

Regular g -measures are not always Gibbsian

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Abstract

Regular g -measures are discrete-time processes determined by conditional expectations with respect to the past. One-dimensional Gibbs measures, on the other hand, are fields determined by simultaneous conditioning on past and future. For the Markovian and exponentially continuous cases both theories are known to be equivalent. Its equivalence for more general cases was an open problem. We present a simple example settling this issue in a negative way: there exist g -measures that are continuous and non-null but are not Gibbsian. Our example belongs, in fact, to a well-studied family of processes with rather nice attributes: It is a chain with variable-length memory, characterized by the absence of phase coexistence and the existence of a visible renewal scheme.

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1 Introduction

Measures on $E^{\mathbb{Z}}$ are the object of two very developed theories. On the one hand, the theory of *chains of complete connections*, started in [16] and developed under a variety of names in slightly non-equivalent frameworks: *chains of infinite order* [10], *g -measures* [12], *uniform martingales* [11], etc. On the other hand, the theory of one-dimensional *Gibbs measures* started in [2, 15] and whose classical reference is the treatise [9]. The former theory interprets \mathbb{Z} as discrete time and measures as discrete-time processes. The building blocks of the theory are, therefore, *transition probabilities*, that is, conditional probabilities *with respect to the past*. Chains, g -measures, etc are defined by their invariance under—or consistency with—these transition probabilities. In contrast, Gibbsianness refers to fields in a spatial setting

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determined by distributions on finite-regions *conditioned on exterior configurations*. In one dimension, if \mathbb{Z} is interpreted as time, this corresponds to conditioning both *with respect to the past and to the future*. Of course, this is only part of the story. Both processes and Gibbs measures are required to satisfy suitable regularity conditions, as reviewed below.

Given this state of affairs it is natural to wonder about the relation of both theories. Are they equivalent? Does one-side conditioning carry the same information as two-side conditioning? Is every regular process a Gibbs measure and vice-versa? To be sure, there is another source of difference related to non-nullness. The standard theory of Gibbs measures deals with systems with no forbidden configurations. Many important instances of processes, on the other hand, include grammars or local exclusion rules (subshifts of finite type). The previous questions should be stated, then, in the common non-null framework. In this set-up, the equivalence of both theories has long been known to be true for Markov processes and fields (see, for instance, [9, Chapter 11]) and when continuity rates are exponentially decreasing [6]. In this note we exhibit a simple example —where all calculations can be explicitly performed— showing that this equivalence is *not* true in general.

In more detail, below we construct a g -measure μ on $\{0, 1\}^{\mathbb{Z}}$ with the following properties:

- μ is non-null: it gives nonzero measure to every cylinder.
- There exists a continuous g function for which μ is the *only* consistent measure.
- μ is left-right symmetric [as proven in the third line of (3.12)].
- μ is a variable-length memory chain which admits a renewal construction with visible renewals (see [8] for example).
- μ can be perfectly simulated (for example by the method of [8]).

Yet, despite all these fine properties, the measure μ is non-Gibbsian.

This example shows that regular g -measures are a different type of creature than Gibbsian measures and respective theories (large-deviations, uniqueness theorems, variational approach) can not, in general, be imported from one to the other. In particular, readers are warned that a g -function of the form $g = e^{\phi}$ with ϕ “nice” does not automatically deserve the qualifier “Gibbsian”. These observations complement previous studies on differences between one-sided and two-sided measurability done in the more general framework of ergodic theory (see, e.g. [4] and references therein).

Another potential source of confusion arises from the traditional use, by people working in dynamical systems, of the word “Gibbsian” to refer to SRB measures (Sinai-Ruelle-Bowen measures, see e.g. [17] or [1]). As briefly reviewed below, the set of non-null SRB measures is *strictly contained* in the set of statistical mechanical Gibbs measures (but the former, unlike the latter, can also incorporate exclusions and subshifts). Hence our g -measure μ is also non-Gibbsian in SRB sense.

