# ON THE VERLINDE FORMULAS FOR SO(3)-BUNDLES 

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#### Abstract

This paper computes the quantization of the moduli space of flat $\mathrm{SO}(3)$-bundles over an oriented surface with boundary, with prescribed holonomies around the boundary circles. The result agrees with the generalized Verlinde formula conjectured by Fuchs and Schweigert.


## 1. Introduction

Let $G$ be a compact, connected Lie group, $\Sigma$ a compact oriented surface of genus $h$ with $r$ boundary components. Given conjugacy classes $\mathcal{C}_{1}, \ldots, \mathcal{C}_{r} \subset G$, denote by

$$
\begin{equation*}
\mathcal{M}\left(\Sigma, \mathcal{C}_{1}, \ldots, \mathcal{C}_{r}\right) \tag{1}
\end{equation*}
$$

the moduli space of flat $G$-bundles over $\Sigma$, with boundary holonomies in prescribed conjugacy classes $\mathcal{C}_{j}$. The choice of an invariant inner product on $\mathfrak{g}$ defines a symplectic structure on the moduli space. Under suitable integrality conditions the moduli space carries a pre-quantum line bundle $L$, and one can define the quantization

$$
\begin{equation*}
\mathcal{Q}\left(\mathcal{M}\left(\Sigma, \mathcal{C}_{1}, \ldots, \mathcal{C}_{r}\right)\right) \in \mathbb{Z} \tag{2}
\end{equation*}
$$

as the index of the $\operatorname{Spin}_{c}$-Dirac operator with coefficients in $L$. (It may be necessary to use a partial desingularization as in [16].) Choosing a complex structure on $\Sigma$ further defines a Kähler structure on the moduli space. If $G$ is simply connected, Kodaira vanishing results [20] show that the above index coincides with the dimension of the space of holomorphic sections of $L$. It is given by the celebrated Verlinde formula $[22,21,7,19,5]$. For symplectic approaches to the Verlinde formulas, much in the spirit of the present paper, see $[11,10,12,6,2]$.

Much less is known for non-simply connected groups. For surfaces without boundary ( $r=0$ ), and taking $G=\mathrm{PU}(n)$, Verlinde-type formulas were obtained by Pantev [17] in the case $n=2$ and by Beauville [4] for $n$ prime. For more general compact, semi-simple connected Lie groups, Fuchs and Schweigert [9] conjectured a generalization of the Verlinde formula, expressed in terms of orbit Lie algebras. Partial results on these conjectures were obtained in [2].

In this article, we will establish Fuchs-Schweigert formulas for the index (2) for the simplest case $G=\mathrm{SO}(3)$. We will use the recently developed quantization procedure [15, 14] for quasi-Hamiltonian actions with group-valued moment map [1]. In order to apply these techniques, we present the moduli spaces (1) as symplectic quotients of quasi-Hamiltonian $\tilde{G}$-spaces for the universal cover $\tilde{G}=\mathrm{SU}(2)$. In more detail, let $\mathcal{D}_{i} \subset \mathrm{SU}(2)$ be conjugacy classes, and consider the quasi-Hamiltonian SU(2)-space

$$
\tilde{M}=\mathcal{D}_{1} \times \cdots \times \mathcal{D}_{s} \times \mathrm{SU}(2)^{2 h}
$$

with moment map the product of holonomies,

$$
\tilde{\Phi}\left(d_{1}, \ldots, d_{s}, a_{1}, b_{1}, \ldots, a_{h}, b_{h}\right)=\prod_{i=1}^{s} d_{i} \prod_{j=1}^{h}\left(a_{j} b_{j} a_{j}^{-1} b_{j}^{-1}\right)
$$

Put $M=\tilde{M} / \Gamma$, where $\Gamma \subset Z^{s+2 h}$ is the subgroup preserving $\tilde{M} \subset \mathrm{SU}(2)^{s+2 h}$ and $\tilde{\Phi}$. Then $M$ is a quasi-Hamiltonian $\mathrm{SU}(2)$-space, and all connected components of moduli spaces (2) are symplectic quotients $M / / \mathrm{SU}(2)$ for suitable choices of $\mathcal{D}_{j}$ (see Section 2.3). Our first main result gives necessary and sufficient conditions under which the space $M$ admits a level $k$ pre-quantization [13]. Using localization, we then compute the corresponding quantization $\mathcal{Q}(M) \in R_{k}(\mathrm{SU}(2))$, an element of the level $k$ fusion ring (Verlinde ring). These results are summarized in Theorem 3.7 . We reformulate the result as an equivariant version of the Fuchs-Schweigert formula (Theorem 4.1); the non-equivariant formula (see (16) in Section 4) is then obtained from a 'quantization commutes with reduction' principle.

Using the results of [14], it is also possible to compute quantizations of moduli spaces for non-simply connected groups of higher rank. However, the determination of the pre-quantization conditions and the evaluation of the fixed point contributions becomes more involved. We will return to these questions in a forthcoming paper; see also the author's abstracts in Oberwolfach Report No. 2011/09.

## 2. Preliminaries

The following notation, consistent with [15], will be used in this paper. For the Lie group $\mathrm{SU}(2)$ let $T$ be the maximal torus given as the image of

$$
j: \mathrm{U}(1) \rightarrow \mathrm{SU}(2), \quad j(z)=\left(\begin{array}{cc}
z & 0 \\
0 & z
\end{array}\right)
$$

Let $\Lambda=\operatorname{ker} \exp _{T} \subset \mathfrak{t}$ denote the integral lattice and $\Lambda^{*} \subset \mathfrak{t}^{*}$ its dual, the (real) weight lattice. Let $\rho \in \Lambda^{*}$ be the generator dual to the generator $\mathrm{d} j(2 \pi i) \in \Lambda$. We will use the basic inner product on $\mathfrak{s u}(2)$,

$$
\xi \cdot \xi^{\prime}:=\frac{1}{4 \pi^{2}} \operatorname{tr}\left(\xi^{\dagger} \xi^{\prime}\right), \quad \xi, \xi^{\prime} \in \mathfrak{s u}(2)
$$

to identify $\mathfrak{s u}(2) \cong \mathfrak{s u}(2)^{*}$. Under this identification, $\|\rho\|^{2}=\frac{1}{2}$, and $\Lambda=2 \Lambda^{*}$ with generator $2 \rho$. The following two elements of $\mathrm{SU}(2)$ will play a special role in this paper:

$$
u_{*}=\left(\begin{array}{cc}
0 & 1 \\
-1 & 0
\end{array}\right), \quad t_{*}=\left(\begin{array}{cc}
i & 0 \\
0 & -i
\end{array}\right)
$$

Observe that $t_{*}=\exp (\rho / 2)$, with square $c=\exp \rho$ the non-trivial element in the center $Z:=Z(\mathrm{SU}(2)) \cong \mathbb{Z}_{2}$. The element $u_{*} \in N(T)$ represents the non-trivial element of the Weyl group $W=N(T) / T \cong \mathbb{Z}_{2}$. Both $u_{*}, t_{*}$ are contained in the conjugacy class $\mathcal{D}_{*} \subset \mathrm{SU}(2)$ of elements of trace 0 . Note that $\mathcal{D}_{*}$ is the unique conjugacy class in $\mathrm{SU}(2)$ that is invariant under multiplication by $Z$. The quotient $\mathcal{C}_{*}=\mathcal{D}_{*} / Z \cong \mathbb{R} P(2)$ is the conjugacy class in $\mathrm{SO}(3)$ consisting of rotations by $\pi$.
2.1. The fusion ring $R_{k}(\mathrm{SU}(2))$. We view the representation ring $R(\mathrm{SU}(2))$ as the subring of $C^{\infty}(\mathrm{SU}(2))$ generated by characters of $\mathrm{SU}(2)$-representations. As a $\mathbb{Z}$-module, it is free with basis $\chi_{0}, \chi_{1}, \chi_{2}, \ldots$, where $\chi_{m}$ is the character of the irreducible $\mathrm{SU}(2)$-representation on the $m$-th symmetric power $S^{m}\left(\mathbb{C}^{2}\right)$. The ring structure is determined by the formula

$$
\chi_{m} \chi_{m^{\prime}}=\chi_{m+m^{\prime}}+\chi_{m+m^{\prime}-2}+\cdots+\chi_{\left|m-m^{\prime}\right|}
$$

For $k=0,1,2, \ldots$ let $I_{k}(\mathrm{SU}(2))$ be the ideal generated by $\chi_{k+1}$ and let

$$
R_{k}(\mathrm{SU}(2))=R(\mathrm{SU}(2)) / I_{k}(\mathrm{SU}(2))
$$

be the level $k$ fusion ring (or Verlinde ring). As a $\mathbb{Z}$-module, $R_{k}(\mathrm{SU}(2))$ is free, with basis $\tau_{0}, \tau_{1}, \ldots, \tau_{k}$ the images of $\chi_{0}, \chi_{1}, \ldots, \chi_{k}$ under the quotient homomorphism. Let $q=e^{\frac{i \pi}{k+2}}$ be the $2 k+4$-th root of unity, and define special points

$$
\begin{equation*}
t_{l}=j\left(q^{l+1}\right), \quad l=0, \ldots, k \tag{3}
\end{equation*}
$$

Then $I_{k}(\mathrm{SU}(2)) \subset R(\mathrm{SU}(2))$ has an alternative description as the ideal of characters vanishing at all special points (3). Hence, the evaluation of characters at the special points descends to evaluations $R_{k}(\mathrm{SU}(2)) \rightarrow \mathbb{C}, \tau \mapsto \tau\left(t_{l}\right)$.

