

The media effect in Axelrod's model explained

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We revisit the problem of introducing an external global field – the mass media – in Axelrod's model of social dynamics, where in addition to their nearest neighbors, the agents can interact with a virtual neighbor whose cultural features are fixed from the outset. The finding that this apparently homogenizing field actually increases the cultural diversity has been considered a puzzle since the phenomenon was first reported more than a decade ago. Here we offer a simple explanation for it, which is based on the pedestrian observation that Axelrod's model exhibits more cultural diversity, i.e., more distinct cultural domains, when the agents are allowed to interact solely with the media field than when they can interact with their neighbors as well. In this perspective, it is the local homogenizing interactions that work towards making the absorbing configurations less fragmented as compared with the extreme situation in which the agents interact with the media only.

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I. INTRODUCTION

Although computational social science – a branch of computer science that aims at predicting the behavior of groups as well as exploring ways of influencing that behavior – is currently in vogue in the scientific media (see, e.g., [1]), it has a long record in the computational physics community [2–4]. In fact, physicists have offered a variety of models of social influence and opinion formation (see, e.g., [5–8]), but they seem to have adopted as a paradigm a model proposed by a political scientist – Axelrod's model [9] – which exhibits the right balance between simplicity and realism, and so succeeds in boiling down a collective social phenomenon to its functional essence [10]. In addition, Axelrod's model has served as inspiration to the proposal of heuristics that use social interactions to solve combinatorial optimization problems [11–13].

Axelrod's model was introduced to explore the mechanisms behind the persistence of cultural and opinion differences in a society, given that the interactions between people tend, on the average, to make them more alike in their beliefs and attitudes. Hinging on a few reasonable assumptions, Axelrod's model offers a non-trivial answer to that issue. In that model, an agent is represented by a string of cultural features, where each feature can adopt a certain number of distinct traits. The interaction between any two agents takes place with probability proportional to their cultural similarity, i.e., proportional to the number of traits they have in common. The result of such interaction is the increase of the similarity between the two agents, as one of them modifies a previously distinct trait to match that of its partner. In spite of this homogenizing assumption, the model exhibits global polarization, i.e., a stable multicultural regime [9]. Subsequent analysis of Axelrod's model by the statistical physics community has revealed a rich dynamic behavior with a nonequilibrium phase transition separating the global polarization regime from the homogeneous regime, where a single culture dominates

the entire population [14–16].

Another counterintuitive result derived from the analysis of Axelrod's model is that the introduction of an external global media (i.e., it is the same for all agents) aiming at influencing the agents opinions actually favors polarization [17]. This result conflicts with the view that mass media, such as newspapers and television, are instruments to control people's opinions and so homogenize society [18]. The effect of the mass media in Axelrod's model has been extensively investigated (see, e.g., [19–22]), but no first-principles explanation for it was offered. In fact, much of that research has focused on the search for a threshold on the intensity of the media influence such that above that threshold, the population is culturally polarized and below it, the population is homogeneous. More recently, it was argued that this threshold is an artifact of finite lattices, and that even a vanishingly small media influence is sufficient to produce cultural polarization in a region of the parameter space where the homogeneous regime is dominant in the absence of the media [23].

Here we offer a simple explanation for the media effect in Axelrod's model, which is based on viewing the problem from a new perspective: rather than focusing on the counterintuitive domain-fragmenting effect of increasing the media strength, we centre at the trite homogenizing effect of the interactions between agents, which become more important as the media strength is weakened. We conclude then that it is the local homogenizing interactions that work towards making the absorbing configurations less heterogeneous as compared with the extreme situation in which the agents can interact with the media only.

II. MODEL

In the original Axelrod's model [9], each agent is characterized by a set of F cultural features which can take on q distinct values. The agents are fixed in the sites of

a square lattice of linear size L with open boundary conditions (i.e., agents in the corners of the lattice interact with two neighbors, agents in the sides with three, and agents in the bulk with four nearest neighbors). The initial configuration is completely random with the features of each agent given by random integers drawn uniformly between 1 and q . At each time we pick an agent at random (this is the target agent) as well as one of its neighbors. These two agents interact with probability equal to their cultural similarity, defined as the fraction of common cultural features. An interaction consists of selecting at random one of the distinct features, and making the selected feature of the target agent equal to its neighbor's corresponding trait. This procedure is repeated until the system is frozen into an absorbing configuration.