Let us conclude with a brief explanation of the non-Gibbsianness argument below. The measure μ is non-null and consistent with a continuous g -measure. This means that, upon conditioning, it becomes asymptotically insensitive to the far past. In order to be Gibbs, the same asymptotic insensitivity must hold but simultaneously with respect to past *and future*. The measure μ is supported on configurations with an infinite number of 1’s. This is because the probability of having a 1 after an infinite sequence of 0’s is a strictly positive number

p_∞ . In addition, by continuity, the probability of having a 1 conditioned on a large string of zeros converges, as the first 1 recedes to $-\infty$, to p_∞ . For the measure to be Gibbsian this same continuity must hold for *two-side* conditioning. A delicate case arises, however, when conditioning on having the all-zero configuration both towards the past and the future. Of course, this is an impossible configuration for μ , so that the actual value of this conditional probability is irrelevant. Rather, the Gibbsianness question refers to whether conditional probabilities converge (to whatever) as both the first 1 to the left and the first 1 to the right move away. This is an essential property in the sense that its absence can not be fixed by measure-zero redefinitions. Theorem 3.1 shows that for some g -functions this two-side continuity is impossible. As an aside, we point out that the “all-0” configuration is the *only* point of discontinuity of the two-side conditional probabilities. The measure μ is, therefore, *almost Gibbsian* and thus *weakly Gibbsian* (see, for instance, [5, Section 4.4] for the corresponding definitions and historical references to these notions).

2 Preliminaries

We consider a measurable space (E, \mathcal{E}) where $E = \{0, 1\}$ is a two-symbol alphabet and \mathcal{E} is the associated discrete σ -algebra. We denote by (Ω, \mathcal{F}) the associated product measurable space, that is $\Omega = E^{\mathbb{Z}}$, $\mathcal{F} = \mathcal{E}^{\mathbb{Z}}$. For each $\Lambda \subset \mathbb{Z}$ we denote $\Omega_\Lambda = E^\Lambda$ and σ_Λ for the restriction of a configuration $\sigma \in \Omega$ to Ω_Λ , namely the family $(\sigma_i)_{i \in \Lambda} \in E^\Lambda$. Also, \mathcal{F}_Λ will denote the sub- σ -algebra of \mathcal{F} generated by cylinders based on Λ (\mathcal{F}_Λ -measurable functions are insensitive to configuration values outside Λ). When Λ is an interval, $\Lambda = [k, n]$ with $k, n \in \mathbb{Z}$ such that $k \leq n$, we use the notation: $\omega_k^n = \omega_{[k, n]} = \omega_k, \dots, \omega_n$, $\Omega_k^n = \Omega_{[k, n]}$ and $\mathcal{F}_k^n = \mathcal{F}_{[k, n]}$. For semi-intervals we denote also $\mathcal{F}_{\leq n} = \mathcal{F}_{(-\infty, n]}$, etc. The concatenation notation $\omega_\Lambda \sigma_\Delta$, where $\Lambda \cap \Delta = \emptyset$, indicates the configuration on $\Lambda \cup \Delta$ coinciding with ω_i for $i \in \Lambda$ and with σ_i for $i \in \Delta$.

We start by briefly reviewing the well-known notions of chains in a shift-invariant setting. In this particular case, chains are also called g -measures (see [12]).

Definition 2.1. A regular g -function P on Ω is a probability kernel $P: E \times \Omega_{-\infty}^{-1} \rightarrow [0, 1]$, i.e.,

$$\sum_{\omega_0 \in E} P(\omega_0 \mid \omega_{-\infty}^{-1}) = 1 \quad \forall \omega_{-\infty}^{-1} \in \Omega_{-\infty}^{-1}, \quad (2.1)$$

such that:

- the function $P(\omega_0 \mid \cdot)$ is **continuous** for each $\omega_0 \in E$, i.e., for all $\epsilon > 0$ there exists $n \geq 0$ so that

$$|P(\omega_0 \mid \omega_{-\infty}^{-1}) - P(\omega_0 \mid \sigma_{-\infty}^{-1})| < \epsilon \quad (2.2)$$

for all $\omega_{-\infty}^0, \sigma_{-\infty}^0 \in \Omega_{-\infty}^0$ with $\omega_{-n}^0 = \sigma_{-n}^0$;

- the function $P(\omega_0 \mid \cdot)$ is **strongly non-null** for each $\omega_0 \in E$, i.e., $P(\omega_0 \mid \cdot) \geq c > 0$.