The product in the complexified fusion ring $R_{k}(\mathrm{SU}(2)) \otimes_{\mathbb{Z}} \mathbb{C}$ can be diagonalized using the $S$-matrix, given by the Kac-Peterson formula

$$
\begin{equation*}
S_{m, l}=\left(\frac{k}{2}+1\right)^{-\frac{1}{2}} \sin \left(\frac{\pi(l+1)(m+1)}{k+2}\right) \tag{4}
\end{equation*}
$$

for $l, m=0,1, \ldots, k$. The $S$-matrix is orthogonal, and the alternative basis elements

$$
\tilde{\tau}_{l}=\sum_{m} S_{0, l} S_{m, l} \tau_{m}
$$

satisfy $\tilde{\tau}_{m}\left(t_{l}\right)=\delta_{m, l}$, hence

$$
\tilde{\tau}_{m} \tilde{\tau}_{m^{\prime}}=\delta_{m, m^{\prime}} \tilde{\tau}_{m}
$$

The basis elements $\left\{\tau_{0}, \ldots, \tau_{k}\right\}$ are expressed in terms of the alternative basis as $\tau_{m}=\sum_{l} S_{0, l}^{-1} S_{m, l} \tilde{\tau}_{l}$.
2.2. Quasi-Hamiltonian $G$-spaces. We recall some basic definitions and facts from [1]. Let $G$ be a compact Lie group with Lie algebra $\mathfrak{g}$, equipped with an invariant inner product, denoted by a dot $\cdot$ Let $\theta^{L}, \theta^{R}$ denote the left-invariant, right-invariant Maurer-Cartan forms on $G$, and let $\eta=\frac{1}{12} \theta^{L} \cdot\left[\theta^{L}, \theta^{L}\right]$ denote the Cartan 3-form on $G$. For a $G$-manifold $M$, and $\xi \in \mathfrak{g}$, let $\xi^{\sharp}$ denote the generating vector field, defined in terms of the action on functions $f \in C^{\infty}(M)$ by $\left(\xi^{\sharp} f\right)(x)=$ $\left.\frac{d}{d t}\right|_{t=0} f(\exp (-t \xi) \cdot x)$. The Lie group $G$ is itself viewed as a $G$-manifold for the conjugation action.

Definition 2.1. A quasi-Hamiltonian $G$-space is a triple $(M, \omega, \Phi)$ consisting of a $G$-manifold $M$, a $G$-invariant 2-form $\omega$ on $M$, and an equivariant map $\Phi: M \rightarrow G$, called the moment map, satisfying:
(1) $d \omega+\Phi^{*} \eta=0$,
(2) $\iota_{\xi^{\sharp}} \omega+\frac{1}{2} \Phi^{*}\left(\left(\theta^{L}+\theta^{R}\right) \cdot \xi\right)=0$ for all $\xi \in \mathfrak{g}$,
(3) at every point $x \in M$, $\operatorname{ker} \omega_{x} \cap \operatorname{kerd} \Phi_{x}=\{0\}$.

The fusion product of two quasi-Hamiltonian $G$-spaces $\left(M_{1}, \omega_{1}, \Phi_{1}\right)$ and $\left(M_{2}, \omega_{2}, \Phi_{2}\right)$ is the product $M_{1} \times M_{2}$, with the diagonal $G$-action, 2-form

$$
\begin{equation*}
\omega=\operatorname{pr}_{1}^{*} \omega_{1}+\operatorname{pr}_{2}^{*} \omega_{2}+\frac{1}{2} \operatorname{pr}_{1}^{*} \Phi_{1}^{*} \theta^{L} \cdot \operatorname{pr}_{2}^{*} \Phi_{2}^{*} \theta^{R} \tag{5}
\end{equation*}
$$

and moment map $\Phi=\Phi_{1} \Phi_{2}$.
The symplectic quotient of a quasi-Hamiltonian $G$-space is the symplectic space $M / / G=\Phi^{-1}(e) / G$. Similar to the theory of Hamiltonian group actions, the group unit $e$ is a regular value of $\Phi$ if and only if $G$ acts locally freely on the level set $\Phi^{-1}(e)$, and in this case the pull-back of the 2 -form to the level set descends to a
symplectic 2 -form on the orbifold $\Phi^{-1}(e) / G$. If $e$ is a singular value, then $M / / G$ is a singular symplectic space as defined in [18].

The conjugacy classes $\mathcal{C} \subset G$ are basic examples of quasi-Hamiltonian $G$-spaces. The moment map is the inclusion into $G$, and the 2 -form $\omega$ is given on generating vector fields by the formula

$$
\begin{equation*}
\omega_{g}\left(\zeta^{\sharp}(g), \xi^{\sharp}(g)\right)=\frac{1}{2}\left(\xi \cdot \operatorname{Ad}_{g} \zeta-\zeta \cdot \operatorname{Ad}_{g} \xi\right) . \tag{6}
\end{equation*}
$$

Together with the double $\mathbf{D}(G)=G \times G$, equipped with diagonal $G$-action and moment map $\Phi(g, h)=g h g^{-1} h^{-1}$, these are the building blocks of the main example appearing in this paper. As shown in [1], the moduli space of flat $G$-bundles over a compact, oriented surface $\Sigma$ of genus $h$ with $s$ boundary components, with boundary holonomies in prescribed conjugacy classes $\mathcal{C}_{j}, j=1, \ldots, s$, is a symplectic quotient of a fusion product:

$$
\begin{equation*}
M\left(\Sigma, \mathcal{C}_{1}, \ldots, \mathcal{C}_{s}\right)=\mathcal{C}_{1} \times \cdots \times \mathcal{C}_{s} \times \mathbf{D}(G)^{h} / / G \tag{7}
\end{equation*}
$$

If the group $G$ is simply connected, then the fibers of the moment map for any compact, connected quasi-Hamiltonian $G$-space are connected. In particular, (7) is connected in that case. If $G$ is non-simply connected, the space (7) may have several components.

To clarify the decomposition into components, we use the following construction. Suppose $p: \breve{G} \rightarrow G$ is a homomorphism of compact, connected Lie groups, with finite kernel $Z$. Then $Z$ is a subgroup of the center of $\check{G}$, and $G=\check{G} / Z$. For any quasi-Hamiltonian $G$-space $(N, \omega, \Phi)$, let $\check{N}$ denote the fiber product defined by the pull-back square


Then $(\check{N}, \check{\omega}, \check{\Phi})$ is a quasi-Hamiltonian $\check{G}$-space, for the diagonal $\check{G}$-action on $\check{N} \subset$ $N \times \check{G}$, and with the 2 -form $\check{\omega}=p_{N}^{*} \omega$. Simple properties of this construction are:

Proposition 2.2. (i) We have a canonical identification of symplectic quotients

$$
\check{N} / / \check{G} \cong N / / G
$$

(ii) For a fusion product $N=N_{1} \times \cdots \times N_{r}$ of quasi-Hamiltonian $G$-spaces, the space $\check{N}$ is a quotient of $\check{N}_{1} \times \cdots \times \check{N}_{r}$ by the group $\left\{\left(c_{1}, \ldots, c_{r}\right) \in\right.$ $\left.Z^{r} \mid \prod_{j=1}^{r} c_{j}=e\right\}$.
(iii) If $\Phi: N \rightarrow G$ lifts to a moment map $\Phi^{\prime}: N \rightarrow \check{G}$, thus turning $N$ into a quasi-Hamiltonian $\check{G}$-space $\left(N, \omega, \Phi^{\prime}\right)$, then

$$
\check{N}=N \times Z
$$

as a fusion product of quasi-Hamiltonian $\check{G}$-spaces. Here $Z$ is viewed as a quasi-Hamiltonian $\check{G}$-space, with trivial action and with moment map the inclusion to $\check{G}$.

Proof. (i) By definition of $\check{N}$, the level sets $\check{\Phi}^{-1}(\check{e})$ and $\Phi^{-1}(e)$ are identified, and the pull-backs of the 2-forms to the level sets coincide. Since central elements in $\check{G}$ act trivially on $\check{N}$, the orbit spaces $\check{\Phi}^{-1}(\check{e}) / \check{G}$ and $\Phi^{-1}(e) / G$ are identified as well.
(ii) Think of the spaces $\check{N}_{i}$ as submanifolds of $N_{i} \times \check{G}$. The canonical map
$\check{N}_{1} \times \cdots \times \check{N}_{r} \rightarrow \check{N},\left(x_{1}, g_{1}, x_{2}, g_{2}, \ldots, x_{r}, g_{r}\right) \mapsto\left(x_{1}, \ldots, x_{r}, g_{1}, \ldots, g_{r}\right)$
is exactly the quotient map by $\left\{\left(c_{1}, \ldots, c_{r}\right) \in Z^{r} \mid \prod_{j=1}^{r} c_{j}=e\right\}$, and it preserves the $\check{G}$-actions and 2 -forms.
(iii) The map $N \times Z \rightarrow \check{N},(x, c) \mapsto\left(x, \Phi^{\prime}(x) c\right)$ is the desired diffeomorphism.
2.3. The moduli space example. Our main interest is the moduli space of flat $\mathrm{SO}(3)$-bundles with prescribed boundary holonomies, i.e. (7) with $G=\mathrm{SO}(3)$. In the notation of the previous Section, we will describe the quasi-Hamiltonian $\mathrm{SU}(2)$ space $\check{N}$ associated to the quasi-Hamiltonian $\mathrm{SO}(3)$-space

$$
N=\mathcal{C}_{1} \times \cdots \times \mathcal{C}_{s} \times \mathbf{D}(\mathrm{SO}(3))^{h}
$$

