# THE EXPANSION IN ULTRASPHERICAL POLYNOMIALS: A SIMPLE PROCEDURE FOR THE FAST COMPUTATION OF THE ULTRASPHERICAL COEFFICIENTS 

ENRICO DE MICHELI<br>Consiglio Nazionale delle Ricerche<br>Via De Marini, 6-16149 Genova, Italy<br>E-mail: enrico.demicheli@cnr.it<br>GIOVANNI ALBERTO VIANO<br>Dipartimento di Fisica - Università di Genova, Istituto Nazionale di Fisica Nucleare - Sezione di Genova, Via Dodecaneso, 33-16146 Genova, Italy<br>E-mail: viano@ge.infn.it


#### Abstract

We present a simple and fast algorithm for the computation of the coefficients of the expansion of a function $f(\cos u)$ in ultraspherical (Gegenbauer) polynomials. We prove that these coefficients coincide with the Fourier coefficients of an Abel-type transform of the function $f(\cos u)$. This allows us to fully exploit the computational efficiency of the Fast Fourier Transform, computing the first $N$ ultraspherical coefficients in just $\mathcal{O}\left(N \log _{2} N\right)$ operations.


## 1. Introduction

In this paper we investigate the generalization to ultraspherical polynomials (also known as Gegenbauer polynomials) of the results obtained in a previous paper [5] regarding the efficient computation of the coefficients of Legendre expansions.

Ultraspherical expansions play a relevant role in various subjects of applied and computational mathematics. These expansions have been used successfully for the solution of linear [8] and nonlinear [6, 7, 15] differential equations, of integral equation [16, and in spectral methods for partial differential equations 3, 19. Gegenbauer filtering (for the suppression of the Gibbs phenomenon) [10, 11, 12] has been proved to provide an exponentially convergent approximation (in the maximum norm) of a piecewise analytic function, starting from the Fourier partial sum of the function itself. The related Gegenbauer reconstruction method has found natural application in the image segmentation problem, in particular of MRI images [2, 13].

[^0]One of the limits of the application of ultraspherical expansions is the high cost of computing the expansion coefficients, this question becoming particularly critical in the case of multivariate functions.

In this paper we give an efficient procedure for computing these coefficients. We first obtain a Dirichlet-Murphy-type integral representation of the ultraspherical polynomials $P_{n}^{(d)}(\cos u)$ of degree $n$ and order $d$. Then we prove that the coefficients of the ultraspherical expansion of a function $f(\cos u)(u \in[0, \pi])$ coincide with the Fourier coefficients (restricted to nonnegative index) of an Abel-type transform of the function $f$.

These results produce straightforwardly an algorithm for the computation of ultraspherical coefficients which is very simple and very fast. In fact, the first $N$ ultraspherical coefficients are obtained in only $\mathcal{O}\left(N \log _{2} N\right)$ operations by a single Fast Fourier Transform of the Abel-type integral function, the latter being easily computable by standard quadrature techniques [14.

Finally, in the Appendix the results obtained for the ultraspherical polynomials are given also in terms of Gegenbauer polynomials $C_{n}^{(\lambda)}(x)$ with $x \in[-1,1]$.

## 2. Connection between Ultraspherical expansions and Fourier series

In dimension $d$, the expansion in ultraspherical polynomials of a function $f=$ $f(\cos u)(u \in[0, \pi])$ reads:

$$
\begin{equation*}
f(\cos u)=\frac{1}{2^{(d-2)} \pi^{\frac{(d-1)}{2}} \Gamma\left(\frac{d-1}{2}\right)} \sum_{n=0}^{\infty} \frac{\Gamma(n+d-2)}{n!}\left(n+\frac{d-2}{2}\right) a_{n}^{(d)} P_{n}^{(d)}(\cos u) \tag{2.1}
\end{equation*}
$$

the coefficients $a_{n}^{(d)}$ being defined by [4]:

$$
\begin{equation*}
a_{n}^{(d)}=\frac{2 \pi^{\frac{(d-1)}{2}}}{\Gamma\left(\frac{d-1}{2}\right)} \int_{0}^{\pi} f(\cos u) P_{n}^{(d)}(\cos u)(\sin u)^{(d-2)} \mathrm{d} u \quad\left(n \in \mathbb{N}_{0}\right) \tag{2.2}
\end{equation*}
$$

where $P_{n}^{(d)}(\cos u)$ denotes the ultraspherical polynomial of degree $n$ and order $d$, which is defined by the integral representation [9, 18]:

$$
\begin{equation*}
P_{n}^{(d)}(\cos u)=\frac{\Gamma\left(\frac{d-1}{2}\right)}{\sqrt{\pi} \Gamma\left(\frac{d-2}{2}\right)} \int_{0}^{\pi}(\cos u+\mathrm{i} \sin u \cos \eta)^{n}(\sin \eta)^{(d-3)} \mathrm{d} \eta \quad\left(n \in \mathbb{N}_{0}\right) . \tag{2.3}
\end{equation*}
$$

Our goal now is to show that the coefficients $a_{n}^{(d)}$ are the Fourier coefficients of a suitable Abel-type transform of $f$. To this end, we first prove the following proposition (see also [4]).

Proposition 1. The following integral representation of the ultraspherical polynomials $P_{n}^{(d)}(\cos u)$ holds:
$P_{n}^{(d)}(\cos u)=\frac{(-\mathrm{i})^{(d-2)}}{\sqrt{\pi}} \frac{\Gamma\left(\frac{d-1}{2}\right)}{\Gamma\left(\frac{d-2}{2}\right)} \frac{1}{(\sin u)^{(d-3)}} \int_{u}^{2 \pi-u} e^{\mathrm{i}\left(n+\frac{d-2}{2}\right) t}[2(\cos u-\cos t)]^{\frac{d-4}{2}} \mathrm{~d} t$.
Proof. In the integral representation (2.3) of $P_{n}^{(d)}(\cos u)$ substitute to $\eta$ the complex integration variable $\tau$ defined by

$$
\begin{equation*}
e^{\mathrm{i} \tau}=\cos u+\mathrm{i} \sin u \cos \eta \tag{2.5}
\end{equation*}
$$

It can be checked that

$$
\begin{equation*}
2 e^{\mathrm{i} \tau}(\cos \tau-\cos u)=\left(e^{\mathrm{i} \tau}-e^{\mathrm{i} u}\right)\left(e^{\mathrm{i} \tau}-e^{-\mathrm{i} u}\right)=\sin ^{2} u \sin ^{2} \eta \tag{2.6}
\end{equation*}
$$

Now, since $e^{\mathrm{i} \tau} \mathrm{d} \tau=-\sin u \sin \eta \mathrm{~d} \eta$, the integrand on the r.h.s. of Eq. (2.3) can be written as follows:

$$
\begin{align*}
& -(\sin u)^{-(d-3)} e^{\mathrm{i}(n+1) \tau}\left[\left(e^{\mathrm{i} \tau}-e^{\mathrm{i} u}\right)\left(e^{\mathrm{i} \tau}-e^{-\mathrm{i} u}\right)\right]^{\frac{d-4}{2}} \mathrm{~d} \tau  \tag{2.7}\\
& \quad=-(\sin u)^{-(d-3)} e^{\mathrm{i}\left(n+\frac{d-2}{2}\right) \tau}[2(\cos \tau-\cos u)]^{\frac{d-4}{2}} \mathrm{~d} \tau
\end{align*}
$$