[Property (2.2) is indeed continuity with respect to the product discrete topology of Ω .]

Definition 2.2. A probability measure μ on (Ω, \mathcal{F}) with underlying process $(X_i)_{i \in \mathbb{Z}}$ on $(\Omega, \mathcal{F}, \mu)$ is said to be a **regular g -measure** if μ is shift-invariant and there exists a regular g -function P such that μ is **consistent** with P , namely,

$$\mu(X_0 = \omega_0 \mid X_{-\infty}^{-1} = \omega_{-\infty}^{-1}) = P(\omega_0 \mid \omega_{-\infty}^{-1}) \quad (2.3)$$

for all $\omega_0 \in E$ and μ -a.e. $\omega_{-\infty}^{-1} \in \Omega_{-\infty}^{-1}$.

Remark 2.3. In the consistency definition (2.3), μ needs only to be defined on $(\Omega_{-\infty}^0, \mathcal{F}_{\leq 0})$. By shift-invariance, μ can be extended in a unique way to (Ω, \mathcal{F}) . Thus, without loss of generality, we identify μ on $(\Omega_{-\infty}^0, \mathcal{F}_{\leq 0})$ with its natural extension on (Ω, \mathcal{F}) .

The previous definition is not very useful to prove the regularity of a measure. For this, the following well known result is often useful. Let us call a measure μ on (Ω, \mathcal{F}) **non-null** if it gives non-zero measure to every cylinder. We then have (see, for instance, [11]):

Theorem 2.4. A probability measure μ on (Ω, \mathcal{F}) with underlying process $(X_i)_{i \in \mathbb{Z}}$ on $(\Omega, \mathcal{F}, \mu)$ is a regular g -measure iff it is non-null, shift-invariant and the sequences

$$\left[\mu(X_0 = \omega_0 \mid X_{-n}^{-1} = \omega_{-n}^{-1}) \right]_{n \geq 1} \quad \text{converge uniformly as } n \rightarrow \infty. \quad (2.4)$$

Let us now review the notions of specification and Gibbs measures. It involves definitions and properties analogous to the preceding ones, but replacing one-side by two-side conditioning. (We specialize to the one-dimensional setting but similar notions and results are valid for any dimension.)

Definition 2.5. A specification γ on (Ω, \mathcal{F}) is a family of probability kernels $\{\gamma_\Lambda: \Lambda \subset \mathbb{Z}, |\Lambda| < \infty\}$, $\gamma_\Lambda: \mathcal{F} \times \Omega \rightarrow [0, 1]$ such that for all finite $\Lambda \subset \mathbb{Z}$:

- $\gamma_\Lambda(A \mid \cdot) \in \mathcal{F}_{\Lambda^c}$, for each $A \in \mathcal{F}$;
- $\gamma_\Lambda(B \mid \omega) = \mathbb{1}_B(\omega)$, for each $B \in \mathcal{F}_{\Lambda^c}$ and $\omega \in \Omega$;
- $\gamma_\Delta \gamma_\Lambda = \gamma_\Delta$, for each finite $\Delta \subset \mathbb{Z}$: $\Delta \supset \Lambda$, i.e.,

$$\iint h(\xi) \gamma_\Lambda(d\xi \mid \sigma) \gamma_\Delta(d\sigma \mid \omega) = \int h(\sigma) \gamma_\Delta(d\sigma \mid \omega) \quad (2.5)$$

for all measurable functions h and configurations $\omega \in \Omega$.

