \int_{\Gamma_{v_{n}}}\left|v_{n}\right|^{p} .
\end{aligned}
$$

Hence we get

$$
c \leq c_{1} t_{n}^{q-2} \int_{\mathbb{R}^{N} \backslash \Gamma_{v_{n}}}\left|v_{n}\right|^{q}+2 c_{1} t_{n}^{p-2} \int_{\Gamma_{v_{n}}}\left|v_{n}\right|^{p}
$$

and by claim 2 of Remark 3 we get the contradiction.
Remark 7. We have that for any given $u \in \mathcal{D}^{1,2} \backslash\{0\}$, there exists a unique real number $t_{u}^{V}>0$ such that $t_{u}^{V} u \in \mathcal{N}_{V}$ and $E_{V}\left(t_{u}^{V} u\right)$ is the maximum for the function

$$
t \rightarrow E_{V}(t u), \quad t>0
$$

Proof. Given $u \neq 0$ we set, for $t \geq 0$ :

$$
\begin{equation*}
g_{u}(t)=g_{u}^{V}(t):=E_{V}(t u)=\int_{\mathbb{R}^{N}} \frac{t^{2}}{2}\left(|\nabla u|^{2}+V u^{2}\right)-f(t u) d x \tag{19}
\end{equation*}
$$

We have:

$$
\begin{align*}
g_{u}^{\prime}(t) & =\int_{\mathbb{R}^{N}} t|\nabla u|^{2}+V t u^{2}-u f^{\prime}(t u) d x  \tag{20}\\
g_{u}^{\prime \prime}(t) & =\int_{\mathbb{R}^{N}}|\nabla u|^{2}+V u^{2}-u^{2} f^{\prime \prime}(t u) d x . \tag{21}
\end{align*}
$$

By hypothesis on $f$, if $g_{u}^{\prime}(\bar{t})=0$ we have

$$
\begin{equation*}
\bar{t}^{2} g_{u}^{\prime \prime}(\bar{t})=\int_{\mathbb{R}^{N}} \bar{t} u f^{\prime}(\bar{t} u)-\bar{t}^{2} u^{2} f^{\prime \prime}(\bar{t} u) d x<0 \tag{22}
\end{equation*}
$$

then $\bar{t}$ is a maximum point for $g_{u}$. Furthermore $0=g_{u}(0)=g_{u}^{\prime}(0)$ and $g_{u}^{\prime \prime}(0)>0$ by the hypothesis on $V$, then 0 is a local minimum point for $g_{u}$. By (3), for $t \geq 1$, we have

$$
\begin{align*}
g_{u}(t) & \leq \int_{\mathbb{R}^{N}} \frac{t^{2}}{2}\left(|\nabla u|^{2}+V u^{2}\right) d x-c_{0} \int_{\{|t u \leq 1|\}}|t u|^{q} d x-c_{0} \int_{\{|t u>1|\}}|t u|^{p} d x \leq \\
& \leq \int_{\mathbb{R}^{N}} \frac{t^{2}}{2}\left(|\nabla u|^{2}+V u^{2}\right) d x-c_{0} \int_{\{|t u>1|\}}|t u|^{p} d x \leq  \tag{23}\\
& \leq \frac{t^{2}}{2} \int_{\mathbb{R}^{N}}\left(|\nabla u|^{2}+V u^{2}\right) d x-c_{0} t^{p} \int_{\{|u>1|\}}|u|^{p} d x
\end{align*}
$$

the last quantity diverges negatively as $t \rightarrow \infty$, since $p>2$, and the claim follows.

We search antisymmetric solutions of $\left(\mathscr{P}_{\tau}\right)$. To do that, we look for critical points of the restriction of $E_{V}$ to $\mathcal{N}_{V}^{\tau}$. In fact, if $\bar{u} \in \mathcal{N}_{V}^{\tau}$ is a critical point of the restriction of $E_{V}$ to $\mathcal{N}_{V}^{\tau}$, then

$$
E_{V}^{\prime}(\bar{u}) \varphi=\left\langle\nabla E_{V}(\bar{u}), \varphi\right\rangle=0 \quad \forall \varphi \in T_{\bar{u}} \mathcal{N}_{V} \cap \mathcal{D}_{\tau}^{1,2}\left(\mathbb{R}^{N}\right)
$$

But $\nabla E_{V}(\bar{u})=\tau \nabla E_{V}(\tau \bar{u})=\tau \nabla E_{V}(\bar{u})$, so we can see that $\nabla E_{V}(\bar{u})=0$.

## 3 A splitting lemma

We recall that a sequence $\left\{u_{n}\right\}_{n} \in \mathcal{D}^{1,2}$ such that $E_{V}\left(u_{n}\right) \rightarrow c$, and there exists a sequence $\varepsilon_{n} \rightarrow 0$ s.t. $\left|E_{V}^{\prime}\left(u_{n}\right) \varphi\right| \leq \varepsilon_{n}\|\varphi\|$, for all $\varphi \in \mathcal{D}^{1,2}$ is a Palais-Smale sequence at level $c$ for $E_{V}$.

In the same way we say that $\left\{u_{n}\right\}_{n} \in \mathcal{N}_{V}^{\tau}$ such that $E_{V}\left(u_{n}\right) \rightarrow c$, and there exists a sequence $\varepsilon_{n} \rightarrow 0$ s.t. $\left|E_{V}^{\prime}\left(u_{n}\right) \varphi\right| \leq \varepsilon_{n}\|\varphi\|$, for all $\varphi \in T_{u_{n}} \mathcal{N}_{V} \cap \mathcal{D}_{\tau}^{1,2}$ is a Palais-Smale sequence at level $c$ for $E_{V}$ restricted to $\mathcal{N}_{V}^{\tau}$.

A functional $f$ satisfies the $(P S)_{c}$ condition if all the Palais-Smale sequences at level $c$ converge.

Unfortunately the functional $E_{V}$ on $\mathcal{N}_{V}^{\tau}$ does not satisfy the $P S$ condition in all the energy range. The following lemma provides a description of the PS sequences in $\mathcal{D}_{\tau}^{1,2}$.

This splitting lemma is a fundamental tool to obtain the claimed results. The main idea of this result spread over a result of M. Struwe [19] that described all the $P S$ sequences for $E_{V}$ on $H^{1}(\Omega)$ when $f(u)=u|u|^{2^{*}-2}$ and $V(x)=\lambda u$

Lemma 8. Let $\left\{u_{n}\right\}_{n}$ a PS sequence at level c for the functional $E_{V}$ restricted to the manifold $\mathcal{N}_{V}^{\tau}$. Then, up to a subsequence, there exist two integers $k_{1}, k_{2} \geq 0, k_{1}+k_{2}$ sequences $\left\{y_{n}^{j}\right\}_{n}$, an antisymmetric solution $u^{0}$ of the problem $-\Delta u+V u=f^{\prime}(u)$, and $k_{1}$ solutions $u^{j}, j=1, \ldots k_{1}$, and $k_{2}$ antisymmetric solutions $u^{j}, j=k_{1}+1, \ldots, k_{1}+k_{2}$, of the problem $-\Delta u=f^{\prime}(u)$ such that

1. if $j=1, \ldots, k_{1}$, then $\tau y_{n}^{j} \neq y_{n}^{j}$, and $\left|y_{n}^{j}\right| \rightarrow \infty$ as $n \rightarrow \infty$
2. if $j=k_{1}+1, \ldots, k_{2}$, then $\tau y_{n}^{j}=y_{n}^{j}$, and $\left|y_{n}^{j}\right| \rightarrow \infty$ as $n \rightarrow \infty$
3. $u_{n}(x)=u^{0}(x)+\sum_{j=1}^{k_{1}}\left[u^{j}\left(x-y_{n}^{j}\right)+\tau u^{j}\left(x-y_{n}^{j}\right)\right]+\sum_{j=k_{1}+1}^{k_{2}} u^{j}\left(x-y_{n}^{j}\right)+o(1)$
4. $E_{V}\left(u_{n}\right) \rightarrow E_{V}\left(u^{0}\right)+2 \sum_{j=1}^{k_{1}} E_{0}\left(u^{j}\right)+\sum_{j=k_{1}+1}^{k_{2}} E_{0}\left(u^{j}\right)$

Proof. Since $u_{n}$ is a $P S$ sequence for the functional $E_{V}$ restricted to the manifold $\mathcal{N}_{V}$, then $u_{n}$ is a $P S$ sequence for the functional $E_{V}$. For Step 1 of [8, Lemma 3.3] we get that $u_{n}$ converges to $u^{0}$ weakly in $\mathcal{D}^{1,2}$ (up to subsequence) and $u^{0}$ solves $-\Delta u+V u=f^{\prime}(u)$.