Choose conjugacy classes $\mathcal{D}_{j} \in \mathrm{SU}(2)$ with $p\left(\mathcal{D}_{j}\right)=\mathcal{C}_{j}$, and define a quasi-Hamiltonian SU(2)-space

$$
\begin{equation*}
\tilde{M}=\mathcal{D}_{1} \times \cdots \times \mathcal{D}_{s} \times \mathbf{D}(\mathrm{SU}(2))^{h} \tag{9}
\end{equation*}
$$

Put

$$
\begin{equation*}
M=\tilde{M} / \Gamma \tag{10}
\end{equation*}
$$

where $\Gamma \subset Z^{s+2 h}$ consists of $\gamma=\left(\gamma_{1}, \ldots, \gamma_{s+2 h}\right)$ with the properties $\prod_{j=1}^{s} \gamma_{j}=e$ and $\gamma_{j} \mathcal{D}_{j}=\mathcal{D}_{j}$ for $j \leq s$. (Equivalently, $\gamma_{j}=e$ for all $\mathcal{D}_{j} \neq \mathcal{D}_{*}$ ). The conditions guarantee that $\gamma \operatorname{acts}$ on $\tilde{M}$, preserving the 2 -form and moment map which hence descend to $M=\tilde{M} / \Gamma$. Let $\mathcal{C}_{*} \cong \mathbb{R} P(2)$ be the $\operatorname{SO}(3)$-conjugacy class consisting of rotations by $\pi$. It is the unique $\mathrm{SO}(3)$-conjugacy class whose pre-image in $\mathrm{SU}(2)$ is connected. This pre-image is the $\mathrm{SU}(2)$-conjugacy class $\mathcal{D}_{*} \cong S^{2}$ of matrices of trace 0 .

Lemma 2.3. With $N$ as above, we have

$$
\check{N} \cong \begin{cases}M & \text { if } \exists j: \mathcal{C}_{j}=\mathcal{C}_{*} \\ M \times Z & \text { if } \forall j: \mathcal{C}_{j} \neq \mathcal{C}_{*}\end{cases}
$$

Proof. The moment map $\mathbf{D}(\mathrm{SU}(2)) \rightarrow \mathrm{SU}(2)$ (given by Lie group commutator) is invariant under the action of $Z \times Z$, hence it descends to a lift $\mathbf{D}(\mathrm{SO}(3)) \rightarrow \mathrm{SU}(2)$ of the commutator map for $\mathrm{SO}(3)$. Thus

$$
\check{\mathbf{D}}(\mathrm{SO}(3))=\mathbf{D}(\mathrm{SO}(3)) \times Z
$$

If $\mathcal{C}_{j} \neq \mathcal{C}_{*}$, the map $\mathcal{D}_{j} \rightarrow \mathcal{C}_{j}$ is a diffeomorphism, and defines a lift of the moment $\operatorname{map} \mathcal{C}_{j} \hookrightarrow \mathrm{SO}(3)$. Hence

$$
\check{\mathcal{C}}_{j}=\mathcal{D}_{j} \times Z
$$

in that case. On the other hand, the conjugacy class $\mathcal{C}_{*}$ satisfies

$$
\check{\mathcal{C}}_{*}=\mathcal{D}_{*} .
$$

With these ingredients, the claim follows from Proposition 2.2.
We may choose the labeling of the conjugacy classes $\mathcal{C}_{1}, \ldots, \mathcal{C}_{s}$ in such a way that $\mathcal{C}_{j}=\mathcal{C}_{*}$ for $j \leq r$ and $\mathcal{C}_{j} \neq \mathcal{C}_{*}$ for $j>r$. The space (10) is then a fusion product

$$
\begin{equation*}
M=M^{\prime} \times \mathcal{D}_{r+1} \times \cdots \times \mathcal{D}_{s} \times \mathbf{D}(\mathrm{SO}(3))^{h} \tag{11}
\end{equation*}
$$

where $\mathbf{D}(\mathrm{SO}(3))$ is viewed as a quasi-Hamiltonian $\mathrm{SU}(2)$-space (using the canonical lift of the $\mathrm{SO}(3)$ moment map, as in the proof of Lemma 2.3), and where

$$
M^{\prime}=\left(\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}\right) / \Gamma^{\prime}
$$

with $r$ factors, and with $\Gamma^{\prime}=\left\{\left(\gamma_{1}, \ldots, \gamma_{r}\right) \in Z^{r} \mid \prod \gamma_{j}=e\right\}$. Let us describe the 2 -form $\omega^{\prime}$ of the space $M^{\prime}$, in terms of its pull-back $\tilde{\omega}^{\prime}$ to the universal cover $\tilde{M}=\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}$. Since the 2 -form on $\mathcal{D}_{*}$ is just zero, only the fusion terms contribute. By iterative use of the formula (5) for the fusion product, one obtains

$$
\begin{equation*}
\tilde{\omega}^{\prime}=\frac{1}{2} \sum_{i<j} g_{i}^{*} \theta^{L} \cdot \operatorname{Ad}_{g_{i+1} \cdots g_{j-1}}\left(g_{j}^{*} \theta^{R}\right), \tag{12}
\end{equation*}
$$

where $g_{i}: \tilde{M} \rightarrow \mathcal{D}_{*} \subset \mathrm{SU}(2)$ denotes projection onto the $i$-th factor.

## 3. Quantization of the moduli space of flat $\operatorname{SO}(3)$-Bundles

In this section we use localization to compute the quantization of the space $M=\left(\mathcal{D}_{1} \times \cdots \times \mathcal{D}_{s} \times \mathbf{D}(\mathrm{SU}(2))^{h}\right) / \Gamma$, as an element of the level $k$ fusion ring $R_{k}(\mathrm{SU}(2))$.
3.1. Pre-quantization. Recall that we fix the inner product $\cdot$ on $\mathfrak{s u}(2)$ to be the basic inner product. Then $\eta \in \Omega^{3}(\mathrm{SU}(2))$ is integral, and represents a generator $x \in H^{3}(\mathrm{SU}(2) ; \mathbb{Z}) \cong \mathbb{Z}$. The condition $d \omega+\Phi^{*} \eta=0$ from the definition of a quasiHamiltonian space says that the pair $(\omega, \eta)$ defines a relative cocycle in $\Omega^{3}(\Phi)$, the algebraic mapping cone of the pull-back map $\Phi^{*}: \Omega^{*}(G) \rightarrow \Omega^{*}(M)$. Let $k \in \mathbb{N}$.
Definition 3.1. [13, 15] A level $k$ pre-quantization of a quasi-Hamiltonian $\mathrm{SU}(2)$ space $(M, \omega, \Phi)$ is an integral lift $\alpha \in H^{3}(\Phi ; \mathbb{Z})$ of the class $k[(\omega, \eta)] \in H^{3}(\Phi ; \mathbb{R})$.

A necessary and sufficient condition for the existence of a level $k$ pre-quantization is that for all smooth singular 2-cycles $\Sigma \in Z_{2}(M)$, and all smooth singular 3-chains $C \in C_{3}(G)$ such that $\partial C=\Phi(\Sigma)$,

$$
k\left(\int_{\Sigma} \omega+\int_{C} \eta\right) \in \mathbb{Z}
$$

We list some basic properties and examples of level $k$ pre-quantizations.
(a) The set of level $k$ pre-quantizations is a torsor under the torsion group $\operatorname{Tor}\left(H^{2}(M, \mathbb{Z})\right)$ of isomorphism classes of flat line bundles.
(b) The level $k$ pre-quantized conjugacy classes of $\mathrm{SU}(2)$ are exactly those of the elements $\exp \left(\frac{m}{k} \rho\right)$ with $m=0, \ldots, k[15$, Proposition 7.3].
(c) The double $\mathbf{D}(\mathrm{SO}(3))$ (viewed as a quasi-Hamiltonian $\mathrm{SU}(2)$-space) admits a level $k$ pre-quantization if and only if $k$ is even [15, Proposition 7.4].
(d) If $M_{1}$ and $M_{2}$ are pre-quantized quasi-Hamiltonian $\mathrm{SU}(2)$-spaces at level $k$, then their fusion product $M_{1} \times M_{2}$ inherits a pre-quantization at level $k$. Conversely, a pre-quantization of the product induces pre-quantizations of the factors. See [13, Proposition 3.8].
(e) A level $k$ pre-quantization of $M$ induces a pre-quantization of the symplectic quotient $M / / \mathrm{SU}(2)$, equipped with the $k$-th multiple of the symplectic form.
(f) The long exact sequence in relative cohomology gives a necessary condition $k \Phi^{*}(x)=0$ for the existence of a level $k$ pre-quantization. If $H^{2}(M ; \mathbb{R})=0$, this condition is also sufficient [13, Proposition 4.2].
(g) The existence of the canonical 'twisted Spin $_{c}$-structure' [15, Section 6] on quasi-Hamiltonian $\mathrm{SU}(2)$-spaces $(M, \omega, \Phi)$ implies that $2 \Phi^{*}(x)=W^{3}(M)$, the third integral Stiefel-Whitney class. Since this is a 2-torsion class, $4 \Phi^{*}(x)=0$. In fact, there is a distinguished element $\beta \in H^{3}(\Phi ; \mathbb{Z})$ whose image in $H^{3}(\mathrm{SU}(2) ; \mathbb{Z})$ is $4 x$. If $H^{2}(M, \mathbb{R})=0$, this element gives a distinguished level 4 pre-quantization.
Given a level $k$ pre-quantization of a quasi-Hamiltonian $\operatorname{SU}(2)$-space ( $M, \omega, \Phi$ ) the construction from [15] produces a quantization $\mathcal{Q}(M) \in R_{k}(\mathrm{SU}(2))$, an element of the level $k$ fusion ring. It is obtained as a push-forward in twisted equivariant $K$-homology, using the Freed-Hopkins-Teleman theorem [8] to identify $R_{k}(\mathrm{SU}(2))$ with the equivariant twisted $K$-homology of $\mathrm{SU}(2)$ at level $k+2$. This is the quasi-Hamiltonian counterpart of the $\operatorname{Spin}_{c}$ quantization of an ordinary compact Hamiltonian $\mathrm{SU}(2)$-space, which produces an element of $R(\mathrm{SU}(2))$ as the equivariant index of a $\mathrm{Spin}_{c}$-Dirac operator with coefficients in an equivariant pre-quantum line bundle. The quantization procedure for quasi-Hamiltonian $G$-spaces satisfies properties similar to its Hamiltonian analog. These include
(1) compatibility with products, $\mathcal{Q}\left(M_{1} \times M_{2}\right)=\mathcal{Q}\left(M_{1}\right) \mathcal{Q}\left(M_{2}\right)$; and
(2) the 'quantization commutes with reduction' principle, $\mathcal{Q}(M / / G)=\mathcal{Q}(M)^{G}$. Here $R_{k}(G) \rightarrow \mathbb{Z}, \tau \mapsto \tau^{G}$ is the trace defined by $\tau_{m}^{G}=\delta_{m, 0}$.
3.2. Pre-quantization of $M$. Let us now consider level $k$ pre-quantizations of the quasi-Hamiltonian $\mathrm{SU}(2)$-space