In order to determine the integration path, consider an intermediate step where $e^{\mathrm{i} \tau}$ is chosen as the integration variable; the original path (corresponding to $\eta \in[0, \pi]$ ) is the (oriented) linear segment $\delta_{0}(u)$ starting at $e^{\mathrm{i} u}$ and ending at $e^{-\mathrm{i} u}$. Since (as shown by (2.7)) the integrand is an analytic function of $e^{\mathrm{i} \tau}$ in the disk $\left|e^{\mathrm{i} \tau}\right|<1$ (since $n \in \mathbb{N}$ ), the integration path $\delta_{0}(u)$ can be replaced by the circular path $\delta_{+}(u)=\left\{e^{\mathrm{i} \tau} ; \tau=t, u \leqslant t \leqslant 2 \pi-u\right\}$ (see Fig. (1). Moreover, by using the fact that $\left[\left(e^{\mathrm{i} \tau}-e^{\mathrm{i} u}\right)\left(e^{\mathrm{i} \tau}-e^{-\mathrm{i} u}\right)\right]^{\frac{d-4}{2}}$ is positive for $e^{\mathrm{i} \tau} \in \delta_{0}(u) \cup \mathbb{R}$ and therefore at $e^{\mathrm{i} \tau}=e^{\mathrm{i} \pi}$, we conclude from the left equality in (2.6) that in the r.h.s. of (2.7) the following specification holds (for $\tau=t ; u \leqslant t \leqslant 2 \pi-u$ ):

$$
\begin{equation*}
[2(\cos t-\cos u)]^{\frac{d-4}{2}}=(-\mathrm{i})^{d-4}[2(\cos u-\cos t)]^{\frac{d-4}{2}} \tag{2.8}
\end{equation*}
$$

Finally, by taking into account the latter expression, the integral representation (2.3) can then be replaced by the integral representation (2.4).

We can now prove the following theorem (see also [4]).
Theorem 2. The ultraspherical coefficients $\left\{a_{n}^{(d)}\right\}_{n=0}^{\infty}$ (see (2.2)) coincide with the Fourier coefficients (restricted to the nonnegative integer) of the form:

$$
\begin{equation*}
a_{n}^{(d)}=\int_{-\pi}^{\pi} \widehat{f}^{(d)}(t) e^{\mathrm{i} n t} \mathrm{~d} t \quad\left(n \in \mathbb{N}_{0}\right) \tag{2.9}
\end{equation*}
$$

where:

$$
\begin{equation*}
\widehat{f}^{(d)}(t)=[-\mathrm{i} \varepsilon(t)]^{d-2} \frac{2 \pi^{\frac{d-2}{2}}}{\Gamma\left(\frac{d-2}{2}\right)} e^{\mathrm{i}\left(\frac{d-2}{2}\right) t} \int_{0}^{t} f(\cos u)[2(\cos u-\cos t)]^{\frac{d-4}{2}} \sin u \mathrm{~d} u \tag{2.10}
\end{equation*}
$$

$\varepsilon(t)$ being the sign function.
Proof. By plugging representation (2.4) into formula (2.2) we have:

$$
\begin{align*}
a_{n}^{(d)}= & (-\mathrm{i})^{d-2} \frac{2 \pi^{\frac{d-2}{2}}}{\Gamma\left(\frac{d-2}{2}\right)} \int_{0}^{\pi} \mathrm{d} u f(\cos u) \sin u \int_{u}^{2 \pi-u} e^{\mathrm{i}\left(n+\frac{d-2}{2}\right) t}[2(\cos u-\cos t)]^{\frac{d-4}{2}} \mathrm{~d} t  \tag{2.11}\\
= & (-\mathrm{i})^{d-2} \frac{2 \pi^{\frac{d-2}{2}}}{\Gamma\left(\frac{d-2}{2}\right)}\left\{\int_{0}^{\pi} \mathrm{d} t e^{\mathrm{i}\left(n+\frac{d-2}{2}\right) t} \int_{0}^{t} \mathrm{~d} u f(\cos u) \sin u[2(\cos u-\cos t)]^{\frac{d-4}{2}}\right. \\
& \left.+\int_{\pi}^{2 \pi} \mathrm{~d} t e^{\mathrm{i}\left(n+\frac{d-2}{2}\right) t} \int_{0}^{2 \pi-t} \mathrm{~d} u f(\cos u) \sin u[2(\cos u-\cos t)]^{\frac{d-4}{2}}\right\}
\end{align*}
$$