Since $\tau u_{n}=u_{n}$, we have $\tau u^{0}=u^{0}$. In fact, if $u_{n} \stackrel{\mathcal{D}^{1,2}}{\sim} u^{0}$, then, for every $R>0$, we have that $u_{n} \xrightarrow{L^{2}\left(B_{R}\right)} u^{0}$, so $u_{n}(x) \rightarrow u^{0}(x)$ for almost all $x$.

We set

$$
\psi_{n}(x)=u_{n}(x)-u^{0}(x) .
$$

Then $\tau \psi_{n}=\psi_{n}$, and $\psi_{n} \rightharpoonup 0$ weakly in $\mathcal{D}^{1,2}$. By Steps 2 and 4 of [8, Lemma 3.3], we get that $\psi_{n}$ is a $P S$ sequence for $E_{0}$. If $\psi_{n} \nrightarrow 0$ strongly in $\mathcal{D}^{1,2}$, then by Steps 3 and 4 of [8, Lemma 3.3] we get that there exists a sequence $\left\{\xi_{n}\right\} \subset \mathbb{R}^{N}$ with $\left|\xi_{n}\right| \rightarrow \infty$ as $n \rightarrow \infty$, and $\psi_{n}\left(x+\xi_{n}\right) \rightharpoonup u^{1}(x)$ where $u^{1}$ is a weak solution of $-\Delta u=f^{\prime}(u)$.

We consider in $\mathbb{R}^{N}=\Gamma \oplus \Gamma^{\perp}$, where $\Gamma:=\left\{x \in \mathbb{R}^{N}: \tau(x)=x\right\}$. We consider $P_{\Gamma}$ the projection on the subspace $\Gamma$. At this point we must distinguish two cases

Case I: if $\left|\xi_{n}-\tau\left(\xi_{n}\right)\right|$ is bounded we define $y_{n}=P_{\Gamma} \xi_{n}$.
Case II: if $\left|\xi_{n}-\tau\left(\xi_{n}\right)\right|$ is unbounded we define $y_{n}=\xi_{n}$.

Case I In this case there exist a solution $\widetilde{u} \in \mathcal{D}_{\tau}^{1,2} \backslash\{0\}$ of $-\Delta u=f^{\prime}(u)$ and a $P S$ sequence $\left\{\widetilde{\psi}_{n}\right\}_{n}$ for $E_{0}$ such that

$$
\begin{equation*}
\psi_{n}=\widetilde{\psi}_{n}+\widetilde{u}\left(x+y_{n}\right) \tag{24}
\end{equation*}
$$

$\widetilde{\psi}_{n} \rightharpoonup 0$ weakly in $\mathcal{D}^{1,2}$, and

$$
\begin{equation*}
E_{0}\left(\widetilde{\psi}_{n}\right)=E_{0}\left(\psi_{n}\right)-E_{0}(\widetilde{u})+o(1) \tag{25}
\end{equation*}
$$

We can assume, without loss of generality, that $\xi_{n}=P_{\Gamma} \xi_{n}+w$, where $w \in \Gamma^{\perp}$. We now consider the sequence $\left\{\psi_{n}\left(x+y_{n}\right)\right\}_{n}$ which is bounded: hence, up to subsequence $\left\{\psi_{n}\left(x+y_{n}\right)\right\}_{n}$ converges to $\widetilde{u}(x)$ weakly in $\mathcal{D}^{1,2}$, $\tilde{u}(x)=u^{1}(x-w) \neq 0$, then $-\Delta \widetilde{u}=f^{\prime}(\widetilde{u})$. Furthermore, because $\tau y_{n}=y_{n}$ we have that $\tau \widetilde{u}=\widetilde{u}$. We define

$$
\begin{equation*}
\widetilde{\psi}_{n}(x):=\psi_{n}(x)-\widetilde{u}\left(x-y_{n}\right) . \tag{26}
\end{equation*}
$$

We will verify that $\widetilde{\psi}_{n}$ is a $P S$ sequence for $E_{0}$. Indeed by Lemma 2.11 of [10] we get

$$
\begin{aligned}
E_{0}\left(\widetilde{\psi}_{n}\right) & =E_{0}\left(\psi_{n}(x)-\widetilde{u}\left(x-y_{n}\right)\right)=E_{0}\left(\psi_{n}\left(x+y_{n}\right)-\widetilde{u}(x)\right)= \\
& =E_{0}\left(\psi_{n}\right)-E_{0}(\widetilde{u})+o(1),
\end{aligned}
$$

and, because $\left\{\psi_{n}\right\}_{n}$ is a $P S$ sequence for $E_{0}$, we have that $E_{0}\left(\psi_{n}\right)$ converges, so $E_{0}\left(\widetilde{\psi}_{n}\right)$ converges, also.

Again, since $\left\{\psi_{n}\right\}$ is a $P S$ sequence, we have that exists an $\varepsilon_{n} \rightarrow 0$ such that, for all $\varphi \in C_{0}^{\infty}$

$$
\begin{aligned}
& \left|\left(E_{0}\right)^{\prime}\left(\widetilde{\psi}_{n}\right)[\varphi]\right|=\left|\int \nabla \psi_{n} \nabla \varphi-\int \nabla \widetilde{u}_{-y_{n}} \nabla \varphi-\int f^{\prime}\left(\psi_{n}-\widetilde{u}_{-y_{n}} \varphi\right)\right|= \\
& =\left|\int\left[f^{\prime}\left(\psi_{n}(x)\right)-f^{\prime}\left(\widetilde{u}\left(x-y_{n}\right)\right)-f^{\prime}\left(\psi_{n}(x)-\widetilde{u}\left(x-y_{n}\right)\right)\right] \varphi(x) d x+\varepsilon_{n} \| \varphi\right|| |= \\
& =\left|\int\left[f^{\prime}\left(\psi_{n}\left(z+y_{n}\right)\right)-f^{\prime}(\widetilde{u}(z))-f^{\prime}\left(\psi_{n}\left(z+y_{n}\right)-\widetilde{u}(z)\right)\right] \varphi_{y_{n}}(z) d z+\varepsilon_{n}\|\varphi\|\right| .
\end{aligned}
$$

Now we can choose an $R>0$ and split this integral as follows.

$$
\begin{aligned}
\left|\left(E_{0}\right)^{\prime}\left(\widetilde{\psi}_{n}\right)[\varphi]\right|= & \mid \int_{B_{R}}\left[f^{\prime}\left(\psi_{n}\left(x+y_{n}\right)\right)-f^{\prime}(\widetilde{u})\right] \varphi_{y_{n}}(z) d z+ \\
& +\int_{\mathbb{R}^{N} \backslash B_{R}}\left[f^{\prime}\left(\psi_{n}\left(x+y_{n}\right)\right)-f^{\prime}\left(\psi_{n}\left(x+y_{n}\right)-\widetilde{u}\right)\right] \varphi_{y_{n}} d z- \\
& -\int_{\mathbb{R}^{N} \backslash B_{R}} f^{\prime}(\widetilde{u}) \varphi_{y_{n}}-\int_{B_{R}} f^{\prime}\left(\psi_{n}\left(x+y_{n}\right)-\widetilde{u}\right) \varphi_{y_{n}} d z+\varepsilon_{n}\|\varphi\| \leq \leq \\
\leq & A_{n} \gamma_{n, R}\|\varphi\|_{\mathcal{D}^{1,2}}+B_{n} M_{R}\|\varphi\|_{\mathcal{D}^{1,2}}+\varepsilon_{n}\|\varphi\|_{\mathcal{D}^{1,2}}
\end{aligned}
$$