$$
M=\left(\mathcal{D}_{1} \times \cdots \times \mathcal{D}_{s} \times \mathbf{D}(\mathrm{SU}(2))^{h}\right) / \Gamma
$$

from (10).
Theorem 3.2. The quasi-Hamiltonian $\mathrm{SU}(2)$-space $M$ carries a level $k$ pre-quantization if and only if the following conditions are satisfied:
(i) The conjugacy classes $\mathcal{D}_{j}$ are of the form $\mathrm{SU}(2) \cdot \exp \left(\frac{m_{j}}{k} \rho\right)$ with $m_{j} \in$ $\{0, \ldots, k\}$,
(ii) if $h \geq 1$, then $k \in 2 \mathbb{N}$,
(iii) if the number of $\mathcal{D}_{*}$-factors is $r \geq 3$, then $k \in 4 \mathbb{N}$.

Note that if at least one $\mathcal{D}_{*}$-factor appears, then the first condition requires that $k \in 2 \mathbb{N}$ since $\mathcal{D}_{*}=\mathrm{SU}(2)$. $\exp \left(\frac{1}{2} \rho\right)$.

Proof. Since a level $k$ pre-quantization of $M$ induces a level $k$ pre-quantization of the universal cover $\tilde{M}$, it is a necessary condition that all $\mathcal{D}_{j}$ be pre-quantizable. That is, $\mathcal{D}_{j}=\mathrm{SU}(2) . \exp \left(\frac{m_{j}}{k} \rho\right)$ with $m_{j} \in\{0, \ldots, k\}$.

Let us enumerate the conjugacy classes in such a way that $\mathcal{D}_{1}=\ldots=\mathcal{D}_{r}=\mathcal{D}_{*}$. Using the decomposition (11) and the known pre-quantization conditions (b), (c) for the conjugacy classes $\mathcal{D}_{j}$ and the double $\mathbf{D}(\mathrm{SO}(3))$, together with the fusion property (d), the proof is reduced to the case $h=0, s=r$. We may thus assume $M=\left(\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}\right) / \Gamma$ with $r$ factors. If $r=1$ then $M=\mathcal{D}_{*}$, which is pre-quantized at level $k$ if and only if $k$ is even. Suppose $r>1$. The non-trivial element $c \in Z$ acts on $H^{2}\left(\mathcal{D}_{*} ; \mathbb{R}\right) \cong \mathbb{R}$ as multiplication by -1 . Hence, $\Gamma$ acts on $H^{2}(M ; \mathbb{R}) \cong \mathbb{R}^{r}$ by componentwise sign changes. In particular, the $\Gamma$-invariant part is trivial. Since $\Gamma$ acts freely, it follows that

$$
H^{2}(M ; \mathbb{R}) \cong H^{2}(\tilde{M} ; \mathbb{R})^{\Gamma}=0
$$

Hence, by Property (f), a level $k$ pre-quantization exists if and only if $k \Phi^{*}(x)=0$. If $r=2$, so that $M=\left(\mathcal{D}_{*} \times \mathcal{D}_{*}\right) / \mathbb{Z}_{2}$, Poincaré duality gives that $H^{3}(M ; \mathbb{Z}) \cong \mathbb{Z}_{2}$; therefore $2 \Phi^{*}(x)=0$. Hence the condition $k \in 2 \mathbb{N}$ is also sufficient if $r=2$.

It remains to consider the case $r \geq 3$. By Property (g), the condition $k \in 4 \mathbb{N}$ is sufficient. Let us show that it is also necessary. Observe that the non-identity component of the normalizer, the circle $T u_{*}=N(T)-T$, is a single conjugacy class inside $N(T)$. Since $u_{*} \in \mathcal{D}_{*}$, it follows that $T u_{*} \subset \mathcal{D}_{*}$. Let $\tilde{X} \subset \tilde{M}=\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}$ be the 2 -torus given as the image of the map

$$
T \times T \rightarrow \tilde{M}, \quad\left(h_{1}, h_{2}\right) \mapsto\left(h_{1} u_{*}, h_{2} u_{*}, h_{1} h_{2} u_{*}, u_{*}, \ldots, u_{*}\right)
$$

and denote by $X$ its image in $M$. Let $\tilde{\omega}_{X}, \omega_{X}$ be the pull-backs of the quasiHamiltonian 2-forms on $\tilde{X}, X$. Since $T u_{*}=u_{*} T$, we have $\tilde{\Phi}(\tilde{X})=\Phi(X) \subset T u_{*}{ }^{r}$. Since the generator $x \in H^{3}(\mathrm{SU}(2), \mathbb{Z})$ pulls back to zero on this circle (for dimension reasons), the existence of a level $k$ pre-quantization of $M$ requires that $k \int_{X} \omega_{X} \in \mathbb{Z}$. Since the projection $\tilde{X} \rightarrow X$ is a 4 -fold covering, $\int_{X} \omega_{X}=\frac{1}{4} \int_{\tilde{X}} \tilde{\omega}_{X}$. Hence it is necessary that $k \int_{\tilde{X}} \tilde{\omega}_{X} \in 4 \mathbb{Z}$.

Let $\theta \in \Omega^{1}(T, \mathfrak{t})$ be the Maurer-Cartan form for $T$. From the general formula (12), and using $\left(h u_{*}\right)^{*} \theta^{L}=-h^{*} \theta,\left(h u_{*}\right)^{*} \theta^{R}=h^{*} \theta$, we obtain

$$
\tilde{\omega}_{X}=\frac{1}{2}\left(-h_{1}^{*} \theta \wedge h_{2}^{*} \theta+h_{1}^{*} \theta \wedge\left(h_{1} h_{2}\right)^{*} \theta-h_{2}^{*} \theta \wedge\left(h_{1} h_{2}\right)^{*} \theta\right)=\frac{1}{2} h_{1}^{*} \theta \wedge h_{2}^{*} \theta
$$

Writing elements of $T$ in the form $h=j\left(e^{2 \pi i v}\right)$, we may take $v \in[0,1]$ as the coordinate on $T \cong \mathbb{R} / \mathbb{Z}$. Since the lattice $\Lambda$ is generated by $2 \rho$, we find $h_{i}^{*} \theta=$ $2 d v_{i} \otimes \rho$, hence

$$
\tilde{\omega}_{X}=2\|\rho\|^{2} d v_{1} \wedge d v_{2}=d v_{1} \wedge d v_{2}
$$

integrates to 1 . This gives the condition $k \in 4 \mathbb{N}$.
3.3. Fixed point components. Suppose $M$ is a level $k$ pre-quantized quasiHamiltonian $\mathrm{SU}(2)$-space, and let $\mathcal{Q}(M) \in R_{k}(\mathrm{SU}(2))$ be its quantization. By [15, Theorem 9.5], the numbers $\mathcal{Q}(M)(t)$ with $t=t_{l}, l=0, \ldots, k$ are given as a sum of contributions from the fixed point manifolds of $t$ :

$$
\begin{equation*}
\mathcal{Q}(M)(t)=\sum_{F \subset M^{t}} \int_{F} \frac{\widehat{A}(F) \operatorname{Ch}\left(\mathcal{L}_{F}, t\right)^{1 / 2}}{D_{\mathbb{R}}\left(\nu_{F}, t\right)} \tag{13}
\end{equation*}
$$

The ingredients of the right hand side will be described below, and explicitly computed in the context of our main example (10). The quantizations of $\mathrm{SU}(2)$ conjugacy classes and of the double $\mathbf{D}(\mathrm{SO}(3)$ ) (viewed as a quasi-Hamiltonian $\mathrm{SU}(2)$-space) were computed in [15].