The introduction of a global media in the standard model follows the ingenious suggestion of adding a virtual agent which interacts with all agents in the lattice and whose cultural traits reflect the media message [17]. In the original version, each cultural feature of the virtual agent has the trait which is the most common in the population – the consensus opinion. To speed up the simulations, here we choose to keep the media message fixed from the outset, so it really models some alien influence impinging on the population. Explicitly, we generate the culture vector of the virtual agent at random and keep it fixed during the dynamics [19, 20]. The interaction of the media (virtual agent) with the real agents is governed by the control parameter $p \in [0, 1]$, which may be interpreted as a measure of the strength of the media influence. As in the original Axelrod's model, we begin by choosing a target agent at random, but now it can interact with the media with probability p or with its neighbors with probability $1 - p$. Since we have defined the media as a virtual agent, the interaction follows exactly the same rules as before. The original model is recovered by setting $p = 0$.

The simulation of the dynamics of Axelrod's model can be greatly accelerated by introducing a list of active agents [16]. An active agent is an agent that has at least one feature in common and at least one feature distinct with at least one of its neighbors, included the virtual agent (media). Clearly, since only active agents can change their cultural features, it is more efficient to select the target agent randomly from the list of active agents rather than from the entire lattice. Note that the randomly selected neighbor of the target agent may not necessarily be an active agent itself. In the case that the cultural features of the target agent are modified by the interaction with its neighbor, we need to re-examine the active/inactive status of the target agent as well as of all its neighbors so as to update the list of active agents. The dynamics is frozen when the list of active agents is empty.

Because the cultural features of the virtual agent are fixed a priori, there are three distinct types of absorbing states of the dynamics. These absorbing states determine the stationary (or frozen) regimes of the model.

First, there is the homogeneous regime where all agents are identical to the media. Second, there are $(q - 1)^F$ distinct homogeneous regimes in which all F cultural traits of the agents are distinct from those of the media. However, we will show below that this type of absorbing configuration is absent for large lattices. Third, there is a heterogeneous regime characterized by the co-existence of domains of agents whose cultural traits are identical to the media with domains of agents whose traits are all distinct from the media.

As pointed out before, the puzzling phenomenon reported first by Shibanaï et al. [17] and then confirmed by many other authors [19–23] is that for a set of control parameters, say $F = 5$ and $q = 5$, the homogeneous regime in the absence of media is replaced by a heterogeneous regime whenever the media is turned on (i.e., for $p > 0$). This goes against the intuitive idea that the effect of the media is to make the culture more homogeneous. Studies based on lattices of small sizes (typically $L = 40$) have suggested the existence of a threshold value p_c on the intensity of the media influence such that for $p > p_c$ the stationary regime is heterogeneous, whereas it is homogeneous for $p \leq p_c$ [19–22]. (We recall that the parameters q and F are chosen such that the population is culturally homogeneous for $p = 0$.) This finding was questioned in a recent study which suggested that p_c is an artifact of finite lattices and hence that even a vanishingly small media influence is sufficient to destabilize the culturally homogeneous regime [23]. However, the problem with that study was that the virtual agent (media) did not take part in the determination of the active status of the agents. As we will argue below, correction of this point does not change the claim that $p_c \rightarrow 0$ when $L \rightarrow \infty$.

III. SIMULATION RESULTS

We consider first the ratio $g \in [0, 1]$ between the number of clusters or cultural domains S and the lattice area L^2 (see fig. 1). A cluster or domain is simply a bounded region of uniform culture. Two or more cultural domains characterized by the same culture are counted separately in the calculation of S . For the parameter set $F = q = 5$, analysis of small lattices indicates a threshold value $p_c \approx 0.03$ [19, 20], but the results exhibited in fig. 1 indicate unequivocally that the heterogeneous regime displaces the homogeneous one even for very small values of the media strength as the lattice size increases. This figure is virtually indistinguishable from that obtained in the case the media does not take part in the determination of the active status of the agents [23]. In this, as well as in the next figures of this paper, the error bars are smaller or at most equal to the symbol sizes. We note that there are many more cultural domains in the case the local interactions are turned off (i.e., $p = 1$) than in any other situation.

A more lively illustration of the effect of increasing the

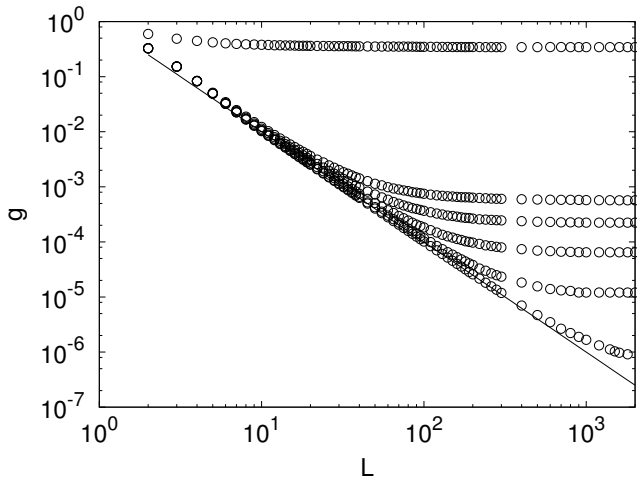


FIG. 1: Logarithmic plot of the ratio g between the number of domains and the lattice area as function of the linear size L of the lattice for (top to bottom) $p = 1, 0.05, 0.04, 0.03, 0.02$ and 0.01 . The solid straight line is $1/L^2$, which corresponds to the value of g in the uniform regime. The parameters are $F = 5$ and $q = 5$.