where

$$
\begin{aligned}
& A_{n}=\left\|f^{\prime \prime}\left(\theta \psi_{n}\left(\cdot+y_{n}\right)+(1-\theta) \widetilde{u}\right)-f^{\prime \prime}\left(\theta \psi_{n}\left(\cdot+y_{n}\right)-\theta \widetilde{u}\right)\right\|_{L^{p / p-2}} ; \\
& \gamma_{n, R}=\left\|\psi_{n}\left(\cdot-y_{n}\right)-\widetilde{u}\right\|_{L^{p}\left(B_{R}\right)} ; \\
& B_{n}=\left\|f^{\prime \prime}\left(\psi_{n}\left(\cdot+y_{n}\right)+\theta \widetilde{u}\right)-f^{\prime \prime}(\theta \widetilde{u})\right\|_{L^{p / p-2} \cap L^{q / q-2}} ; \\
& M_{R}=\|\widetilde{u}\|_{L^{p}+L^{q}\left(\mathbb{R}^{N} \backslash B_{R}\right)},
\end{aligned}
$$

for some $0<\theta<1$.
By [10, Lemma 2.11] we have that both $A_{n}$ and $B_{n}$ are bounded. Since $M_{R} \rightarrow 0$ as $R \rightarrow+\infty$ and, given $R, \gamma_{n, R} \rightarrow 0$ as $n \rightarrow+\infty$, we get the claim. Case II In this case there exist a solution $u^{1} \neq 0$ of $-\Delta u=f^{\prime}(u)$ and a $P S$ sequence $\left\{\widetilde{\psi}_{n}\right\}_{n} \subset \mathcal{D}_{\tau}^{1,2}$ for $E_{0}$ such that

$$
\begin{align*}
& \psi_{n}(x)=\widetilde{\psi}_{n}(x)+u^{1}\left(x-y_{n}\right)-u^{1}\left(\tau x-y_{n}\right)+o(1)  \tag{27}\\
& E_{0}\left(\widetilde{\psi}_{n}\right)=E_{0}\left(\psi_{n}\right)-2 E_{0}\left(u^{1}\right)+o(1) \tag{28}
\end{align*}
$$

We define $\widetilde{\psi}_{n}=\psi_{n}-\gamma_{n}$,

$$
\begin{equation*}
\gamma_{n}(x)=u^{1}\left(x-y_{n}\right) \chi\left(\frac{\left|x-y_{n}\right|}{\rho_{n}}\right)-u^{1}\left(\tau x-y_{n}\right) \chi\left(\frac{\left|x-\tau y_{n}\right|}{\rho_{n}}\right), \tag{29}
\end{equation*}
$$

where $\rho_{n}:=\frac{\left|y_{n}-\tau y_{n}\right|}{2} \rightarrow \infty$ for $n \rightarrow \infty$, and, as usual, $\chi: \mathbb{R}_{0}^{+} \rightarrow[0,1]$ is a $C^{\infty}$ function such that $\chi(s) \equiv 0$ for all $s \geq 2, \chi(s) \equiv 1$ for all $s \leq 1$ and $\left|\chi^{\prime}(s)\right| \leq 2$ for all $s$.

It is trivial that $\tau \gamma_{n}=\gamma_{n}$, so $\tau \widetilde{\psi}_{n}=\widetilde{\psi}_{n}$. Furthermore, easily we have

$$
\psi_{n}(x)=\widetilde{\psi}_{n}(x)+u^{1}\left(x-y_{n}\right)-u^{1}\left(\tau x-y_{n}\right)+o(1) .
$$

Now we have to prove (28), and to show that $\tilde{\psi}_{n}$ is a $P S$ sequence.
At first we prove that

$$
\begin{equation*}
\left\|\tilde{\psi}_{n}\right\|_{\mathcal{D}^{1,2}}^{2}=\left\|\psi_{n}-\gamma_{n}\right\|^{2}=\left\|\psi_{n}\right\|^{2}+2\left\|u^{1}\right\|^{2}+o(1) \tag{30}
\end{equation*}
$$

In fact we have that $\left\|\psi_{n}-\gamma_{n}\right\|^{2}=\left\|\psi_{n}\right\|^{2}+\left\|\gamma_{n}^{2}\right\|^{2}-2\left(\psi_{n}, \gamma_{n}\right)_{\mathcal{D}^{1,2}}$, and it is easy to see that $\left\|\gamma_{n}^{2}\right\|^{2} \rightarrow 2\left\|u^{1}\right\|^{2}$. Furthermore

$$
\begin{align*}
\left(\psi_{n}, \gamma_{n}\right)_{\mathcal{D}^{1,2}}= & \int \nabla \psi_{n} \nabla\left(u^{1}\left(x-y_{n}\right) \chi\left(\frac{\left|x-y_{n}\right|}{\rho_{n}}\right)\right)+  \tag{31}\\
& +\int \nabla \psi_{n} \nabla\left(u^{1}\left(\tau x-y_{n}\right) \chi\left(\frac{\left|\tau x-y_{n}\right|}{\rho_{n}}\right)\right),
\end{align*}
$$

and the first term converges to $\int\left|\nabla u^{1}\right|$. For the second term we have

$$
\begin{aligned}
& \int \nabla \psi_{n} \nabla\left(u^{1}\left(\tau x-y_{n}\right) \chi\left(\frac{\left|\tau x-y_{n}\right|}{\rho_{n}}\right)\right)= \\
& =\int\left(\nabla \psi_{n} \nabla u^{1}\left(\tau x-y_{n}\right)\right) \chi(\cdot)+\int\left(\nabla \psi_{n} \nabla \chi(\cdot)\right) u^{1}\left(\tau x-y_{n}\right) .
\end{aligned}
$$

The last term of the equation vanishes when $n \rightarrow \infty$, while, remembering that $\psi_{n}$ is symmetric, and setting $z=\tau x-y_{n}$, we have

$$
\begin{aligned}
& \int\left(\nabla \psi_{n}(x) \nabla u^{1}\left(\tau x-y_{n}\right)\right) \chi(\cdot) d x= \\
& \int\left(\tau \nabla \psi_{n}\left(z+y_{n}\right) \nabla u^{1}(z)\right) \chi\left(\frac{z}{\rho_{n}}\right) d z \rightarrow-\int\left|\nabla u^{1}\right|^{2},
\end{aligned}
$$

so we have proved (30).
We want now to estimate

$$
\int f\left(\widetilde{\psi}_{n}\right)=\int f\left(\psi_{n}-\gamma_{n}\right)
$$

Set

$$
\begin{aligned}
& I_{1}=\int_{\left|x-y_{n}\right|<2 \rho_{n}} f\left(\psi_{n}(x)-u^{1}\left(x-y_{n}\right) \chi\left(\frac{\left|x-y_{n}\right|}{\rho_{n}}\right)\right) \\
& I_{2}=\int_{\left|\tau x-y_{n}\right|<2 \rho_{n}} f\left(\psi_{n}(x)+u^{1}\left(\tau x-y_{n}\right) \chi\left(\frac{\left|\tau x-y_{n}\right|}{\rho_{n}}\right)\right) ; \\
& I_{3}=\int_{\left\{B\left(y_{n}\right)_{2 \rho_{n}} \cup B\left(\tau y_{n}\right) 2 \rho_{n}\right\}^{C}} f\left(\psi_{n}(x)\right),
\end{aligned}
$$

we have

$$
\begin{equation*}
\int f\left(\widetilde{\psi}_{n}\right)=I_{1}+I_{2}+I_{3} \tag{32}
\end{equation*}
$$

We have that

$$
\begin{equation*}
\left[\psi_{n}\left(z+y_{n}\right)-u^{1}(z) \chi\left(\frac{|z|}{\rho_{n}}\right)\right] \chi\left(\frac{|z|}{\rho_{n}}\right) \rightharpoonup 0 \text { in } \mathcal{D}^{1,2} . \tag{33}
\end{equation*}
$$

By [10, Lemma 2.11], then we have that

$$
\begin{aligned}
I_{1} & =\int_{|z|<2 \rho_{n}} f\left(\psi_{n}\left(z+y_{n}\right)-u^{1}(z) \chi\left(\frac{|z|}{\rho_{n}}\right)\right)= \\
& =\int_{|z|<2 \rho_{n}} f\left(\psi_{n}\left(z+y_{n}\right)\right)-\int_{|z|<2 \rho_{n}} f\left(u^{1}(z) \chi\left(\frac{|z|}{\rho_{n}}\right)\right)+o(1)= \\
& =\int_{|z|<2 \rho_{n}} f\left(\psi_{n}\left(z+y_{n}\right)\right)-\int_{\mathbb{R}^{n}} f\left(u^{1}(z)\right)+o(1) .
\end{aligned}
$$