For the remainder of this section, we therefore focus on the case $h=0, s=r \geq 2$, i.e. $M=\left(\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}\right) / \Gamma$.
3.3.1. Fixed point sets of $M$. We need to determine the components $F \subset M^{t}$ of the fixed point manifold for $t=t_{l}, l=0, \ldots, k$, and describe various aspects of $F$ and its normal bundle $\nu_{F}$. Consider first a general regular element $t \in T^{\mathrm{reg}}$. Define the following two submanifolds of $\mathcal{D}_{*}$, labeled by the elements of the center $Z=\{e, c\}$ as follows:

$$
Y^{(e)}=\mathcal{D}_{*} \cap T=\left\{t_{*}, t_{*}^{-1}\right\}, \quad Y^{(c)}=T u_{*}
$$

Thus $Y^{(e)}$ is the fixed point set of $\operatorname{Ad}\left(t_{*}\right)$, while $Y^{(c)}$ consists of elements satisfying $\operatorname{Ad}\left(t_{*}\right)(g)=c g$. Note that both are $Z$-invariant. For $\gamma=\left(\gamma_{1}, \ldots, \gamma_{r}\right) \in \Gamma$, consider the $\Gamma$-invariant submanifold

$$
\tilde{F}^{(\gamma)}=Y^{\left(\gamma_{1}\right)} \times \cdots \times Y^{\left(\gamma_{r}\right)}
$$

and put $F^{(\gamma)}=\tilde{F}^{(\gamma)} / \Gamma$. Let $\mathrm{I}(\gamma)$ be the number of $\gamma_{i}$ 's that are equal to $c$. Then $\tilde{F}^{(\gamma)}$ is a disjoint union of $2^{r-\mathrm{I}(\gamma)}$ tori of dimension $\mathrm{I}(\gamma)$. Let $\varepsilon=(e, \ldots, e)$ denote the group unit in $\Gamma$. If $\gamma \neq \varepsilon$, then $\Gamma$ acts transitively on the set of components of $\tilde{F}^{(\gamma)}$. Hence $F^{(\gamma)}$ is a (connected) torus, and since $|\Gamma|=2^{r-1}$, it follows that the projection restricts to a $2^{1(\gamma)-1}$-fold covering on each component of $\tilde{F}^{(\gamma)}$. If $\gamma=\varepsilon$, $\tilde{F}^{(\varepsilon)}$ consists of $2^{r}$ points, and hence $F^{(\varepsilon)}$ consists of two points.

Proposition 3.3. The fixed point set of $t \in T^{\mathrm{reg}}$ in $M$ is

$$
M^{t}= \begin{cases}F^{(\varepsilon)} & \text { if } t \notin\left\{t_{*}, t_{*}^{-1}\right\} \\ \coprod_{\gamma \in \Gamma} F^{(\gamma)} & \text { if } t \in\left\{t_{*}, t_{*}^{-1}\right\}\end{cases}
$$

Proof. An element $\left(g_{1}, \ldots, g_{r}\right) \in \tilde{M}$ maps to a point in $M^{t}$ if and only if there exists $\gamma=\left(\gamma_{1}, \ldots, \gamma_{r}\right) \in \Gamma$ with $\operatorname{Ad}(t) g_{i}=g_{i} \gamma_{i}$, for $i=1, \ldots, r$. If $\gamma_{i}=e$, this condition gives $g_{i} \in T$, since $t$ is regular. If $\gamma_{i}=c$, the condition says that $\operatorname{Ad}\left(g_{i}^{-1}\right)(t)=\gamma_{i} t$. Since $t$ is regular, this happens if and only if $t \in\left\{t_{*}, t_{*}{ }^{-1}\right\}$, with $g_{i} \in N(T)$ representing the non-trivial Weyl group element.
3.3.2. The symplectic volume of the components of the fixed point set. Each $F^{(\gamma)} \subset$ $M^{t}$ is a quasi-Hamiltonian $T$-space, with moment map the restriction of $\Phi$. (See e.g. [14, Proposition 3.1].) In particular, they are symplectic.

Lemma 3.4. The symplectic volume of each component of $\tilde{F}^{(\gamma)}$ is equal to 1. Thus

$$
\operatorname{vol}\left(F^{(\gamma)}\right)=2^{1-\mathrm{I}(\gamma)}
$$

Proof. The construction from [3] associates to any quasi-Hamiltonian $G$-space (with $G$ compact, but possibly disconnected) a Liouville volume, in such a way that the volume of a fusion product is the product of the volumes. If $G=T$, so that the space is symplectic, the Liouville volume coincides with the symplectic volume. For a $G$-conjugacy class $\mathcal{C} \cong G / G_{g}$, the Liouville volume is given by the formula [3, Proposition 3.6]

$$
\operatorname{vol} \mathcal{C}=\left|\operatorname{det}_{\mathfrak{g}_{g}^{\perp}}\left(1-\operatorname{Ad}_{g}\right)\right|^{1 / 2} \frac{\operatorname{vol}(G)}{\operatorname{vol}\left(G_{g}\right)}
$$

involving the Riemannian volumes of $G$ and of the stabilizer group $G_{g}$. The spaces $Y^{(z)}$ for $z \in Z$ can be viewed as conjugacy classes for the group $N(T)$, of elements $t_{*}$ if $z=e$ and $u_{*}$ if $z=c$. Application of the formula gives

$$
\operatorname{vol}\left(Y^{(z)}\right)= \begin{cases}2 & \text { if } z=e \\ 1 & \text { if } z=c\end{cases}
$$

This is obvious for $z=e$, while for $z=c$ (so that $g=u_{*}, N(T)_{g}=\mathbb{Z}_{4}$ ) we have $\left|\operatorname{det}_{\mathfrak{t}}\left(1-\operatorname{Ad}_{u_{*}}\right)\right|^{1 / 2}=\sqrt{2}\left(\right.$ since $\operatorname{Ad}_{u_{*}}$ acts as -1 on $\left.\mathfrak{t}\right), \operatorname{vol}(N(T))=2 \operatorname{vol}(T)=$ $2\|\alpha\|=2 \sqrt{2}$, and $\operatorname{vol}\left(N(T)_{g}\right)=4$. It follows that

$$
\operatorname{vol}\left(\tilde{F}^{(\gamma)}\right)=\prod_{i=1}^{r} \operatorname{vol}\left(Y^{\left(\gamma_{i}\right)}\right)=2^{r-\mathrm{I}(\gamma)}
$$

Since the moment map for the quasi-Hamiltonian $N(T)$-space $\tilde{F}^{(\gamma)}$ takes values in $T$, this coincides with the symplectic volume. Since $2^{r-\mathrm{I}(\gamma)}$ is also the number of components of $\tilde{F}^{(\gamma)}$, it follows that each component has volume 1.
3.4. Fixed point contributions. In this Section, we assume that $M=\left(\mathcal{D}_{*} \times \cdots \times\right.$ $\left.\mathcal{D}_{*}\right) / \Gamma$ carries a level $k$ pre-quantization. Thus $k \in 2 \mathbb{N}$ if $r=2$ and $k \in 4 \mathbb{N}$ if $r>2$. Our aim is to compute the fixed point contributions to $\mathcal{Q}(M)(t)$, as described in formula (13), for $t=t_{l}, l=0, \ldots, k$.

If $t \neq t_{*}$, Proposition 3.3 shows that $M^{t}=F^{(\varepsilon)}$ consists of just two points, covered by the set $\tilde{M}^{t}=\tilde{F}^{(\varepsilon)}$ (consisting of $2^{r}$ points). The fixed point contribution of $F^{(\varepsilon)}$ is just that for $\tilde{F}^{(\varepsilon)}$, divided by $|\Gamma|=2^{r-1}$. Hence

$$
\mathcal{Q}(M)(t)=2^{1-r} \mathcal{Q}\left(\tilde{M}^{t}\right)=2^{1-r} \mathcal{Q}\left(\mathcal{D}_{*}\right)^{r}(t)
$$

with $\mathcal{Q}\left(\mathcal{D}_{*}\right)=\tau_{k / 2}[15$, Proposition 11.2].
If $t=t_{*}, \mathcal{Q}(M)\left(t_{*}\right)$ is a sum over the contributions from all $F^{(\gamma)}, \gamma \in \Gamma$. The contribution from $F^{(\varepsilon)}$ is $2^{1-r}\left(\mathcal{Q}\left(\mathcal{D}_{*}\right)\left(t_{*}\right)\right)^{r}$, as before. Calculation of the contributions from $F=F^{(\gamma)}, \gamma \neq \varepsilon$ requires more work:
Proposition 3.5. The contribution of the fixed point manifold $F=F^{(\gamma)}, \gamma \neq \varepsilon$ to $\mathcal{Q}(M)\left(t_{*}\right)$ is

$$
\int_{F} \frac{\widehat{A}(F) \operatorname{Ch}\left(\mathcal{L}_{F}, t_{*}\right)^{1 / 2}}{D_{\mathbb{R}}\left(\nu_{F}, t_{*}\right)}=2^{1-r}\left(\frac{k}{2}+1\right)^{\mathrm{I}(\gamma) / 2} \varphi^{(\gamma)}
$$

where the scalar $\varphi^{(\gamma)}=\mu_{F(\gamma)}\left(t_{*}\right) \in \mathrm{U}(1)$ is the action of $t_{*}$ on the pre-quantum line bundle over $F^{(\gamma)}$.
Proof. Since $F=F^{(\gamma)}$ is a torus, $\widehat{A}(F)=1$. To compute the $D_{\mathbb{R}}$-class, note that the normal bundle of $T u_{*}$ in $\mathcal{D}_{*}$ is an orientable real line bundle, hence it is trivializable. Consequently, the normal bundle $\nu_{\tilde{F}(\gamma)}$ to $\tilde{F}^{(\gamma)}$ in $\tilde{M}$ is trivializable, and thus the normal bundle $\nu_{F}=\nu_{\tilde{F}(\gamma)} / \Gamma$ to $F$ in $M$ is a flat Euclidean vector bundle of rank $2 r-\mathrm{I}(\gamma)$. The element $t_{*}$ acts by multiplication by -1 on the fibers of $\nu_{F}$, since $\operatorname{Ad}\left(t_{*}\right)$ has order 2 and cannot act trivially. By definition of the $D_{\mathbb{R}}$-class (see [2, Section 2.3] or [14, Section 5.3]), it follows that

$$
D_{\mathbb{R}}\left(\nu_{F}, t_{*}\right)=i^{\operatorname{rank}\left(\nu_{F}\right) / 2} \operatorname{det}_{\mathbb{R}}^{1 / 2}(1-(-1))=(2 i)^{r-\frac{\mathrm{l}(\gamma)}{2}}
$$