lattice size is offered in fig. 2, which shows the fraction of runs ξ_h which ended up in a homogeneous absorbing configuration. Typically we carried out from 10^3 to 10^4 independent runs of the stochastic dynamics for a given value of the media strength p . This figure highlights the somewhat odd predominance of the homogeneous absorbing configurations for lattices of intermediate size, which has led to the suggestion of a nonzero threshold value p_c . For $p = 1$ and $L > 5$ we find that the absorbing configurations are always heterogeneous. In addition, we note that, except for small lattice sizes, the homogeneous absorbing configurations are identical to the media. In fact, we have estimated the fraction of runs ξ_m which ended up in the media configuration and plotted in fig. 3 the fraction of homogeneous absorbing configurations distinct from the media, i.e., the quantity $1 - \xi_m/\xi_h$. As shown in the figure, this fraction vanishes exponentially fast with the increase of the lattice size L^2 and so, for large lattices, there is essentially a single homogeneous absorbing configuration – the one where all agents assume the media cultural traits.

The results exhibited in figs. 1 and 2 merely reassert the claim that the media promotes cultural diversity, since the number of cultural domains increases with the strength of the media, measured by the parameter p . However, this is a rather odd view of the situation. In fact, in all runs and for all parameter settings considered, we found that the largest domain is composed of agents identical to the media. In addition, the size of this domain is much larger than the sum of the sizes of the other domains, as shown in fig. 4, which exhibits the fraction η of the lattice sites (agents) that do not belong to the largest domain. We note that although some of these small domains are also characterized by the media traits,

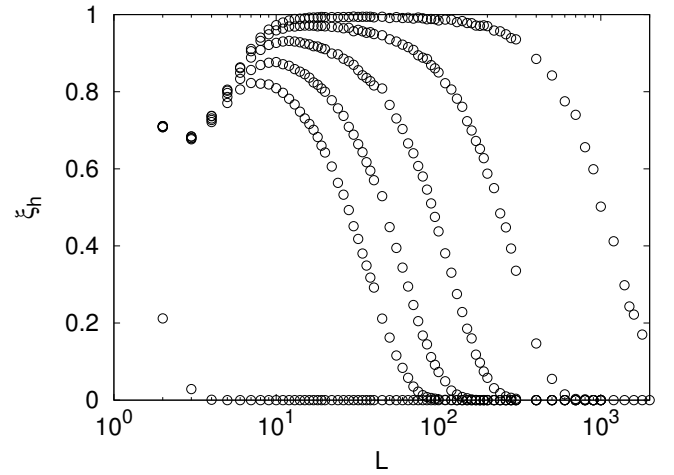


FIG. 2: Fraction of runs trapped into homogeneous absorbing configuration ξ_h as function of the linear size L of the lattice for (top to bottom) $p = 0.01, 0.02, 0.03, 0.04, 0.05$ and 1 . The parameters are $F = 5$ and $q = 5$.

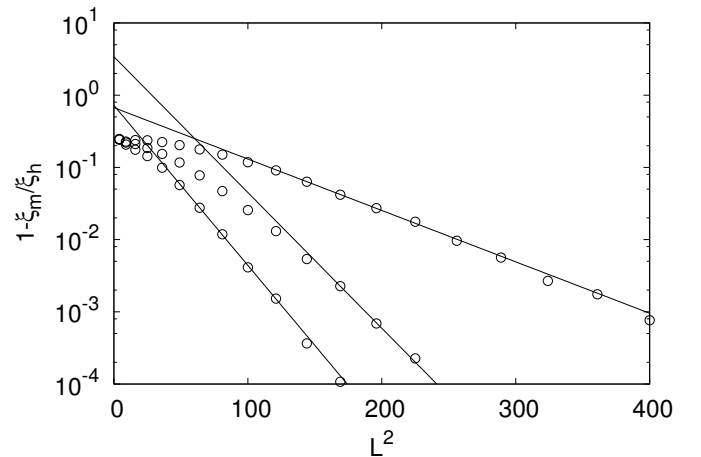


FIG. 3: Fraction of the homogeneous absorbing configurations which are distinct from the media as function of the lattice size L^2 for (top to bottom) $p = 0.01, 0.02$ and 0.03 . The solid lines are the fittings $a_p \exp(-b_p L^2)$. The parameters are $F = 5$ and $q = 5$.

they count as distinct domains as they are disconnected from the largest cluster – a situation known as cultural diaspora [24]. In agreement with the conclusion drawn from the previous figures, the domains different from the media occupy a finite, though very small, portion of the lattice in the thermodynamic limit. More pointedly, in this limit the average size of the domains, excluded the largest one, can be obtained by calculating the ratio η/g and yields about 3 sites, regardless of the value of p .