In the same way, because $\psi_{n}$ is symmetric,

$$
\begin{equation*}
\left[\psi_{n}\left(\tau z+\tau y_{n}\right)-u^{1}(z) \chi\left(\frac{|z|}{\rho_{n}}\right)\right] \chi\left(\frac{|z|}{\rho_{n}}\right) \rightharpoonup 0 \text { in } \mathcal{D}^{1,2}, \tag{34}
\end{equation*}
$$

and

$$
I_{2}=\int_{\left|\tau x-y_{n}\right|<2 \rho_{n}} f\left(\psi_{n}(x)\right)-\int_{\mathbb{R}^{n}} f\left(u^{1}(x)\right)+o(1) .
$$

At last we have

$$
\begin{align*}
\int f\left(\widetilde{\psi}_{n}\right)= & \int_{\left|x-y_{n}\right|<2 \rho_{n}} f\left(\psi_{n}(x)\right)+\int_{\left|\tau x-y_{n}\right|<2 \rho_{n}} f\left(\psi_{n}(x)\right)+ \\
& +\int_{\left\{B\left(y_{n}\right)_{2 \rho_{n}} \cup B\left(\tau y_{n}\right)_{2 \rho_{n}}\right\}^{C}} f\left(\psi_{n}(x)\right)-2 \int_{\mathbb{R}^{n}} f\left(u^{1}(x)\right)+o(1)=  \tag{35}\\
= & \int_{\mathbb{R}^{N}} f\left(\psi_{n}(x)\right)-2 \int_{\mathbb{R}^{n}} f\left(u^{1}(x)\right)+o(1) .
\end{align*}
$$

From (30) and (35) we obtain, as claimed

$$
\begin{equation*}
E_{0}\left(\tilde{\psi}_{n}\right)=E_{0}\left(\psi_{n}-\gamma_{n}\right)=E_{0}\left(\psi_{n}\right)-2 E_{0}\left(u^{1}\right)+o(1) \tag{36}
\end{equation*}
$$

furthermore, because $\left\{\psi_{n}\right\}_{n}$ is a $P S$ sequence for $E_{0}$, we have that $E_{0}\left(\widetilde{\psi}_{n}\right) \rightarrow c$ for some $c \in \mathbb{R}$.

To complete the proof we must show that

$$
\begin{equation*}
\left|\left(E_{0}\right)^{\prime}\left(\widetilde{\psi}_{n}\right) \varphi\right| \leq \varepsilon_{n}\|\varphi\|_{\mathcal{D}^{1,2}}, \tag{37}
\end{equation*}
$$

where $\varepsilon_{n} \rightarrow 0$. Set

$$
\begin{aligned}
I_{n}^{1}= & \int_{\left|x-y_{n}\right|<2 \rho}\left[f^{\prime}\left(\psi_{n}(x)\right)-f^{\prime}\left(\psi_{n}(x)-u^{1}\left(x-y_{n}\right) \chi\left(\frac{\left|x-y_{n}\right|}{\rho_{n}}\right)\right)\right] \varphi(x) d x- \\
& -\int \nabla\left(u^{1}\left(x-y_{n}\right) \chi\left(\frac{\left|x-y_{n}\right|}{\rho_{n}}\right)\right) \nabla \varphi(x) ; \\
I_{n}^{2}= & \int_{\left|\tau x-y_{n}\right|<2 \rho}\left[f^{\prime}\left(\psi_{n}(x)\right)-f^{\prime}\left(\psi_{n}(x)-u^{1}\left(\tau x-y_{n}\right) \chi\left(\frac{\left|\tau x-y_{n}\right|}{\rho_{n}}\right)\right)\right] \varphi(x) d x- \\
& -\int \nabla\left(u^{1}\left(\tau x-y_{n}\right) \chi\left(\frac{\left|\tau x-y_{n}\right|}{\rho_{n}}\right)\right) \nabla \varphi(x) ; \\
I_{n}^{3}= & \int_{\left\{B\left(y_{n}\right)_{2 \rho_{n}} \cup B\left(\tau y_{n}\right)_{\left.2 \rho_{n}\right\}^{C}}\right.}\left[\nabla \psi_{n}-f^{\prime}\left(\psi_{n}\right)\right] \varphi,
\end{aligned}
$$

we have that

$$
\begin{equation*}
\left|\left(E_{0}\right)^{\prime}\left(\widetilde{\psi}_{n}\right) \varphi\right|=I_{n}^{1}+I_{n}^{2}+I_{n}^{3} . \tag{38}
\end{equation*}
$$

Immediately we have that $I_{n}^{3} \leq \varepsilon_{n}\|\varphi\|$; furthermore, we can estimate $I_{n}^{1}$ as before, obtaining

$$
\begin{aligned}
I_{n}^{1}= & \int_{\left|z<2 \rho_{n}\right|}\left[f^{\prime}\left(\psi_{n}\left(z+y_{n}\right)\right)-f^{\prime}\left(\psi_{n}\left(z+y_{n}\right)-u^{1}(z) \chi\left(\frac{|z|}{\rho_{n}}\right)\right)\right] \varphi\left(z+y_{n}\right) d z- \\
& -\int_{\left|z<2 \rho_{n}\right|} f^{\prime}\left(u^{1} \chi\right) \varphi\left(z+y_{n}\right) d z+\varepsilon_{n}\|\varphi\| .
\end{aligned}
$$

Setting

$$
\begin{equation*}
\alpha_{n}(z):=\psi_{n}\left(z+y_{n}\right)-u^{1}(z) \chi\left(\frac{|z|}{\rho_{n}}\right), \tag{39}
\end{equation*}
$$

we have

$$
\left|I_{n}^{1}\right|=\left|\int_{B \rho_{n}}\left[f^{\prime}\left(u^{1} \chi+\alpha_{n}\right)-f^{\prime}\left(\alpha_{n}\right)-f^{\prime}\left(u^{1} \chi\right)\right] \varphi_{y_{n}} d z+\varepsilon_{n}\|\varphi\|\right|,
$$

and, chosen an $R>0$,

$$
\begin{aligned}
\left|I_{n}^{1}\right|= & \mid \int_{B_{\rho_{n}}}\left[f^{\prime}\left(u^{1} \chi+\alpha_{n}\right)-f^{\prime}\left(u^{1} \chi\right)\right] \varphi_{y_{n}} d z+ \\
& +\int_{\left(\mathbb{R}^{N} \backslash B_{R}\right) \cap B_{2 \rho_{n}}}\left[f^{\prime}\left(u^{1} \chi+\alpha_{n}\right)-f^{\prime}\left(\alpha_{n}\right)\right] \varphi_{y_{n}} d z- \\
& -\int_{\left(\mathbb{R}^{N} \backslash B_{R}\right) \cap B_{2 \rho_{n}}} f^{\prime}\left(u^{1} \chi\right) \varphi_{y_{n}} d z-\int_{B_{\rho_{n}}} f^{\prime}\left(\alpha_{n}\right) \varphi_{y_{n}} d z+\varepsilon_{n}\|\varphi\| \mid .
\end{aligned}
$$

Using that $f^{\prime}(0)=0$ at last we have

$$
\begin{aligned}
& \left|I_{n}^{1}\right| \leq\left\|f^{\prime \prime}\left(u^{1} \chi+\theta_{1} \alpha_{n}\right)-f^{\prime \prime}\left(\theta_{1} \alpha_{n}\right)\right\|_{L^{p / p-2}\left(\mathbb{R}^{N}\right)}\|\varphi\|_{L^{p}\left(\mathbb{R}^{N}\right)}\left\|\alpha_{n}\right\|_{L^{p}\left(B_{R}\right)}+ \\
& +\left\lvert\,\left\|f^{\prime \prime}\left(\alpha_{n}+\theta u^{1} \chi\right)-f^{\prime \prime}\left(\theta_{1} u^{1} \chi\right)\right\|_{L^{p-2} \cap L^{\frac{p}{q-2}}\left(\mathbb{R}^{N}\right)}\|\varphi\|_{L^{p}+L^{q}}\left\|u^{1} \chi\right\|_{L^{p}+L^{q}\left(B_{R}^{C}\right)}+\right. \\
& +\varepsilon_{n}\|\varphi\|,
\end{aligned}
$$

where $0<\theta, \theta_{1}<1$. By Remark 4 we get

$$
\begin{equation*}
\left|I_{n}^{1}\right| \leq \varepsilon_{n}\|\varphi\|_{\mathcal{D}^{1,2}} . \tag{40}
\end{equation*}
$$

In the same way we can estimate $\left|I_{n}^{2}\right|$, and this concludes the proof.