By [15, Proposition 9.3], the restriction $\left.T M\right|_{F}$ inherits a distinguished $\operatorname{Spin}_{c^{-}}$ structure (depending on the choice of level $k$ pre-quantization), equivariant for the action of $t_{*}$. The line bundle $\mathcal{L}_{F} \rightarrow F$ is the $\operatorname{Spin}_{c}$-line bundle associated to this Spin $_{c}$-structure, and

$$
\operatorname{Ch}\left(\mathcal{L}_{F}, t_{*}\right)^{1 / 2}=\sigma\left(\mathcal{L}_{F}\right)\left(t_{*}\right)^{1 / 2} \exp \left(\frac{1}{2} c_{1}\left(\mathcal{L}_{F}\right)\right)
$$

is the square root of its equivariant Chern character, with $\sigma\left(\mathcal{L}_{F}\right)\left(t_{*}\right) \in \mathrm{U}(1)$ the action of $t_{*}$ the $\operatorname{Spin}_{c}$-line bundle. As discussed in [2, Section 2.3] (see also [14, Section 5.3]), the sign of the square root is determined as follows. Since $\Phi$ restricts to a surjective map $F \rightarrow T$, the fixed point set $F$ meets $\Phi^{-1}(e)$. Pick any $x \in$ $F \cap \Phi^{-1}(e)$. Observe that $\omega$ is non-degenerate at points of $\Phi^{-1}(e)$, and choose a $t_{*}$-invariant compatible complex structure to view $T_{x} M$ as a Hermitian vector space. Let $A \in \mathrm{U}\left(T_{x} M\right)$ be the transformation defined by $t_{*}$ and $A^{1 / 2}$ its unique square root for which all eigenvalues are of the form $e^{i u}$ with $0 \leq u<\pi$. Then

$$
\sigma\left(\mathcal{L}_{F}\right)\left(t_{*}\right)^{1 / 2}=\varphi^{(\gamma)} \operatorname{det}_{\mathbb{C}}\left(A^{1 / 2}\right)
$$

Since $t_{*}$ acts trivially on $T_{m} F$ and as -1 on the normal bundle, the transformation $A^{1 / 2}$ acts trivially on $T_{x} F$ and as $i$ on the normal bundle. Thus $\operatorname{det}_{\mathbb{C}}\left(A^{1 / 2}\right)=$ $i^{r-I(\gamma) / 2}$, which cancels a similar factor in the expression for the $D_{\mathbb{R}}$-class.

It remains to find the integral $\int_{F} \exp \left(\frac{1}{2} c_{1}\left(\mathcal{L}_{F}\right)\right)$. To this end, we interpret $\mathcal{L}_{F}$ as a pre-quantum line bundle. By the same argument as in Property (g) of Section 3.1, (see also [15, Section 11.1]), the level $k$ pre-quantization and the canonical twisted $\operatorname{Spin}_{c}$-structure on $M$ combine to give an element of $H^{3}(\Phi ; \mathbb{Z})$ at level $2 k+4$. Since $H^{2}(M ; \mathbb{R})=0$, this element defines a pre-quantization at level $2 k+4$. Pull-back to $F$ defines a level $2 k+4$ pre-quantization of $F$, with $\mathcal{L}_{F}$ as the pre-quantum line bundle. Hence $c_{1}\left(\mathcal{L}_{F}\right)$ is the $2 k+4$-th multiple of the class of the symplectic form on $F$. It follows that

$$
\int_{F} \exp \left(\frac{1}{2} c_{1}\left(\mathcal{L}_{F}\right)\right)=(k+2)^{\mathrm{I}(\gamma)} \operatorname{vol}(F)=2^{1-\mathrm{I}(\gamma) / 2}\left(\frac{k}{2}+1\right)^{\mathrm{I}(\gamma) / 2}
$$

where we have used Lemma 3.4.
The phase factors $\varphi^{(\gamma)}$ depend on the choice of pre-quantization. Recall again that the set of pre-quantizations of a quasi-Hamiltonian $\mathrm{SU}(2)$-space is a torsor under the group of isomorphism classes of flat line bundles. In our case this is the group

$$
\operatorname{Tor}\left(H^{2}(M ; \mathbb{Z})\right) \cong \operatorname{Hom}(\Gamma, \mathrm{U}(1))
$$

The homomorphism $\psi: \Gamma \rightarrow \mathrm{U}(1)$ defines the flat line bundle $\tilde{M} \times_{\Gamma} \mathbb{C}_{\psi}$, where $\mathbb{C}_{\psi}$ is the 1-dimensional $\Gamma$-representation defined by $\psi$. Changing the pre-quantization by such a flat line bundle changes $\varphi^{(\gamma)}$ for $F=F^{(\gamma)}$ to $\psi(\gamma) \varphi^{(\gamma)}$. By Property (g) of Section 3.1, and since $H^{2}(M ; \mathbb{R})=0$, there is a distinguished pre-quantization at any level $k \in 4 \mathbb{N}$. Hence, the inequivalent pre-quantizations at level $k \in 4 \mathbb{N}$ are labeled by $\operatorname{Hom}(\Gamma, U(1))$.

Lemma 3.6. If $r \geq 3$ and $k \in 4 \mathbb{N}$, the phase factor for the pre-quantization labeled by $\psi \in \operatorname{Hom}(\Gamma, \mathrm{U}(1))$ is given by

$$
\varphi^{(\gamma)}=(-1)^{\frac{k}{4}(r-\mathrm{I}(\gamma) / 2)} \psi(\gamma)
$$

Proof. The phase factor $\varphi^{(\gamma)}$ for the distinguished pre-quantization at level 4 is given by $\operatorname{det}_{\mathbb{C}}(A)=(-1)^{r-I(\gamma) / 2}$, in the notation from the proof of Proposition 3.5. For the distinguished pre-quantization at level $k \in 4 \mathbb{N}$, we have to take the $\frac{k}{4}$-th power of this number, and changing the pre-quantization by $\psi$ we have to multiply by $\psi(\gamma)$.

If $r=2$, there are $|\Gamma|=2$ distinct pre-quantizations at all even levels $k \in 2 \mathbb{N}$, related by elements $\psi \in \operatorname{Hom}(\Gamma, \mathrm{U}(1))$. Aside from the discrete fixed point set $F^{(\varepsilon)}$, there is a single non-discrete fixed point component $F^{(\gamma)}$ of $t_{*}$, given by $\gamma=(c, c)$. The non-trivial homomorphism $\psi \in \operatorname{Hom}(\Gamma, \mathrm{U}(1)) \cong \mathbb{Z}_{2}$ satisfies $\psi(c, c)=-1$, hence the weight $\varphi^{(\gamma)}$ is equal to 1 for one of the pre-quantizations and -1 for the other.
3.5. Quantization of $M$. We are now ready to summarize our computation of $\mathcal{Q}(M)$ for $M=\left(\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}\right) / \Gamma$. Assuming that $k$ is even, recall that $\mathcal{D}_{*}$ has a unique pre-quantization at level $k$, and $\mathcal{Q}\left(\mathcal{D}_{*}\right)=\tau_{k / 2}$. Define an element

$$
\chi=\tau_{0}-\tau_{2}+\tau_{4}-\cdots+(-1)^{k / 2} \tau_{k} \in R_{k}(\mathrm{SU}(2))
$$

By the orthogonality relations for $R_{k}(\mathrm{SU}(2))$, this element satisfies $\chi\left(t_{*}\right)=\left(\frac{k}{2}+1\right)$ and $\chi(t)=0$ for $t=t_{l}, l \neq k / 2$. Hence we may write the sum over the fixed point contributions as follows:

$$
\mathcal{Q}(M)(t)=2^{1-r}\left(\tau_{k / 2}(t)^{r}+\chi(t) \sum_{\gamma \in \Gamma \backslash\{\varepsilon\}}\left(\frac{k}{2}+1\right)^{1(\gamma) / 2-1} \varphi^{(\gamma)}\right)
$$

