WAVELET ANALYSIS ON ADELES AND PSEUDO-DIFFERENTIAL OPERATORS

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ABSTRACT. This paper is devoted to wavelet analysis on adele ring \mathbb{A} and the theory of pseudo-differential operators. We develop the technique which gives the possibility to generalize finite-dimensional results of wavelet analysis to the case of adeles \mathbb{A} by using infinite tensor products of Hilbert spaces. The adele ring is roughly speaking a subring of the direct product of all possible (*p*-adic and Archimedean) completions \mathbb{Q}_p of the field of rational numbers \mathbb{Q} with some conditions at infinity. Using our technique, we prove that $L^2(\mathbb{A}) = \bigotimes_{e,p \in \{\infty,2,3,5,\dots\}} L^2(\mathbb{Q}_p)$ is the infinite tensor product of the spaces $L^2(\mathbb{Q}_p)$ with a stabilization $e = (e_p)_p$, where $e_p(x) = \Omega(|x|_p) \in L^2(\mathbb{Q}_p)$, and Ω is a characteristic function of the unit interval $[0, 1], \mathbb{Q}_p$ is the field of *p*-adic numbers, p = $2, 3, 5, \ldots; \mathbb{Q}_{\infty} = \mathbb{R}$. This description allows us to construct an infinite family of Haar wavelet bases on $L^2(\mathbb{A})$ which can be obtained by shifts and multi-delations. The adelic multiresolution analysis (MRA) in $L^2(\mathbb{A})$ is also constructed. In the framework of this MRA another infinite family of Haar wavelet bases is constructed. We introduce the adelic Lizorkin spaces of test functions and distributions and give the characterization of these spaces in terms of wavelet functions. One class of pseudo-differential operators (including the fractional operator) is studied on the Lizorkin spaces. A criterion for an adelic wavelet function to be an eigenfunction for a pseudo-differential operator is derived. We prove that any wavelet function is an eigenfunction of the fractional operator. These results allow one to create the necessary prerequisites for intensive using of adelic wavelet bases and pseudo-differential operators in applications.

1. INTRODUCTION

1.1. *p*-Adic and adelic analysis. During a few hundred years theoretical physics has been developed on the basis of real and, later, complex numbers. The *p*-adic numbers were described by K. Hensel in 1897 to transfer the ideas and techniques of power series methods to number theory. According to the well-known Ostrovsky theorem, any nontrivial valuation on the field \mathbb{Q} of rational numbers is equivalent either to the real valuation $|\cdot|$ or to one of the *p*-adic valuations $|\cdot|_p$. The corresponding completions of \mathbb{Q} give the fields \mathbb{R} or \mathbb{Q}_p . The theory of *p*-adic numbers has already penetrated intensively into several areas of mathematics and its applications. In the last 20 years the field of *p*-adic numbers \mathbb{Q}_p (as well as its algebraic extensions, including the field of complex *p*-adic numbers \mathbb{C}_p) has been intensively used in theoretical and mathematical physics, *p*-adic string theory, gravity and cosmology, the theory of stochastic differential equations over the field of *p*-adic numbers, Feynman path integration over *p*-adics, the

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theory of *p*-adic valued probabilities and dynamical systems, in theory of disordered systems (spin glasses) (see [8], [26], [27], [30], [38], [40], [64] and the references therein). Applications were, however, not only restricted to physics. *p*-Adic models were also proposed in psychology, cognitive and social sciences, and, e.g., in biology, image analysis (see [27], [28]).

These applications induced and stimulated a development of new branches of p-adic analysis, in particular, the theory of *p*-adic wavelets. Recall that nowadays wavelets are applied in a lot of branches of modern mathematics and engineering area. The first real wavelet basis (3.5), (3.6) was introduced by Haar in 1910. However, for almost a century nobody could find another wavelet function (a function whose shifts and delations form an orthogonal basis). Only in the early nineties a method for a more general construction of the wavelet functions appeared [46], [47], [50], [51]. This method is based on the notion of *multiresolution analysis* (MRA in the sequel). In the *p*-adic setting the situation was the following. In 2002 S. V. Kozyrev [39] constructed a compactly supported p-adic wavelet basis (3.21) in $L^2(\mathbb{Q}_p)$, which is an analog of the real Haar basis (3.5), (3.6). J. J. Benedetto and R. L. Benedetto [10], R. L. Benedetto [11] suggested a method for constructing wavelet bases on locally compact abelian groups with compact open subgroups. This method is applicable for the *p*-adic setting. It is based on a *theory* of wavelet sets and only allows the construction of wavelet functions whose Fourier transforms are the characteristic functions of some sets (see [10, Proposition 5.1.]). Moreover, these authors doubted that the development of the MRA approach is possible. In spite of the above opinions and arguments [10], [11], in [57], the p-adic MRA theory in $L^2(\mathbb{Q}_p)$ was developed and new *p*-adic wavelet bases were constructed. Some important results in p-adic wavelet theory were obtained in [1], [2], [34], [35]. It turned out that the theory of p-adic wavelets plays an important role in the study of p-adic pseudo-differential operators and equations [3], [4], [5], [33], [39], [40], [57]. This theory gives a powerful technique to deal with p-adic pseudo-differential equations. Recall that on complexvalued functions defined on \mathbb{Q}_p , the operation of differentiation is not defined. As a result, a large number of *p*-adic models use pseudo-differential equations instead of differential equations. The p-adic multidimensional fractional operator D^{α} was introduced by M. Taibleson [58] (see also [59]) in the space of distributions $\mathcal{D}'(\mathbb{Q}_n^n)$. The spectral theory of this fractional operator was developed by V.S. Vladimirov in [62], in particular, explicit formulas for the eigenfunctions of this operator were constructed (see also [64]). In [63] (see also [64]) V.S. Vladimirov constructed the spectral theory of the Schrödingertype operator $D^{\alpha} + V(x)$, which was further developed by A.N. Kochubei [36], [37] (see also [38]).

The adele ring A is some subring of the direct product of all possible (*p*-adic and Archimedean) completions \mathbb{Q}_p of the field of rational numbers \mathbb{Q} . The group of ideles was introduced by Chevalley in 1936 [16] as a part of his program to formulate class field theory so that it worked for infinite-degree extensions. The adeles were introduced by Weil in the late 1930s as an additive analogue of ideles. The ring A is often used in advanced parts of number theory (for example, see [13], [15], [17], [48], [49], [65]). For example, Tate's proof (see in [15, J. T. Tate, pp.305–347]) of the functional equation Z(s) = Z(1-s), where $Z(s) = \pi^{-s/2}\Gamma(s/2)\zeta(s)$ for the Riemann zeta function $\zeta(s) = \sum_{n \in \mathbb{N}} \frac{1}{n^s}$, $\Re(s) > 1$, is based on the Fourier analysis on adeles A and ideles (see also [17, 4.7]). Recently the theory of adeles has been successfully applied in various parts of contemporary mathematical and theoretical physics. Namely, there are adelic constructions, e.g., in statistical mechanics, stochastics, string theory, quantum cosmology, and quantum mechanics (see [9], [14], [19]– [21], [25], [38, 5.8] [26], [27], [60], [61], [64], [66], and the references therein).

However it should be noted that in contrast to the real and p-adic analysis, the adelic analysis practically (in particular, the adelic wavelet theory and the theory of adelic pseudo-differential operators and equations) has not been developed so far. In particular, the necessity of the development the wavelet theory on adeles was mentioned in [10].

1.2. Contents of the paper. In this paper some problems of the adelic harmonic analysis are studied. The points to be considered are, first, the theory of adelic Haar wavelets; and second, the theory of simplest adelic pseudo-differential operators in connection with wavelets.

In Sec. 2, some facts from the *p*-adic and adelic analysis are given. In Sec. 3, we recall definitions of the real MRA and present some results on *p*-adic MRA and wavelet bases from the papers [35], [57].

In Sec. 4, the theory of infinite tensor products of Hilbert spaces [52], [12] is used to generalize finite-dimensional results to the case of adeles. We recall the constructions of infinite tensor product of Hilbert spaces $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} \mathcal{H}_n$, the complete von Neumann product of infinitely many Hilbert spaces. The space \mathcal{H}_e can be obtained as the closure of the union of some subspaces. We observe *certain stability* of the space \mathcal{H}_e with respect to the variation of the corresponding subspaces (see Lemma 4.2). In Subsec. 4.4, using Lemma 4.2 for the special stabilizing sequence $e = (e_p)_p$, $e_p(x_p) = \phi_p(x_p) = \Omega(|x_p|_p)$ (here $\Omega(t)$ is the characteristic function of the segment $[0, 1] \subset \mathbb{R}$) we show that $L^2(\mathbb{A})$ coincides with the infinite tensor product of the Hilbert spaces $L^2(\mathbb{Q}_p)$ over all possible completions of the field \mathbb{Q} : $L^2(\mathbb{A}) = \bigotimes_{e,p} L^2(\mathbb{Q}_p)$ (see Lemma 4.3).

In Sec. 5, we apply the above scheme to construct adelic wavelet bases (5.5) on $L^2(\mathbb{A}, dx)$ generated by the tensor product of one-dimensional Haar wavelet bases (3.5), (3.22), (3.25). We would like to stress that to construct adelic wavelet bases, we need the *p*-adic wavelet bases that contain functions $\phi_p(x_p) = \Omega(|x_p|_p)$. Thus, instead of the Haar wavelet basis (3.20) we will use modified Haar basis (3.22). According to [1], there are no orthogonal *p*-adis MRA based wavelet bases except for those described in Theorem 3.1 and Corollary 3.1. Thus, formula (5.5) gives all adelic Haar wavelet bases generated by the tensor product of one-dimensional Haar wavelet bases. Here infinite tensor product depends on the special stabilization $e = (e_p)_p$, where $e_p(x_p) = \Omega(|x_p|_p)$. Recall that in [20, Sec.4, formula (4.1)], a basis in the space $L^2(\mathbb{A})$ associated with eigenfunctions of harmonic oscillators was constructed (see Remark 5.1). To construct adelic wavelet bases, we use the standard Haar measures on \mathbb{Q}_p and on the adele ring A. For measures on \mathbb{Q}_p different from Haar measure, see [3, Appendix D], based on [42].

In Sec. 6, using the idea of constructing separable multidimensional MRA by means of the tensor product of one-dimensional MRAs suggested by Y. Meyer [50] and S. Mallat [46], [47], we construct adelic wavelet bases in $L^2(\mathbb{A})$. In a general situation we show how using some system of closed subspaces $V_j^{(k)}$, $j \in \mathbb{Z}$, in a Hilbert spaces H_n , $n \in \mathbb{N}$, with properties (a)–(c) of the MRA (see Definitions 3.1, 3.2) one can construct various systems of subspaces with the same properties in the infinite tensor product of the spaces H_n . In Theorem 6.2 adelic separable MRA is constructed. The refinable function of this MRA given by (6.14) is an infinite product refinable functions of all one-dimensional MRAs. In the framework of this MRA an infinite family of adelic Haar wavelet bases (6.25) are constructed.

In Sec. 7, we introduce the adelic Lizorkin spaces of test functions $\Phi(\mathbb{A})$ and distributions $\Phi'(\mathbb{A})$. These spaces are constructed by using the original real Lizorkin spaces introduced in [43]– [45] and the *p*-adic Lizorkin spaces introduced in [4] (see also [3, Ch. 7]). A basic motivation for using Lizorkin spaces rather than the Schwartz and Bruhat–Schwartz spaces of distributions $\mathcal{S}'(\mathbb{R})$ and $\mathcal{D}'(\mathbb{Q}_p)$ is due to the fact that the latter spaces are not invariant under the fractional operators. Next, in Sec. 8, the characterization of the adelic Lizorkin spaces in terms of wavelets is given. Namely, it is proved that any test function from $\Phi(\mathbb{A})$ can be represented in the form of a *finite* combination of adelic wavelet functions (5.6), and any distribution from $\Phi'(\mathbb{A})$ can be represented as an *infinite* linear combination of adelic wavelet functions (5.6) (in [7], assertions of these types were stated for ultrametric Lizorkin spaces).

In the framework of our constructions the following three facts seem to have the same deep reason: (1) functions $\phi^{(p)}(x_p) = \Omega(|x_p|_p)$ are stabilization functions in the adelic Bruhat–Schwartz space (see Definition 2.3), (2) we use the stabilization sequence $e = (e_p)_p$ where $e_p(x_p) = \Omega(|x_p|_p)$ in proving the fact that $L^2(\mathbb{A}) = \bigotimes_{e,p} L^2(\mathbb{Q}_p)$, (3) under the projection of the space $L^2(\mathbb{Q}_p)$ onto $L^2(\mathbb{Z}_p)$ some elements of wavelet basis (3.20) and Kozyrev's wavelet basis (3.21) are transformed into functions which are proportional to the same function $\Omega(|x_p|_p)$ (see Propositions 3.1–3.3).

In Sec. 9, by Definition (9.17), (9.18) a class of pseudo-differential operators on the adelic Lizorkin spaces is introduced. The fractional operators $D^{\widehat{\gamma}}, \widehat{\gamma} \in \mathbb{C}^{\infty}$ (see Definition (9.9), (9.10), (9.13)), and $D^{\gamma}, \gamma \in \mathbb{C}$ (see Definition (9.14), (9.15)) belong to this class. We prove that the Lizorkin spaces of test functions $\Phi(\mathbb{A})$ and distributions $\Phi'(\mathbb{A})$ are *invariant* under the above-mentioned pseudo-differential operators. Moreover, a family of fractional operators on the space of distributions $\Phi'(\mathbb{A})$ forms an abelian group. Thus the Lizorkin spaces constitute "natural" domains of definition for this class of pseudo-differential operators. Note, that in [31] the fractional operator was considered in $L^2(\mathbb{A})$. In Subsec. 9.4, the spectral theory of adelic pseudo-differential operators is developed. By Theorem 9.1, we derive a criterion for an adelic wavelet function to be an eigenfunction for a pseudo-differential operator. It is proved that any wavelet function is an eigenfunction of a fractional operator. Thus the *adelic wavelet analysis is closely* connected with the spectral analysis of pseudo-differential operators. Using results of Sec. 9, similarly to the *p*-adic case, one can develop the "variable separation method" (an analog of the classical Fourier method) to reduce solving adelic pseudo-differential equations to solving ordinary differential equations with respect to the real variable t(for details, see [3, Ch. 10], [6]).

2. Preliminary results

2.1. *p*-Adic numbers. We shall systematically use the notation and results from the book [64]. Let \mathbb{N} , \mathbb{Z} , \mathbb{Q} , \mathbb{R} , \mathbb{C} be the sets of positive integers, integers, rational, real, and complex numbers, respectively.

According to the well-known Ostrovsky theorem, any nontrivial valuation on the field of rational numbers \mathbb{Q} is equivalent either to the real valuation $|\cdot|$ or to one of the padic valuations $|\cdot|_p$. This p-adic norm $|\cdot|_p$ is defined as follows: if an arbitrary rational number $x \neq 0$ is represented as $x = p^{\gamma} \frac{m}{n}$, where $\gamma = \gamma(x) \in \mathbb{Z}$ and the integers m, n are not divisible by p, then The norm $|\cdot|_p$ is non-Archimedean and satisfies the strong triangle inequality $|x+y|_p \leq \max(|x|_p, |y|_p)$. The completion of \mathbb{Q} with respect to the usual absolute value $|\cdot|$ gives the field of real numbers \mathbb{R} . The field \mathbb{Q}_p of p-adic numbers is defined as the completion of the field of rational numbers \mathbb{Q} with respect to the norm $|\cdot|_p$. Next we will denote $|x|_{\infty} = |x|, \mathbb{Q}_{\infty} = \mathbb{R}$ and $\mathbb{Z}_{\infty} = \mathbb{Z}$. By

(2.2)
$$V_{\mathbb{Q}} = \{\infty, 2, 3, 5, \dots\}$$

we denote the set of indices for all valuations on the field \mathbb{Q} .

Any p-adic number $x \in \mathbb{Q}_p$, $x \neq 0$, is represented in the *canonical form*

(2.3)
$$x = \sum_{k=\gamma}^{\infty} x_k p^k$$

where $\gamma = \gamma(x) \in \mathbb{Z}$, $x_k \in \mathbb{F}_p = \{0, 1, \dots, p-1\}, x_\gamma \neq 0, \gamma \leq k < \infty$. The series is convergent in the *p*-adic norm $|\cdot|_p$, and one has $|x|_p = p^{-\gamma}$. The *fractional part* of the number $x \in \mathbb{Q}_p$ (given by (2.3)) is defined as follows

(2.4)
$$\{x\}_p = \begin{cases} 0, & \text{if } \gamma(x) \ge 0 & \text{or } x = 0, \\ x_{\gamma} p^{\gamma} + \dots + x_{-1} p^{-1}, & \text{if } \gamma(x) < 0 \end{cases}$$

The set $\mathbb{Q}_p^{\times} = \mathbb{Q}_p \setminus \{0\}$ is the multiplicative group of the field \mathbb{Q}_p . p-Adic numbers $\mathbb{Z}_p = \{x \in \mathbb{Q}_p : |x|_p \leq 1\}$ are called *integer p-adic numbers*. In view of (2.3), \mathbb{Z}_p consists of p-adic numbers

(2.5)
$$x = \sum_{k=0}^{\infty} x_k p^k.$$

 \mathbb{Z}_p is a subring of the ring \mathbb{Q}_p . The multiplicative group of the ring \mathbb{Z}_p is the set

$$\mathbb{Z}_{p}^{\times} = \left\{ x \in \mathbb{Z}_{p} : |x|_{p} = 1 \right\} = \left\{ x \in \mathbb{Z}_{p} : x = \sum_{k=0}^{\infty} x_{k} p^{k}, \quad x_{0} \neq 0 \right\}.$$

Denote by $B_{\gamma}(a) = \{x \in \mathbb{Q}_p : |x - a|_p \leq p^{\gamma}\}$ the ball of radius p^{γ} with the center at the point $a \in \mathbb{Q}_p$ and by $S_{\gamma}(a) = \{x \in \mathbb{Q}_p : |x - a|_p = p^{\gamma}\} = B_{\gamma}(a) \setminus B_{\gamma-1}(a)$ the corresponding sphere, $\gamma \in \mathbb{Z}$. For a = 0 we set $B_{\gamma} = B_{\gamma}(0)$ and $S_{\gamma} = S_{\gamma}(0)$.

2.2. *p*-Adic distributions. A complex-valued function f defined on \mathbb{Q}_p is called *locally-constant* if for any $x \in \mathbb{Q}_p$ there exists an integer $l(x) \in \mathbb{Z}$ such that

$$f(x+y) = f(x), \quad y \in B_{l(x)}.$$

Let $\mathcal{E}(\mathbb{Q}_p)$ and $\mathcal{D}(\mathbb{Q}_p)$ be the linear spaces of locally-constant \mathbb{C} -valued functions on \mathbb{Q}_p and locally-constant \mathbb{C} -valued compactly supported functions (so-called Bruhat– Schwartz test functions), respectively [64, VI.1.,2.]. If $\varphi \in \mathcal{D}(\mathbb{Q}_p)$, then according to Lemma 1 from [64, VI.1.], there exists $l \in \mathbb{Z}$, such that

$$\varphi(x+y) = \varphi(x), \quad y \in B_l, \quad x \in \mathbb{Q}_p.$$

The largest of such numbers $l = l(\varphi)$ is called the *parameter of constancy* of the function φ . Let us denote by $\mathcal{D}_N^l(\mathbb{Q}_p)$ the finite-dimensional space of test functions $\mathcal{D}_N^l(\mathbb{Q}_p)$ from $\mathcal{D}(\mathbb{Q}_p)$ having supports in the ball B_N and with *parameters of constancy* $\geq l$. The following embedding holds:

$$\mathcal{D}_N^l(\mathbb{Q}_p) \subset \mathcal{D}_{N'}^{l'}(\mathbb{Q}_p), \quad N \le N', \quad l \ge l'.$$

We have $\mathcal{D}(\mathbb{Q}_p) = \lim_{N \to \infty} \operatorname{ind} \mathcal{D}_N(\mathbb{Q}_p)$, where $\mathcal{D}_N(\mathbb{Q}_p) = \lim_{l \to -\infty} \operatorname{ind} \mathcal{D}_N^l(\mathbb{Q}_p)$ (see [64, VI.2]). These representations give us the *inductive limit topology* on the corresponding spaces.