Remark 9. If $u_{n} \in \mathcal{N}_{V}^{\tau}$ is a Palais-Smale sequence of the restriction of $E_{V}$ to $\mathcal{N}_{V}^{\tau}$, that is $E_{V}\left(u_{n}\right)$ converges and

$$
\left|E_{V}^{\prime}\left(u_{n}\right) w\right| \leq \varepsilon_{n}\|w\|_{\mathcal{D}^{1,2}} \quad \forall w \in T_{u_{n}} \mathcal{N}_{V} \cap \mathcal{D}_{\tau}^{1,2}
$$

where $\varepsilon_{n} \rightarrow 0$, then $u_{n}$ is a Palais-Smale sequence for the functional $E_{V}$.
This remark, combined with the splitting lemma, provides a complete description of the $P S$ sequences in our case.

## 4 The main result

At this point we prove some technical lemmas.
Let $u \in \mathcal{N}_{V}$, then $u^{+}$and $u^{-}$belong to $\mathcal{N}_{V}$. Furthermore, if $u$ is antisymmetric, we have $E_{V}\left(u^{+}\right)=E_{V}\left(u^{-}\right)$. So, if $u \in \mathcal{N}_{V}^{\tau}$, we get

$$
E_{V}(u)=E_{V}\left(u^{+}\right)+E_{V}\left(u^{-}\right)=2 E_{V}\left(u^{+}\right) \geq 2 \inf _{\mathcal{N}_{V}} E_{V}=2 \mu_{V}
$$

This implies that

$$
\begin{align*}
\mu_{V}^{\tau} & \geq 2 \mu_{V}  \tag{41}\\
\mu_{0}^{\tau} & \geq 2 \mu_{0} \tag{42}
\end{align*}
$$

Remark 10. We have $\mu_{0}^{\tau}=2 \mu_{0}$
Proof. We have to proof that $\mu_{0}^{\tau} \leq 2 \mu_{0}$. It is possible to find a sequence $\left\{u_{k}\right\}_{k} \subset \mathcal{N}_{0}^{\tau}$ such that $E_{0}\left(u_{k}\right) \rightarrow 2 \mu_{0}$. So

$$
\mu_{0}^{\tau} \leq \inf _{k} E_{0}\left(u_{k}\right) \leq 2 \mu_{0} .
$$

The construction of $\left\{u_{k}\right\}_{k}$ is quite similar to the construction of $\left\{z_{k}\right\}_{k}$ in the next theorem. So also the proof that $E_{0}\left(u_{k}\right) \rightarrow 2 \mu_{0}$. Therefore, for the sake of simplicity, we omit the detailed proof of this result.

We are ready now to prove the main lemma of this section
Lemma 11. We have that $\mu_{V}^{\tau} \leq \mu_{0}^{\tau}$
Proof. We prove it by steps
Step I We know that $w$ exists such that $\mu_{0}=E_{0}(w)$. Let $\chi(x)$ a smooth, real function such that

$$
\chi= \begin{cases}1 & B(0,1) ; \\ 0 & \mathbb{R}^{N} \backslash B(0,3) .\end{cases}
$$

We also ask that $\chi(x)=\chi(|x|)$ and that $|\nabla \chi| \leq 1$.
Let $\left\{y_{k}\right\} \subset \mathbb{R}^{N}$ s.t. $\left|y_{k}\right| \rightarrow \infty$ and $\left|\tau\left(y_{k}\right)-y_{k}\right| \rightarrow \infty$. Let $\rho_{k}$ be defined as

$$
\rho_{k}:=\frac{\left|\tau\left(y_{k}\right)-y_{k}\right|}{6} .
$$

At last we define a function in $\mathcal{D}^{1,2}$

$$
\begin{equation*}
z_{k}=z_{k}^{1}+z_{k}^{2}, \tag{43}
\end{equation*}
$$

where

$$
\begin{align*}
z_{k}^{1} & =w\left(x-y_{k}\right) \chi\left(\frac{x-y_{k}}{\rho_{k}}\right)  \tag{44}\\
z_{k}^{2} & =-w\left(\tau(x)-y_{k}\right) \chi\left(\frac{x-\tau\left(y_{k}\right)}{\rho_{k}}\right) . \tag{45}
\end{align*}
$$

Obviously we have that $\tau z_{k}^{1}=z_{k}^{2}$ and $\tau z_{k}^{2}=z_{k}^{1}$, so $z_{k} \in \mathcal{D}_{\tau}^{1,2} \forall k$. Furthermore $z_{k}^{1}$ and $z_{k}^{2}$ have disjoint supports, so

$$
\begin{equation*}
E_{V}\left(t \cdot z_{k}\right)=E_{V}\left(t \cdot z_{k}^{1}\right)+E_{V}\left(t \cdot z_{k}^{2}\right) \forall t>0 . \tag{46}
\end{equation*}
$$

We know, from Remark 7, that it exists a $t_{k}>0$ s.t. $t_{k} \cdot z_{k}^{1} \in \mathcal{N}_{V}$. It's easy to see that, for such $t_{k}$ we have that $t_{k} \cdot z_{k}^{2} \in \mathcal{N}_{V}$ and $t_{k} \cdot z_{k} \in \mathcal{N}_{V}^{\tau}$.

In the next we will prove that $E_{V}\left(t_{k} z_{k}\right) \rightarrow 2 \mu_{0}$, when $k \rightarrow \infty$.
Step II We prove that $\left\|z_{k}^{1}(x)-w\left(x-y_{k}\right)\right\|_{\mathcal{D}^{1,2}} \rightarrow 0$ for $k \rightarrow \infty$.
Set $w_{k}:=w\left(x-y_{k}\right)$, and $\gamma_{k}:=\left(1-\chi\left(\frac{x-y_{k}}{\rho_{k}}\right)\right)$, we have

$$
\begin{aligned}
\left\|z_{k}^{1}(x)-w_{k}\right\|_{\mathcal{D}^{1,2}}^{2} & =\int\left|\nabla\left[\gamma_{k} w_{k}\right]\right|^{2} \leq 2 \int \gamma_{k}^{2}\left|\nabla w_{k}\right|^{2}+2 \int\left|\nabla \gamma_{k}\right|^{2} w_{k}^{2} \leq \\
& \leq \int_{\left|x-y_{k}\right|>3 \rho_{k}}\left|\nabla w_{k}\right|^{2}+\frac{2}{\rho_{k}^{2}} \int_{\rho_{k}<\left|x-y_{k}\right|<3 \rho_{k}} w_{k}^{2} \rightarrow 0,
\end{aligned}
$$

as $k \rightarrow \infty$.
Step III We prove that it exists $c, C>0$ such that $c<t_{k}<C$ for all $k$.
By Remark 6, we know that, if $t_{k} z_{k}^{1} \in \mathcal{N}_{V}$, then it exists $M>0$ such that, for all $k, M \leq\left\|t_{k} z_{k}^{1}\right\|$. Furthermore, for the above step

$$
\begin{equation*}
\left\|z_{k}^{1}\right\| \rightarrow\left\|w_{k}\right\|=\|w\| . \tag{47}
\end{equation*}
$$

This implies that $c$ exists such that $t_{k}>c>0$ for all $k$.