Theorem 3.7. Consider the quasi-Hamiltonian $\mathrm{SU}(2)$-space $M=\left(\mathcal{D}_{*} \times \cdots \times \mathcal{D}_{*}\right) / \Gamma$ with $r \geq 2$ factors, where $\Gamma \subset Z^{r}$ consists of all $\gamma=\left(\gamma_{1}, \ldots, \gamma_{r}\right)$ with $\prod_{i=1}^{r} \gamma_{i}=e$.
(1) If $r \geq 3$, the space $M$ is pre-quantized at level $k$ if and only if $k \in 4 \mathbb{N}$. The different pre-quantizations are indexed by the elements $\psi \in \operatorname{Hom}(\Gamma, \mathrm{U}(1))$, and the corresponding level $k$ quantization is given by the formula,

$$
\mathcal{Q}_{\psi}(M)=2^{1-r}\left(\left(\tau_{k / 2}\right)^{r}+\chi \sum_{\gamma \in \Gamma \backslash\{\varepsilon\}} \psi(\gamma)\left(\frac{k}{2}+1\right)^{\frac{1(\gamma)}{2}-1}(-1)^{\frac{k}{4}\left(r-\frac{1(\gamma)}{2}\right)}\right)
$$

(2) If $r=2$, the space $M$ is pre-quantized at level $k$ if and only if $k \in 2 \mathbb{N}$. At any such level, there are two distinct pre-quantizations indexed by the action $\pm 1$ of $t_{*}$ on the pre-quantum line bundle over $F^{(\gamma)}$, for $\gamma=(c, c)$. The corresponding level $k$ quantizations of $M$ are

$$
\mathcal{Q}_{ \pm}(M)=\frac{1}{2}\left(\left(\tau_{k / 2}\right)^{2} \pm \chi\right)
$$

3.6. Multiplicity computations. Being elements of $R_{k}(\mathrm{SU}(2))$, the coefficients of $\mathcal{Q}(M)$ in its decomposition with respect to the basis $\tau_{0}, \ldots, \tau_{k}$ must be integers. In this Section, we will compute these multiplicities for small $r$.
3.6.1. $r=2$ factors. Assume $k \in 2 \mathbb{N}$, and let $\mathcal{Q}_{ \pm}(M)$ be the quantizations corresponding to the pre-quantizations labeled by $\pm 1$. The multiplication rules for level $k$ characters give

$$
\left(\tau_{k / 2}\right)^{2}=\tau_{0}+\tau_{2}+\ldots+\tau_{k}
$$

Hence, if $k \in 4 \mathbb{N}$ we obtain

$$
\begin{aligned}
& \mathcal{Q}_{+}(M)=\tau_{0}+\tau_{4}+\ldots+\tau_{k} \\
& \mathcal{Q}_{-}(M)=\tau_{2}+\tau_{6}+\ldots+\tau_{k-2}
\end{aligned}
$$

while for $k \in 4 \mathbb{N}-2$,

$$
\begin{aligned}
& \mathcal{Q}_{+}(M)=\tau_{0}+\tau_{4}+\ldots+\tau_{k-2} \\
& \mathcal{Q}_{-}(M)=\tau_{2}+\tau_{6}+\ldots+\tau_{k}
\end{aligned}
$$

3.6.2. $r=3$ factors. Let $\mathcal{Q}_{\psi}(M)$ denote the level $k \in 4 \mathbb{N}$ pre-quantization indexed by $\psi \in \operatorname{Hom}(\Gamma, \mathrm{U}(1))$. Since $r=3, \mathrm{I}(\gamma)=2$ for any $\gamma \neq \varepsilon$ and the quantization formula simplifies to:

$$
\mathcal{Q}_{\psi}(M)=\frac{1}{4}\left(\tau_{2 m}^{3}+\chi \sum_{\gamma \neq \varepsilon} \psi(\gamma)\right)
$$

For the trivial homomorphism $\psi=1$, we have $\sum_{\gamma \neq \varepsilon} \psi(\gamma)=3$, while for a nontrivial homomorphism $\psi \neq 1, \sum_{\gamma \neq \varepsilon} \psi(\gamma)=-1$. We have,

$$
\left(\tau_{k / 2}\right)^{3}=\tau_{0}+3 \tau_{2}+\ldots+\left(\frac{k}{2}+1\right) \tau_{k / 2}+\ldots+3 \tau_{k-2}+\tau_{k}
$$

We therefore obtain

$$
\begin{array}{rlrl}
\mathcal{Q}_{\psi}(M)= & \left(\tau_{0}+2 \tau_{4}+3 \tau_{8}+\ldots 3 \tau_{k-8}+2 \tau_{k-4}+\tau_{k}\right) & & \text { if } \psi=1 \\
& +\left(\tau_{6}+2 \tau_{10}+\ldots+2 \tau_{k-10}+\tau_{k-6}\right) & \\
\mathcal{Q}_{\psi}(M)= & \left(\tau_{0}+\tau_{4}+2 \tau_{8}+\ldots+3 \tau_{k-8}+2 \tau_{k-4}+\tau_{k}\right) & \\
& +\left(\tau_{2}+2 \tau_{6}+3 \tau_{10}+\ldots+3 \tau_{k-10}+2 \tau_{k-6}+\tau_{k-2}\right) & & \text { if } \psi \neq 1
\end{array}
$$

Note that the coefficients are symmetric about the midpoint $\frac{k}{2}$ of the interval $[0, k]$. In closed form, $\mathcal{Q}_{\psi}(M)=\sum_{j=0}^{k / 2} a_{2 j} \tau_{2 j}$, where

$$
a_{2 j}= \begin{cases}\frac{1}{4}\left(2 j+1+\left(4 \delta_{\psi, 1}-1\right)(-1)^{j}\right) & : 2 j \leq k / 2 \\ \frac{1}{4}\left(k-2 j+1+\left(4 \delta_{\psi, 1}-1\right)(-1)^{j}\right) & : 2 j \geq k / 2\end{cases}
$$

3.6.3. $r=4$ factors. If $r=4$ we have $|\Gamma|=8$. There is a unique element $\gamma^{\prime} \in \Gamma$ with $\mathrm{I}\left(\gamma^{\prime}\right)=4$, and $\mathrm{I}(\gamma)=2$ for $\gamma \neq \gamma^{\prime}, \varepsilon$. Hence we may write the quantization formula for levels $k \in 4 \mathbb{N}$ as:

$$
\mathcal{Q}_{\psi}(M)=\frac{1}{8}\left(\tau_{k / 2}^{4}+\left(\psi\left(\gamma^{\prime}\right)\left(\frac{k}{2}+1\right)+(-1)^{k / 4} \sum_{\mathrm{I}(\gamma)=2} \psi(\gamma)\right) \chi\right)
$$

One finds that there are 4 homomorphisms $\psi$ with $\sum_{\mathrm{I}_{( }(\gamma)=2} \psi(\gamma)=0, \psi\left(\gamma^{\prime}\right)=-1$ and 3 homomorphisms with $\sum_{\mathrm{I}(\gamma)=2} \psi(\gamma)=-2, \psi\left(\gamma^{\prime}\right)=1$. Of course, $\sum_{\mathrm{I}(\gamma)=2} \psi(\gamma)=$ $6, \psi\left(\gamma^{\prime}\right)=1$ for $\psi=1$. Therefore, we have

$$
\mathcal{Q}_{\psi}(M)= \begin{cases}\frac{1}{8}\left(\tau_{k / 2}^{4}+\left(6(-1)^{k / 4}+\left(\frac{k}{2}+1\right)\right) \chi\right) & : \psi=1 \\ \frac{1}{8}\left(\tau_{k / 2}^{4}-\left(\frac{k}{2}+1\right) \chi\right) & : \sum_{\mathrm{I}(\gamma)=2} \psi(\gamma)=0 \\ \frac{1}{8}\left(\tau_{k / 2}^{4}+\left(2(-1)^{k / 4+1}+\left(\frac{k}{2}+1\right)\right) \chi\right) & : \sum_{\mathrm{l}(\gamma)=2} \psi(\gamma)=-2\end{cases}
$$

with

$$
\left(\tau_{k / 2}\right)^{4}=\sum_{j=0}^{k / 2}\left(\frac{k}{2}+1-2 j^{2}+j k\right) \tau_{2 j}
$$

One may verify that the multiplicities of $\tau_{2 j}$ in $\mathcal{Q}_{\psi}(M)$ are integers, as required.

## 4. Fuchs-Schweigert

The formulas appearing in Theorem 3.7 may be rewritten in terms of the socalled $S$-matrix. For $z \in Z$, define $S_{m, l}^{(z)}$ by

$$
S_{m, l}^{(z)}= \begin{cases}1 & \text { if } z=c \\ \mathrm{~S}_{m, l} & \text { if } z=e\end{cases}
$$

In the terminology of [9], $S_{m, l}^{(z)}$ is the $S$-matrix of the orbit Lie algebra associated to the central element $z$. (This interpretation may seem obscure for $\mathrm{SU}(2)$, but becomes natural for higher rank groups.) Consider once again the space $M=\tilde{M} / \Gamma$ from (10). Recall that $\Gamma$ consists of elements $\gamma=\left(\gamma_{1}, \ldots, \gamma_{s+2 h}\right) \in Z^{s+2 h}$ such that $\prod_{j=1}^{s} \gamma_{j}=e$, and $\gamma_{j}=e$ for all $j \leq s$ with $\mathcal{C}_{j} \neq \mathcal{C}_{*}$. In particular $|\Gamma|=2^{2 h+r-1}$ if $r \geq 1$, while $|\Gamma|=2^{2 h}$ if $r=0$. To write the Fuchs-Schweigert formula, it is convenient to use the following notation. For $\gamma \in \Gamma$, let $\sum_{l}^{(\gamma)}$ denote the full sum $\sum_{l=0}^{k}$ if all $\gamma_{i}=e$, and consisting of the single term $l=\frac{k}{2}$ if at least one $\gamma_{i} \neq e$. (For higher rank groups, this becomes a sum over level $k$ weights that are fixed
under the action of all $\gamma_{i} \in Z$ on the set of level $k$ weights.) We will prove the following equivariant analogue to the Fuchs-Schweigert formula:
Theorem 4.1. Suppose the quasi-Hamiltonian $\mathrm{SU}(2)$-space

$$
M=\left(\mathcal{D}_{1} \times \cdots \times \mathcal{D}_{s} \times \mathbf{D}(\mathrm{SU}(2))^{h}\right) / \Gamma
$$

is pre-quantized at level $k$. Then

$$
\begin{equation*}
\mathcal{Q}(M)=\frac{1}{|\Gamma|} \sum_{\gamma \in \Gamma} \varphi^{\prime}(\gamma) \sum_{l}^{(\gamma)} \frac{S_{m_{1}, l}^{\left(\gamma_{1}\right)} \cdots S_{m_{s}, l}^{\left(\gamma_{s}\right)}}{\left(S_{0, l}\right)^{s+2 h}} \tilde{\tau}_{l} \tag{14}
\end{equation*}
$$

where $\varphi^{\prime}(\gamma) \in \mathrm{U}(1)$ are phase factors depending on the choice of pre-quantization, with $\varphi^{\prime}(\varepsilon)=1$.