Denote by $\mathcal{D}'(\mathbb{Q}_p)$ the set of all linear functionals (Bruhat–Schwartz distributions) on $\mathcal{D}(\mathbb{Q}_p)$ [64, VI.3.].

The Fourier transform of $\varphi \in \mathcal{D}(\mathbb{Q}_p)$ is defined by the formula

$$\widehat{\varphi}(\xi) = F[\varphi](\xi) = \int_{\mathbb{Q}_p} \chi_p(\xi x) \varphi(x) \, dx, \qquad \xi \in \mathbb{Q}_p,$$

where dx is the Haar measure on \mathbb{Q}_p such that $\int_{|x|_p \leq 1} dx = 1$, and

(2.6)
$$\chi_p(x) = e^{2\pi i \{x\}_p}$$

is the additive character on \mathbb{Q}_p (see [64, III.1.]), $\{x\}_p$ is the fractional part (2.4) of the number $x \in \mathbb{Q}_p$. The Fourier transform is a linear isomorphism $\mathcal{D}(\mathbb{Q}_p)$ onto $\mathcal{D}(\mathbb{Q}_p)$ [59, III,(3.2)], [64, VII.2.]. Moreover,

(2.7)
$$\varphi \in \mathcal{D}_{N}^{l}(\mathbb{Q}_{p}) \quad \text{iff} \quad F[\varphi] \in \mathcal{D}_{-l}^{-N}(\mathbb{Q}_{p}).$$

The Fourier transform of a distribution $f \in \mathcal{D}'(\mathbb{Q}_p)$ is the distribution $\widehat{f} = F[f]$ defined by the relation

$$\langle F[f], \varphi \rangle = \langle f, F[\varphi] \rangle, \quad \forall \varphi \in \mathcal{D}(\mathbb{Q}_p).$$

Here and in the sequel $\langle f, \varphi \rangle$ denotes the action of a distribution f on a test function φ . If $f \in \mathcal{D}'(\mathbb{Q}_p), \ a \in \mathbb{Q}_p^{\times}, b \in \mathbb{Q}_p$, then [64, VII,(3.3)]:

(2.8)
$$F[f(ax+b)](\xi) = |a|_p^{-1}\chi_p\left(-\frac{b}{a}\xi\right)F[f(x)]\left(\frac{\xi}{a}\right), \quad x,\xi \in \mathbb{Q}_p.$$

According to [64, IV, (3.1)],

(2.9)
$$F[\Omega(p^{-k}|\cdot|_p)](x) = p^k \Omega(p^k|x|_p), \quad k \in \mathbb{Z}, \qquad x \in \mathbb{Q}_p$$

where $\Omega(t)$ is the characteristic function of the segment $[0,1] \subset \mathbb{R}$. In particular,

(2.10)
$$F[\Omega(|\xi|_p)](x) = \Omega(|x|_p).$$

2.3. Adeles. We use the notation and results from [24, Ch. III,§1,2] and [17].

Definition 2.1. The *adeles* of \mathbb{Q} are

$$\mathbb{A} = \mathbb{A}_{\mathbb{Q}} = \left\{ (x_p)_{p \in V_{\mathbb{Q}}} \in \prod_{p \in V_{\mathbb{Q}}} \mathbb{Q}_p : x_p \in \mathbb{Z}_p \text{ for almost all } p \neq \infty \right\},\$$

and the *ideles* of \mathbb{Q} are

$$\mathbb{J} = \mathbb{J}_{\mathbb{Q}} = \big\{ (x_p)_{p \in V_{\mathbb{Q}}} \in \prod_{p \in V_{\mathbb{Q}}} \mathbb{Q}_p^{\times} : x_p \in \mathbb{Z}_p^{\times} \text{ for almost all } p \neq \infty \big\},\$$

where $V_{\mathbb{Q}} = \{\infty, 2, 3, 5, \dots\}$ is the set of indices.

The adele ring A is the restricted direct product of \mathbb{R} and \mathbb{Q}_p for all p = 2, 3, ... with respect to the integer rings \mathbb{Z}_p (see [17]). If componentwise operations of addition and multiplication are introduced, A is a ring of adeles and J is a multiplicative group. The additive group of the ring A is called the *adelic group*.

There is a natural imbedding $\mathbb{Q} \mapsto \mathbb{A}$ given by

$$\mathbb{Q} \ni r \mapsto (r, r, \dots, r, \dots) \in \mathbb{A}.$$

Indeed, any constant sequence $(r, r, \ldots, r, \ldots)$ is an adele since $r \in \mathbb{Z}_p$ for any p not dividing the denominator of r. The adeles (ideles) of the form $(r, r, \ldots, r, \ldots)$, where $r \in \mathbb{Q}$, are called *principal adeles* (respectively, *ideles*).

To define a topology on \mathbb{A} , we show that \mathbb{A} is a union of locally compact groups. Note that the direct product $\prod_{p \in V_{\mathbb{Q}}} \mathbb{Q}_p$ is not a locally compact group.

Definition 2.2. Let S be a finite subset of $V_{\mathbb{Q}}$ such that $\infty \in S \subset V_{\mathbb{Q}}$. Define the S-adeles of \mathbb{Q} as

$$\mathbb{A}_S = \mathbb{A}_{\mathbb{Q},S} = \prod_{p \in S} \mathbb{Q}_p \times \prod_{p \notin S} \mathbb{Z}_p.$$

For an arbitrary S, the space \mathbb{A}_S of S-adeles is locally compact (in the Tikhonov product topology) as the product of a finite product of locally compact spaces \mathbb{Q}_p , $p \in S$ by an infinite product of compact spaces \mathbb{Z}_p , $p \notin S$. Since we have $\mathbb{A} = \bigcup_S \mathbb{A}_S$ [17, Theorem 2.15], the adele group \mathbb{A} is locally compact (i.e. any neighborhood of a point contains a compact neighborhood of this point). A fundamental system of open neighborhoods of zero in \mathbb{A} is the following set (see [17, Theorem 2.17])

(2.11)
$$\prod_{p \in S} U_p \times \prod_{p \notin S} \mathbb{Z}_p,$$

where S is a finite subset of $V_{\mathbb{Q}}$ that contains ∞ and U_p is an open set in \mathbb{Q}_p containing $0 \in \mathbb{Q}_p$ for $p \in S$.

The non-Archimedean part $\widetilde{\mathbb{A}}$ of the adele ring \mathbb{A} is defined (see [31]) as the set of infinite sequences

$$x' = (x_2, \dots, x_p, \dots), \text{ where } x_p \in \mathbb{Q}_p, p = 2, 3, \dots,$$

and there exists a prime number P = P(x') such that $x_p \in \mathbb{Z}_p$ for $p \ge P$. The sequence of adeles $\{x^{(n)} \mid n \in \mathbb{N}\}$ converges to the adele $x \mid x^{(n)} \to x \mid n \in \mathbb{N}\}$

The sequence of adeles
$$\{x^{(n)}, n \in \mathbb{N}\}$$
 converges to the adele $x \ (x^{(n)} \to x, n \to \infty)$ if

(1) $x_p^{(n)} \to x_p, n \to \infty$, for any $p \in V_{\mathbb{Q}}$, where $V_{\mathbb{Q}}$ is defined by (2.2);

(2) there exists N such that $x_p^{(n)} - x_p \in \mathbb{Z}_p$ for all $n \ge N$.

Since the adele group A is a locally compact commutative group, *it possesses the* Haar measure which will be denoted by dx, where $x = (x_{\infty}, x_2, \ldots, x_p, \ldots)$. The Haar measure dx can be expressed in terms of the measures dx_p on the groups \mathbb{Q}_p as follows:

$$(2.12) dx = dx_{\infty} dx_2 \cdots dx_p \cdots,$$

where

$$\int_0^1 dx_\infty = 1, \qquad \int_{\mathbb{Z}_p} dx_p = 1, \quad p \in V_{\mathbb{Q}}.$$

Here formula (2.12) is understood in the following sense: if

$$f(x) = f_{\infty}(x_{\infty})f_2(x_2)\cdots f_p(x_p)$$

is a cylindrical function, then

$$\int_{\mathbb{A}} f(x) \, dx = \int_{\mathbb{R}} f_{\infty}(x_{\infty}) \, dx_{\infty} \int_{\mathbb{Q}_2} f_2(x_2) \, dx_2 \cdots \int_{\mathbb{Q}_p} f_p(x_p) \, dx_p.$$

Any additive character $\chi(x)$ on the adelic ring has the form

$$\chi(x) = \chi_0(ax), \quad x = (x_\infty, x_2, \dots x_p, \dots) \in \mathbb{A},$$

for some $a = (a_{\infty}, a_2, \dots, a_p, \dots) \in \mathbb{A}$ (see [24, Ch. III,§1.5,(3)]). Here

(2.13)
$$\chi_0(a) = \prod_{p \in V_{\mathbb{Q}}} \chi_p(a_p), \quad a \in \mathbb{A},$$

where $\chi_{\infty}(a_{\infty}) = e^{2\pi i a_{\infty}}$, and $\chi_p(a_p)$ is defined by (2.6), $p = 2, 3, \ldots$. It is clear that for any $a \in \mathbb{A}$ there exists a prime number P = P(a) such that $\chi_p(a_p) = 1$ for $p \ge P$, i.e., in fact, the product (2.13) is finite. In other words, $\chi_0(a) = \exp(2\pi i \sigma(a))$, where $\sigma(a) = \sum_{p \in V_0} a_p \pmod{1}$.

Definition 2.3. (see [24, Ch. III,§2.1]) Let $\mathcal{S}(\mathbb{R})$ be the real Schwartz space of tempered test functions. The space of Bruhat–Schwartz adelic test functions $\mathcal{S}(\mathbb{A})$ consists of finite linear combinations of elementary functions of the form

(2.14)
$$\varphi(x) = \prod_{p \in V_{\mathbb{Q}}} \varphi_p(x_p), \quad x \in \mathbb{A},$$

where the factors $\varphi_p(x_p)$ are such that:

(i) $\varphi_{\infty}(x_{\infty}) \in \mathcal{S}(\mathbb{R});$

(ii)
$$\varphi_p(x_p) \in \mathcal{D}(\mathbb{Q}_p), p = 2, 3, \dots$$

(ii) $\varphi_p(x_p) \subset \mathcal{P}(\varphi_p)$, p = 2, 3, ...,(iii) there exists $P = P(\varphi)$ such that $\varphi_p(x_p) = \Omega(|x_p|_p)$ for all p > P (the number $P(\varphi)$ is called the *parameter of finiteness* of an elementary function φ).

In view of condition (*iii*) the space of test functions $\mathcal{S}(\mathbb{A})$ admits a natural representation

(2.15)
$$\mathcal{S}(\mathbb{A}) = \lim_{m \in V_{\mathbb{Q}} \setminus \infty} \operatorname{ind} \mathcal{S}^{[m]}(\mathbb{A}),$$

where $\mathcal{S}^{[m]}(\mathbb{A})$ is the subspace of all test functions with the *parameter of finiteness* m, $m \in V_{\mathbb{Q}} \setminus \infty$. The representation (2.15) equips the space $\mathcal{S}(\mathbb{A})$ with the *inductive limit topology*.

The Bruhat–Schwartz adelic test functions are continuous on \mathbb{A} . The Bruhat–Schwartz space of *adelic distributions* $\mathcal{S}'(\mathbb{A})$ was studied in [31], [22].

The spaces of test functions and distributions connected with the non-Archimedean part of adeles $\widetilde{\mathbb{A}}$ we denote by $\mathcal{S}(\widetilde{\mathbb{A}})$ and $\mathcal{S}'(\widetilde{\mathbb{A}})$, correspondingly (see [31]).

The Fourier transform of $\varphi \in \mathcal{S}(\mathbb{A})$ is defined by the formula

$$\widehat{\varphi}(\xi) = F[\varphi](\xi) \stackrel{def}{=} \int_{\mathbb{A}} \chi_0(\xi x) \varphi(x) \, dx, \qquad \xi \in \mathbb{A},$$

where χ_0 is defined by (2.13). It is clear that $F[\mathcal{S}(\mathbb{A})] = \mathcal{S}(\mathbb{A})$. If $f \in \mathcal{S}'(\mathbb{A})$, then

$$\langle F[f], \varphi \rangle \stackrel{def}{=} \langle f, F[\varphi] \rangle, \quad \forall \varphi \in \mathcal{S}(\mathbb{A}),$$

and $F[f] \in \mathcal{S}'(\mathbb{A})$.

For $f, g \in L^2(\mathbb{A})$ we have

(2.16)
$$(f,g) = (F[f], F[g]), \qquad ||f||_2 = ||F[f]||_2,$$

where

$$(f,g) = \int_{\mathbb{A}} f(x)\overline{g(x)} \, dx$$

is the scalar product of the functions f and g in $L^2(\mathbb{A})$.

It is clear that the space $\mathcal{S}(\mathbb{A})$ is dense in $L^2(\mathbb{A})$.

3. Real and *p*-adic MRAs and wavelet bases

3.1. Real MRA.. Now we recall the definitions of the real and p-adic multiresolution analysis.

Definition 3.1. (for example, see [53, 1.2]) A collection of closed spaces $V_j \subset L^2(\mathbb{R})$, $j \in \mathbb{Z}$, is called a *multiresolution analysis (MRA) in L*²(\mathbb{R}) if the following axioms hold

- (a) $V_i \subset V_{i+1}$ for all $j \in \mathbb{Z}$;
- (b) $\bigcup V_i$ is dense in $L^2(\mathbb{R})$;
- (c) $\bigcap_{j \in \mathbb{Z}}^{j \in \mathbb{Z}} V_j = \{0\};$
- (d) $f(\cdot) \in V_j \iff f(2\cdot) \in V_{j+1}$ for all $j \in \mathbb{Z}$;

(e) there exists a function $\phi \in V_0$ such that the system $\{\phi(\cdot - n), n \in \mathbb{Z}\}$ is an orthonormal basis for V_0 .

The function ϕ from axiom (e) is called *refinable* or *scaling*. One also says that a MRA is generated by its scaling function. The function ϕ is a solution of a special kind of functional equations which are called *refinement equations*. Their solutions are called refinable functions.

According to the standard MRA-scheme (see, e.g., [53, $\S1.3$]) for each j, we define (wavelet spaces)

$$(3.1) W_j = V_{j+1} \ominus V_j, \quad j \in \mathbb{Z}.$$

It is easy to see that

$$f \in W_j \iff f(2 \cdot) \in W_{j+1}, \text{ for all } j \in \mathbb{Z}$$

and $W_j \perp W_k, j \neq k$. Taking into account axioms (b) and (c), we obtain

(3.2)
$$L^{2}(\mathbb{R}) = \bigoplus_{j \in \mathbb{Z}} W_{j} = V_{0} \oplus \left(\bigoplus_{j \in \mathbb{Z}_{+}} W_{j}\right),$$

where $\mathbb{Z}_+ = \{0\} \cup \mathbb{N}$.

It is well known that the Haar refinement equation has the following form

(3.3)
$$\phi^{H}(t) = \phi^{H}(2t) + \phi^{H}(2t-1), \quad t \in \mathbb{R}$$

Its solution (the characteristic function $\chi_{[0,1]}(t)$ of the unit interval [0,1])

(3.4)
$$\phi^{H}(t) = \chi_{[0,1]}(t) = \begin{cases} 1, & t \in [0,1], \\ 0, & t \notin [0,1], \end{cases}, \quad t \in \mathbb{R},$$

generates the Haar MRA. In the framework of the Haar MRA one can construct the well-known Haar wavelet basis in $L^2(\mathbb{R})$

(3.5)
$$\psi_{jn}^{H}(t) = 2^{j/2} \psi^{H} \left(2^{j} t - n \right), \quad t \in \mathbb{R}, \quad j \in \mathbb{Z}, \quad n \in \mathbb{Z},$$

where

(3.6)
$$\psi^{H}(t) = \begin{cases} 1, & 0 \le t < \frac{1}{2}, \\ -1, & \frac{1}{2} \le t < 1, \\ 0, & t \notin [0, 1), \end{cases} = \chi_{[0, \frac{1}{2})}(t) - \chi_{[\frac{1}{2}, 1)}(t), \quad t \in \mathbb{R},$$

is called the *Haar wavelet function* (whose dyadic shifts and delations form the *Haar basis* (3.5)).

Using the second decomposition (3.2), instead of the Haar wavelet basis (3.5) in $L^2(\mathbb{R})$, one can consider the following wavelet basis

(3.7)
$$\begin{aligned} \phi^H(t-n) \in V_0, & n \in \mathbb{Z}, \\ \psi^H_{jn}(t) &= 2^{j/2} \psi^H (2^j t - n) \in W_j, & j \in \mathbb{Z}_+, n \in \mathbb{Z}, \quad t \in \mathbb{R}, \end{aligned}$$

where the wavelet function ψ^{H} is given by (3.6). Basis (3.7) will be called *modified Haar* basis.

3.2. *p*-adic MRA.. In the *p*-adic case a "natural" set of shifts for \mathbb{Q}_p is the following (see [39] and [57]):

(3.8)
$$I_p = \{a \in \mathbb{Q}_p : \{a\}_p = a\} = \{a = \frac{a_{-\gamma}}{p^{\gamma}} + \dots + \frac{a_{-1}}{p} : a_j \in \mathbb{F}_p, \ -\gamma \le j \le -1\}.$$

In [57], similarly to Definition 3.1, the following definition was introduced.

Definition 3.2. ([57]) A collection of closed spaces $V_j \subset L^2(\mathbb{Q}_p), j \in \mathbb{Z}$, is called a multiresolution analysis (MRA) in $L^2(\mathbb{Q}_p)$ if the following axioms hold

- (a) $V_j \subset V_{j+1}$ for all $j \in \mathbb{Z}$;
- (b) $\bigcup V_j$ is dense in $L^2(\mathbb{Q}_p)$;
- (c) $\bigcap_{j\in\mathbb{Z}} V_j = \{0\};$

(d) $f(\cdot) \in V_j \iff f(p^{-1} \cdot) \in V_{j+1}$ for all $j \in \mathbb{Z}$; (e) there exists a function $\phi \in V_0$ such that the system $\{\phi(\cdot - a), a \in I_p\}$ is an orthonormal basis for V_0 .

It follows immediately from axioms (d) and (e) that the functions $p^{j/2}\phi(p^{-j}\cdot -a)$, $a \in I_p$, form an orthonormal basis for $V_j, j \in \mathbb{Z}$.

According to the standard scheme (see, e.g., [53, §1.3]) for the construction of MRAbased wavelets, for each j, we define a space W_i (wavelet space) as the orthogonal complement of V_j in V_{j+1} , i.e.,

(3.9)
$$V_{j+1} = V_j \oplus W_j, \qquad j \in \mathbb{Z}.$$

It is not difficult to see that

(3.10)
$$f \in W_j \iff f(p^{-1} \cdot) \in W_{j+1}, \text{ for all } j \in \mathbb{Z}$$

and $W_j \perp W_k, j \neq k$. Taking into account axioms (b) and (c), we obtain

(3.11)
$$L^{2}(\mathbb{Q}_{p}) = \bigoplus_{j \in \mathbb{Z}} W_{j} = V_{0} \oplus \left(\bigoplus_{j \in \mathbb{Z}_{+}} W_{j}\right)$$

In view of (3.9) and axiom (a), we have

(3.12)
$$V_j = V_0 \oplus \left(\bigoplus_{0 \le k \le j-1} W_k\right) \quad j \in \mathbb{N}$$

If we now find a finite number of functions $\psi_{\nu} \in W_0, \nu \in A$, such that the system $\{\psi_{\nu}(x-a): a \in I_p, \nu \in A\}$ forms an orthonormal basis for W_0 , then, due to (3.10), (3.11), the system

$$\{p^{j/2}\psi_{\nu}(p^{-j}\cdot -a): a\in I_p, j\in\mathbb{Z}, \nu\in A\}$$

is an orthonormal basis for $L^2(\mathbb{Q}_p)$. Such functions $\psi_{\nu}, \nu \in A$, are called *wavelet* functions and the corresponding basis is called a *wavelet basis*.