In words, a specification is almost a *regular system of conditional probabilities* for finite regions conditioned on the external sigma algebras. The only difference lies in the fact that conditional probabilities satisfy (2.5) for *almost all* ω with respect to some measure known beforehand. In the definition of specification there is no previously known measure, thus condition (2.5) is asked *for all* ω (see e.g. [3] for more details). We can now define regularity, as expected, as a property of all the kernels γ_Λ . In our setting, however, we are interested only in non-null specifications and they, in fact, are entirely determined by their single-site part $\{\gamma_{\{i\}}: i \in \mathbb{Z}\}$ (see [7] and references therein). Relevant definitions need only be done, thus, at the level of these single-site kernels. That is what we do now, to emphasize the parallelism with the treatment of g -measures. For brevity we denote $\gamma_{\{i\}}(\omega_i \mid \cdot)$ the kernel $\gamma_{\{i\}}$ applied to the cylinder of base ω_i .

Definition 2.6. A specification γ on (Ω, \mathcal{F}) is **regular** if, for each $i \in \mathbb{Z}$,

- the function $\gamma_{\{i\}}(\omega_i \mid \cdot)$ is **continuous** for each $\omega_i \in E$, i.e., for all $\epsilon > 0$, there exists $n, m \geq 0$ so that

$$\left| \gamma_{\{i\}}(\omega_i \mid \omega_{\{i\}^c}) - \gamma_{\{i\}}(\omega_i \mid \sigma_{\{i\}^c}) \right| < \epsilon \quad (2.6)$$

for all $\omega, \sigma \in \Omega$ with $\omega_{-n}^m = \sigma_{-n}^m$;

- the function $\gamma_{\{i\}}(\omega_i | \cdot)$ is **strongly non-null** for each $\omega_i \in E$, i.e., $\gamma_{\{i\}}(\omega_i | \cdot) \geq c > 0$.

Definition 2.7. A probability measure μ on (Ω, \mathcal{F}) with underlying field $(X_i)_{i \in \mathbb{Z}}$ on $(\Omega, \mathcal{F}, \mu)$ is said to be a **Gibbs measure** if μ is **consistent** with some regular specification γ , namely,

$$\mu(X_i = \omega_i | X_{\{i\}^c} = \omega_{\{i\}^c}) = \gamma_{\{i\}}(\omega_i | \omega_{\{i\}^c}) \quad (2.7)$$

for all $\omega_i \in E$ and μ -a.e. $\omega_{\{i\}^c} \in \Omega_{\{i\}^c}$.

The following is the analogues of Theorem 2.4.

Theorem 2.8. A shift-invariant probability measure μ on (Ω, \mathcal{F}) with underlying field $(X_i)_{i \in \mathbb{Z}}$ on $(\Omega, \mathcal{F}, \mu)$ is a Gibbs measure iff it is non-null and the sequences

$$\left[\mu(X_0 = \omega_0 | X_{-m}^{-1} = \omega_{-m}^{-1}, X_1^m = \omega_1^m) \right]_{n, m \geq 1} \text{ converge uniformly as } n, m \rightarrow \infty. \quad (2.8)$$

A measure violating (2.8) is, thus, a *non-Gibbsian measure* (see, e.g. [5] and references therein). More specifically, the absence of convergence of a sequence (2.8) corresponds to an *essential discontinuity* at $\omega \in \Omega$, that is, a discontinuity in the conditional expectations that can not be removed by a redefinition on a zero-measure set (a more detailed discussion of these issues can be found in [5, Section 5.3]).

The link between Definition 2.7 and the usual notion in classical statistical mechanics is provided by a theorem due to Kozlov [14] that states that a specification is Gibbsian if and only if it has the Boltzmann form

$$\gamma_\Lambda(\omega_i | \omega_{\{i\}^c}) = \exp\left\{-\sum_{A \ni i} \phi_A(\omega_A)\right\} / \text{Norm.}, \quad (2.9)$$

where the functions ϕ_A (*interaction*) satisfy the summability condition

$$\sup_{i \in \mathbb{Z}} \sum_{A \ni i} \|\phi_A\|_\infty < \infty. \quad (2.10)$$

In the theory of dynamical systems, often Gibbsianness is associated with the SRB measures. These are measures μ for which there exists a function $\psi: \{-1, 1\}^{\mathbb{Z}^+} \rightarrow \mathbb{R}$, a constant $\Theta = \Theta(\psi)$, and some finite positive constants \underline{c}, \bar{c} such that

$$\underline{c} \leq \frac{\mu(\omega_0^n)}{\exp\left(\sum_{i=0}^n \psi(\tau^i \omega) - (n+1)\Theta\right)} \leq \bar{c} \quad (2.11)$$