For the other inequality we must prove that

$$
\begin{equation*}
\int V(x) w_{k}^{2}(x) d x \rightarrow 0 \text { when } k \rightarrow \infty \tag{48}
\end{equation*}
$$

in fact, fixed an $R>0$ we have

$$
\begin{aligned}
\left|\int V w_{k}^{2}\right| & \leq \int_{B_{R}}|V| w_{k}^{2}+\int_{\mathbb{R}^{N} \backslash B_{R}}|V| w_{k}^{2} \leq \\
& \leq\|V\|_{L^{\frac{N}{2}}\left(B_{R}\right)}\left\|w_{k}\right\|_{L^{2^{*}}\left(B_{R}\right)}^{2}+\|V\|_{L^{\frac{N}{2}}\left(\mathbb{R}^{N} \backslash B_{R}\right)}\left\|w_{k}\right\|_{L^{2^{*}}\left(\mathbb{R}^{N}\right)}^{2} .
\end{aligned}
$$

We have that $\|V\|_{L^{\frac{N}{2}}\left(\mathbb{R}^{N} \backslash B_{R}\right)} \rightarrow 0$ as $R \rightarrow \infty$; furthermore

$$
\begin{equation*}
\int_{B_{R}} w^{2^{*}}\left(x-y_{k}\right) d x=\int_{B_{R}\left(-y_{k}\right)} w^{2^{*}}(s) d s \leq \int_{\mathbb{R}^{N} \backslash B_{\left(\left|y_{k}\right|-R\right)}} w^{2^{*}}(s) d s, \tag{49}
\end{equation*}
$$

thus $\left\|w_{k}\right\|_{L^{2^{*}}}^{2} \rightarrow 0$ as $k \rightarrow \infty$, that proves (48).
Now set a function

$$
\begin{equation*}
g_{z_{k}^{1}}^{V}(t):=E_{V}\left(t z_{k}^{1}\right) . \tag{50}
\end{equation*}
$$

Obviously $g_{w}^{0}(t)=E_{0}(t w)=\frac{1}{2} t^{2} \int|\nabla w|^{2}-\int f(t w)$. For Remark 7 we know that there exists a $\bar{t}$ such that $g_{w_{k}}^{0}(\bar{t})=g_{w}^{0}(\bar{t})<0$ for all $k$. We want to prove that, for $k$ sufficiently big, we have also $g_{z_{k}^{1}}^{V}(\bar{t})<0$.

$$
\begin{aligned}
g_{z_{k}^{1}}^{V}(\bar{t})-g_{w}^{0}(\bar{t})= & g_{z_{k}^{1}}^{V}(\bar{t})-g_{w_{k}}^{0}(\bar{t})=\frac{1}{2} \int\left|\nabla \bar{t} z_{k}^{1}\right|^{2}+\frac{1}{2} \int V(x) \bar{t}^{2}\left(z_{k}^{1}\right)^{2}- \\
& -\int f\left(\bar{t} z_{k}^{1}\right)-\frac{1}{2} \int\left|\nabla \bar{t} w_{k}\right|^{2}+\int f\left(\bar{t} w_{k}\right)= \\
= & \frac{\bar{t}^{2}}{2}\left[\int\left|\nabla z_{k}^{1}\right|^{2}-\left|\nabla w_{k}\right|^{2}+V(x)\left(z_{k}^{1}\right)^{2}\right]-\int f\left(\bar{t} z_{k}^{1}\right)-f\left(\bar{t} w_{k}\right) .
\end{aligned}
$$

By (47) and by (48) we have that the first integral of the right hand side of the equation vanishes when $k \rightarrow \infty$. We estimate the last term.

$$
\begin{aligned}
\int f\left(\bar{t} z_{k}^{1}\right)-f\left(\bar{t} w_{k}\right) & =\int_{\left|x-y_{k}\right|>\rho_{k}} f\left(\bar{t} w_{k} \chi\left(\frac{x-y_{k}}{\rho_{k}}\right)\right)-f\left(\bar{t} w_{k}\right)= \\
& =\int_{|s|>\rho_{k}} f\left(\bar{t} w \chi\left(\frac{s}{\rho_{k}}\right)\right)-f(\bar{t} w)= \\
& =\int_{|s|>\rho_{k}} f^{\prime}\left(\left[\theta \chi\left(\frac{s}{\rho_{k}}\right)+(1-\theta)\right] \bar{t} w\right)\left(\chi\left(\frac{s}{\rho_{k}}\right)-1\right) \bar{t} w= \\
& \leq \bar{t} \int_{|s|>\rho_{k}} w f^{\prime}\left(\left[\theta \chi\left(\frac{s}{\rho_{k}}\right)+(1-\theta)\right] \bar{t} w\right)\left(\chi\left(\frac{s}{\rho_{k}}\right)-1\right),
\end{aligned}
$$

$$
\left|\int f\left(\bar{t} z_{k}^{1}\right)-f\left(\bar{t} w_{k}\right)\right| \leq \bar{t} \int_{|s|>\rho_{k}}\left|f^{\prime}\left(\left[\theta \chi\left(\frac{s}{\rho_{k}}\right)+(1-\theta)\right] \bar{t} w\right)\right|\left|\chi\left(\frac{s}{\rho_{k}}\right)-1\right| w
$$

and the last term vanishes when $k \rightarrow \infty$. In fact, for Remark 4, $f^{\prime}$ is bounded in $L^{\frac{p}{p-1}} \cap L^{\frac{q}{q-1}}$ and $\left|\chi\left(\frac{s}{\rho_{k}}\right)-1\right| w \rightarrow 0$ in $L^{2^{*}}$. So, for $k_{0}$ big enough, we have that

$$
\exists c, C \in \mathbb{R}^{+} \text {s.t. } 0<c<t_{k}<C, \forall k>k_{0}
$$

Step IV We want to prove that $E_{V}\left(t_{k} z_{k}^{1}\right) \rightarrow \mu_{0}$. We have

$$
\left|E_{V}\left(t_{k} z_{k}^{1}\right)-E_{V}\left(w_{k}\right)\right|=\left|E_{V}^{\prime}\left(\theta t_{k} z_{k}^{1}+(1-\theta) w_{k}\right)\left(z_{k}^{1}-w_{k}\right)\right|
$$

We know, for Step 2, that $\left\|z_{k}^{1}-w_{k}\right\|_{\mathcal{D}^{1,2}} \rightarrow 0$.
Furthermore

$$
\begin{equation*}
\left\|\theta t_{k} z_{k}^{1}+(1-\theta) w_{k}\right\| \leq\left\|z_{k}^{1}\right\| t_{k}+\left\|w_{k}\right\| \tag{51}
\end{equation*}
$$

that is bounded because $t_{k}$ is bounded and by Step II.
At this point by Remark 4 we get the claim.
We know also that

$$
\begin{aligned}
E_{V}\left(w_{k}\right)-\mu_{0} & =E_{V}\left(w_{k}\right)-E_{0}(w)=E_{V}\left(w_{k}\right)-E_{0}\left(w_{k}\right)= \\
& =\int V(x) w_{k}^{2} \xrightarrow{k \rightarrow \infty} 0
\end{aligned}
$$

for (48). Then

$$
\begin{equation*}
\left|E_{V}\left(t_{k} z_{k}^{1}\right)-\mu_{0}\right| \leq\left|E_{V}\left(t_{k} z_{k}^{1}\right)-E_{V}\left(w_{k}\right)\right|+\left|E_{V}\left(w_{k}\right)-\mu_{0}\right| \rightarrow 0 \tag{52}
\end{equation*}
$$

as we wanted to prove.
Conclusion We know that $t_{k} z_{k} \in \mathcal{N}_{V}^{\tau}$. Then

$$
E_{V}\left(t_{k} z_{k}\right) \geq \mu_{V}^{\tau}:=\inf _{u \in \mathcal{N}_{V}^{\tau}} E_{V}(u)
$$

Hence

$$
\mu_{V}^{\tau} \leq E_{V}\left(t_{k} z_{k}\right)=E_{V}\left(t_{k} z_{k}^{1}\right)+E_{V}\left(t_{k} z_{k}^{2}\right) \rightarrow 2 \mu_{0}=\mu_{0}^{\tau}
$$

that gives us the proof.
We are ready, now, to prove the first result claimed in the introduction.