In [32], the following conjecture was proposed: to construct a p-adic analog of the real Haar MRA, we can use the following *refinement equation*

(3.13)
$$\phi(x) = \sum_{r=0}^{p-1} \phi\left(\frac{1}{p}x - \frac{r}{p}\right), \quad x \in \mathbb{Q}_p,$$

whose solution (a refinable function) ϕ is the characteristic function $\Omega(|x|_p)$ of the unit disc. The above refinement equation is natural and reflects the fact that the unit disc $B_0 = \{x : |x|_p \leq 1\}$ is the union of p mutually disjoint discs of radius p^{-1} :

$$B_{0} = \bigcup_{r=0}^{p-1} B_{-1}(r), \quad \text{where} \quad B_{-1}(r) = \left\{ x : \left| x - r \right|_{p} \le \frac{1}{p} \right\}, \quad r \in \mathbb{F}_{p}.$$

This geometric fact is the result of the ultrametric structure of the p-adic field \mathbb{Q}_p .

The term *p*-adic Haar MRA is connected with the fact that for p = 2 the equation (3.13) is a 2-adic refinement equation

$$\phi(x) = \phi\left(\frac{1}{2}x\right) + \phi\left(\frac{1}{2}x - \frac{1}{2}\right), \quad x \in \mathbb{Q}_2,$$

which is a direct analog of the *refinement equation* (3.3) generating the Haar MRA and the Haar wavelet basis (3.5), (3.6) in the real case.

In [57], using the above relation (3.13) as a refinement equation, the Haar MRA was constructed for p = 2 (for the case $p \neq 2$, see [3, 8.4]). Here

(3.14)
$$V_j = \overline{\operatorname{span}\{p^{j/2}\phi(p^{-j}\cdot -a): a \in I_p\}}, \quad j \in \mathbb{Z},$$

where $\phi(x) = \Omega(|x|_p)$. In contrast to the Haar MRA in $L^2(\mathbb{R})$ (which generates only one wavelet basis (3.5), (3.6)), in the *p*-adic setting there exist *infinity many different Haar orthogonal bases* for $L^2(\mathbb{Q}_p)$ generated by the same MRA. Explicit formulas for generating *p*-adic wavelet functions were obtained in [57] for p = 2 and later in [35] for $p \neq 2$. Let us recall these results.

Theorem 3.1. ([35]) The set of all compactly supported wavelet functions is given by

(3.15)
$$\psi_{\mu}(x) = \sum_{\nu=1}^{p-1} \sum_{k=0}^{p^s-1} \alpha_{\nu;k}^{\mu} \psi_{\nu}^{(0)} \left(x - \frac{k}{p^s} \right), \quad \mu \in \mathbb{F}_p^{\times} = \{1, 2, \dots, p-1\},$$

(here supp $\psi_{\mu} \subset B_s(0), s \ge 0$), where

(3.16)
$$\psi_{\nu}^{(0)}(x) = \chi_p \Big(\frac{\nu}{p} x\Big) \Omega\big(|x|_p\big), \quad \nu \in \mathbb{F}_p^{\times}, \qquad x \in \mathbb{Q}_p,$$

are Kozyrev's wavelet functions, $s = 0, 1, 2, \ldots$, and

$$\alpha^{\mu}_{\nu;k} =$$

(3.17)
$$= \begin{cases} -p^{-s} \sum_{m=0}^{p^{s}-1} e^{-2\pi i \frac{-\frac{\nu}{p}+m}{p^{s}}} k \sigma_{\mu m} u_{\mu \mu}, & \mu = \nu, \\ -p^{-2s} \sum_{m=0}^{p^{s}-1} \sum_{n=0}^{p^{s}-1} e^{-2\pi i \frac{-\frac{\nu}{p}+m}{p^{s}}} k \frac{1-e^{2\pi i \frac{\mu-\nu}{p}}}{1-e^{2\pi i \frac{\mu-\nu}{p^{s}}}} \sigma_{\nu m} u_{\nu \mu}, & \mu \neq \nu, \end{cases}$$

 $|\sigma_{\mu m}| = 1, u_{\mu \nu}$ are entries of an arbitrary unitary $(p-1) \times (p-1)$ matrix $U; \mu, \nu \in \mathbb{F}_p^{\times}; k = 0, 1, \ldots, p^s - 1.$

Corollary 3.1. ([57]) Let p = 2. For every $s = 0, 1, 2, \ldots$ the function

(3.18)
$$\psi(x) = \sum_{k=0}^{2^s - 1} \alpha_k \psi^{(0)} \left(x - \frac{k}{2^s} \right),$$

is a compactly supported wavelet function (supp $\psi \subset B_s(0)$) for the Haar MRA if and only if

(3.19)
$$\alpha_k = 2^{-s} \sum_{r=0}^{2^s - 1} \gamma_r e^{-i\pi \frac{2r - 1}{2^s}k},$$

where the wavelet function $\psi^{(0)} := \psi_1^{(0)}$ is given by (3.16), $\gamma_r \in \mathbb{C}$ is an arbitrary constant such that $|\gamma_r| = 1$; $k, r = 0, 1, \ldots, 2^s - 1$.

According to the general wavelet theory, all dilations and shifts of wavelet functions (3.18), (3.19) or (3.15), (3.17) form a *p*-adic orthonormal Haar wavelet basis in $L^2(\mathbb{Q}_p)$:

(3.20)
$$\psi_{k;ja}(x) = p^{j/2}\psi_k(p^{-j}x - a), \quad k \in \mathbb{F}_p^{\times}, \ j \in \mathbb{Z}, \ a \in I_p.$$

In particular, Kozyrev's wavelet basis in $L^2(\mathbb{Q}_p)$ generated by wavelet functions (3.16) is the following

$$\psi_{k;ja}^{(0)}(x) = p^{j/2} \psi_k^{(0)}(p^{-j}x - a)$$

$$(3.21) \qquad = p^{j/2} \chi_p \Big(\frac{k}{p}(p^{-j}x - a)\Big) \Omega\Big(|p^{-j}x - a|_p\Big), \quad k \in \mathbb{F}_p^{\times}, \ j \in \mathbb{Z}, \ a \in I_p.$$

The wavelet functions (3.16) can be expressed in terms of the refinable function $\phi(x) = \Omega(|x|_p)$:

$$\psi_k^{(0)}(x) = \sum_{r=0}^{p-1} e^{2\pi i \frac{kr}{p}} \phi\left(\frac{1}{p}x - \frac{r}{p}\right), \quad k \in \mathbb{F}_p^{\times}, \quad x \in \mathbb{Q}_p.$$

It was proved in [1] that there are no orthogonal MRA based wavelet bases except for those described in Theorem 3.1 and Corollary 3.1.

To construct adelic wavelet bases, we will need the *p*-adic wavelet bases that contain the function $\phi(x)$. Therefore, taking into account the second decomposition (3.11), instead of the *p*-adic wavelet basis (3.20) in $L^2(\mathbb{Q}_p)$, we consider the following wavelet basis

(3.22)
$$\begin{aligned} \phi(x-a) \in V_0, & a \in I_p, \\ \psi_{k;ja}(x) = p^{j/2} \psi_k(p^{-j}x-a) \in W_j, & k \in \mathbb{F}_p^{\times}, \ j \in \mathbb{Z}_+, \ a \in I_p. \end{aligned}$$

where the wavelet functions ψ_k , $k \in \mathbb{F}_p^{\times}$, are given by (3.15)–(3.17). In particular, we have the following wavelet basis (3.23)

$$\phi(x-a) \in V_0, \qquad a \in I_p, \\ \psi_{k;ja}^{(0)}(x) = p^{j/2} \chi_p \Big(\frac{k}{p} (p^{-j}x-a)\Big) \Omega\Big(|p^{-j}x-a|_p\Big) \in W_j, \quad k \in \mathbb{F}_p^{\times}, \ j \in \mathbb{Z}_+, \ a \in I_p$$

The fact that bases (3.22) and (3.23) contain the refinable function $\phi(x) = \Omega(|x|_p)$ plays a crucial role in our construction (see below Sec. 5). Bases (3.22), (3.23) will be called modified Haar type bases.

Taking into account the second decomposition (3.11), we obtain the following statement.

Proposition 3.1. The restriction of basis (3.23) on \mathbb{Z}_p constitutes an orthonormal basis in $L^2(\mathbb{Z}_p, dx_p^0)$:

(3.24)
$$\begin{aligned} \phi(x) &= \Omega(|x|_p), \\ \widetilde{\psi}_{k;ja}^{(0)}(x) &= p^{j/2}\psi_k(p^{-j}x-a)\Omega(|p^{-j}x-a|_p)|_{\mathbb{Z}_p}, \, k \in \mathbb{F}_p^{\times}, \, a \in I_p^j, \, j \in \mathbb{Z}_+, \end{aligned}$$

where $dx_p^0 = dx_p \big|_{\mathbb{Z}_p}$ and $I_p^j = \{a \in I_p : p^j a \in \mathbb{Z}_p\}.$

Proof. It is easy to see that the restriction of Kozyrev's wavelets (3.21) is the following

$$\begin{split} \widetilde{\psi}_{k;ja}^{(0)}(x) &= p^{j/2} \chi_p \Big(\frac{k}{p} (p^{-j} x - a) \Big) \Omega \Big(|p^{-j} x - a|_p \Big) \Big|_{\mathbb{Z}_p} \\ &= \begin{cases} p^{j/2} \chi_p \Big(\frac{k}{p} (p^{-j} x - a) \Big) \Omega \Big(|p^{-j} x - a|_p \Big), & k \in \mathbb{F}_p, j \in \mathbb{Z}_+, a \in I_p^j, \\ p^{j/2} \phi(x), & j \leq -1, a = 0, \\ 0, & j \leq -1, a \neq 0. \end{cases} \end{split}$$

For $j \in \mathbb{Z}_+$, we have $\widetilde{\psi}_{k;ja}^{(0)}(x)|_{\mathbb{Z}_p} \neq 0$ only if $|p^{-j}x-a|_p \leq 1$. Let $a = a_{-\gamma}p^{-\gamma} + \cdots + a_{-1}p^{-1}$ and $x = \sum_{k=0}^{\infty} x_k p^k$. Then we have

$$p^{-j}x - a = (x_0p^{-j} + \dots + x_{j-1}p^{-1}) - (a_{-\gamma}p^{-\gamma} + \dots + a_{-1}p^{-1}) + \sum_{k=j}^{\infty} x_k p^{k-j}.$$

Since $p^{-j}x - a \in \mathbb{Z}_p$, the latter relation implies that $x_0p^{-j} + \cdots + x_{j-1}p^{-1} = a_{-\gamma}p^{-\gamma} + \cdots + a_{-1}p^{-1}$. Hence $p^ja \in \mathbb{Z}_p$.

Let us calculate the scalar product

$$\begin{split} \widetilde{\psi}_{k;ja}^{(0)}, \widetilde{\psi}_{k';j'a'}^{(0)} \\ &= p^{(j+j')/2} \int_{\mathbb{Z}_p} \chi_p \Big(\frac{k}{p} (p^{-j}x - a) \Big) \chi_p \Big(-\frac{k'}{p} (p^{-j'}x - a) \Big) \times \\ &\times \Omega \Big(|p^{-j}x - a|_p \Big) \Omega \Big(|p^{-j'}x - a'|_p \Big) \, dx, \quad k, k' \in \mathbb{F}_p, j, j' \in \mathbb{Z}_+, a, a' \in I_p. \end{split}$$

Here $(\widetilde{\psi}_{k;ja}^{(0)}, \widetilde{\psi}_{k';j'a'}^{(0)}) \neq 0$ only if $|p^{-j}x - a|_p \leq 1$ and $|p^{-j'}x - a'|_p \leq 1$. Just as above, we conclude that $a, a' \in I_p$ are such that $p^j a \in \mathbb{Z}_p$, $p^{j'}a' \in \mathbb{Z}_p$. Thus for $a \in I_p^j$, $a' \in I_p^{j'}$, we have supp $\Omega(|p^{-j}x - a|_p)$, supp $\Omega(|p^{-j'}x - a'|_p) \in \mathbb{Z}_p$. Now, taking into account orthonormality of wavelet functions $\psi_{k;ja}^{(0)}, \psi_{k';j'a'}^{(0)}$ in \mathbb{Q}_p , one can conclude that $(\widetilde{\psi}_{k;ja}^{(0)}, \widetilde{\psi}_{k';j'a'}^{(0)}) = \delta_{kk'}\delta_{jj'}\delta_{aa'}$, where $\delta_{ss'}$ is the Kronecker symbol.

Using Proposition 3.1 and Theorem 3.1, one can prove the following statement.

Proposition 3.2. The restriction of the basis (3.22) to \mathbb{Z}_p constitutes in $L^2(\mathbb{Z}_p, dx_p^0)$ the orthonormal basis

(3.25)
$$\begin{aligned} \phi(x) &= \Omega(|x|_p), \\ \widetilde{\psi}_{k;ja}(x) &= p^{j/2}\psi_k(p^{-j}x-a)\big|_{\mathbb{Z}_p}, \, k \in \mathbb{F}_p^{\times}, \, a \in I_p^{(j)}, \, j \in \mathbb{Z}_+ \end{aligned}$$

where the wavelet functions ψ_k , $k \in \mathbb{F}_p^{\times}$, are given by (3.15)–(3.17), and the set $I_p^{(j)}$ is defined in Remark 3.1.

Remark 3.1. We note that the set $I_p^{(j)}$, for which the restricted functions $\widetilde{\psi}_{k;ja}(x)$, $a \in I_p^{(j)}$ are nonzero, is finite and its description is similar to the description of the set I_p^j , though more complicated. For other wavelet bases, we will have a similar situation.

Consider the orthogonal projection

(3.26)
$$P^{[0]}: L^2(\mathbb{Q}_p, dx_p) \mapsto L^2(\mathbb{Z}_p, dx_p^0),$$

given by

(3.27)
$$L^2(\mathbb{Q}_p, dx_p) \ni f(x) \mapsto (P^{[0]}f)(x) = f(x)\Omega(|x|_p) \in L^2(\mathbb{Z}_p, dx_p^0).$$

In our case, the projections of *modified Haar* bases (3.22) and (3.23) to the subspace $L^2(\mathbb{Z}_p, dx_p^0)$ are also bases in $L^2(\mathbb{Z}_p, dx_p^0)$, where dx_p^0 is the restriction of the Haar measure dx_p on \mathbb{Q}_p to \mathbb{Z}_p . We note that the projections of some basis elements become zeros.

Proposition 3.3. The restriction of the p-adic Haar MRA (given by formula (3.14)) to the space $L^2(\mathbb{Z}_p)$ consists of the spaces

(3.28)
$$\widetilde{V}_j = V_j \big|_{\mathbb{Z}_p} = \overline{\operatorname{span}\{p^{j/2}\phi(p^{-j}\cdot -a), x \in \mathbb{Z}_p : a \in I_p\}}, \quad j \in \mathbb{Z}_+,$$

where
$$V_j|_{\mathbb{Z}_p} = V_0 = \operatorname{span}\{\phi(\cdot)\}$$
 for all $j \leq -1$, $\phi(x) = \Omega(|x|_p)$, and
 $\widetilde{V}_0 \subset \widetilde{V}_1 \subset \widetilde{V}_2 \subset \cdots$.

Here according to (3.9)–(3.11),

(3.29)
$$L^{2}(\mathbb{Z}_{p}) = \widetilde{V}_{0} \bigoplus \left(\bigoplus_{j \in \mathbb{Z}_{+}} \widetilde{W}_{j} \right),$$

where

$$\widetilde{W}_{j} = \widetilde{V}_{j+1} \ominus \widetilde{V}_{j}$$
$$= \overline{\operatorname{span}\left\{p^{j/2}\psi_{k}(p^{-j}x-a)\big|_{\mathbb{Z}_{p}} : k \in \mathbb{F}_{p}^{\times}, a \in I_{p}\right\}}, \quad j \in \mathbb{Z}_{+}$$

and the wavelet functions $\psi_k, k \in \mathbb{F}_p^{\times}$, are given by (3.15)–(3.17). In particular, we can use the wavelet functions (3.23).

4. Basis in the space $L^2(\mathbb{A}, dx)$

4.1. Infinite tensor product of Hilbert spaces. We recall [52] (see also [12, Ch.1, §2.3]) the definition of *infinite tensor product* $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} H_n$ of Hilbert spaces H_n , $n \in \mathbb{N}$.

Fix the sequence

(4.1)
$$e = (e^{(n)})_{n=1}^{\infty}, e^{(n)} \in H_n, ||e^{(n)}||_{H_n} = 1$$

called a stabilizing sequence. Denote by E the set of all stabilizing sequences. Fix an orthonormal basis (o.n.b) $(e_k^{(n)})_{k\in\mathbb{N}}$ in the Hilbert space H_n such that $e^{(n)} = e_1^{(n)}$, $n \in \mathbb{N}$. Let Λ be the set of multi-indices $\alpha = (\alpha_n)_{n\in\mathbb{N}}$, $\alpha_n \in \mathbb{N}$, $n \in \mathbb{N}$ such that $\alpha_n = 1$ for sufficiently big n depending on α .

By definition, the o.n.b. of the space $\otimes_{e,n\in\mathbb{N}}H_n$ with the stabilizing sequence e consists of all vectors $(e_{\alpha})_{\alpha\in\Lambda}$ of the following form

(4.2)
$$e_{\alpha} = e_{\alpha_1}^{(1)} \otimes e_{\alpha_2}^{(2)} \otimes \cdots \otimes e_{\alpha_n}^{(n)} \otimes e^{(n+1)} \otimes \cdots, \ \alpha \in \Lambda$$

where $\alpha_n = 1$ for sufficiently big *n* depending on α . An arbitrary element *f* of the space \mathcal{H}_e has the following form:

(4.3)
$$f = \sum_{\alpha \in \Lambda} f_{\alpha} e_{\alpha}, \text{ with } \|f\|_{\mathcal{H}_{e}}^{2} = \sum_{\alpha \in \Lambda} |f|_{\alpha}^{2} < \infty$$

and the scalar product $(f, g)_{\mathcal{H}_e}$ of two vectors $f, g \in \mathcal{H}_e$ has the following form:

(4.4)
$$(f,g)_{\mathcal{H}_e} = \sum_{\alpha \in \Lambda} f_\alpha \overline{g_\alpha}.$$

It will often be convenient for us (see [12, Ch.1,§ 2.3]) to represent the set Λ as the union of disjoint sets, each consisting of "finite" sequences. Namely, for an $\alpha \in \Lambda$ let $\nu(\alpha)$ denote the minimal $m = 1, 2, \ldots$ such that $\alpha_{m+1} = \alpha_{m+2} = \cdots = 1$. Let

(4.5)
$$\Lambda_n = \{ \alpha \in \Lambda : \nu(\alpha) = n \}, \quad n \in \mathbb{N}$$

Obviously, $\Lambda_n \bigcap \Lambda_m = \emptyset \ (n \neq m)$ and $\Lambda = \bigcup_{n=1}^{\infty} \Lambda_n$.

Example 4.1. ([12, Ch.1, §2.3, example 1]) Let $L^2(X_k, \mu_k)$, $k \in \mathbb{N}$ be the space of square integrable complex functions on the measurable space X_k with a probability measure μ_k . Choose the stabilizing sequence $e = (e^{(k)})_{k=1}^{\infty}$, where $e^{(k)}(x) \equiv 1$, $x \in X_k$, $k \in \mathbb{N}$. In this case we have

Theorem 4.1. The following two spaces are isomorphic:

(4.6)
$$\bigotimes_{e,k\in\mathbb{N}} L^2(X_k,\mu_k) \cong L^2\Big(\prod_{k\in\mathbb{N}} X_k,\otimes_{k\in\mathbb{N}}\mu_k\Big).$$

Example 4.2. If in the previous example the first *m* measures μ_k are not necessarily probability, i.e., $\mu_k(X_k) = \infty$, then we get following statement.