Now, changing the variables $(t, u) \rightarrow(t+2 \pi, u)$ in the second term inside the parentheses on the rightmost side of (2.11), the latter becomes:

$$
\begin{equation*}
e^{\mathrm{i} \pi(d-2)} \int_{-\pi}^{0} \mathrm{~d} t e^{\mathrm{i}\left(n+\frac{d-2}{2}\right) t} \int_{0}^{-t} \mathrm{~d} u f(\cos u) \sin u[2(\cos u-\cos t)]^{\frac{d-4}{2}} \tag{2.12}
\end{equation*}
$$



Figure 1. Integration path for evaluating the integral representation (2.4) of the ultraspherical polynomials.

Finally, from (2.11) and (2.12) we have:

$$
\begin{align*}
a_{n}^{(d)}= & (-\mathrm{i})^{d-2} \frac{2 \pi^{\frac{d-2}{2}}}{\Gamma\left(\frac{d-2}{2}\right)}\left\{\int_{0}^{\pi} \mathrm{d} t e^{\mathrm{i} n t}\left[e^{\mathrm{i} \frac{d-2}{2} t} \int_{0}^{t} \mathrm{~d} u f(\cos u) \sin u[2(\cos u-\cos t)]^{\frac{d-4}{2}}\right]\right.  \tag{2.13}\\
& \left.+e^{\mathrm{i} \pi(d-2)} \int_{-\pi}^{0} \mathrm{~d} t e^{\mathrm{i} n t}\left[e^{\mathrm{i} \frac{d-2}{2} t} \int_{0}^{t} \mathrm{~d} u f(\cos u) \sin u[2(\cos u-\cos t)]^{\frac{d-4}{2}}\right]\right\} \\
= & \int_{-\pi}^{\pi} \widehat{f}^{(d)}(t) e^{\mathrm{i} n t} \mathrm{~d} t
\end{align*}
$$

with $\widehat{f}^{(d)}(t)$ given by (2.10).
It is easy to check that the $2 \pi$-periodic $\widehat{f}^{(d)}(t)$ function enjoys the following symmetry properties:

$$
\begin{equation*}
\widehat{f}^{(d)}(t)=(-1)^{d} e^{\mathrm{i}(d-2) t} \widehat{f}^{(d)}(-t) \quad(t \in \mathbb{R}) \tag{2.14}
\end{equation*}
$$

and, consequently,

$$
\begin{equation*}
a_{n}^{(d)}=(-1)^{d} a_{-(n+d-2)}^{(d)} \quad(n \in \mathbb{Z}) \tag{2.15}
\end{equation*}
$$

## Appendix A. From ultraspherical to Gegenbauer polynomials

A.1. The Gegenbauer polynomials $\boldsymbol{C}_{\boldsymbol{n}}^{(\boldsymbol{\lambda})}$. For the convenience of the reader we summarize hereafter the main properties of the Gegenbauer polynomials $C_{n}^{(\lambda)}(x)$, which are strictly related to the ultraspherical polynomials $P_{n}^{(d)}(x)$ that we introduced in Section 2, The following formulae are listed in Ref. 20 and can also be found in refs. [1, Chapter 22] and [8, 11, 17].

The Gegenbauer polynomials of order $\lambda$ can be defined in terms of their generating function:

$$
\begin{equation*}
\sum_{n=0}^{\infty} C_{n}^{(\lambda)}(x) t^{n}=\left(1-2 x t+t^{2}\right)^{-\lambda} \tag{A.1}
\end{equation*}
$$

From (A.1) we see that the Legendre polynomials $P_{n}(x)$ are the particular case of Gegenbauer polynomials with $\lambda=\frac{1}{2}$, i.e.: $C_{n}^{\left(\frac{1}{2}\right)}(x)=P_{n}(x)$.
They satisfy the recurrence relation:

$$
\begin{align*}
& n C_{n}^{(\lambda)}(x)=2(n+\lambda-1) x C_{n-1}^{(\lambda)}(x)-(n+2 \lambda-2) C_{n-2}^{(\lambda)}(x) \\
& C_{0}^{(\lambda)}(x)=1, \quad C_{1}^{(\lambda)}(x)=2 \lambda x \tag{A.2}
\end{align*}
$$