Proof of Theorem 1. First, we prove that

$$
\begin{equation*}
\forall u \in \mathcal{N}_{0} \exists t_{u}^{V} u \in \mathcal{N}_{V} \text { s.t } E_{V}\left(t_{u}^{V} u\right) \geq E_{0}(u) . \tag{53}
\end{equation*}
$$

In fact, by Remark 7, we have that for every $u \in \mathcal{N}_{0}$, there exist $t_{u}^{V}>0$ such that $t_{u}^{V} u \in \mathcal{N}_{V}$. Then we have:

$$
0=g_{u}^{\prime}\left(t_{u}^{V} u\right)=\left\langle\nabla E_{0}\left(t_{u}^{V} u\right), u\right\rangle+t_{u}^{V} \int_{\mathbb{R}^{N}} V u^{2}
$$

Since $V>0$ we have that $\int V u^{2}>0$ and $\left\langle\nabla E_{0}\left(t_{u}^{V} u\right), u\right\rangle<0$. Hence $t_{u}^{V}>$ $t_{u}^{0}=1$. Let us observe that by (22) the function $s \rightarrow \int \frac{1}{2} f^{\prime}(s u) s u-f(s u) d x$ is strictly increasing, then, remembering that $t_{u}^{V} u \in \mathcal{N}_{V}$, we have:

$$
\begin{aligned}
E_{V}\left(t_{u}^{V} u\right) & =\frac{1}{2} \int f^{\prime}\left(t_{u}^{V} u\right) t_{u}^{V} u-f\left(t_{u}^{V} u\right) d x \geq \\
& \geq \frac{1}{2} \int f^{\prime}(u) u-f(u) d x=E_{0}(u)
\end{aligned}
$$

If $u \in \mathcal{N}_{0}^{\tau}$, we can prove in the same way that $t_{u}^{V} u \in \mathcal{N}_{V}^{\tau}$ and that

$$
\begin{equation*}
E_{V}\left(t_{u}^{V} u\right) \geq E_{0}(u) \geq \mu_{0}^{\tau} \tag{54}
\end{equation*}
$$

So

$$
\begin{equation*}
\inf _{w \in \mathcal{N}_{V}^{\tau}} E_{V}(w)=\inf _{u \in \mathcal{N}_{0}^{\tau}} E_{V}\left(t_{u}^{V} u\right) \geq E_{0}(u) \geq \mu_{0}^{\tau} . \tag{55}
\end{equation*}
$$

Theorem 11 provides us the other inequality.
Suppose now that there exists $v \in \mathcal{N}_{V}^{\tau}$ such that $\mu_{V}^{\tau}=E_{V}(v)$. We know that $\int V(x) v(x)^{2}>0$ and

$$
0=\left\langle\nabla E_{0}(v), v\right\rangle+\int V(x) v(x)^{2} d x
$$

so, consequently $\left\langle\nabla E_{0}(v), v\right\rangle<0$. Then, by Remark 7 , we get $t_{v}^{0}<t_{v}^{V}=1$. As said before, the function $s \mapsto \int \frac{1}{2} f^{\prime}(s v) s v-f(s v) d x$ is strictly increasing, so we have

$$
E_{0}\left(v t_{v}^{0}\right)=\int \frac{1}{2} f^{\prime}\left(t_{v}^{0}\right) t_{v}^{0} v-f\left(t_{v}^{0} v\right) d x<\int \frac{1}{2} f^{\prime}(v) v-f(v) d x=E_{V}(v)=\mu_{V}^{\tau}
$$

and we get a contradiction.

Now we prove the following preliminary result.

Proposition 12. There exists a class of potential $V(x)$ such that $\mu_{V}^{\tau}<\mu_{0}^{\tau}$.
Proof. We consider the class of potentials defined in (10). We want to show that, when $|y| \rightarrow \infty$ and $|y-\tau y| \rightarrow \infty$, then $\mu_{V_{y}}^{\tau}<\mu_{0}^{\tau}$. We prove it by steps.

Take $w \in \mathcal{N}_{0}$ such that $E_{0}(w)=\mu_{0}, w$ radially symmetric and $w>0$ (see [10, 11, 12]). By means of $w$, we define

$$
\begin{equation*}
z_{y}(x)=w(x-y)-w(x-\tau y) . \tag{56}
\end{equation*}
$$

Step I We prove that, for $|y-\tau y| \rightarrow \infty$,

$$
g_{z_{y}}^{V_{y}}(t) \rightarrow t^{2} \int|\nabla w(x)|^{2} d x+a t^{2} \int[|x|-1] w^{2}(x) d x-2 \int f(t w(x)) d x
$$

where

$$
\begin{equation*}
g_{z_{y}}^{V_{y}}(t)=E_{V_{y}}\left(t z_{y}\right)=\int_{\mathbb{R}^{N}} \frac{t^{2}}{2}\left(\left|\nabla z_{y}\right|^{2}+V_{y} z_{y}^{2}\right)-f\left(t z_{y}\right) d x \tag{57}
\end{equation*}
$$

Now

$$
\begin{equation*}
\frac{t^{2}}{2} \int\left|\nabla z_{y}\right|^{2}=\frac{t^{2}}{2} \int|\nabla w(x-y)|^{2}+|\nabla w(\tau x-y)|^{2}+\nabla w(x-y) \nabla w(x-\tau y) \tag{58}
\end{equation*}
$$

After a change of variables, the first two terms are equals to $\frac{t^{2}}{2} \int|\nabla w|^{2}$, and the last term vanishes. So we have that, for all $t$,

$$
\begin{equation*}
\frac{t^{2}}{2} \int\left|\nabla z_{y}\right|^{2} \rightarrow t^{2} \int|\nabla w|^{2} \quad \text { when }|y-\tau y| \rightarrow \infty \tag{59}
\end{equation*}
$$

In a similar way consider

$$
\begin{align*}
\int V_{y} z_{y}^{2}= & a \int_{|x-y|<1}(|x-y|-1) z_{y}^{2}+a \int_{|x-\tau y|<1}(|x-\tau y|-1) z_{y}^{2}= \\
= & a \int_{|x-y|<1}(|x-y|-1)[w(x-y)-w(x-\tau y)]^{2}+  \tag{60}\\
& +a \int_{|x-\tau y|<1}(|x-\tau y|-1)[w(x-y)-w(x-\tau y)]^{2} .
\end{align*}
$$

By means of a change of variables we obtain

$$
\int_{|x-y|<1}(|x-y|-1) z_{y}^{2}=\int_{|s|<1}(|s|-1)[w(s)+w(s+y-\tau y)]^{2}
$$

It is not difficult to prove that

$$
\begin{align*}
\int_{|s|<1}(|s|-1) w^{2}(s+y-\tau y) & \rightarrow 0  \tag{61}\\
\int_{|s|<1}(|s|-1) w(s) w(s+y-\tau y) & \rightarrow 0 . \tag{62}
\end{align*}
$$

In the same way we proceed for the second term of the (60), obtaining

$$
\begin{equation*}
\frac{t^{2}}{2} \int V_{v} z_{y}^{2} \rightarrow a t^{2} \int_{|x|<1}(|x|-1) w^{2}(x) d x \text { when }|y| \rightarrow \infty \tag{63}
\end{equation*}
$$

We have to estimate now $\int f(t w)$. Fixed an $R>0$, we have

$$
\begin{equation*}
\int f\left(t z_{y}\right)=\int_{B_{R}(y)} f\left(t z_{y}\right)+\int_{B_{R}(\tau y)} f\left(t z_{y}\right)+\int_{\left(B_{R}(y) \cup B_{R}(\tau y)\right)^{C}} f\left(t z_{y}\right) . \tag{64}
\end{equation*}
$$

For the first term we have

$$
\begin{aligned}
\int_{B_{R}(y)} f\left(t z_{y}\right)= & \int_{B_{R}} f(t w(s)-t w(s+y-\tau y))=\int_{B_{R}} f(t w)+ \\
& +\int_{B_{R}} f^{\prime}(\theta t w(s)+(1-\theta) t w(s+y-\tau y))[t w(s+y-\tau y)]
\end{aligned}
$$

for some $\theta \in[0,1]$.
Now, for Remark (4, we have that $f^{\prime}(\theta t w(x-y)+(1-\theta) \operatorname{tw}(\tau x-y))$ is bounded in $L^{p^{\prime}} \cap L^{q^{\prime}}$, in fact $\theta t w(x-y)+(1-\theta) t w(\tau x-y)$ is bounded in $\mathcal{D}^{1,2}$ and so in $L^{p}+L^{q}$. Furthermore, $w(s+y-\tau y) \rightarrow 0$ strongly in $L^{P}\left(B_{R}\right)$ when $|y-\tau y| \rightarrow+\infty$.