$[\tau^i$ is the i th iterate of the shift on $\Omega]$. In the non-null case, these SRB measures form a strict subset of the one-dimensional Gibbs measures. This can be seen in two ways. First, general Gibbs measures satisfy (2.11) but with the constants \underline{c}, \bar{c} substituted by $o(n)$ -functions. Second, the corresponding interactions for the non-null SRB measures must satisfy the condition

$$\sup_{i \in \mathbb{Z}} \sum_{A \ni i} \text{diam}(A) \|\phi_A\|_\infty < \infty \quad (2.12)$$

which is stronger than (2.10).

3 Main result

For any $\underline{\omega} = \omega_{-\infty}^{-1} \in \Omega_{-\infty}^{-1}$ and $\bar{\omega} = \omega_1^{+\infty} \in \Omega_1^{+\infty}$, let

$$\ell(\underline{\omega}) = \min\{j \geq 0: \omega_{-j-1} = 1\} \quad \text{and} \quad m(\bar{\omega}) = \min\{j \geq 0: \omega_{j+1} = 1\} \quad (3.1)$$

denote the number of 0's before finding the first 1 when looking backward in $\underline{\omega}$ and forward in $\bar{\omega}$, respectively [$\ell(\underline{0}) = m(\bar{0}) = \infty$].

Our g -function is defined by a converging sequence $\{p_i\}_{i \geq 0}$ of numbers with values in $(0, 1)$, satisfying

$$\inf_{i \geq 0} p_i = \epsilon > 0 \quad , \quad p_\infty = \lim_{i \rightarrow +\infty} p_i . \quad (3.2)$$

The kernel P is defined on $E \times \Omega_{-\infty}^{-1}$ by

$$P(1 | \underline{\omega}) = p_{\ell(\underline{\omega})} \quad \forall \underline{\omega} \in \Omega_{-\infty}^{-1} . \quad (3.3)$$

Note that the continuity of the kernel P follows from the fact that

$$\begin{aligned} \sup_{\substack{\underline{\omega}, \underline{\sigma} \in \Omega_{-\infty}^{-1} \\ \omega_{-k}^{-1} = \sigma_{-k}^{-1}}} \left| P(1 | \underline{\omega}) - P(1 | \underline{\sigma}) \right| &= \sup_{\omega_{-k}^{-k-1}, \sigma_{-k}^{-k-1} \in \Omega_{-\infty}^{-k-1}} \left| P(1 | 0_{-k}^{-1} \omega_{-\infty}^{-k-1}) - P(1 | 0_{-k}^{-1} \sigma_{-\infty}^{-k-1}) \right| \\ &= \sup_{l, m \geq k} |p_l - p_m| \end{aligned} \quad (3.4)$$

which goes to 0 since $\{p_k\}_{k \geq 0}$ converges.

This g -function is, therefore, regular and, furthermore, in [8] it is proven that there exists a unique stationary chain μ compatible with P which is the renewal chain with infinitely many 1's separated by intervals of 0's of random length and having exponential tail distribution. For all practical purposes, this chain is as regular as it can be. Nevertheless, it is not necessary Gibbsian.

Theorem 3.1. *There exist choices of the sequence $\{p_i\}_{i \geq 0}$ satisfying (3.2) for which*

$$\mu(X_0 = 0 | \cdot) \text{ is essentially discontinuous at } 0_{-\infty}^{+\infty} , \quad (3.5)$$

where μ is the (unique and non-null) g -measure compatible with the kernel defined by (3.3).

By Kozlov's theorem [14] this means that the resulting g -measure is non-Gibbsian in the statistical mechanical sense, and hence neither in SRB sense.