Figure 4 hints that there is something trivial about the nature of the cultural domains distinct from the media. There are two points which are worth emphasizing

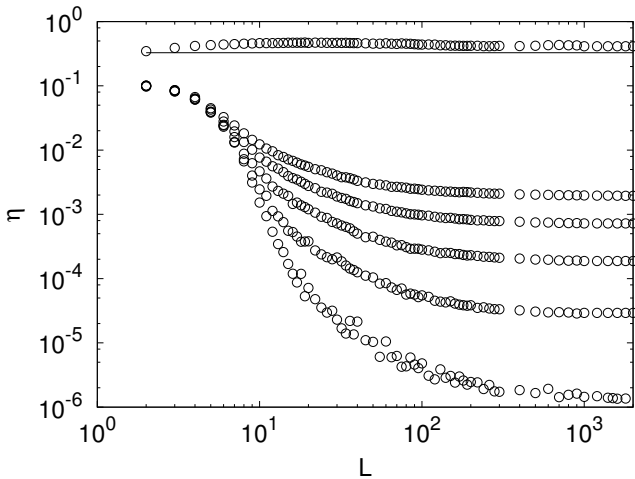


FIG. 4: Fraction of the lattice sites that are not part of the largest domain (media) as function of the linear size L of the lattice for (top to bottom) $p = 1, 0.05, 0.04, 0.03, 0.02$ and 0.01 . The horizontal line is $(1 - 1/q)^F$. The parameters are $F = 5$ and $q = 5$.

here. First, when the agents are allowed to interact solely with the fixed external media, the absorbing configurations are heterogeneous (see fig. 2), since agents whose F cultural traits are all distinct from the media at the onset will be unaffected by the dynamic rules. For instance, the horizontal line in fig. 4 gives the fraction of sites distinct from the media regardless of the domain organization in the case the local interactions are turned off ($p = 1$), i.e., $(1 - 1/q)^F$. Those sites are frozen from the very beginning of the simulation. Second, there are many more agents different from the media for $p = 1$ than for $p < 1$, i.e., the local homogenizing interactions contribute to increase the number of agents with the media traits. Summing up, due to the non-zero similarity restriction for the occurrence of interactions, the media in Axelrod's model does not promote uniformity; on the contrary, there are more cultural diversity when the media is the only partner the agents can interact with than when local interactions are allowed.

IV. CONCLUSION

There is nothing extraordinary about the effect of the media in Axelrod's model, provided we interpret the results summarized in fig. 4 from top to bottom: when assembling the random initial configuration a fraction $(1 - 1/q)^F$ of the agents have their F cultural traits distinct from the media and so they form the cultural domains in the case the local interactions are neglected

($p \rightarrow 1$). As the media strength p decreases, the local interactions have the chance to flip some of those traits, resulting in the increase of the size of the media cluster. Finally, when $p \rightarrow 0$ the media domain takes over the entire lattice, leading to a unique homogeneous absorbing configuration. For $p = 0$ the homogeneous absorbing configuration is not unique – there are q^F different possibilities – but, as in the ferromagnetic Ising model, a vanishingly small external field suffices to break this symmetry and direct the relaxation to a unique absorbing configuration. Although we have restricted our analysis to a single setting of the control parameters, namely, $F = q = 5$, due mainly to the computational demand of the simulations of large lattices, our conclusion is valid in general since it provides an intuitive first-principles explanation for the effect of the media in Axelrod's model, which is not affected by the tuning of the control parameters of the model.

The somewhat trivial nature of the cultural domains of Axelrod's model in presence of a external media is reminiscent of the organization of the domains of the extended majority model [25–27], which is completely determined by the initial random assignment of cultural traits. As in that model, relaxation to an absorbing configuration is very fast since the random clusters formed at the beginning are practically unaltered by the update rules of the model.

In summary, we show here that the effect of external global media in Axelrod's model can be described much more prosaically by pointing out that the local homogenizing interactions destabilize the domains of agents (usually, isolated sites) which have zero similarity with the media and hence are not affected by it. As a result, decreasing the media strength and, consequently, increasing the role of the local interactions will favor less fragmented absorbing configurations. Reversely, increasing the media strength will favor more heterogeneous absorbing configurations. It seems to us that the cause of the ruckus about the media effect in Axelrod's model was the failure to realize that the absorbing configurations can already be very fragmented due to the interactions with the media only. In fact, the mere inclusion of the data for $p = 1$ in figs. 1 and 4 turns the media effect from a surprising phenomenon to a trivial one.

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