Theorem 4.2. The following two spaces are isomorphic:

(4.7)
$$L^{2}\left(\prod_{k=1}^{m} X_{k}, \otimes_{k=1}^{m} \mu_{k}\right) \otimes \left(\bigotimes_{e,k=m+1}^{\infty} L^{2}(X_{k},\mu_{k})\right) \cong L^{2}\left(\prod_{k\in\mathbb{N}} X_{k}, \otimes_{k\in\mathbb{N}} \mu_{k}\right)$$

Remark 4.1. Let us consider in the infinite tensor product space $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} H_n$ the infinite tensor product $A = \bigotimes_{n \in \mathbb{N}} A_n$ of bounded operators A_n acting in the space H_n , $n \in \mathbb{N}$. By definition, an operator A acts on the total set of well-defined vectors

$$f = \bigotimes_{n \in \mathbb{N}} f_n \text{ in the space } \mathcal{H}_e \text{ (see Lemma 4.1) in the following way:}$$

$$(4.8) \qquad Af := \bigotimes_{n \in \mathbb{N}} A_n f_n = A_1 f_1 \otimes \cdots \otimes A_m f_m \otimes \cdots,$$

if the latter expression is well-defined in \mathcal{H}_e . In the sequel we shall use the projections $P^{[m]}: \mathcal{H}_e \to H^{(m)}$, where the subspaces $H^{[m]} \subset \mathcal{H}_e$ are defined on the total family of vectors $f = \bigotimes_{n \in \mathbb{N}} f_n$ as follows:

$$\mathcal{H}_e \ni f = \bigotimes_{n \in \mathbb{N}} f_n \mapsto P^{[m]} f := \bigotimes_{n=1}^m f_n \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots \in H^{[m]},$$

and

$$H^{[m]} = \bigotimes_{n=1}^{m} H_n \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots$$

We shall denote by $A_{(m)}$ the projections $P^{[m]}AP^{[m]}$ of the operator A in \mathcal{H}_e onto the space $H^{[m]}$, i.e. define $A_{(m)} := P^{[m]}AP^{[m]}$. It is clear that

$$A_{(m)} = \bigotimes_{n=1}^{m} A_n \otimes \bigotimes_{n>m} Id_n,$$

where Id_n is the identity operator in H_n . We have the strong convergence $A_{(m)} \rightarrow A$ as $m \rightarrow \infty$ on a suitable set of vectors $f = \bigotimes_{n \in \mathbb{N}} f_n$. It is sufficient to estimate $\|(A - A_{(m)})f\|_{\mathcal{H}_e}$. For details, see [12, Ch. I, §2.7.].

As usual, we denote by $A|_X$ the restriction of an operator A acting in a Hilbert space H to the invariant subspace $X \subset H$.

4.2. Complete von Neumann product of infinitely many Hilbert spaces. The complete von Neumann tensor product $\mathcal{H} = \bigotimes_{k \in \mathbb{N}} H_k$ of Hilbert spaces H_k , $k \in \mathbb{N}$ is by definition the orthogonal sum of the spaces \mathcal{H}_e ([52], see also [12, Ch.1,§ 2.10])

(4.9)
$$\mathcal{H} = \bigoplus_{e \in E_{/\sim}} \mathcal{H}_e$$

over all possible equivalence classes $E_{/\sim}$ of stabilizing sequences e.

To be more precise, fix the space $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} H_n$. We define the vector $f = \bigotimes_{n \in \mathbb{N}} f^{(n)}$,

where $f^{(n)} \in H_n$, as the week limit (if it exists) in \mathcal{H}_e of the vectors

(4.10)
$$f[m] = f^{(1)} \otimes \cdots \otimes f^{(m)} \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots$$

as $m \to \infty$. Since the set span $\{e_{\alpha} : \alpha \in \Lambda\}$ is dense in \mathcal{H}_e , the week limit of the vectors f[m] exists if and only if: 1) the norms $||f[m]||_{\mathcal{H}_e}$ are uniformly bounded with respect to $m = 1, 2, \ldots$, and 2) $\lim_{m\to\infty} (f[m], e_{\alpha})_{\mathcal{H}_e}$ exists for each $\alpha \in \Lambda$. The following statements are proved in [12, Ch.I,§ 2.10].

Lemma 4.1. The strong limit of vectors (4.10) exists in \mathcal{H}_e , as $m \to \infty$, if and only if the product $\prod_{n=1}^{\infty} \|f^{(n)}\|_{H_n}$ and $\prod_{n=q}^{\infty} (f^{(n)}, e^{(n)})_{H_n}$ (q = 1, 2, ...) converge to finite numbers, and we have $\prod_{n=1}^{\infty} \|f^{(n)}\|_{H_n} = 0$ when $\prod_{n=q}^{\infty} (f^{(n)}, e^{(n)})_{H_n} = 0$ for each q.

Corollary 4.1. If $f^{(n)}$ in (4.10) are taken to be unit vectors, then the strong limit $\lim_{m\to\infty} f[m]$ exists if and only if for some $q = 1, 2, \ldots$ the product $\prod_{n=q}^{\infty} (f^{(n)}, e^{(n)})_{H_n}$ converges to a finite nonzero number.

Definition 4.1. (see [12, Ch. I, § 2.10, Theorem 2.9]) Consider the set E of all stabilizing sequences $e = (e^{(n)})_{n=1}^{\infty}$ of the form (4.1). A stabilizing sequence $l \in E$ is said to be equivalent to a stabilizing sequence $e \in E$ $(l \sim e)$, if each strong limit

$$l^{(1)'} \otimes l^{(2)'} \otimes \cdots = \lim_{m \to \infty} l^{(1)'} \otimes \cdots \otimes l^{(m)'} \otimes e^{(m+2)} \otimes e^{(m+2)} \otimes \cdots,$$

exists in \mathcal{H}_e , where $(l^{(n)'})_{k=1}^{\infty}$ is the sequence $(l^{(n)})_{k=1}^{\infty}$, "diluted" by the vectors $e^{(n)}$, i.e. each $l^{(n)'}$ is equal either to $l^{(n)}$ or to $e^{(n)}$. The relation \sim is an equivalence relation and we denote by $E_{/\sim}$ the set of all equivalent classes of E.

Theorem 4.3. Two infinite tensor products $\mathcal{H}_e = \bigotimes_{e,n\in\mathbb{N}} \mathcal{H}_n$ and $\mathcal{H}_l = \bigotimes_{l,n\in\mathbb{N}} \mathcal{H}_n$ corresponding to two equivalent stabilizing sequences $e = (e^{(n)})_{n=1}^{\infty}$ and $l = (l^{(n)})_{n=1}^{\infty}$ are isomorphic. For $e \not\sim l$, the spaces \mathcal{H}_e and \mathcal{H}_l are orthogonal.

4.3. On some subspaces of the infinite tensor products. Let X_n be some subspaces in the Hilbert spaces H_n , $n \in \mathbb{N}$ and let $(e_k^{(n)})_{k \in \mathbb{Z}}$ be an orthonormal basis in H_n such that the orthonormal basis in X_n is $(e_k^{(n)})_{k \in \mathbb{N}}$, and let $(e^{(n)})_{n \in \mathbb{N}}$ be the stabilizing sequence $e^{(n)} = e_1^{(n)}$, $n \in \mathbb{N}$. Note that

(4.11)
$$e^{(n)} = e_1^{(n)} \in X_n \subset H_n, \quad n \in \mathbb{N}.$$

Consider two spaces \mathcal{H}_e and $\mathcal{H}_e^l(X), l \in \mathbb{N}$, where

$$\mathcal{H}_e = \bigotimes_{e,n\in\mathbb{N}} H_n = H_1 \otimes H_2 \otimes \cdots \otimes H_l \otimes H_{l+1} \otimes H_{l+2} \otimes \cdots,$$

(4.12)
$$\mathcal{H}_{e}^{l}(X) = \bigotimes_{k=1}^{l} H_{k} \otimes \bigotimes_{e,k=l+1}^{\infty} X_{k} = H_{1} \otimes H_{2} \otimes \cdots \otimes H_{l} \otimes X_{l+1} \otimes \cdots$$

In the particular case $X_n = \mathbb{C}e^{(n)}$ we have

$$\mathcal{H}_{e}^{l}(X) = H_{1} \otimes H_{2} \otimes \cdots \otimes H_{l} \otimes \mathbb{C}e^{(l+1)} \otimes \mathbb{C}e^{(l+2)} \otimes \cdots, \quad l \in \mathbb{N}.$$

Lemma 4.2. For arbitrary subspaces X_n in H_n and a stabilizing sequence $(e^{(n)})_{n \in \mathbb{N}}$ with properties (4.11), we have

(4.13)
$$\mathcal{H}_e = \overline{\bigcup_{l \ge 1} \mathcal{H}_e^l(X)}.$$

Proof. The basis in the space \mathcal{H}_e is the following (see (4.2)):

(4.14)
$$e_{\alpha} = e_{\alpha_1}^{(1)} \otimes e_{\alpha_2}^{(2)} \otimes \cdots \otimes e_{\alpha_n}^{(n)} \otimes e^{(n+1)} \otimes \cdots, \ \alpha \in \Lambda = \bigcup_{n \in \mathbb{N}} \Lambda_n,$$

where Λ_n is defined by (4.5), and the basis in the space $\mathcal{H}_e^l(X)$ is

$$e_{\alpha}^{l} = e_{\alpha_{1}}^{(1)} \otimes e_{\alpha_{2}}^{(2)} \otimes \cdots \otimes e_{\alpha_{l}}^{(l)} \otimes e_{\alpha_{l+1}}^{(l+1)} \otimes \cdots \otimes e_{\alpha_{l+k}}^{(l+k)} \otimes e^{(l+k+1)} \otimes \cdots,$$

$$(4.15) \qquad \qquad \alpha \in \Lambda_{l} = \bigcup_{k \in \mathbb{N} \cup \{0\}} \Lambda_{l,k},$$

where

$$\Lambda_{l,k} = \{ \alpha \in \Lambda : \alpha_i \in \mathbb{Z}, \ 1 \le i \le l, \ \alpha_i \in \mathbb{N}, \ l+1 \le i \le l+k, \ \alpha_{l+k+i} = 1, \ i > 1 \}.$$

Obviously $\mathcal{H}_e \supseteq \mathcal{H}_e^l(X)$ for all $l \in \mathbb{N}$, hence $\mathcal{H}_e \supseteq \overline{\bigcup_{l \ge 1} \mathcal{H}_e^l(X)}$.

We show that $\mathcal{H}_e \subseteq \overline{\bigcup_{l\geq 1} \mathcal{H}_e^l(X)}$. It is sufficient to show that all vectors e_α of the form (4.14) are contained in the family of vectors e_α^l of the form (4.15). Indeed, if we take in (4.15) l = n and k = 0, we obtain all vectors e_α of the form (4.14).

4.4. Infinite tensor product of Hilbert spaces $L^2(\mathbb{Q}_p, dx_p)$. Using the results of the previous section we can define the infinite tensor product

(4.16)
$$\mathcal{H}_e(\mathbb{Q}) = \bigotimes_{e, p \in V_{\mathbb{Q}}} L^2(\mathbb{Q}_p, dx_p)$$

of Hilbert spaces $L^2(\mathbb{Q}_p, dx_p)$ with an arbitrary stabilizing sequence $e = (e^{(p)})_{p \in V_{\mathbb{Q}}} e^{(p)} \in L^2(\mathbb{Q}_p, dx_p)$.

Let $\{e_{\alpha_{\infty}}^{(\infty)}\}_{\alpha_{\infty}\in\mathbb{I}_{\infty}}$ and $\{e_{\alpha_{p}}^{(p)}\}_{\alpha_{p}\in\mathbb{I}_{p}}$ be arbitrary orthonormal bases in $L^{2}(\mathbb{Q}_{\infty}, dx_{\infty})$ and $L^{2}(\mathbb{Q}_{p}, dx_{p})$, respectively, where \mathbb{I}_{∞} and \mathbb{I}_{p} are the corresponding multi-indices. Then the orthonormal basis in the space $\mathcal{H}_{e}(\mathbb{Q})$ is the following:

(4.17)
$$e_{\alpha} = e_{\alpha_{\infty}}^{(\infty)} \otimes \bigotimes_{p \le m} e_{\alpha_p}^{(p)} \otimes \bigotimes_{m < p} e^{(p)}, \quad \alpha = (\alpha_{\infty}, \alpha_2, \alpha_3, \dots) \in \Lambda,$$

where $(e^{(p)})_{p \in V_{\mathbb{Q}}}$, $e^{(p)} \in L^2(\mathbb{Q}_p, dx_p)$, is some stabilizing sequence and Λ is defined below.

Definition 4.2. Define Λ as the set of multi-indices $\alpha = (\alpha_p)_{p \in V_Q}, \alpha_p \in \mathbb{I}_p$ such that $e_{\alpha_p}^{(p)} = e^{(p)}$ for sufficiently big p depending on α .

Fix the stabilizing sequence of the form

(4.18)
$$e = (e^{(p)})_{p \in V_{\mathbb{Q}}}$$
 $e^{(p)}(x_p) = \phi_p(x_p) = \Omega(|x_p|_p) \in L^2(\mathbb{Z}_p, dx_p^0), \quad p \in V_{\mathbb{Q}}.$

Using Lemma 4.2, for the stabilizing sequence (4.18), we obtain the following description of the space $L^2(\mathbb{A}, dx)$.

Lemma 4.3. We have

(4.19)
$$L^{2}(\mathbb{A}, dx) = \mathcal{H}_{e}(\mathbb{Q}) = \bigotimes_{e, p \in V_{\mathbb{Q}}} L^{2}(\mathbb{Q}_{p}, dx_{p}).$$

Proof. It is sufficient to use Lemma 4.2 and set $X_p = L^2(\mathbb{Z}_p)$, $H_p = L^2(\mathbb{Q}_p)$, $\mathcal{H}_e^p = L^2(\mathbb{A}^p, \mu^p) = \bigotimes_{q \leq p} L^2(\mathbb{Q}_q) \otimes \bigotimes_{q > p} L^2(\mathbb{Z}_q)$, $p \in V_{\mathbb{Q}}$, where $\mathbb{A}^p = \mathbb{Q}_{\infty} \times \prod_{p' \leq p} \mathbb{Q}_{p'} \times \prod_{p' > p} \mathbb{Z}_{p'}$, $p \in V_{\mathbb{Q}} \setminus \infty$ (see (4.27)), μ^p is the restriction on \mathbb{A}^p of the Haar measure dx on \mathbb{A} .

4.5. Some remarks about the stabilizations. Consider the space (4.16)

(4.20)
$$\mathcal{H}_e(\mathbb{Q}) = \bigotimes_{e, p \in V_{\mathbb{Q}}} L^2(\mathbb{Q}_p, dx_p),$$

where the stabilizing sequence $e = (e^{(p)})_{p \in V_{\mathbb{Q}}}$ has the special form $e^{(p)}(x_p) = \phi_p(x_p) \in L^2(\mathbb{Q}_p, dx_p)$. We show that the sequence e and $M^j_{V_{\mathbb{Q}}}e$ for $j \in \mathbb{Z}$ are not equivalent (see Definition 4.1), where

(4.21)
$$M_{V_{\mathbb{Q}}}^{j}e := (M_{p}^{j}\phi^{(p)})_{p\in V_{\mathbb{Q}}}$$
 and $M_{p}^{j}\phi^{(p)}(x_{p}) := p^{-j/2}\phi^{(p)}(p^{j}x_{p}).$

Indeed, we have

$$(\phi^{(p)}, M_p^j \phi^{(p)})_{L^2(\mathbb{Q}_p)} = \int_{\mathbb{Q}_p} \phi^{(p)}(x_p) p^{-j/2} \phi^{(p)}(p^j x_p) dx = \begin{cases} p^{-j/2}, & j \ge 0, \\ p^{j/2}, & j \le -1. \end{cases}$$

Set $c_{\infty} = (\phi^{(\infty)}, M^j_{\infty} \phi^{(\infty)})_{L^2(\mathbb{Q}_p \infty)}$, then we get

$$c_{\infty}^{-1}(e, M_{V_{\mathbb{Q}}}^{j}e) = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} (\phi^{(p)}, M_{p}^{j}\phi^{(p)})_{L^{2}(\mathbb{Q}_{p})} = \begin{cases} \left(\prod_{p \in V_{\mathbb{Q}} \setminus \infty} p\right)^{-j/2}, & j \ge 0, \\ \left(\prod_{p \in V_{\mathbb{Q}} \setminus \infty} p\right)^{j/2}, & j \le -1 \end{cases}$$

Thus in both cases the product is divergent, i.e., $(e, M_{V_Q}^j e) = 0$ hence the sequences e and $M_{V_Q}^j e$ are not equivalent!

A similar argument holds if we take as the stabilizing sequence some elements of Kozurev's basis $\psi_{\alpha_p}^{(p)}$ in the space $L^2(\mathbb{Q}_p)$, i.e., $e = (\psi_{\alpha_p}^{(p)})_{p \in V_{\mathbb{Q}}}$. In this case we immediately get $(\psi_{\alpha_p}^{(p)}, M_p^j \psi_{\alpha_p}^{(p)})_{L^2(\mathbb{Q}_p)} = 0$ for all $j \in \mathbb{Z}$. Hence the sequences $e = (\psi_{\alpha_p}^{(p)})_{p \in V_{\mathbb{Q}}}$ and $M_{V_{\mathbb{Q}}}^j e = (M_p^j \psi_{\alpha_p}^{(p)})_{p \in V_{\mathbb{Q}}}$ are not equivalent.

The above considerations force us to consider the following space

(4.22)
$$\mathcal{H}_{e,\mathbb{Z}}(\mathbb{Q}) = \bigoplus_{j \in \mathbb{Z}} \mathcal{H}_{M^j_{V_{\mathbb{Q}}}e}(\mathbb{Q})$$

to be sure that the operator $M_{V_{\mathbb{Q}}}^{j}$ is well-defined. Further we set

(4.23)
$$W_k = \mathcal{H}_{M_{V_Q}^k}(\mathbb{Q})$$

(4.24)
$$V_j = \bigoplus_{k=-\infty}^j W_k = \bigoplus_{k=-\infty}^j \mathcal{H}_{M_{V_{\mathbb{Q}}}^k}(\mathbb{Q}).$$

Theorem 4.4. The collection of closed spaces $V_j \subset \mathcal{H}_{e,\mathbb{Z}}(\mathbb{Q}), j \in \mathbb{Z}$, defined by (4.24) is a multiresolution analysis in $\mathcal{H}_{e,\mathbb{Z}}(\mathbb{Q})$, i.e., the following properties hold:

(a)
$$V_j \subset V_{j+1}$$
 for all $j \in \mathbb{Z}$;
(b) $\bigcup_{j \in \mathbb{Z}} V_j$ is dense in $\mathcal{H}_{e,\mathbb{Z}}(\mathbb{Q})$;
(c) $\bigcap_{j \in \mathbb{Z}} V_j = \{0\}$;
(d) $f(j) \in V_j$ (e) $f(M^j) \in V_j$

(d)
$$f(\cdot) \in V_j \iff f(M_{V_0}^j \cdot) \in V_{j+1}$$
 for all $j \in \mathbb{Z}$;

(e) there exists a basis $(e_i)_{i \in I}$ in the space W_0 such that $(M_{V_Q}^k e_i)_{i \in I}$ is a basis in the space $W_k, k \in \mathbb{Z}$.

4.6. Complete von Neumann product of Hilbert spaces $L^2(\mathbb{Q}_p, dx_p)$. It is natural to consider also the complete von Neumann product (see (4.25)) of Hilbert spaces $L^2(\mathbb{Q}_p, dx_p), p \in V_{\mathbb{Q}}$

(4.25)
$$\mathcal{H}(\mathbb{Q}) = \bigoplus_{e \in E_{/\sim}} \mathcal{H}_e(\mathbb{Q}),$$

where $\mathcal{H}_e(\mathbb{Q})$ is defined by (4.20). Perhaps this space could be useful in further development of analysis on the adele space \mathbb{A} . The space $\mathcal{H}_{e,\mathbb{Z}}(\mathbb{Q})$ defined by (4.22) certainly, contains the space $L^2(\mathbb{A})$ and is roughly speaking an infinite direct sum of non isomorphic copies of the space similar to $L^2(\mathbb{A})$.

4.7. **Basis on** $L^2(\mathbb{A}, dx)$. We can construct a basis in $L^2(\mathbb{A}, dx)$ using description (4.19) of the space $L^2(\mathbb{A}, dx)$ in Lemma 4.3. Without using this description we can proceed as follows.

We present A as the union of some subgroups \mathbb{A}^p : $\mathbb{A} = \bigcup_{p \in V_{\mathbb{Q}}} \mathbb{A}^p$ (see (4.27)), then we realize $L^2(\mathbb{A}^p, \mu^p)$ as the infinite tensor product of Hilbert spaces (see (4.30)) and use the following considerations.