In terms of Gaussian hypergeometric function they can be written as:

$$
\begin{equation*}
C_{n}^{(\lambda)}(x)=\frac{2^{1-2 \lambda} \sqrt{\pi} \Gamma(n+2 \lambda)}{n!\Gamma(\lambda)}{ }_{2} F_{1}\left(-n, 2 \lambda+n ; \lambda+\frac{1}{2} ; \frac{1-x}{2}\right) . \tag{A.3}
\end{equation*}
$$

They can be written explicitly as

$$
\begin{equation*}
C_{n}^{(\lambda)}(x)=\sum_{k=0}^{\lfloor n / 2\rfloor}(-1)^{k} \frac{\Gamma(\lambda+n-k)}{k!(n-2 k)!\Gamma(\lambda)}(2 x)^{n-2 k}, \tag{A.4}
\end{equation*}
$$

and have the following integral representation:

$$
\begin{equation*}
C_{n}^{(\lambda)}(\cos u)=\frac{2^{1-2 \lambda} \Gamma(n+2 \lambda)}{n![\Gamma(\lambda)]^{2}} \int_{0}^{\pi}(\cos u+\mathrm{i} \sin u \cos \eta)^{n}(\sin \eta)^{2 \lambda-1} \mathrm{~d} \eta \tag{A.5}
\end{equation*}
$$

The Gegenbauer polynomials can be computed by the Rodrigues formula:

$$
\begin{equation*}
C_{n}^{(\lambda)}(x)=\frac{(-2)^{n}}{n!} \frac{\Gamma(n+\lambda) \Gamma(n+2 \lambda)}{\Gamma(\lambda) \Gamma(2 n+2 \lambda)}\left(1-x^{2}\right)^{-\lambda+\frac{1}{2}} \frac{\mathrm{~d}^{n}}{\mathrm{~d} x^{n}}\left[\left(1-x^{2}\right)^{n+\lambda-\frac{1}{2}}\right] \tag{A.6}
\end{equation*}
$$

which follows by induction from

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} x} C_{n}^{(\lambda)}(x)=2 \lambda C_{n-1}^{(\lambda+1)}(x) \tag{A.7}
\end{equation*}
$$

For fixed $\lambda$, the Gegenbauer polynomials are orthogonal on the interval $[-1,1]$ with respect to the weight function:

$$
\begin{equation*}
w^{(\lambda)}(x)=\left(1-x^{2}\right)^{\lambda-\frac{1}{2}}, \tag{A.8}
\end{equation*}
$$

that is, for $n \neq m$ :

$$
\begin{equation*}
\int_{-1}^{1} C_{n}^{(\lambda)}(x) C_{m}^{(\lambda)}(x)\left(1-x^{2}\right)^{\lambda-\frac{1}{2}} \mathrm{~d} x=0 \tag{A.9}
\end{equation*}
$$

and are normalized by:

$$
\begin{equation*}
\int_{-1}^{1}\left[C_{n}^{(\lambda)}(x)\right]^{2}\left(1-x^{2}\right)^{\lambda-\frac{1}{2}} \mathrm{~d} x=\frac{\pi 2^{1-2 \lambda} \Gamma(n+2 \lambda)}{n!(n+\lambda)[\Gamma(\lambda)]^{2}} \tag{A.10}
\end{equation*}
$$

A.2. Relation between Gegenbauer polynomials $C_{n}^{(\lambda)}$ and ultraspherical polynomials $\boldsymbol{P}_{\boldsymbol{n}}^{(d)}$. Comparing the integral representations (2.4) and (A.5), it is easily seen that the relation between ultraspherical $P_{n}^{(d)}(\cos u)$ and Gegenbauer polynomials $C_{n}^{(\lambda)}(\cos u)$ is:

$$
\begin{align*}
C_{n}^{(\lambda)}(\cos u) & =\frac{\Gamma(n+2 \lambda)}{n!\Gamma(2 \lambda)} P_{n}^{(2 \lambda+2)}(\cos u)  \tag{A.11a}\\
P_{n}^{(d)}(\cos u) & =\frac{n!\Gamma(d-2)}{\Gamma(n+d-2)} C_{n}^{\left(\frac{d-2}{2}\right)}(\cos u) \tag{A.11b}
\end{align*}
$$