An explicit description of the phase factors $\varphi^{\prime}(\gamma)$ will be given during the course of the proof.

Proof of Theorem 4.1. The space $M$ is a fusion product of the space $\widetilde{\left(\mathcal{C}_{*}\right)^{r}}$, conjugacy classes $\mathcal{D}_{j} \neq \mathcal{D}_{*}$, and $h$ factors of $\mathbf{D}(\mathrm{SO}(3))$ (viewed as a quasi-Hamiltonian $\mathrm{SU}(2)$-space). Since the fusion product in the basis $\tilde{\tau}_{m}$ is diagonalized, we may verify the formula separately for factors of these three types.

We begin with the case $h=0, s=r$, with $r \geq 3$ (thus necessarily $k \in 4 \mathbb{N}$ ). We re-write the right hand side of (14), separating the term $\gamma=\varepsilon$ from the sum over terms $\gamma \neq \varepsilon$. The right hand side of (14) becomes

$$
\begin{equation*}
\mathcal{Q}(M)=\frac{1}{|\Gamma|}\left(\varphi^{\prime}(\varepsilon) \sum_{l} \frac{\left(S_{k / 2, l}\right)^{r}}{\left(S_{0, l}\right)^{r}} \tilde{\tau}_{l}+\sum_{\gamma \neq \varepsilon} \varphi^{\prime}(\gamma) \frac{\left(S_{k / 2, k / 2}\right)^{r-\mathrm{I}(\gamma)}}{\left(S_{0, k / 2}\right)^{r}} \tilde{\tau}_{k / 2}\right) . \tag{15}
\end{equation*}
$$

The sum over $l$ is just $\left(\tau_{k / 2}\right)^{r}$. The element $\chi \in R_{k}(\mathrm{SU}(2))$ considered in Section 3.5 satisfies $\chi\left(t_{l}\right)=\left(\frac{k}{2}+1\right) \delta_{l, k / 2}$ for $l=0, \ldots, k$, hence

$$
\tilde{\tau}_{\frac{k}{2}}=\left(\frac{k}{2}+1\right)^{-1} \chi .
$$

Furthermore, by definition of the $S$-matrix,

$$
S_{0, k / 2}=\left(\frac{k}{2}+1\right)^{-\frac{1}{2}}, \quad S_{k / 2, k / 2}=\left(\frac{k}{2}+1\right)^{-\frac{1}{2}}(-1)^{\frac{k}{4}}
$$

Equation (15) becomes

$$
\mathcal{Q}(M)=\frac{1}{2^{r-1}}\left(\varphi^{\prime}(\varepsilon)\left(\tau_{k / 2}\right)^{r}+\sum_{\gamma \neq \varepsilon} \varphi^{\prime}(\gamma)(-1)^{\frac{k}{4}(r-\mathrm{I}(\gamma))}\left(\frac{k}{2}+1\right)^{\frac{\mathrm{I}(\gamma)}{2}-1} \chi\right)
$$

which agrees with Theorem 3.7 for $\varphi^{\prime}(\gamma)=\psi(\gamma)(-1)^{\frac{k 1(\gamma)}{8}}$.
The calculation is similar for the case $h=0, s=r=2, k \in 2 \mathbb{N}$. Here, $|\Gamma|=2$, and the generator $\gamma=(c, c) \in \Gamma$ has $\mathrm{I}(\gamma)=2$. We hence obtain

$$
\mathcal{Q}(M)=\frac{1}{2}\left(\varphi^{\prime}(e, e)\left(\tau_{k / 2}\right)^{r}+\varphi^{\prime}(c, c) \chi\right)
$$

which agrees with Theorem 3.7 if we put $\varphi^{\prime}(e, e)=1$, and $\varphi^{\prime}(c, c)= \pm 1$. If $h=0$ and $s=r=1, k \in 2 \mathbb{N}$, then $\Gamma=\{e\}$, and the formula becomes $\mathcal{Q}(M)=$ $\varphi^{\prime}(e) \tau_{k / 2}$, which is the correct expression for $\mathcal{Q}\left(\mathcal{D}_{*}\right)$ for $\varphi^{\prime}(e)=1$. Similarly, if $h=r=0, s=1$ so that $M$ is a conjugacy class $\mathcal{D}_{j} \neq \mathcal{D}_{*}$, the formula reduces to $\mathcal{Q}(M)=\tau_{m_{j}}=\mathcal{Q}\left(\mathcal{D}_{j}\right)$.

Consider finally the case $h=1, s=0$ so that $M=\mathbf{D}(\mathrm{SO}(3))$. Pre-quantizability of this space requires $k \in 2 \mathbb{N}$, and as shown in [15] the distinct pre-quantizations
are indexed by $\varphi \in \operatorname{Hom}(\Gamma, \mathrm{U}(1))$, with $\Gamma=Z \times Z$. Separating off the term $(e, e)$, (14) becomes

$$
\mathcal{Q}(M)=\frac{1}{4}\left(\varphi^{\prime}(\epsilon) \sum_{l} \frac{1}{S_{0, l}^{2}} \tilde{\tau}_{l}+\sum_{\gamma \neq(e, e)} \varphi^{\prime}(\gamma) \frac{1}{S_{0, k / 2}^{2}} \tilde{\tau}_{k / 2}\right) .
$$

We have $\frac{1}{S_{0, k / 2}^{2}} \tilde{\tau}_{k / 2}=\chi$, and

$$
\mathcal{Q}(\mathbf{D}(\mathrm{SU}(2)))=\sum_{m} \tau_{m}^{2}=\sum_{l, m} \frac{S_{m, l}^{2}}{S_{0, l}^{2}} \tilde{\tau}_{l}=\sum_{l} \frac{1}{S_{0, l}^{2}} \tilde{\tau}_{l}
$$

where we use the symmetry and orthogonality of the $S$-matrix. Thus the formula may be re-written

$$
\mathcal{Q}(M)=\frac{1}{4}\left(\varphi^{\prime}(\epsilon) \mathcal{Q}(\mathbf{D}(\mathrm{SU}(2)))+\sum_{\gamma \neq(e, e)} \varphi^{\prime}(\gamma) \chi\right)
$$

This agrees with the formula for $\mathcal{Q}(\mathbf{D}(\mathrm{SO}(3))$ given in [15, Section 11.4] if one puts $\varphi^{\prime}(e, e)=1$ and $\varphi^{\prime}(\gamma)=(-1)^{k / 2} \varphi(\gamma)$ for $\gamma \neq(e, e)$.

By combining this result with the 'quantization commutes with reduction' theorem for quasi-Hamiltonian spaces [15, Theorem 10.1], and since the coefficient of $\tau_{0}$ in $\tilde{\tau}_{l}$ is $S_{0, l}^{2}$, we obtain the Fuchs-Schweigert formula [9] for the $\mathrm{SO}(3)$ moduli space $\mathcal{M}\left(\Sigma, \mathcal{C}_{1}, \ldots, \mathcal{C}_{s}\right)$, where $\Sigma$ is of genus $h$ with $s$ boundary components. Recall that this moduli space has up two 2 connected components, of the form $M / / \mathrm{SU}(2)$ for suitable choice of lifts $\mathcal{D}_{j}$. We have,

$$
\begin{equation*}
\mathcal{Q}(M / / \mathrm{SU}(2))=\frac{1}{|\Gamma|} \sum_{\gamma \in \Gamma} \varphi^{\prime}(\gamma) \sum_{l}^{(\gamma)} \frac{S_{m_{1}, l}^{\left(\gamma_{1}\right)} \cdots S_{m_{s}, l}^{\left(\gamma_{s}\right)}}{\left(S_{0, l}\right)^{s+2 h-2}} \tag{16}
\end{equation*}
$$

Remark 4.2. The above Fuchs-Schweigert type formula computes the quantization of the moduli space of $\mathrm{SO}(3)$-bundles interpreted as the index of a pre-quantum line bundle, while the original conjecture in [9] concerns the dimension of the space of conformal blocks. It is expected that, just as in the case of simply-connected groups, the space of conformal blocks can be re-interpreted as the space of holomorphic sections, and that a Kodaira vanishing result can further identify its dimension with the index considered here. We are not aware of a reference addressing such questions in generality for non-simply connected groups.

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