Let (X, μ) be some measurable space and let X be the union of the measurable sets $X_n, X = \bigcup_{n \in \mathbb{N}} X_n$. Set $\mu_n = \mu|_{X_n}, \nu_n = \mu|_{X_n \setminus X_{n-1}}, n \ge 2, \quad \nu_1 = \mu_1$ and $X_0 = \emptyset$. In this case we have

(4.26)
$$L^2(X,\mu) = \bigoplus_{n \in \mathbb{N}} L^2(X_n \setminus X_{n-1},\nu_n) = \overline{\bigcup_{n \in \mathbb{N}} L^2(X_n,\mu_n)}$$

For any prime p we denote by p_{-} and p_{+} the previous and the following primes.

Define the subgroup \mathbb{A}^p of the group \mathbb{A} as follows:

(4.27)
$$\mathbb{A}^p = \mathbb{Q}_{\infty} \times \prod_{p' \le p} \mathbb{Q}_{p'} \times \prod_{p' > p} \mathbb{Z}_{p'}, \quad p \in V_{\mathbb{Q}}.$$

We have $\mathbb{A} = \bigcup_{p \in V_{\mathbb{O}}} \mathbb{A}^p$.

Let dx_p^0 be the restriction of the Haar measure dx_p on \mathbb{Q}_p to the subgroup \mathbb{Z}_p , let $d\mu^p$ be the restriction to the subgroup \mathbb{A}^p of the measure dx on \mathbb{A} , and let $d\nu_p$ be the restriction of the measure dx to the subset $\mathbb{A}^p \setminus \mathbb{A}^{p-}$. Then we have

(4.28)
$$d\mu^p = dx_{\infty} \otimes \left(\otimes_{p' \le p} dx_{p'} \right) \otimes \left(\otimes_{p' > p} dx_{p'}^0 \right)$$

and using (4.26) we get.

Lemma 4.4. For an arbitrary prime p we have the following description:

(4.29)
$$L^{2}(\mathbb{A}, dx) = L^{2}(\mathbb{A}^{p}, \mu^{p}) \oplus \left(\bigoplus_{p' > p} L^{2}(\mathbb{A}^{p'} \setminus \mathbb{A}^{p'_{-}}, \nu_{p'})\right) = \overline{\bigcup_{p \in V_{\mathbb{Q}}} L^{2}(\mathbb{A}^{p}, \mu^{p})}.$$

To construct a basis in the space $L^2(\mathbb{A}, dx)$ using the latter description, it is sufficient to construct a basis in the space $L^2(\mathbb{A}^p, \mu^p)$. The latter space $L^2(\mathbb{A}^p, \mu^p)$ is the infinite tensor product of the spaces $L^2(\mathbb{Q}_p, dx_p)$ and $L^2(\mathbb{Z}_p, dx_p^0)$:

(4.30)
$$L^{2}(\mathbb{A}^{p},\mu^{p}) = L^{2}(\mathbb{Q}_{\infty},dx_{\infty}) \otimes \left(\bigotimes_{p' \leq p} L^{2}(\mathbb{Q}_{p'},dx_{p'})\right) \otimes \left(\bigotimes_{e,p'>p} L^{2}(\mathbb{Z}_{p'},dx_{p'}^{0})\right)$$

Remark 4.2. Since the measures dx_p^0 on \mathbb{Z}_p are *probability measures*, decomposition (4.30) allows us to use an explicit description of the basis in the infinite tensor product $\otimes_{k \in \mathbb{N}} L^2(X_k, \mu_k)$ of Hilbert spaces $L^2(X_k, \mu_k)$ with probability measures $\mu_k, k \in \mathbb{N}$ (for example, see [12]) in order to construct a basis in the spaces $L^2(\mathbb{A}^p, \mu^p)$ and $L^2(\mathbb{A}, dx)$.

Suppose that $\{e_{\alpha_{\infty}}^{(\infty)}\}_{\alpha_{\infty}\in\mathbb{I}_{\infty}}$, and $\{e_{\alpha_{p}}^{(p)}\}_{\alpha_{p}\in\mathbb{I}_{p}}$ are *arbitrary* orthonormal bases in the spaces $L^{2}(\mathbb{Q}_{\infty})$ and $L^{2}(\mathbb{Q}_{p})$, respectively, \mathbb{I}_{∞} and \mathbb{I}_{p} are the corresponding indices. Fix the stabilizing sequence $(e^{(p)})_{p\in V_{\mathbb{Q}}}$, where $e^{(p)}$ is some element of the basis $\{e_{\alpha_{p}}^{(p)}\}_{\alpha_{p}\in\mathbb{I}_{p}}$ for all $p \in V_{\mathbb{Q}}$ such that $e^{(p)}(x) \in L^{2}(\mathbb{Z}_{p}, dx_{p}^{0})$. We can construct a basis in the space $L^{2}(\mathbb{A}^{p}, \mu^{p})$ for all $p \in V_{\mathbb{Q}}$ using decomposition (4.30). In such a way we can construct a basis in $L^{2}(\mathbb{A}, dx)$ using (4.29). As it was mentioned before, using Lemma 4.3 and the basis in the infinite tensor product we obtain.

Theorem 4.5. The vectors

(4.31)
$$e_{\alpha} = e_{\alpha_{\infty}}^{(\infty)} \otimes \bigotimes_{q \le m} e_{\alpha_{q}}^{(q)} \otimes \bigotimes_{m < q} e^{(q)}, \quad \alpha \in \Lambda = \bigcup_{m \in V_{\mathbb{Q}}} \Lambda_{m},$$

form the orthonormal basis in the space $L^2(\mathbb{A}, dx)$, where Λ is the set of multi-indices $\alpha = (\alpha_p)_{p \in V_{\mathbb{Q}}}$, such that $e_{\alpha_p}^{(p)} = e^{(p)}$ for sufficiently big p depending on α (see (4.2)), $\Lambda_p = \{\alpha \in \Lambda : \nu_a(\alpha) = p\}, p \in V_{\mathbb{Q}}, and \nu_a(\alpha)$ denote the minimal $p \in V_{\mathbb{Q}}$ such that $e_{\alpha_q}^{(q)} = e^{(q)}$ for q > p.

Corollary 4.2. The vectors

(4.32)
$$\widetilde{e}_{\alpha} = \bigotimes_{q \le m} e_{\alpha_q}^{(q)} \otimes \bigotimes_{m < q} e^{(q)}, \quad \alpha \in \widetilde{\Lambda},$$

form an orthonormal basis in the space $L^2(\widetilde{\mathbb{A}}, dx)$, where $\widetilde{\Lambda}$ is the set of multi-indices $\alpha = (\alpha_p)_{p \in V_{\mathbb{Q}} \setminus \{\infty\}}$ such that $\psi_{\alpha_p}^{(p)} = e^{(p)}$ for sufficiently big p depending on α , i.e.,

(4.33)
$$\widetilde{\Lambda} = \bigcup_{m \in V_{\mathbb{Q}} \setminus \infty} \widetilde{\Lambda}_m \quad and \quad \widetilde{\Lambda}_m = \{ \alpha \in \widetilde{\Lambda} : \nu_a(\alpha) = m \}, \quad m \in V_{\mathbb{Q}} \setminus \infty.$$

5. Adelic wavelet bases generated by tensor product of one-dimensional wavelet bases

To construct an adelic wavelet basis, we apply the above scheme from Subsec. 4.7 to the one-dimensional bases (3.5), (3.22), (3.25). In particular, one can use the Haar wavelet bases (3.5), (3.23), (3.24). Let

(5.1)
$$\psi_{\alpha_{\infty}}^{(\infty)}(x_{\infty}) = \psi_{j_{\infty}a_{\infty}}^{(\infty)}(x_{\infty}) = \psi_{j_{\infty}a_{\infty}}^{H}(x_{\infty}), \quad \alpha_{\infty} = (j_{\infty}, a_{\infty}) \in \mathbb{I}_{\infty} := (\mathbb{Z}, \mathbb{Z});$$

be the real Haar wavelet basis (3.5) in $L^2(\mathbb{R})$,

(5.2)
$$\begin{aligned} \psi_{\alpha_p}^{(p)}(x_p) &= \psi_{k_p;j_p a_p}^{(p)}(x_p), \quad \alpha_p = (k_p, j_p, a_p) \in \mathbb{I}_p^+ := (\mathbb{F}_p^{\times}, \mathbb{Z}_+, I_p), \\ \psi_{\alpha_p}^{(p)}(x_p) &= \psi_{a_p}^{(p)}(x_p) = \phi^{(p)}(x_p - a_p), \quad \alpha_p = (0, 0, a_p) \equiv a_p \in I_p; \end{aligned}$$

be the *p*-adic modified Haar wavelet basis (3.22) in $L^2(\mathbb{Q}_p)$, where I_p is defined by (3.8), and let

(5.3)
$$\begin{aligned} \psi_{\alpha_p}^{(p)}(x_p) &= \psi_{k_p; j_p a_p}^{(p)}(x_p), \quad \alpha_p = (k_p, j_p, a_p) \in \mathbb{I}_p^+ = (\mathbb{F}_p^{\times}, \mathbb{Z}_+, I_p), \\ \psi_{\alpha_p}^{(p)}(x_p) &= \psi_0^{(p)}(x_p) = \phi^{(p)}(x_p), \quad \alpha_p = (0, 0, 0) = 0; \end{aligned}$$

be the *p*-adic Haar wavelet basis (3.25) in $L^2(\mathbb{Z}_p)$, which is the restriction of the basis (5.2) to \mathbb{Z}_p . We recall that the restrictions of some basis elements (5.2) are equal to zero (see Remark 3.1 and Proposition 3.1). *Here and in what follows, for a stabilizing sequence we take* (4.18).

Hence we have the following bases in $L^2(\mathbb{Q}_p)$ and in $L^2(\mathbb{Z}_p)$, respectively,

(5.4)
$$\psi_{\alpha_p}^{(p)}, \quad \alpha_p \in \mathbb{I}_p = \mathbb{I}_p^+ \bigcup I_p; \qquad \psi_{\alpha_p}^{(p)}, \quad \alpha_p \in \mathbb{I}_p^{(0)} := \mathbb{I}_p^+ \bigcup \{0\} \subset \mathbb{I}_p,$$

where $\{\psi_{\alpha_p}^{(p)}\}_{\alpha_p \in \mathbb{I}_p^{(0)}}$ is the projection of $\{\psi_{\alpha_p}^{(p)}\}_{\alpha_p \in \mathbb{I}_p}$ on $L^2(\mathbb{Z}_p)$. Now, using Lemma 4.3, we obtain the following orthonormal wavelet basis in the space $L^2(\mathbb{A}, dx)$ (see (4.31)):

(5.5)
$$\Psi_{\alpha} = \Psi_{(\widehat{k},\widehat{j},\widehat{a})}(x) = \psi_{\alpha_{\infty}}^{(\infty)} \otimes \bigotimes_{q \le m} \psi_{\alpha_{q}}^{(q)} \otimes \bigotimes_{m < q} \phi^{(q)}, \quad \alpha \in \Lambda = \bigcup_{m \in V_{\mathbb{Q}}} \Lambda_{m}.$$

and the orthonormal wavelet basis in the space $L^2(\widetilde{\mathbb{A}}, dx)$ (see (4.32))

(5.6)
$$\widetilde{\Psi}_{\alpha} = \widetilde{\Psi}_{(\widehat{k},\widehat{j},\widehat{a})}(x) = \bigotimes_{2 \le q \le m} \psi_{\alpha_q}^{(q)} \otimes \bigotimes_{m < q} \phi^{(q)}, \quad \alpha \in \widetilde{\Lambda},$$

with the stabilization $\phi = (\phi^{(p)})_{p \in V_{\mathbb{Q} \setminus \infty}}$. Here $\hat{k} = (k_p)_p$, $\hat{j} = (j_p)_p$, $\hat{a} = (a_p)_p$ (see (5.2)).

Remark 5.1. We note that in [20, Sec.4, formula (4.1)] the basis in the space $L^2(\mathbb{A}, dx)$ is constructed as follows

$$\psi_{\alpha,\beta} = \psi_{n_0}^{\infty}(x_{\infty}) \prod_{p=2,3,\dots} \psi_{\alpha_p,\beta_p}(x_p),$$

where $\psi_{n_0}^{\infty}(x_{\infty})$ and $\psi_{\alpha_p,\beta_p}(x_p)$ are orthonormal eigenfunctions (of harmonic oscillators) in the real and p-adic cases, respectively, where $\psi_{\alpha_p,\beta_p}(x_p) = \Omega(|x_p|_p)$ for almost all p.

Let us introduce the adelic set of shifts $I_{\mathbb{A}}$ and dilations $\mathbb{Z}_{\mathbb{A}}$:

$$I_{\mathbb{A}} = \left\{ \widehat{a} = (a_{\infty}, a_2, \dots, a_p, \dots) : a_{\infty} \in I_{\infty} = \mathbb{Z}, \ a_p \in I_p, \right\}$$

(5.7) there exists n depending on \hat{a} such that $a_p = 0$ for all p > n.

$$\mathbb{Z}_{\mathbb{A}} = \{ \hat{j} = (j_{\infty}, j_2, \dots, j_p, \dots) : j_{\infty} \in \mathbb{Z}, \, j_p \in \mathbb{Z}, \, d_p \in$$

(5.8) there exists *n* depending on \hat{j} such that $j_p = 0$ for all p > n.

Set $m(\hat{a}) = \min\{p : a_q = 0 \text{ for all } q > p\}$ and $m(\hat{j}) = \min\{p : j_q = 0 \text{ for all } q > p\}$.

On the space $L^2(\mathbb{A})$ we define the operators of shifts $T_{\hat{a}}$, $\hat{a} \in \mathbb{Z}_{\mathbb{A}}$, and multi-dilation $M^{\hat{j}}, \hat{j} \in \mathbb{Z}_{\mathbb{A}}$ which are defined as the infinite tensor product of one-dimensional operators (4.21):

(5.9)
$$M^{\hat{j}} = M_{\infty}^{-j_{\infty}} \otimes \bigotimes_{2 \le q \le m} M_q^{j_q} \otimes \bigotimes_{q > m} Id_q,$$

where m = m(j) and Id_q is the identity operator on $L^2(\mathbb{Q}_q)$. We suppose that the operators M^j_{∞} and M^j_p in the spaces $L^2(\mathbb{Q}_{\infty})$ and $L^2(\mathbb{Q}_p)$ (on the functions $f^{(\infty)} \in L^2(\mathbb{Q}_{\infty})$ and $f^{(p)} \in L^2(\mathbb{Q}_p)$) act as follows

$$(M_{\infty}^{j}f^{(\infty)})(x_{\infty}) = 2^{-j/2}f^{(\infty)}(2^{-j}x_{\infty}), \quad (M_{p}^{j}f^{(p)})(x) = p^{-j/2}f^{(p)}(p^{j}x_{p}).$$

Let $f(x) = \bigotimes_{q \in V_{\mathbb{Q}}} f_q(x_q) \in L^2(\mathbb{A})$, where $f^{(\infty)} \in L^2(\mathbb{Q}_{\infty})$, $f^{(q)} \in L^2(\mathbb{Q}_q)$, $q \ge 2$, and $f^{(q)} = \phi^{(q)}$ for almost all q > m. Then

(5.10)
$$(T_{\widehat{a}}f)(x) \stackrel{def}{=} f(x-\widehat{a}) = f_{\infty}(x_{\infty}-a_{\infty}) \otimes \bigotimes_{2 \leq q} f_q(x_q-a_q),$$

$$(M^{\hat{j}}f)(x) \stackrel{def}{=} 2^{-j_{\infty}/2} f^{(\infty)}(2^{-j_{\infty}}x_{\infty}) \otimes$$

(5.11)
$$\otimes \Big(\bigotimes_{2 \le q \le m} q^{-j_q/2} f^{(q)}(q^{j_q} x_q)\Big) \otimes \Big(\bigotimes_{m < q} f^{(q)}(x_q)\Big),$$

$$\left(M^{\widehat{j}}T_{\widehat{a}}f\right)(x) \stackrel{def}{=} 2^{-j_{\infty}/2}f^{(\infty)}(2^{-j_{\infty}}x_{\infty}-a_{\infty})\otimes$$

(5.12)
$$\otimes \left(\bigotimes_{2 \le q \le m} q^{-j_q/2} f^{(q)}(q^{j_q} x_q - a_q)\right) \otimes \left(\bigotimes_{m < q} f^{(q)}(x_q)\right).$$

In the latter relation, we assume that $m = m(\hat{j}) = m(\hat{a})$.

Now one can obtain the adelic wavelet basis functions Ψ_{α} given in (5.5) by all *adelic* shifts and dilations of wavelet functions

(5.13)
$$\Psi_{(\hat{k},0,0)} := \psi_{(0,0)}^{(\infty)}(x_{\infty}) \otimes \bigotimes_{q \le m} \psi_{(k_q,0,0)}^{(q)}(x_q) \otimes \bigotimes_{m < q} \phi^{(q)}(x_q).$$

Namely,

$$\Psi_{\alpha}(x) =: \Psi_{(\widehat{k},\widehat{j},\widehat{a})}(x) = \left(M^{\widehat{j}}T_{\widehat{a}}\Psi_{(\widehat{k},0,0)}\right)(x)$$

where $\hat{k} = (0, k_2, ..., k_m, 0, 0, ...), k_q \in \mathbb{F}_q^{\times}; \ \hat{a} = (a_{\infty}, a_2, ..., a_m, 0, 0, ...), a_{\infty} \in \mathbb{Z}, a_q \in I_q; \ \hat{j} = (j_{\infty}, j_2, ..., j_m, 0, 0, ...), j_{\infty} \in \mathbb{Z}, j_q \in \mathbb{Z}_+; m \in V_{\mathbb{Q}} \setminus \infty.$ We stress that here $m(\hat{a}) = m(\hat{j}) = m(\hat{k}) = m.$

The wavelet systems (5.5) described above is of considerable interest and can be useful in various situations. Nevertheless, they do not possess all the advantages of one-dimensional wavelet bases, in particular, the localization property, which is of great value for applications. In the one-dimensional case, the real Haar basis functions with large indices $j_{\infty} \in \mathbb{Z}$ and the *p*-adic Haar basis functions with large indices $j_p \in \mathbb{Z}_+$, $p = 2, 3, \ldots$, have small supports. Thus the support of the multidimensional basis function may be large along one or several directions and small along some other directions. To avoid these drawbacks, we shall use a different approach. Its main idea is using the tensor product of the MRAs generating these bases instead of the tensor product of available wavelet bases.

6. Separable adelic MRA generated by tensor product of one-dimensional MRAs

Using the idea of constructing separable multidimensional MRA by means of the tensor product of one-dimensional MRAs (suggested by Y. Meyer [50] and S. Mallat [46], [47] (see, e.g., [53, §2.1]), we construct wavelet bases for the space $L^2(\mathbb{A}, dx)$.