Then, rephrasing Proposition 1 in terms of Gegenbauer polynomials, from (2.4) and (A.11) it follows that the polynomials $C_{n}^{(\lambda)}(\cos u)(u \in[0, \pi])$ have the following integral representation:
$C_{n}^{(\lambda)}(\cos u)=\frac{\Gamma(2 \lambda+n) \Gamma\left(\lambda+\frac{1}{2}\right)}{\sqrt{\pi} n!\Gamma(2 \lambda) \Gamma(\lambda)} \frac{(-\mathrm{i})^{2 \lambda}}{(\sin u)^{2 \lambda-1}} \int_{u}^{2 \pi-u} e^{\mathrm{i}(n+\lambda) t}[2(\cos u-\cos t)]^{\lambda-1} \mathrm{~d} t$.
A.3. Expansions in Gegenbauer polynomials $\boldsymbol{C}_{\boldsymbol{n}}^{(\boldsymbol{\lambda})}$. The Gegenbauer expansion of a function $f(x)$, defined in $x \in[-1,1]$, reads [11]:

$$
\begin{equation*}
f(x)=\sum_{n=0}^{\infty} f_{n}^{(\lambda)} C_{n}^{(\lambda)}(x) \quad(x \in[-1,1] ; \lambda>0) \tag{A.13}
\end{equation*}
$$

where the Gegenbauer coefficients are given by

$$
\begin{equation*}
f_{n}^{(\lambda)}=\frac{n!\Gamma(2 \lambda) \Gamma(\lambda)(n+\lambda)}{\sqrt{\pi} \Gamma(n+2 \lambda) \Gamma\left(\lambda+\frac{1}{2}\right)} \int_{-1}^{1} f(x) C_{n}^{(\lambda)}(x)\left(1-x^{2}\right)^{\left(\lambda-\frac{1}{2}\right)} \mathrm{d} x \quad\left(n \in \mathbb{N}_{0}\right) \tag{A.14}
\end{equation*}
$$

Now, from (A.14), (A.11) and using definition (2.2) of the coefficients $a_{n}^{(d)}$ we see that the Gegenbauer coefficients are related to the ultraspherical coefficients by

$$
\begin{equation*}
f_{n}^{(\lambda)}=\frac{\Gamma(\lambda)(n+\lambda)}{2 \pi^{(\lambda+1)}} a_{n}^{(2 \lambda+2)} \tag{A.15}
\end{equation*}
$$

We can now rewrite Theorem 2 in terms of Gegenbauer coefficients $f_{n}^{(\lambda)}$.
Theorem 2'. The coefficients $\left\{\phi_{n}^{(\lambda)}\right\}_{n=0}^{\infty}$, defined as $\phi_{n}^{(\lambda)} \doteq \frac{f_{n}^{(\lambda)}}{(n+\lambda)}$ (see (A.14)) coincide with the following Fourier coefficients (restricted to the nonnegative integer):

$$
\begin{equation*}
\phi_{n}^{(\lambda)} \doteq \frac{f_{n}^{(\lambda)}}{(n+\lambda)}=\int_{-\pi}^{\pi} \widehat{\phi}^{(\lambda)}(t) e^{\mathrm{i} n t} \mathrm{~d} t \quad\left(n \in \mathbb{N}_{0}\right) \tag{A.16}
\end{equation*}
$$

where:

$$
\begin{equation*}
\widehat{\phi}^{(\lambda)}(t)=\frac{[-\mathrm{i} \varepsilon(t)]^{2 \lambda}}{\pi} e^{\mathrm{i} \lambda t} \int_{\cos t}^{1} f(x)[2(x-\cos t)]^{\lambda-1} \mathrm{~d} x \tag{A.17}
\end{equation*}
$$

$\varepsilon(t)$ being the sign function.