Proof. The proof consists in the observation that if ω has $\ell(\underline{\omega}) = i$ and $m(\bar{\omega}) = j$, then, as we will see in (3.12) below, $\mu(X_0 = 0 | X_{-m}^{-1} = \omega_{-m}^{-1}, X_1^n = \omega_1^n)$ is determined by the ratio

$$\prod_{k=0}^{i-1} \frac{1 - p_k}{1 - p_{k+j}} . \quad (3.6)$$

Thus, the discontinuity at $0_{-\infty}^{+\infty}$ is equivalent to the existence of a sequence of p_k for which this ratio oscillates with i and j . The most economical way of achieving this is to define

$$p_k = 1 - (1 - p_\infty) \xi^{v_k} , \quad (3.7)$$

so that

$$\prod_{k=0}^{i-1} \frac{1-p_k}{1-p_{k+j}} = \xi^{\sum_{k=0}^{i-1} (v_k - v_{k+j})}. \quad (3.8)$$

The discontinuity is obtained by choosing a sequence v_k converging to 0 when $k \rightarrow \infty$, but such that $\sum_{k=0}^i v_k$ oscillates.

To formalize this idea, let us first provide an explicit expression for the conditional probabilities. By construction, the measure μ compatible with P has the property that, if $\ell(\underline{\omega}) = i \leq m < \infty$ and $m(\overline{\omega}) = j \leq n < \infty$

$$\begin{aligned} \mu(X_{-m}^n = \omega_{-m}^n) \\ = \mu(X_{-m}^{-i-1} = \omega_{-m}^{-i-1}) \mu(X_{-i}^{j+1} = \omega_{-i}^{j+1} \mid X_{-i-1} = 1) \mu(X_{j+2}^n = \omega_{j+2}^n \mid X_{j+1} = 1) \end{aligned} \quad (3.9)$$

which becomes

$$\mu(X_{-m}^{-i-1} = \omega_{-m}^{-i-1}) \left(\prod_{k=0}^{i+j} (1-p_k) p_{i+j+1} \right) \mu(X_{j+2}^n = \omega_{j+2}^n \mid X_{j+1} = 1) \quad (3.10)$$

when $\omega_0 = 0$, and

$$\mu(X_{-m}^{-i-1} = \omega_{-m}^{-i-1}) \left(\prod_{k=0}^{i-1} (1-p_k) p_i \right) \left(\prod_{k=0}^{j-1} (1-p_k) p_j \right) \mu(X_{j+2}^n = \omega_{j+2}^n \mid X_{j+1} = 1) \quad (3.11)$$

when $\omega_0 = 1$. Thus, the finite-volume 2-sided conditional probability equals

$$\begin{aligned} \mu(X_0 = 0 \mid X_{-m}^{-1} = \omega_{-m}^{-1}, X_1^n = \omega_1^n) \\ = \frac{\mu(X_{-m}^n = \omega_{-m}^n)}{\mu(X_{-m}^n = \omega_{-m}^n) + \mu(X_{-m}^n = \omega_{-m}^{-1} 1_0 \omega_1^n)} \\ = \frac{\prod_{k=0}^{i+j} (1-p_k) p_{i+j+1}}{\prod_{k=0}^{i-1} (1-p_k) p_i \prod_{k=0}^{j-1} (1-p_k) p_j + \prod_{k=0}^{i+j} (1-p_k) p_{i+j+1}} \\ = \left(1 + \frac{p_i p_j}{(1-p_{i+j}) p_{i+j+1}} \prod_{k=0}^{i-1} \frac{1-p_k}{1-p_{k+j}} \right)^{-1} \end{aligned} \quad (3.12)$$

for all $n \geq i$ and $m \geq j$. Notice that the intermediate equality is symmetric under the interchange $i \leftrightarrow j$; this proves the left-right symmetry of the conditional expectations. Writing the probabilities p_k in the form (3.7) this becomes

$$\begin{aligned} \mu(X_0 = 0 \mid X_{-i-1}^{-1} = 1_{-i-1} 0_{-i}^{-1}, X_1^{j+1} = 0_1^j 1_{j+1}) \\ = \left(1 + \frac{[1 - (1-p_\infty)\xi^{v_i}][1 - (1-p_\infty)\xi^{v_j}]}{(1-p_\infty)\xi^{v_{i+j}}[1 - (1-p_\infty)\xi^{v_{i+j+1}}]} \xi^{\sum_{k=0}^{i-1} (v_k - v_{k+j})} \right)^{-1}. \end{aligned} \quad (3.13)$$