We start with some general facts. Let we have a collection of closed subspaces $V_j^{(k)}$, $j \in \mathbb{Z}$, in a Hilbert space H_k , k = 1, 2, ..., m, having the properties (a), (b) and (c) of the Definition 3.1: (a) $V_j^{(k)} \subset V_{j+1}^{(k)}$ for all $j \in \mathbb{Z}$, k = 1, 2, ..., m; (b) $H_k = \bigcup_{j \in \mathbb{Z}} V_j^{(k)}$, (c) $\bigcap_{j \in \mathbb{Z}} V_j^{(k)} = \{0\}$. We use the following notations for the *finite tensor product*:

(6.1)
$$H^{[m]} = \bigotimes_{k=1}^{m} H_k, \ V_{j+1}^{(k)} = V_j^{(k)} \oplus W_j^{(k)}, \quad j \in \mathbb{Z}, \quad k = 1, 2, \dots, m,$$
$$V_{j+1}^{[m]} = \bigotimes_{k=1}^{m} V_{j+1}^{(k)} = \bigotimes_{k=1}^{m} (V_j^{(k)} \oplus W_j^{(k)}) = V_j^{[m]} \oplus W_j^{[m]},$$

where

$$V_j^{[m]} = \bigotimes_{k=1}^m V_j^{(k)}, \qquad W_j^{[m]} = \bigoplus_{(i_1, i_2, \dots, i_m; j) \in \{1, 2\}^m \setminus \{1, 1, \dots, 1\}} W_{(i_1, i_2, \dots, i_m; j)},$$

(6.2)
$$W_{(i_1,i_2,\dots,i_m;j)} = Z_{i_1,j}^{(1)} \otimes Z_{i_2,j}^{(2)} \otimes \dots \otimes Z_{i_m,j}^{(m)},$$

and

(6.3)
$$Z_{1,j}^{(k)} = V_j^{(k)}, \ Z_{2,j}^{(k)} = W_j^{(k)}, \ k = 1, 2, \dots, m.$$

For the *infinite tensor product* we keep the similar notations:

$$H = \mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} H_n, \quad H^{[m]} = \bigotimes_{k=1}^m H_k \bigotimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots, \quad m \in \mathbb{N}$$
$$V_{j+1}^{[m]} = \bigotimes_{k=1}^m V_{j+1}^{(k)} \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots$$
$$= \bigotimes_{k=1}^m (V_j^{(k)} \oplus W_j^{(k)}) \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots = V_j^{[m]} \oplus W_j^{[m]},$$

where

(6.4)
$$V_j^{[m]} = \bigotimes_{k=1}^m V_j^{(k)} \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots,$$

$$W_{j}^{[m]} = \bigoplus_{(i_{1},i_{2},\ldots,i_{m})\in\{1,2\}^{m}\backslash\{1,1,\ldots,1\}} W_{(i_{1},i_{2},\ldots,i_{m};j)}$$

(6.5)
$$W_{(i_1,i_2,\dots,i_m;j)} = Z_{i_1,j}^{(1)} \otimes Z_{i_2,j}^{(2)} \otimes \dots \otimes Z_{i_m,j}^{(m)} \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \dots$$

Define the space V_j as an appropriate limit of the spaces $V_j^{[m]}$

(6.6)
$$V_j = \overline{\operatorname{span}(V_j^{[m]} : m \in \mathbb{N})}.$$

Theorem 6.1. Let subspaces $\{V_j^{(n)}\}_{j\in\mathbb{Z}}$ of the space H_n , $n \in \mathbb{N}$ satisfy the following properties:

(a) $V_j^{(n)} \subset V_{j+1}^{(n)}$ for all $j \in \mathbb{Z}$ and $n \in \mathbb{N}$;

(b) $\bigcup_{j \in \mathbb{Z}} V_j^{(n)}$ is dense in H_n ; (c) $\bigcap_{j \in \mathbb{Z}} V_j^{(n)} = \{0\}$ in H_n . Then the subspaces $V_j, j \in \mathbb{Z}$ of the space $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} H_n$ defined by (6.6) satisfy the following properties:

- (aH) $V_j \subset V_{j+1}$ for all $j \in \mathbb{Z}$; (bH) $\bigcup_{j \in \mathbb{Z}} V_j$ is dense in $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} H_n$;
- (cH) $\bigcap_{i \in \mathbb{Z}} V_i = \{0\}.$

Proof. (aH) Since $V_j^{(n)} \subset V_{j+1}^{(n)}$ for all $n \in \mathbb{N}$, we conclude that $V_j^{[m]} \subset V_{j+1}^{[m]}$, $m \in \mathbb{N}$ (see (6.4)), hence $V_j = \operatorname{span}(V_j^{[m]} : m \in \mathbb{N}) \subset \operatorname{span}(V_{j+1}^{[m]} : m \in \mathbb{N}) = V_{j+1}$.

(bH) It is clear that for any fixed $m \in \mathbb{N}$ the space $\bigcup_{j \in \mathbb{Z}} V_j^{[m]}$ is dense in the space

$$H^{[m]} = \bigotimes_{k=1}^{m} H_k \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots$$

hence all vectors e_{α} , $\alpha \in \Lambda_m$ of the basis (4.14) are contained in $\bigcup_{j \in \mathbb{Z}} V_j^{[m]}$. To finish the proof, we note that all elements of the basis $e_{\alpha}, \alpha \in \Lambda$ in $\mathcal{H}_e = \bigotimes_{e,n \in \mathbb{N}} \mathcal{H}_n$ are given by (4.14), where $\Lambda = \bigcup_{m \in \mathbb{N}} \Lambda_m$.

(cH) 1. Let us suppose that for any $m \in \mathbb{N}$, we have

$$(6.7) \qquad \qquad \cap_{j \in \mathbb{Z}} V_j^{[m]} = \{0\}.$$

2. In this case we get

(6.8)
$$\bigcap_{j \in \mathbb{Z}} V_j = \bigcap_{j \in \mathbb{Z}} \operatorname{span}(V_j^{[m]} : m \in \mathbb{N}) \subset \operatorname{span}(\bigcap_{j \in \mathbb{Z}} V_j^{[m]} : m \in \mathbb{N}) = \{0\}.$$

1. To prove (6.7), we shall use (c). We set $W_i^{(n)} = V_{i+1}^{(n)} \ominus V_i^{(n)}$ for $j \in \mathbb{Z}$ and $n \in \mathbb{N}$. Then we get

(6.9)
$$H_n = \bigoplus_{k \in \mathbb{Z}} W_k^{(n)}, \text{ and } V_j^{(n)} = \bigoplus_{k=-\infty}^{j-1} W_k^{(n)}.$$

Let $(e_{\alpha_n}^{(n)})_{\alpha_n \in I_k^{(n)}}$ be an orthonormal basis in W_k , then by construction, $(e_{\alpha_n}^{(n)})_{\alpha_n \in I^{(n)}}$ is an orthonormal basis in H_n , where $I^{(n)} = \bigcup_{k \in \mathbb{Z}} I_k^{(n)}$. Define $\Lambda(I)$ (see Definition (4.2)) as the set of multi-indices $\alpha = (\alpha_n)_{n \in \mathbb{N}}, \ \alpha_n \in I^{(n)}$ such that $e_{\alpha_n}^{(n)} = e^{(n)}$ for sufficiently big *n* depending on α , here $e = (e^{(n)})_{n \in \mathbb{N}}$ is a stabilizing sequence. Using (6.9) and (6.4), we get

$$V_j^{[m]} = \left(\bigotimes_{k=1}^m V_j^{(k)}\right) \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots$$
$$= \bigotimes_{k=1}^m \left(\bigoplus_{n=-\infty}^{j-1} W_n^{(k)}\right) \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots$$

so the basis in the space $V_i^{[m]}$ can be chosen as e_{α} , $\alpha \in \Lambda_{m,j}(I)$, where

$$e_{\alpha} = e_{\alpha_1}^{(1)} \otimes e_{\alpha_2}^{(2)} \otimes \cdots \otimes e_{\alpha_m}^{(m)} \otimes e^{(m+1)} \otimes \cdots,$$

$$\Lambda_{m,j}(I) = \left(\alpha \in \Lambda(I) : \alpha = (\alpha_n)_{n \in \mathbb{N}}, \quad (\alpha_1, \dots, \alpha_m) \in \left(\bigcup_{k=-\infty}^{j-1} I_k^{(1)}, \dots, \bigcup_{k=-\infty}^{j-1} I_k^{(m)}\right)\right).$$

Note that $\bigcap_{j \in \mathbb{Z}} \Lambda_{m,j}(I) = \lim_{j \to -\infty} \left(\bigcup_{k=-\infty}^{j-1} I_k^{(1)}, \dots, \bigcup_{k=-\infty}^{j-1} I_k^{(m)} \right) = \emptyset$, this imply (6.7) (compare also with Lemma 6.1 below, where $\Lambda_{1,j}(I) = (-\infty, j-1] \cap \mathbb{Z}$). Indeed, for $f \in V_j^{[m]}$ we have

$$f = \sum_{\alpha \in \Lambda_{m,j}(I)} c_{\alpha} e_{\alpha} \quad \text{with} \quad \|f\|^2 = \sum_{\alpha \in \Lambda_{m,j}(I)} |c_{\alpha}|^2 < \infty$$

If $f \in \bigcap_{j \in \mathbb{Z}} V_j^{[m]}$ then

$$||f||^2 = \lim_{j \to -\infty} \sum_{\alpha \in \Lambda_{m,j}(I)} |c_{\alpha}|^2 = \sum_{\alpha \in \varnothing} |c_{\alpha}|^2 = 0.$$

2. To prove (6.8), using (6.4) we get

$$V_j^{[m]} = \otimes_{k=1}^m V_j^{(k)} \otimes e^{(m+1)} \otimes e^{(m+2)} \otimes \cdots, \quad m \in \mathbb{N},$$

$$V_j^{[1]} = V_j^{(1)} \otimes e^{(2)} \otimes \cdots \otimes e^{(m)} \otimes e^{(m+1)} \otimes \cdots,$$

$$V_j^{[2]} = V_j^{(1)} \otimes V_j^{(2)} \otimes \cdots \otimes e^{(m)} \otimes e^{(m+1)} \otimes \cdots,$$

$$V_j^{[m]} = V_j^{(1)} \otimes V_j^{(2)} \otimes \cdots \otimes V_j^{(m)} \otimes e^{(m+1)} \otimes \cdots.$$

Denote by $V_j^{(k)}(e) = \operatorname{span}(e^{(k)}, V_j^{(k)})$ for $k - 1 \in \mathbb{N}$, then we get

$$\operatorname{span}(V_j^{[1]}, V_j^{[2]}) = V_j^{(1)} \otimes V_j^{(2)}(e) \otimes e^{(3)} \otimes \cdots \otimes e^{(m)} \otimes e^{(m+1)} \otimes \cdots,$$

$$\operatorname{span}(V_j^{[1]}, V_j^{[3]}, V_j^{[2]}) = V_j^{(1)} \otimes V_j^{(2)}(e) \otimes V_j^{(3)}(e) \otimes \cdots \otimes e^{(m)} \otimes e^{(m+1)} \otimes \cdots,$$

$$\operatorname{span}(V_j^{[k]}, k = 2, \dots, m) = V_j^{(1)} \otimes V_j^{(2)}(e) \otimes V_j^{(3)}(e) \otimes \cdots \otimes V_j^{(m)}(e) \otimes e^{(m+1)} \otimes \cdots.$$

Since $V_j^{(n)}(e) = \operatorname{span}(e^{(n)}, V_j^{(n)}) \subset V_0^{(n)}$ for $j \leq 0$, we conclude that

$$V_{j} = \overline{\operatorname{span}(V_{j}^{[k]} : k \in \mathbb{N})} = \overline{V_{j}^{(1)} \otimes \bigotimes_{e,n \ge 2} V_{j}^{(n)}(e)} = V_{j}^{(1)} \otimes \bigotimes_{e,n \ge 2} V_{j}^{(n)}(e) \subset V_{j}^{(1)} \bigotimes_{e,n \ge 2} V_{0}^{(n)}.$$

 So

$$\bigcap_{j \in \mathbb{Z}} V_j = \bigcap_{j \in \mathbb{Z}} \overline{\operatorname{span}(V_j^{[k]} : k \in \mathbb{N})} \subset \\ \subset \bigcap_{j \in \mathbb{Z}} \left(V_j^{(1)} \otimes \bigotimes_{e,n \ge 2} V_0^{(n)} \right) = \bigcap_{j \in \mathbb{Z}} \left(V_j^{(1)} \right) \otimes \bigotimes_{e,n \ge 2} V_0^{(n)} = \{0\}.$$

Remark 6.1. We can give another definition of the spaces V_j , $j \in \mathbb{Z}$, namely, for fixed $n \in \mathbb{N}$ denote by $V_j(n)$, $j \in \mathbb{Z}$ the following family of spaces

(6.10)
$$V_j(n) = \overline{\operatorname{span}(V_j^{[m]} : m \in \mathbb{N}, \ m \ge n)}.$$

For this system of spaces, an analog of Theorem 6.1 holds. Moreover, in the proof of (cH) we have the following formula (see the end of the proof of the Theorem 6.1):

$$\bigcap_{j\in\mathbb{Z}}V_j(n)\subset\bigcap_{j\in\mathbb{Z}}\left(\bigotimes_{k=1}^nV_j^{(k)}\right)\otimes\bigotimes_{e,k>n}V_0^{(k)}=\{0\}.$$

We would like to present the following elementary lemma.

Lemma 6.1. Let we have the decomposition of the Hilbert space $H = \bigoplus_{k \in \mathbb{Z}} W_k$ and let $(e_{\alpha})_{\alpha \in I_k}$ be an orthonormal basis in the space W_k . Set $V_j := \bigoplus_{k=-\infty}^{j-1} W_k$, then $\bigcap_{j \in \mathbb{Z}} V_j = \{0\}$.

Proof. By construction, $(e_{\alpha})_{\alpha \in J}$ is an orthonormal basis in H, where $J = \bigcup_{k \in \mathbb{Z}} J_k$. We have for any $f \in H$

(6.11)
$$f = \sum_{\alpha \in J} c_{\alpha} e_{\alpha} = \sum_{k \in \mathbb{Z}} \left(\sum_{\alpha \in J_k} c_{\alpha} e_{\alpha} \right) \text{ and } \|f\|^2 = \sum_{k \in \mathbb{Z}} \left(\sum_{\alpha \in J_k} c_{\alpha}^2 \right) < \infty.$$

If $f \in V_n = \bigoplus_{k=-\infty}^{n-1} W_k$ for all $n \in \mathbb{Z}$, using (6.11), we conclude that

$$||f||^2 = \lim_{n \to -\infty} \sum_{k=-\infty}^{n-1} \left(\sum_{\alpha \in J_k} |c_{\alpha}|^2 \right) = 0.$$

Let $\{V_j^{(\nu)}\}_{j\in\mathbb{Z}}$ be one-dimensional MRA (see Definitions 3.1, 3.2), and let $\phi^{(\nu)}$ be its refinable function, $\nu \in V_{\mathbb{Q}}$. We introduce the spaces (see (6.4))

(6.12)
$$V_j^{[m]} = \bigotimes_{\nu \in \{\infty, 2, 3, \dots, m\}} V_j^{(\nu)} \otimes \phi^{(m_+)} \otimes \dots \subset L^2(\mathbb{A}, dx), \quad m \in V_{\mathbb{Q}}, \quad j \in \mathbb{Z}.$$

Now we introduce in $L^2(\mathbb{A}, dx)$ the delation operator M^j , which is defined by its projection $M_{(m)}^j$ on any subspace $V_j^{[m]}$ (for details, see Remark 4.1):

(6.13)
$$M_{(m)}^{j} = M_{\infty}^{-j} \otimes \bigotimes_{2 \le q \le m} M_{q}^{j} \otimes \bigotimes_{q > m} Id_{q}, \quad m \in V_{\mathbb{Q}} \setminus \infty,$$

where Id_q is the identity operator in $L^2(\mathbb{Q}_q)$. For $f \in L^2(\mathbb{A})$ we define $M^j_{(m)}f$ and $(M^{-j}_{(m)}T_{\widehat{a}})f$ similarly to (5.11) and (5.12), respectively, where $I_{\mathbb{A}}$ is given by (5.7). As before, we assume that $m = m(\widehat{a})$.

Theorem 6.2. Let $\{V_j^{(\nu)}\}_{j\in\mathbb{Z}}$ be the one-dimensional MRA (see Definitions 3.1, 3.2), and let $\phi^{(\nu)}$ be its refinable function, $\nu \in V_{\mathbb{Q}}$. Let

(6.14)
$$\Phi(x) = \phi^{(\infty)}(x_{\infty}) \otimes \phi^{(2)}(x_2) \otimes \cdots \otimes \phi^{(p)}(x_p) \otimes \cdots, \quad x \in \mathbb{A}$$

and

$$V_j = \overline{\operatorname{span}(V_j^{[m]} : m \in V_{\mathbb{Q}})}$$

(6.15)
$$= \overline{\operatorname{span}}\{\left(M_{(m)}^{-j}T_{\widehat{a}}\Phi\right)(\cdot): \widehat{a} \in I_{\mathbb{A}}, \ m(\widehat{a}) = m \in V_{\mathbb{Q}}\}, \quad j \in \mathbb{Z}.$$

Then

(6.16)
$$V_j = \bigotimes_{e;\nu \in V_{\mathbb{Q}}} V_j^{(\nu)} \subset L^2(\mathbb{A}, dx), \quad j \in \mathbb{Z}_+$$

(where $e = (\phi^{(\nu)})_{\nu \in V_{\mathbb{Q}} \setminus \infty}$ is the stabilization sequence (4.18) and the subspaces (6.15) satisfy the following properties:

(a) $V_j \subset V_{j+1}$ for all $j \in \mathbb{Z}$; (b) $\cup_{j \in \mathbb{Z}} V_j$ is dense in $L^2(\mathbb{A}, dx)$; (c) $\cap_{j \in \mathbb{Z}} V_j = \{0\}$;

(d) $f(\cdot) \in V_j^{[m]} \iff (M_{(m)}^{-1}f)(\cdot) \in V_{j+1}^{[m]}$ for all $j \in \mathbb{Z}$, for any $m \in V_{\mathbb{Q}}$; (e) the function $\Phi \in V_0$ is such that the system $\{\Phi(x - \hat{a}) : \hat{a} \in I_{\mathbb{A}}\}$ is an orthonormal basis for V_0 .

Proof. 1. For each $\nu \in V_{\mathbb{Q}}$, the system of functions $\{\phi^{(\nu)}(\cdot - a_{\nu})\}_{a_{\nu} \in I_{\nu}}$ is an orthonormal basis of the space $V_0^{(\nu)}$ (see axioms (e) in Definitions 3.1, 3.2). Using the construction of Sec. 4 and taking into account that any shift $\hat{a} = (a_{\infty}, a_2, \dots, a_p, \dots) \in I_{\mathbb{A}}$ is *finite* and $\phi^{(\nu)} \in V_0^{(\nu)} \subset L^2(\mathbb{Q}_{\nu}, dx_{\nu})$ is a stabilization sequence (4.18) in $L^2(\mathbb{A}, dx)$, we conclude that the system $\{\Phi(x - \hat{a}) : \hat{a} \in I_{\mathbb{A}}\}$ is an orthonormal basis in V_0 and

$$V_0 = \bigotimes_{e; \nu \in V_{\mathbb{Q}}} V_0^{(\nu)} \subset L^2(\mathbb{A}, dx)$$

Using definition (5.12) we get for $m(\hat{a}) = m$.

$$\left(M_{(m)}^{-j}T_{\widehat{a}}\Phi\right)(x) = 2^{j/2}\phi^{(\infty)}(2^{j}x_{\infty} - a_{\infty})\otimes$$

(6.17)
$$\otimes \left(\bigotimes_{\nu=2}^{m} \nu^{j/2} \phi^{(\nu)}(\nu^{-j} x_{\nu} - a_{\nu})\right) \otimes \phi^{(m_+)}(x_{m_+}) \otimes \dots \in V_j^{[m]} \subset V_j$$

for any $\hat{a} \in I_{\mathbb{A}}, m \in V_{\mathbb{Q}}$.

Since $\{\nu^{j/2}\phi^{(\nu)}(\nu^{-j}x_{\nu}-a_{\nu}): a_{\nu} \in I_{\nu}\}$ is the basis in $V_{i}^{(\nu)}$ and $\phi^{(\nu)} \in V_{0}^{(\nu)} \subset V_{i}^{(\nu)}$ for $j \in \mathbb{N}$, using (6.15) and (6.17), we conclude that

$$V_j = \bigotimes_{e; \nu \in V_{\mathbb{Q}}} V_j^{(\nu)} \subset L^2(\mathbb{A}, dx), \quad j \in \mathbb{N}.$$

Taking into account axioms (d) in Definitions 3.1, 3.2, one can see that Definitions (6.15), (6.13) imply that $f \in V_0^{[m]}$ if and only if $M_{(m)}^{-j} f(\cdot) \in V_j^{[m]}$, $j \in \mathbb{N}$, $m \in V_{\mathbb{Q}}$.

Thus representation (6.16) and properties (e) and (d) hold.

2. Due to [12, Ch. I, §2.2, Theorem 2.2], we immediately obtain that property (a)holds for $j \in \mathbb{N}$. For any $j \in \mathbb{Z}$ the property (a) holds by Theorem 6.1. Here we give an independent proof. According to (6.15), $V_0 = \overline{\operatorname{span}\{\Phi(x - \hat{a}) : \hat{a} \in I_{\mathbb{A}}\}},$ where $\Phi(x-\hat{a}) = \phi^{(\infty)}(x_{\infty}-a_{\infty}) \otimes \phi^{(2)}(x_2-a_2) \otimes \cdots \otimes \phi^{(p)}(x_p-a_p) \otimes \phi^{(p_+)}(x_{p_+}) \otimes \cdots$ Applying the refinement equation (3.13) to any factor of the above product and taking into account that the numbers $\frac{a_{\nu}}{\nu}, \frac{a_{\nu}}{\nu} + \frac{r}{\nu} \in I_{\nu}$ for all $a_{\nu} \in I_{\nu}$ and $r = 1, 2, ..., \nu - 1$, $\nu = 2, 3, ...,$ we conclude that $V_0 \subset V_1$. By Definition (6.15) of the spaces V_j and property (d), this yields property (a).