## References

[1] M. Abramowitz, I. Stegun, Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables, New York, Dover, 1965.
[2] R. Archibald, K. Chen, A. Gelb and R. Renaut, Improving tissue segmentation of human brain MRI through preprocessing by the Gegenbauer reconstruction method, NeuroImage 20, 489-502, (2003)
[3] G. Ben-Yu, Gegenbauer approximation and its applications to differential equations with rough asymptotic behaviors at infinity, App. Num. Maths. 38, 403-425, (2001).
[4] J. Bros and G. A. Viano, Connection between the harmonic analysis on the sphere and the harmonic analysis on the one-sheeted hyperboloid: an analytic continuation viewpoint III, Forum Math. 9, 165-191, (1997).
[5] E. De Micheli and G. A. Viano, A new and efficient method for the computation of Legendre coefficients, submitted to Math. Comp., (2011).
[6] H. G. Denman and J. E. Howard, Application of ultraspherical polynomials to nonlinear oscillations. I. Free oscillations of the pendulum, Quart. Appl. Math. 21, 325-330, (1964).
[7] H. G. Denman and Y. K. Liu, Application of ultraspherical polynomials to nonlinear oscillations. II. Free oscillations, Quart. Appl. Math. 22, 273-291, (1965).
[8] D. Elliott, The expansion of functions in Ultraspherical polynomials, J. Australian Math. Soc. 1, 428-438, (1960).
[9] J. Faraut, Analyse Harmonique et Fonctions Spéciales, Ecole d'Été d'Analyse Harmonique de Tunis, (1984).
[10] A. Gelb and J. Tanner, Robust reprojection methods for the resolution of the Gibbs phenomenon, Appl. Comput. Harmon. Anal. 20, 3-25, 2006.
[11] D. Gottlieb, Chi-Wang Shu, A. Solomonoff and H. Vandeven, On the Gibbs phenomenon I: recovering exponential accuracy from the Fourier partial sum of a nonperiodic analytic function, J. Comp. Appl. Math. 43, 81-98, (1992).
[12] D. Gottlieb and C.-W. Shu, On the Gibbs phenomenon and its resolution, SIAM Rev. 39, 644-668, (1997).
[13] X. Huang and W. Chen, A fast algorithm to reduce Gibbs ringing artifact in MRI, Proc. of 2005 IEEE Engineering in Medicine and Biology 27th Annual Conference, Shanghai, China, (2005).
[14] G. Monegato and L. Scuderi, Numerical integration of functions with boundary singularities, J. Comp. Appl. Math., Special Issue: "Numerical Evaluation of Integrals", D. Laurie, R. Cools (eds.), 112 (1999), 201-214.
[15] S. C. Sinha and P. Srinivasan, Application of ultraspherical polynomials to non-linear autonomous systems, J. Sound Vibration 18, 55-60, (1971).
[16] K. N. Srivastava, A class of integral equations involving ultraspherical polynomials as kernel, Proc. Amer. Math. Soc. 14, 932-940, (1963).
[17] G. Szegö, Orthogonal Polynomials, Providence, RI: American Mathematical Society, 1975.
[18] N. I. Vilenkin, Special Functions and the Theory of Group Representations, Transl. Math. Monogr. 22, Amer. Math. Soc., Providence, R.I., 1968.
[19] L. Vozovoi, A. Weill and M. Israeli, Spectrally accurate solution of nonperiodic differential equations by the Fourier-Gegenbauer method, SIAM J. Numer. Anal. 34, 1451-1471, (1997).
[20] http://functions.wolfram.com/HypergeometricFunctions/GegenbauerC3General/


[^0]:    2010 Mathematics Subject Classification. 42C10, 65T50.
    Key words and phrases. Ultraspherical expansions, Gegenbauer coefficients, Abel transform.