To conclude we must choose a sequence v_k with an oscillating sum, as proposed after (3.8). We choose $\xi \in (1, (1-p_\infty)^{-2})$ and consider the sequence $\{v_k\}_{k \geq 0}$ with

$$v_k = \frac{(-1)^{r_k}}{r_k} \quad \text{with} \quad r_k = \inf \left\{ i \geq 1 : \sum_{j=1}^i j \geq k+1 \right\}. \quad (3.14)$$

The first terms of this sequence are as follows:

$$-1, \frac{1}{2}, \frac{1}{2}, -\frac{1}{3}, -\frac{1}{3}, -\frac{1}{3}, \frac{1}{4}, \frac{1}{4}, \frac{1}{4}, \frac{1}{4}, -\frac{1}{5}, -\frac{1}{5}, -\frac{1}{5}, -\frac{1}{5}, -\frac{1}{5}, \dots \quad (3.15)$$

Clearly, v_k converges to 0 when $k \rightarrow \infty$, while $\sum_{k=0}^i v_k$ does not converge as i diverges because it oscillates inside $[-1, 0]$. Due to (3.14), there are strictly increasing subsequences

- $\{i_n^{(1)}\}_{n \geq 1}$ such that $\sum_{k=0}^{i_n^{(1)}-1} v_k = -1$ for any $n \geq 1$;
- $\{j_n^{(1)}\}_{n \geq 1}$, where $j_n^{(1)}$ depends on $i_n^{(1)}$, such that $\sum_{k=0}^{i_n^{(1)}-1} v_{j_n^{(1)}+k} = 0$ for any $n \geq 1$;
- $\{i_n^{(2)}\}_{n \geq 1}$ such that $\sum_{k=0}^{i_n^{(2)}-1} v_k = 0$ for any $n \geq 1$;
- $\{j_n^{(2)}\}_{n \geq 1}$, where $j_n^{(2)}$ depends on $i_n^{(2)}$, such that $\sum_{k=0}^{i_n^{(2)}-1} v_{j_n^{(2)}+k} = 0$ for any $n \geq 1$.

Therefore, we have that

$$\sum_{k=0}^{i_n^{(1)}-1} (v_k - v_{j_n^{(1)}+k}) = -1 \quad (3.16)$$

and

$$\sum_{k=0}^{i_n^{(2)}-1} (v_k - v_{j_n^{(2)}+k}) = 0. \quad (3.17)$$

Since

$$\lim_{i,j \rightarrow \infty} \frac{[1 - (1 - p_\infty)\xi^{v_i}][1 - (1 - p_\infty)\xi^{v_j}]}{(1 - p_\infty)\xi^{v_i+j}[1 - (1 - p_\infty)\xi^{v_i+j+1}]} = \frac{p_\infty}{(1 - p_\infty)}, \quad (3.18)$$

it follows from (3.13)–(3.18) that

$$\begin{aligned} \lim_{n \rightarrow \infty} \mu \left(X_0 = 0 \mid X_{-i_n^{(\ell)}-1}^{-1} = 1_{-i_n^{(\ell)}-1} 0_{-i_n^{(\ell)}}^{-1}, X_1^{j_n^{(\ell)}+1} = 0_1^{j_n^{(\ell)}} 1_{j_n^{(\ell)}+1} \right) \\ = \begin{cases} \frac{(1 - p_\infty)\xi}{(1 - p_\infty)\xi + p_\infty} & \text{if } \ell = 1 \\ (1 - p_\infty) & \text{if } \ell = 2. \end{cases} \end{aligned} \quad (3.19)$$

Hence $\lim_{i,j \rightarrow \infty} \mu(X_0 = 0 \mid X_{-i-1}^{-1} = 1_{-i-1} 0_{-i}^{-1}, X_1^{j+1} = 0_1^j 1_{j+1})$ does not exist. This contradicts (2.8) and, more specifically, proves the essential discontinuity (3.5). \blacksquare

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