3. The property (b) holds by Theorem 6.1. Here we give an independent proof. It is clear that any element $g \in L^2(\mathbb{A}, dx)$ can be approximated by a finite linear combination of the basis elements $\Psi_{\alpha}(x) = \Psi_{(\hat{k};\hat{j}|\hat{a})}(x)$ of $L^2(\mathbb{A}, dx)$ given by (5.5). At the same time any element $\Psi_{\alpha}(x) = \Psi_{(\hat{k};\hat{i}|\hat{a})}(x)$ is a finite product of the one-dimensional basis elements (5.1) - (5.4):

$$\psi_{\alpha_{\infty}}^{(\infty)}(x_{\infty}) = \psi_{j_{\infty}a_{\infty}}^{(\infty)}(x_{\infty}) \in L^{2}(\mathbb{Q}_{\infty}), \qquad \alpha_{\infty} \in (\mathbb{Z}, \mathbb{Z}),$$

$$\psi_{\alpha_{q}}^{(q)}(x_{q}) = \psi_{(k_{q};j_{q}a_{q})}^{(q)}(x_{q}) \in L^{2}(\mathbb{Q}_{q}), \qquad \alpha_{p} \in \mathbb{I}_{p} = (\mathbb{F}_{p}^{\times}, \mathbb{Z}_{+}, I_{p}) \bigcup I_{p},$$

$$\psi_{\alpha_{q}}^{(q)}(x_{q}) = \psi_{(k_{q};j_{q}a_{q})}^{(q)}(x_{q}) \in L^{2}(\mathbb{Z}_{q}), \qquad \alpha_{p} \in \mathbb{I}_{p}^{(0)} = (\mathbb{F}_{p}^{\times}, \mathbb{Z}_{+}, I_{p}) \bigcup \{0\}.$$

In turn, it follows from the completeness property for each MRA $\{V_i^{(\nu)}\}_{i \in \mathbb{Z}}$ (see axioms (b) in Definitions 3.1, 3.2) that any element $f_{\infty} \in L^2(\mathbb{Q}_{\infty})$ and $f_{\nu} \in L^2(\mathbb{Q}_{\nu})$ can be approximated by a finite linear combination of the functions $2^{j/2}\phi^{(\infty)}(2^j \cdot -a_{\infty}), a_{\infty} \in I_{\infty}$, and $\nu^{j/2}\phi^{(\nu)}(\nu^{-j} \cdot -a_{\nu}), a_{\nu} \in I_{\nu}$, respectively, for sufficiently large j $(j \in \mathbb{Z}_+), \nu \in \{2, 3, \ldots, p\}$. Hence any basis element $\Psi_{\alpha}(x) = \Psi_{(\widehat{k};\widehat{j}|\widehat{a})}(x) \in L^2(\mathbb{A}, dx)$ can be approximated by an element of the space V_j for some sufficiently large j $(j \in \mathbb{Z}_+)$. This implies that the element g can be also approximated by an element of the space V_j for some j, i.e., the collection of all the spaces $V_j, j \in \mathbb{Z}$, is dense in $L^2(\mathbb{A}, dx)$.

4. The property (c) holds by Theorem 6.1. Here we give an independent proof. Let us prove that the intersection of all V_j does not contain a nonzero element. Assume that there exists an element $g \in L^2(\mathbb{A}, dx)$ such that $g \in V_j$ for all $j \in \mathbb{Z}$ and $||g|| \neq 0$. According to Sec. 5, the wavelet functions $\Psi_{\alpha} = \Psi_{(\widehat{k};\widehat{j}|\widehat{a})} \in L^2(\mathbb{A}, dx)$ given by (5.5) form an orthonormal basis in $L^2(\mathbb{A}, dx)$, then g can be represented as $g = \sum_{\alpha \in \Lambda} c_\alpha \Psi_\alpha \in V_j$ for all $j \in \mathbb{Z}$, where the coefficients c_α are independent of j. Here $\alpha = (\alpha_p)_{p \in V_Q} = (\widehat{k}; \widehat{j}|\widehat{a}),$ $\alpha_{\infty} \in (\mathbb{Z}, \mathbb{Z}), \alpha_p \in \mathbb{I}_p = (\mathbb{F}_p^{\times}, \mathbb{Z}_+, I_p) \bigcup I_p, p = 2, 3, \ldots$

Consider a basis element $\Psi_{\alpha} = \Psi_{(\widehat{k};\widehat{j}\ \widehat{a})}, \ \widehat{j} = (j_{\infty}, j_2, \dots, j_m, 0, 0, \dots)$. Let us take the index $j_{\alpha} = \min(0, j_{\infty}, j_2, j_3, \dots, j_m)$. For any wavelet function given by the first formula in (5.2), we have $\psi_{\alpha_p}^{(p)} = \psi_{k_p;j_p a_p}^{(p)} \perp V_0^{(p)}, \ j_p \in \mathbb{Z}_+, \text{ and } V_j^{(p)} \subset V_0^{(p)}, \ p = 2, 3, \dots$ In view of (3.1), (3.2) we have $\psi_{\alpha_{\infty}}^{(\infty)} = \psi_{j_{\infty}}^{(\infty)} \perp V_j^{(\infty)}$ for $j < j_{\infty}$. Consequently, the basis element $\Psi_{\alpha} = \Psi_{(\widehat{k};\widehat{j}\ \widehat{a})} \perp V_j, \ j < j_{\alpha}, \text{ i.e., } g \perp \Psi_{\alpha}$. In such a way we prove that $g \perp \Psi_{\alpha}$ for all indices $\alpha \in \Lambda$. The fact that the system $\{\Psi_{\alpha} : \alpha \in \Lambda\}$ is an orthonormal basis in $L^2(\mathbb{A}, dx)$ contradicts the assumption $||g|| \neq 0$.

The collection of the spaces V_j , $j \in \mathbb{Z}_+$, satisfying conditions (a)-(e) of Theorem 6.2 will be called the *adelic separable MRA in the space* $L^2(\mathbb{A})$. The function (6.14) is a refinable function of this MRA.

Next, following the standard finite-dimensional scheme (see [18, Ch. 10.1], [53, §2.1]), we define the wavelet space W_j as the orthogonal complement of V_j in V_{j+1} :

(6.18)
$$W_j = V_{j+1} \ominus V_j, \quad j \in \mathbb{Z}_+.$$

It follows from properties (a) and (b) of Theorem 6.2 that (cf. (3.2) and (3.11))

$$L^2(\mathbb{A}) = V_0 \oplus \bigoplus_{j \in \mathbb{Z}_+} W_j.$$

Let us present some evident formulas in $\bigotimes_{k=1}^{n} H_k$. Namely, we have

(6.19)
$$(a_{11} \oplus a_{12}) \otimes \cdots \otimes (a_{n1} \oplus a_{n2}) = \bigoplus_{(i_1, \dots, i_n) \in \{1, 2\}^n} a_{1i_1} \otimes \cdots \otimes a_{ni_n},$$

and similarly,

(6.20)
$$(\bigoplus_{k=1}^{m} a_{1k}) \otimes \cdots \otimes (\bigoplus_{k=1}^{m} a_{nk}) = \bigoplus_{(i_1, \dots, i_n) \in \{1, \dots, m\}^n} a_{1, i_1} \otimes \cdots \otimes a_{n, i_n}$$

Since according to (3.1), (3.2), (3.9)–(3.11), we have $V_{j+1}^{(\nu)} = V_j^{(\nu)} \oplus W_j^{(\nu)}$, setting $Z_{1,j}^{(\nu)} = V_j^{(\nu)}$ and $Z_{2,j}^{(\nu)} = W_j^{(\nu)}$, using (6.19), (6.20), and taking into account (6.17), we obtain

$$V_{j+1} = \bigotimes_{\nu \in V_{\mathbb{Q}}} V_{j+1}^{(\nu)} = \bigotimes_{\nu \in V_{\mathbb{Q}}} \left(V_{j}^{(\nu)} \oplus W_{j}^{(\nu)} \right)$$

= $V_{j} \oplus \left(\bigoplus_{p \in V_{\mathbb{Q}}} \bigoplus_{(i_{\infty}, i_{2}, \dots, i_{p}) \in \{1, 2\}^{\sharp(p)+1} \setminus \{1, 1, \dots, 1\}} W_{(i_{\infty}, i_{2}, \dots, i_{p}), j} \right), \quad j \in \mathbb{Z}_{+},$

where we set $\sharp(p) = \sharp\{2, 3, \dots, p\}$ ($\sharp(A)$ is the number of elements of the set A), and

(6.21)
$$W_{\hat{i}_{p},j}^{[p]} = W_{(i_{\infty},i_{2},\dots,i_{p}),j}^{[p]} = Z_{i_{\infty},j}^{(\infty)} \otimes Z_{i_{2},j}^{(2)} \otimes \dots \otimes Z_{i_{p},j}^{(p)} \otimes \phi^{(p_{+})} \otimes \dots$$

 $\hat{i}_p = (i_{\infty}, i_2, \dots, i_p) \in \{1, 2\}^{\sharp(p)+1} \setminus \{1, 1, \dots, 1\}$. Thus the space W_j given by (6.18) is a direct sum of infinite number of subspaces (6.21).

Thus we have

(6.22)
$$L^{2}(\mathbb{A}) = V_{0} \oplus \bigoplus_{j \in \mathbb{Z}_{+}} W_{j} = \bigoplus_{j \in \mathbb{Z}_{+}} \left(\bigoplus_{p \in V_{\mathbb{Q}} \ \widehat{i}_{p} \in \{1,2\}^{\sharp(p)+1} \setminus \{1,1,\dots,1\}} W_{\widehat{i}_{p},j}^{[p]} \right).$$

It is clear that an orthonormal basis for the space

$$W_0 = \bigoplus_{p \in V_{\mathbb{Q}}} \bigoplus_{\hat{i}_p = (i_{\infty}, i_2, \dots, i_p) \in \{1, 2\}^{\sharp(p)+1} \setminus \{1, 1, \dots, 1\}} W_{(i_{\infty}, i_2, \dots, i_p), 0}^{[p]}$$

where

$$W_{\hat{i}_{p},0}^{[p]} = W_{(i_{\infty},i_{2},\dots,i_{p}),0}^{[p]} = Z_{i_{\infty},0}^{(\infty)} \otimes Z_{i_{2},0}^{(2)} \otimes \dots \otimes Z_{i_{p},0}^{(p)} \otimes \phi^{(p_{+})} \otimes \dots$$

is formed by the shifts (with respect to $\hat{a} \in I_{\mathbb{A}}$) of the functions

(6.23)
$$\Psi_{\widehat{k},\widehat{i}_p} := \Psi_{\widehat{k},(i_{\infty},i_2,\ldots,i_p)} = \vartheta_{i_{\infty}}^{(\infty)} \otimes \vartheta_{i_2}^{(2)} \otimes \cdots \otimes \vartheta_{i_p}^{(p)} \otimes \phi^{(p_+)} \otimes \cdots,$$

where all factors in this product starting from a certain prime *p* constitute a *stabilization* sequence $(\phi^{(\nu)})_{\nu>p}$; $\vartheta_1^{(\infty)} = \psi^H$, $\vartheta_2^{(\infty)} = \phi^H$ (see (3.5), (3.4)); $\vartheta_1^{(\nu)} = \psi_{k_\nu}^{(\nu)}$, $k_\nu \in \mathbb{F}_{\nu}^{\times}$, $\vartheta_2^{(\nu)} = \phi^{(\nu)}$, (see (3.22), (3.25) or (3.23), (3.24)), $\nu = 2, 3, \ldots, p; \ p = 2, 3, \ldots; \ \hat{i_p} = (i_{\infty}, i_2, \ldots, i_p) \in \{1, 2\}^{\sharp(p)+1} \setminus \{1, 1, \ldots, 1\}.$

According to (6.22), an orthonormal Haar wavelet basis in $L^2(\mathbb{A})$ is formed by all shifts $\hat{a} \in I_{\mathbb{A}}$ of the refinable function (6.14)

$$\phi^{(\infty)}(x_{\infty})\phi^{(2)}(x_{2})\cdots\phi^{(p_{-})}(x_{p_{-}})\phi^{(p)}(x_{p})\phi^{(p_{+})}(x_{p_{+}})\cdots$$

and all shifts $\hat{a} \in I_{\mathbb{A}}$ and dilations $j \in \mathbb{Z}_+$ of wavelet functions (6.23):

(6.24)

$$\psi^{(\infty)}(x_{\infty})\psi^{(2)}_{k_{2}}(x_{2})\cdots\psi^{(p_{-})}_{k_{p_{-}}}(x_{p_{-}})\psi^{(p)}_{k_{p}}(x_{p})\phi^{(p_{+})}(x_{p_{+}})\cdots, \\
\psi^{(\infty)}(x_{\infty})\psi^{(2)}_{k_{2}}(x_{2})\cdots\psi^{(p_{-})}_{k_{p_{-}}}(x_{p_{-}})\phi^{(p)}(x_{p})\phi^{(p_{+},0)}_{k_{p_{+}}}(x_{p_{+}})\cdots, \\
\vdots \\
\phi^{(\infty)}(x_{\infty})\psi^{(2)}_{k_{2}}(x_{2})\cdots\psi^{(p_{-})}_{k_{p_{-}}}(x_{p_{-}})\psi^{(p)}_{k_{p}}(x_{p})\psi^{(p_{+},0)}_{k_{p_{+}}}(x_{p_{+}})\cdots, \\
\vdots \\
\psi^{(\infty)}(x_{\infty})\phi^{(2)}(x_{2})\cdots\phi^{(p_{-})}(x_{p_{-}})\phi^{(p)}(x_{p})\phi^{(p_{+})}(x_{p_{+}})\cdots, \\
\vdots \\
\phi^{(\infty)}(x_{\infty})\phi^{(2)}(x_{2})\cdots\phi^{(p_{-})}(x_{p_{-}})\psi^{(p)}_{k_{p}}(x_{p})\phi^{(p_{+})}(x_{p_{+}})\cdots, \\
\vdots \\
\phi^{(\infty)}(x_{\infty})\phi^{(2)}(x_{2})\cdots\phi^{(p_{-})}(x_{p_{-}})\psi^{(p)}_{k_{p}}(x_{p})\phi^{(p_{+})}(x_{p_{+}})\cdots,$$

where $p = 2, 3, \ldots$; the real Haar wavelet function $\psi^{(\infty)}(x_{\infty})$ and refinable function $\phi^{(\infty)}(x_{\infty}) = \Omega(x_{\infty})$ are given by (3.6) and (3.4), $x_{\infty} \in \mathbb{Q}_{\infty}$; the *r*-adic Haar wavelet functions $\psi_{k_r}^{(r)}(x_r)$, $k_r \in \mathbb{F}^{\times}$, are given by formulas (3.15)–(3.17), the refinable function $\phi^{(r)}(x_r) = \Omega(|x_r|_r)$, $x_r \in \mathbb{Q}_r$; $r = 2, 3, \ldots$

Thus, this wavelet basis has the form

(6.25)
$$T_{\widehat{a}}\Phi, \quad \Psi_{\widehat{k},\widehat{i}_p;j\,\widehat{a}} := (2 \cdot 2 \cdot 3 \cdots p)^{j/2} \left(M_{(p)}^{-j} T_{\widehat{a}} \right) \Psi_{\widehat{k},\widehat{i}_p}$$

where $\hat{i}_p = (i_{\infty}, i_2, \dots, i_p) \in \{1, 2\}^{\sharp(p)+1} \setminus \{1, 1, \dots, 1\}, \ \hat{k} = (k_2, \dots, k_p), \ k_s \in \mathbb{F}_s^{\times}, \ s = 2, 3, \dots, p; \ j \in \mathbb{Z}_+; \ \hat{a} = (a_{\infty}, a_2, \dots, a_p) \in I_{\mathbb{A}}, \ p \in V_{\mathbb{Q}}.$ We stress that here $p(\hat{a}) = p(j) = p(\hat{k}) = p$.

Due to Theorem 3.1 and [1], by formula (6.25) all Haar bases generated by adelic Haar MRA are described.

Corollary 6.1. Let (see (6.23))

(6.26)
$$\widetilde{\Psi}_{\widehat{k},\widehat{i}'_p}(x') := \widetilde{\Psi}_{\widehat{k},(i_2,\dots,i_m)} = \vartheta_{i_2}^{(2)} \otimes \dots \otimes \vartheta_{i_m}^{(m)} \otimes \phi^{(m_+)} \otimes \dots$$

The vectors (see (6.25))

(6.27)
$$T_{\widehat{a}'}\Phi, \quad \widetilde{\Psi}_{\widehat{k},\widehat{i}'_p;j\,\widehat{a}'} := (2\cdot 3\cdots p)^{j/2} \left(M^{-j}_{(p)}T_{\widehat{a}}\right) \Psi_{\widehat{k},\widehat{i}'},$$

form the orthonormal basis in the space $L^2(\widetilde{\mathbb{A}})$, where $\hat{i}'_p = (i_2, \ldots, i_p) \in \{1, 2\}^{\sharp(p-1)+1} \setminus \{1, \ldots, 1\}, \ \hat{k} = (k_2, \ldots, k_p), \ k_s \in \mathbb{F}_s^{\times}, \ s = 2, 3, \ldots, m; \ j \in \mathbb{Z}_+; \ \hat{a}' = (a_2, \ldots, a_p) \in I_{\mathbb{A}}, \ p \in V_{\mathbb{Q}} \setminus \infty.$ We stress that here $p(\hat{a}') = p(j) = p(\hat{k}') = p$.

7. LIZORKIN SPACES ON ADELES

7.1. **Real and** *p***-adic Lizorkin spaces.** Recall some facts from [43]–[45], [54, 2.], [55, §25.1]. Consider the following space

$$\Psi(\mathbb{R}) = \{\psi(\xi) \in \mathcal{S}(\mathbb{R}) : \psi^{(j)}(0) = 0, \, j = 0, 1, 2, \dots\},\$$

where $\mathcal{S}(\mathbb{R})$ is the real Schwartz space of tempered test functions. The space of functions

$$\Phi(\mathbb{R}) = \{ \phi : \phi = F[\psi], \ \psi \in \Psi(\mathbb{R}) \}.$$

is called the *real Lizorkin space of test functions*. The Lizorkin space can be equipped with the topology of the space $\mathcal{S}(\mathbb{R})$, which makes Φ a complete space [54, 2.2.], [55, §25.1.]. Since the Fourier transform is a linear isomorphism $\mathcal{S}(\mathbb{R})$ onto $\mathcal{S}(\mathbb{R})$, this space admits the following characterization: $\phi \in \Phi(\mathbb{R})$ if and only if $\phi \in \mathcal{S}(\mathbb{R})$ is orthogonal to polynomials, i.e.,

(7.1)
$$\int_{\mathbb{R}} x^n \phi(x) \, dx = 0, \quad n = 0, 1, 2, \dots$$

The space $\Phi'(\mathbb{R})$ is called the *real Lizorkin space of distributions*.

According to [4], [3, Ch. 7] the *p*-adic Lizorkin space of test functions is defined as $\Phi(\mathbb{Q}_p) = \{\phi : \phi = F[\psi], \psi \in \Psi(\mathbb{Q}_p)\}$, where $\Psi(\mathbb{Q}_p) = \{\psi(\xi) \in \mathcal{D}(\mathbb{Q}_p) : \psi(0) = 0\}$. The space $\Phi(\mathbb{Q}_p)$ equipped with the topology of the space $\mathcal{D}(\mathbb{Q}_p)$ is a complete space.