Concluding we get

$$
\begin{equation*}
\int_{B_{R}(y)} f\left(t z_{y}\right)=\int_{B_{R}} f(t w(s)) d s+I_{1}(R, y), \tag{65}
\end{equation*}
$$

where, given $R>0, I_{1}(R, y) \rightarrow 0$ when $|y-\tau y| \rightarrow \infty$. In the same way we can conclude that

$$
\begin{equation*}
\int_{B_{R}(\tau y)} f\left(t z_{y}\right)=\int_{B_{R}} f(t w(s)) d s+I_{2}(R, y) \tag{66}
\end{equation*}
$$

where, again, given $R>0, I_{2}(R, y) \rightarrow 0$ when $|y-\tau y| \rightarrow \infty$.
For the last term we have that there exist a $\theta \in[0,1]$ such that

$$
\begin{aligned}
\int_{\left(B_{R}(y) \cup B_{R}(\tau y)\right)^{C}} f\left(t z_{y}\right)= & \int_{\left(B_{R}(y) \cup B_{R}(\tau y)\right)^{C}} f(t(w(x-y)))+ \\
& +\int_{\left(B_{R}(y) \cup B_{R}(\tau y)\right)^{C}} f^{\prime}(\theta t(w(x-y))-(1-\theta) w(x-\tau y)) t w(x-\tau y)= \\
= & \int_{\left(B_{R} \cup B_{R}(y-\tau y)\right)^{C}} f(t(w(s)))+ \\
& \quad+\int_{\left(B_{R} \cup B_{R}(\tau y-y)\right)^{C}} f^{\prime}(\theta t(w(\xi+\tau y-y))-(1-\theta) w(\xi)) t w(\xi) .
\end{aligned}
$$

Now,

$$
\begin{equation*}
\left|\int_{\left(B_{R} \cup B_{R}(\tau y-y)\right)^{C}} f^{\prime}(\cdot) t w(\xi)\right| \leq t| | f^{\prime}(\cdot)\left\|_{L^{p^{\prime} \cap L^{q^{\prime}}\left(\mathbb{R}^{N}\right)}}\right\| w \|_{L^{p}+L^{q}\left(\mathbb{R}^{N} \backslash B_{R}\right)}, \tag{67}
\end{equation*}
$$

and we use that $\|w\|_{\mathcal{D}^{1,2}\left(\mathbb{R}^{N} \backslash B_{R}\right)}$ goes to zero when $R \rightarrow \infty$ and that $\left\|f^{\prime}(\cdot)\right\|_{L^{p^{\prime}} \cap L^{q^{\prime}}\left(\mathbb{R}^{N}\right)}$ is bounded by Remark 4 .

At this point we have that

$$
\begin{equation*}
\int f\left(t z_{y}\right) \rightarrow 2 \int f(w) \quad \text { when }|y-\tau y| \rightarrow \infty \tag{68}
\end{equation*}
$$

and we get the claim.
Step II There exists a $\bar{t}<1$ such that

$$
\begin{equation*}
t_{z_{y}}^{V_{y}} \rightarrow \bar{t} \quad \text { when }|y-\tau y| \rightarrow \infty, \tag{69}
\end{equation*}
$$

where $t_{z_{y}}^{V_{y}}$ is the maximum point of $g_{z y}^{V_{y}}(t)$.
We set

$$
\begin{equation*}
\varphi(t)=t^{2} \int|\nabla w(x)|^{2} d x+a t^{2} \int[|x|-1] w^{2}(x) d x-2 \int f(t w(x)) d x \tag{70}
\end{equation*}
$$

By Remark 7 there exists a unique maximizer $\bar{t}>0$ for the function $\varphi(t)$.
We know that

$$
\begin{equation*}
g_{w}^{0}(t)=\frac{1}{2} t^{2} \int \nabla w^{2}-\int f(t w) \tag{71}
\end{equation*}
$$

reach its maximum for $t=1$. Thus the maximum of the function

$$
\begin{equation*}
\varphi(t)=2 g_{w}^{0}(t)+a t^{2} \int[|x|-1] w^{2}(x) d x \tag{72}
\end{equation*}
$$

is achieved for $\bar{t}$, with $0<\bar{t}<1$.
Given $t_{1}<\bar{t}<t_{2}$, we can choose a $\delta>0$ such that

$$
\begin{equation*}
\max \left\{\varphi\left(t_{1}\right), \varphi\left(t_{2}\right)\right\}+\delta<\varphi(\bar{t})-\delta \tag{73}
\end{equation*}
$$

By Step I, for $|y-\tau y|$ sufficiently large, we obtain

$$
\begin{equation*}
g_{z_{y}}^{V_{y}}\left(t_{i}\right)<\varphi\left(t_{i}\right)+\delta<\varphi(\bar{t})-\delta<g_{z_{y}}^{V_{y}}(\bar{t}) . \tag{74}
\end{equation*}
$$

By Remark 7, we know that $g_{z_{y}}^{V_{y}}(t)$ has an unique maximum point $t_{z_{y}}^{V_{y}}$, thus we conclude that

$$
\begin{equation*}
t_{1}<t_{z_{y}}^{V_{y}}<t_{2} . \tag{75}
\end{equation*}
$$

Since $t_{1}$ and $t_{2}$ are arbitrarily chosen, we get the claim.
Step III For $|y-\tau y|$ sufficiently large we have

$$
\begin{equation*}
\mu_{V_{y}}^{\tau}<\mu_{0}^{\tau} . \tag{76}
\end{equation*}
$$

We know that

$$
\begin{equation*}
E_{V_{y}}\left(t_{z_{y}}^{V_{y}} z_{y}\right)=g_{z_{y}}^{V_{y}}\left(t_{z_{y}}\right) \rightarrow \varphi(\bar{t}) \quad \text { for }|y-\tau y| \rightarrow \infty \tag{77}
\end{equation*}
$$

in fact, for all $\varepsilon>0$ we have that, for $|y-\tau y|$ sufficiently large,

$$
\begin{aligned}
\left|g_{z_{y}}^{V_{y}}\left(t_{z_{y}}\right)-\varphi(\bar{t})\right| & =\left|g_{z_{y}}^{V_{y}}\left(t_{z_{y}}\right)-g_{z_{y}}^{V_{y}}(\bar{t})+g_{z_{y}}^{V_{y}}(\bar{t})-\varphi(\bar{t})\right| \leq \\
& \leq\left|g_{z_{y}}^{V_{y}}\left(t_{z_{y}}\right)-g_{z_{y}}^{V_{y}}(\bar{t})\right|+\left|g_{z_{y}}^{V_{y}}(\bar{t})-\varphi(\bar{t})\right| \leq \varepsilon .
\end{aligned}
$$

By Step I the second term goes to zero when $|y-\tau y| \rightarrow \infty$. By Step II, $t_{z_{y}}^{V_{y}} \rightarrow \bar{t}$, so, arguing as in Step I, we get the claim. We observe that

$$
\begin{aligned}
\varphi(\bar{t}) & =\bar{t} \int|\nabla w|^{2}+a \bar{t}^{2} \int[|x|-1] w^{2}-2 \int f(\bar{t} w)< \\
& <2 E_{0}(\bar{t} w)<2 \mu_{0}
\end{aligned}
$$

because $\bar{t}<1$ and $E_{0}(w)=\mu_{0}$. By (77) we get

$$
\begin{equation*}
\mu_{V_{y}}^{\tau} \leq E_{V_{y}}\left(t_{z_{y}}^{V_{y}} z_{y}\right)<2 \mu_{0}=\mu_{0}^{\tau} \quad \text { for }|y-\tau y| \text { large enough, } \tag{78}
\end{equation*}
$$

that concludes the proof

Now we are ready to prove the second result claimed in the introduction.
Proof of theorem 园. By the Splitting Lemma and the above Proposition, we get the existence of a minimizer for $E_{V_{y}}$, for the class of potential $V_{y}$ defined by (10), when $|y-\tau y|$ large enough.

Let $\omega$ be this minimizer. We know that $\omega$ changes sign, because it is antisymmetric by construction. We have to prove that $\omega$ changes sign exactly once. Suppose that the set $\left\{x \in \mathbb{R}^{N}: \omega(x)>0\right\}$ has $k$ connected components $\Omega_{1}, \ldots, \Omega_{k}$. Set

$$
\omega_{i}= \begin{cases}\omega(x) & x \in \Omega_{i} \cup \tau \Omega_{i} ;  \tag{79}\\ 0 & \text { elsewhere }\end{cases}
$$

For all $i, \omega_{i} \in \mathcal{N}_{V_{y}}^{\tau}$. Furthermore we have

$$
\begin{equation*}
E_{V_{y}}(\omega)=\sum_{i} E_{V_{y}}\left(\omega_{i}\right), \tag{80}
\end{equation*}
$$

thus

$$
\begin{equation*}
\mu_{V_{y}}^{\tau}=E_{V_{y}}(\omega)=\sum_{i=1}^{k} E_{V_{y}}\left(\omega_{i}\right) \geq k \mu_{V_{y}}^{\tau}, \tag{81}
\end{equation*}
$$

so $k=1$, that concludes the proof.

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