Lemma 7.1. ([4]) We have (a) $\phi \in \Phi(\mathbb{Q}_p)$ iff $\phi \in \mathcal{D}(\mathbb{Q}_p)$ and

(7.2)
$$\int_{\mathbb{Q}_p} \phi(x) \, dx = 0;$$

(b)
$$\phi \in \mathcal{D}_N^l(\mathbb{Q}_p) \cap \Phi(\mathbb{Q}_p)$$
 iff $\psi = F^{-1}[\phi] \in \mathcal{D}_{-l}^{-N}(\mathbb{Q}_p) \cap \Psi(\mathbb{Q}_p)$.

The topological dual of the space $\Phi(\mathbb{Q}_p)$ is the space of *p*-adic Lizorkin distributions $\Phi'(\mathbb{Q}_p)$. The space $\Phi'(\mathbb{Q}_p)$ can be obtained from $\mathcal{D}'(\mathbb{Q}_p)$ by "sifting out" constants (see [4]).

7.2. Adelic Lizorkin spaces. Consider the subspace $\Psi(\mathbb{A})$ of the space of test functions $\mathcal{S}(\mathbb{A})$ (see Subsec. 2.3) consisting of finite linear combinations of elementary functions (2.14) $\eta(\xi) = \prod_{p \in V_{\mathbb{Q}}} \eta_p(\xi_p)$ (where $V_{\mathbb{Q}}$ is defined by (2.2)) with properties (i) - (iii)of Definition 2.3, where $\eta_p(\xi_p) = \Omega(|\xi_p|_p)$ for all p > P) and such that

$$(iv) \left. \frac{d^s}{d\xi_{\infty}^s} \eta(\xi_{\infty}, \xi_2, \xi_3, \dots, \xi_{r_-}, \xi_r, \xi_{r_+}, \dots, \xi_P, \xi_{P_+}, \dots) \right|_{\xi_r = 0} = 0, \quad s = 0, 1, 2, \dots,$$

for all $r \in \{\infty, 2, 3, ..., P\}$, where $P = P(\eta)$ is the *parameter of finiteness* of an elementary function η .

Let $\Psi(\widetilde{\mathbb{A}})$ be the space of test functions $\mathcal{S}(\widetilde{\mathbb{A}})$ connected with the *non-Archimedean* part of adeles $\widetilde{\mathbb{A}}$ (see Subsec. 2.3) consisting of finite linear combinations of elementary functions $\eta(\xi') = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \eta_p(\xi_p)$ with properties (i) - (iii) of Definition 2.3 (here $\eta_p(\xi_p) = \Omega(|\xi_p|_p)$ for all p > P) and such that

$$(iv*) \eta(\xi_2,\xi_3,\ldots,\xi_{r_-},\xi_r,\xi_{r_+},\ldots,\xi_P,\xi_{P_+},\ldots)\big|_{\xi_r=0} = 0, \qquad r \in \{2,3,\ldots,P\},$$

where $P = P(\eta)$ is the parameter of finiteness of an elementary function η .

Definition 7.1. The spaces

$$\Phi(\mathbb{A}) = \left\{ \zeta : \zeta = F[\eta], \, \eta \in \Psi(\mathbb{A}) \right\} \quad \text{and} \quad \Phi(\widetilde{\mathbb{A}}) = \left\{ \zeta : \zeta = F[\eta], \, \eta \in \Psi(\widetilde{\mathbb{A}}) \right\}$$

are called the *adelic Lizorkin spaces of test functions*.

 $\Phi(\mathbb{A})$ can be equipped with the topology of the space $\mathcal{S}(\mathbb{A})$ (see (2.15)). Since the Fourier transform is a linear isomorphism $\mathcal{S}(\mathbb{A})$ onto $\mathcal{S}(\mathbb{A})$ (see [24, Ch. III,§2.2]), we have $\Phi(\mathbb{A}) \subset \mathcal{S}(\mathbb{A})$.

The Lizorkin spaces $\Phi(\mathbb{A})$ and $\Phi(\widetilde{\mathbb{A}})$ admit the following characterizations.

Lemma 7.2. (1) Any elementary function

$$\zeta(x) = \zeta_{\infty}(x_{\infty}) \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \zeta_p(x_p) \in \Phi(\mathbb{A})$$

(here $\zeta_p(x_p) = \Omega(|x|_p)$ for all p > P) if and only if $\zeta \in \mathcal{S}(\mathbb{A})$ and

(7.3)
$$\int_{\mathbb{Q}_r} x_r^s \zeta(x_\infty, x_2, x_3, \dots, x_{r_-}, x_r, x_{r_+}, \dots, x_P, x_{P_+}, \dots) \, dx_r = 0,$$

where if $r = \infty$, then s = 0, 1, 2, ..., and if r = 2, 3, ..., P, then s = 0. (2) Any elementary function

$$\widetilde{\zeta}(x') = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \zeta_p(x_p) \in \Phi(\widetilde{\mathbb{A}})$$

if and only if $\widetilde{\zeta} \in \mathcal{S}(\widetilde{\mathbb{A}})$ and

(7.4)
$$\int_{\mathbb{Q}_r} \widetilde{\zeta}(x_2, x_3, \dots, x_{r_-}, x_r, x_{r_+}, \dots, x_P, x_{P_+}, \dots) \, dx_r = 0, \quad r = 2, 3, \dots, P.$$

Remark 7.1. In view of condition (*iii*), similarly to (2.15), the spaces $\Phi(\mathbb{A})$ and $\Phi(\widehat{\mathbb{A}})$ admit a representation in the form

(7.5)
$$\Phi(\mathbb{A}) = \lim_{m \in V_{\mathbb{Q}} \setminus \infty} \operatorname{ind} \Phi^{[m]}(\mathbb{A}), \qquad \Phi(\widetilde{\mathbb{A}}) = \lim_{m \in V_{\mathbb{Q}} \setminus \infty} \operatorname{ind} \Phi^{[m]}(\widetilde{\mathbb{A}}),$$

where $\Phi^{[m]}(\mathbb{A}) \subset \Phi(\mathbb{A})$ and $\Phi^{[m]}(\widetilde{\mathbb{A}}) \subset \Phi(\widetilde{\mathbb{A}})$ are subspaces of the corresponding test functions with the *parameter of finiteness* $m, m \in V_{\mathbb{Q}}$. The representation (7.5) equip the spaces $\Phi(\mathbb{A})$ and $\Phi(\widetilde{\mathbb{A}})$ with the *inductive limit topology*.

The spaces $\Phi'(\mathbb{A})$ and $\Phi'(\widetilde{\mathbb{A}})$ are called the *adelic Lizorkin spaces of distributions*.

We define the Fourier transform of distributions $f \in \Phi'(\mathbb{A})$ and $g \in \Psi'(\mathbb{A})$ by the relations

(7.6)
$$\langle F[f], \eta \rangle = \langle f, F[\eta] \rangle, \quad \forall \eta \in \Psi(\mathbb{A}), \\ \langle F[g], \zeta \rangle = \langle g, F[\zeta] \rangle, \quad \forall \zeta \in \Phi(\mathbb{A}).$$

By definition, $F[\Phi(\mathbb{A})] = \Psi(\mathbb{A})$, $F[\Psi(\mathbb{A})] = \Phi(\mathbb{A})$, i.e., (7.6) give well defined objects. Moreover, we have $F[\Phi'(\mathbb{A})] = \Psi'(\mathbb{A})$, $F[\Psi'(\mathbb{A})] = \Phi'(\mathbb{A})$.

7.3. Adelic Lizorkin spaces and wavelets. It is well known that *p*-adic Haar wavelet functions (3.20) from $L^2(\mathbb{Q}_p)$ satisfy the condition (see [3, Ch. 8])

(7.7)
$$\int_{\mathbb{Q}_p} \psi_{k;ja}(x) \, dx = 0, \quad k \in \mathbb{F}_p^{\times}, \, j \in \mathbb{Z}, \, a \in I_p,$$

where $\psi_{k;ja}(x) = p^{j/2}\psi_k(p^{-j}x - a)$ and wavelet functions $\psi_k, k \in \mathbb{F}_p^{\times}$, are given by Theorem 3.1.

In view of the formula (see [3, Ch. 8] or [39])

$$\left(\Omega(|x|_p), \psi_{k;ja}(x)\right) = \begin{cases} p^{j/2}, & a = 0, \ j \le -1, \\ 0, & \text{otherwise}, \end{cases}$$

p-adic Haar wavelet functions (3.24) from $L^2(\mathbb{Z}_p)$ satisfy the following condition:

(7.8)
$$\int_{\mathbb{Z}_p} \widetilde{\psi}_{k;ja}^{(0)}(x) \, dx = \left(\Omega(|x|_p), \psi_{k;ja}(x)\right) = 0, \quad k \in \mathbb{F}_p^{\times}, \, j \in \mathbb{Z}_+, \, a \in I_p^j,$$

where I_p^j is defined in Proposition 3.1. It is easy to see that for the *p*-adic Haar wavelet functions (3.25) from $L^2(\mathbb{Z}_p)$ we also have

(7.9)
$$\int_{\mathbb{Z}_p} \widetilde{\psi}_{k;ja}(x) \, dx = 0, \quad k \in \mathbb{F}_p^{\times}, \, j \in \mathbb{Z}_+, \, a \in I_p^{(j)}$$

where $I_p^{(j)}$ is described in Remark 3.1.

Thus, taking into account relations (7.7)–(7.9), $\int_{\mathbb{Q}_p} \phi(x) dx = 1$, and Lemma 7.2, we conclude that the adelic wavelet functions (5.5) and (6.25) belong to the Lizorkin space $\Phi(\widetilde{\mathbb{A}})$ only if products (5.6) and (6.27) do not contain the functions $\phi^{(p)}(x_p - a_p)$ as factors for $p \leq m$.

8. CHARACTERIZATION OF THE ADELIC LIZORKIN SPACES IN TERMS OF WAVELETS

To construct adelic wavelet bases in $L^2(\mathbb{A})$, we used the real Haar basis (3.5) in $L^2(\mathbb{R})$, for which we have $\int_{\mathbb{R}} \psi_{jn}^H(t) dt = 0$. Since $\int_{\mathbb{R}} t^n \psi_{jn}^H(t) dt = 0$ does not hold for $n \in \mathbb{N}$, it is clear that $\psi_{jn}^H \notin \Phi(\mathbb{R})$. Therefore, in contrast to the *p*-adic case (see [3, 8.14]), the characterization of the adelic Lizorkin spaces in terms of wavelets is possible only for the space $\Phi(\widetilde{\mathbb{A}})$.

For simplicity here and in what follows we will suppose that the adelic wavelets (5.6) are constructed by using the one-dimensional wavelet basis (3.23) and its restriction (3.24) to \mathbb{Z}_p .

Lemma 8.1. Any test function $\widetilde{\zeta} \in \Phi(\widetilde{\mathbb{A}})$ can be represented in the form of a finite sum (8.1) $\widetilde{\zeta}(x') = \sum_{\alpha \in \widetilde{\mathbb{A}}} c_{\alpha} \widetilde{\Psi}_{\alpha}(x'), \quad x' = (x_2, x_3, \dots) \in \widetilde{\mathbb{A}},$

where $\widetilde{\Lambda}$ is the set of indices (4.33); c_{α} are constants; $\widetilde{\Psi}_{\alpha}(x')$ are wavelet functions (5.6)

$$\widetilde{\Psi}_{\alpha} = \widetilde{\Psi}_{(\widehat{k};\widehat{j}\ \widehat{a})} = \bigotimes_{2 \le q \le m} \psi_{\alpha_q}^{(q)} \otimes \bigotimes_{m < q} \phi^{(q)}, \quad \alpha \in \widetilde{\Lambda},$$

which do not contain the factors $\psi_{\alpha_q}^{(q)}(x_q) = \phi^{(q)}(x_q - a_q), a_p \in I_p$, (see (5.2), (5.3)) for $q \leq m$.

Proof. By definition of the space $\Phi(\widetilde{\mathbb{A}})$ (see Subsec. 7.2), it is sufficient to prove this lemma for the case of an elementary function $\widetilde{\zeta}(x') = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \zeta_p(x_p)$, where $\zeta_p(x_p) = \phi^{(p)}(x_p) = \Omega(|x_p|_p)$ for all p > P, and $P = P(\widetilde{\zeta})$ is the parameter of finiteness of an elementary function $\widetilde{\zeta}$. It is clear that $\widetilde{\zeta} \in L^2(\widetilde{\mathbb{A}})$. Then (8.1) holds. We will prove that only the finite number of $c_{\alpha} \neq 0$.

According to Subsec. 7.3, $\widetilde{\Psi}_{\alpha}(x') \in \Phi(\widetilde{\mathbb{A}})$ only if the product (5.6) does not contain functions $\phi^{(p)}(x_p - a_p)$ as factors for $p \leq m$. Now we need to calculate the coefficients

(8.2)
$$c_{\alpha} = c_{(\widehat{k};\widehat{j}\ \widehat{a})} = \left(\widetilde{\zeta}(x'), \widetilde{\Psi}_{\alpha}(x')\right) = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \left(\zeta_q(x_q), \psi_{\alpha_q}^{(q)}(x_q)\right)$$

where one-dimensional wavelet functions $\psi_{\alpha_q}^{(q)} = \psi_{k_q;j_q a_q}^{(q)}(x_q), \ \alpha_q = (k_q, j_q, a_q) \in \mathbb{I}_p^+ = (\mathbb{F}_q^{\times}, \mathbb{Z}_+, I_q)$, are given by (3.23), (3.24), (5.2), (5.3), and among them there are no wavelet functions $\phi^{(q)}(x_q - a_q)$.

1. Suppose that m > P. Then we have

$$c_{\alpha} = c_{(\widehat{k};\widehat{j}\ \widehat{a})} = \left(\widetilde{\zeta}(x'), \widetilde{\Psi}_{\alpha}(x')\right) = \prod_{2 \le q \le P} \left(\zeta_q(x_q), \psi_{\alpha_q}^{(q)}(x_q)\right) \times \prod_{P < q \le m} \left(\phi^{(q)}(x_q), \psi_{\alpha_q}^{(q)}(x_q)\right) \prod_{m < q} \left(\phi^{(q)}(x_q), \phi^{(q)}(x_q)\right).$$

According to (7.7) and (7.8), (7.9), we have

$$\int_{\mathbb{Z}_q} \psi_{\alpha_q}^{(q)}(x_q) \, dx_q = 0, \quad \text{for} \quad P < q \le m,$$

i.e., $c_{\alpha} = 0$.

2. Let m < P. Consider the part of product (8.2)

$$\prod_{n < q \le P} \left(\zeta_q(x_q), \, \phi^{(q)}(x_q) \right) = \prod_{m < q \le P} \left(\zeta_q^0(x_q), \, \phi^{(q)}(x_q) \right),$$

where $\zeta_q^0(x_q) := \zeta_q(x_q) |_{\mathbb{Z}_q} \in \Phi(\mathbb{Z}_q)$ and $\Phi(\mathbb{Z}_q)$ is the Lizorkin space of test functions with support in \mathbb{Z}_q (see Sec. 7.1). Using the Parseval-Steklov theorem, we obtain

$$\prod_{n < q \le P} \left(\zeta_q(x_q), \, \phi^{(q)}(x_q) \right) = \prod_{m < q \le P} \left(F[\zeta_q^0](\xi_q), \, F[\phi^{(q)}(x_q)](\xi_q) \right).$$

It is clear that $\zeta_q^0(x_q) \in \Phi(\mathbb{Z}_q)$ belongs to one of the spaces $\mathcal{D}_{N_q}^{l_q}(\mathbb{Z}_q), l_q \leq N_q \leq 0$, and satisfies condition (7.2). Then in view of (2.7) and Lemma 7.1, $F[\zeta_q^0] \in \Psi(\mathbb{Q}_q) \cap \mathcal{D}_{-l_q}^{-N_q}(\mathbb{Z}_q)$, and supp $F[\zeta_q^0] \subset B_{-l_q} \setminus B_{-N_q}, 0 \leq -N_q \leq -l_q$. At the same time, $F[\phi^{(q)}](\xi_q) = \phi^{(q)}(\xi_q)$. Thus supp $F[\zeta_q^0] \cap$ supp $\phi^{(q)} = \emptyset$ and, consequently, $\left(F[\zeta_q^0](\xi_q), F[\phi^{(q)}(x_q)](\xi_q)\right) = 0$, $m < q \leq P$. That is in this case $c_{\alpha} = 0$.

3. Let m = P. In this case, using the Parseval-Steklov theorem, one can rewrite the product (8.2) in the form

$$c_{\alpha} = c_{(\widehat{k};\widehat{j}\ \widehat{a})} = \left(\widetilde{\zeta}(x'), \widetilde{\Psi}_{\alpha}(x')\right)$$
$$= \prod_{2 \le q \le m} \left(\zeta_q(x_q), \psi_{\alpha_q}^{(q)}(x_q)\right) = \prod_{2 \le q \le m} \left(F[\zeta_q](\xi_q), F[\psi_{\alpha_q}^{(q)}](\xi_q)\right).$$

Let $\zeta_q \in \Phi(\mathbb{Q}_q)$ for $2 \leq q \leq p$. According to Definition 7.1, Lemma 7.1, and relation (2.7), any function $\zeta_q(x_q) \in \Phi(\mathbb{Q}_q)$ belongs to one of the spaces $\mathcal{D}_{N_q}^{l_q}(\mathbb{Q}_q)$ and satisfies condition (7.2), i.e., $F[\zeta_q] \in \Psi(\mathbb{Q}_q) \cap \mathcal{D}_{-l_q}^{-N_q}(\mathbb{Q}_q)$, and supp $F[\zeta_q] \subset B_{-l_q} \setminus B_{-N_q}$, $2 \leq q \leq p$. For the wavelet function (5.2),

$$\psi_{\alpha_q}^{(q)}(x_q) = \psi_{k_q; j_q a_q}^{(q)}(x_q)$$

(8.3)
$$= q^{j_q/2} \chi_q \Big(\frac{k_q}{q} (q^{-j_q} x_q - a_q) \Big) \Omega \Big(|q^{-j_q} x_q - a_q|_q \Big), \quad x_q \in \mathbb{Q}_q,$$

we have

(8.4)
$$F[\psi_{k_q;j_q a_q}^{(q)}](\xi_q) = q^{-j_q/2} \chi_q (q^{j_q} a_q \xi_q) \Omega \left(\left| \frac{k_q}{q} + q^{j_q} \xi_q \right|_q \right), \quad j_q \in \mathbb{Z}_+.$$

Taking into account that supp $F[\zeta_q] \subset B_{-l_q} \setminus B_{-N_q}$ and using formula (8.4), one can conclude that

(8.5)
$$\left(F[\zeta_q](\xi_q), F[\psi_{\alpha_q}^{(q)}](\xi_q)\right) \neq 0,$$

only if $q^{-N_q} \leq |\xi_q|_q \leq q^{-l_q}$ and $\frac{k_q}{q} + q^{j_q}\xi_q = \eta_q \in \mathbb{Z}_q$ for any $2 \leq q \leq p$. Since $\xi_q = q^{-j_q} \left(\eta_q - \frac{k_q}{q}\right)$ and $|\xi_q|_q = q^{j_q} \left|\eta_q - \frac{k_q}{q}\right|_q = q^{j_q+1}$, one can see that the product (8.5) is nonzero only for *finite number of indices j_q* such that

$$q^{-N_q} < |\xi_q|_q = q^{j_q+1} \le q^{-l_q}, \quad 2 \le q \le p.$$

Now we consider the scalar product $(F[\zeta_q](\xi_q), F[\psi_{\alpha_q}^{(q)}](\xi_q))$, where $\zeta_q \in \Phi(\mathbb{Z}_q)$ and

$$\psi_{\alpha_q}^{(q)}(x_q) = \psi_{k_q; j_q a_q}^{(q)}(x_q) = q^{j_q/2} \chi_q \Big(\frac{k_q}{q} (q^{-j_q} x_q - a_q) \Big) \Omega \Big(|q^{-j_q} x_q - a_q|_q \Big), \quad x_q \in \mathbb{Z}_q,$$

is a wavelet function given by (5.3), $p < q \le m$. Using the identity [64, VII.1], [39] (8.6) $\Omega(|p^j x_p - a_p|_p)\Omega(|p^{j'} x_p - a'_p|_p) = \Omega(|p^j x_p - a_p|_p)\Omega(|p^{j'-j} a_p - a'_p|_p), \quad j \le j',$ by explicit calculation we obtain

(8.7)
$$F[\psi_{\alpha_{q}}^{(q)}\big|_{\mathbb{Z}_{q}}](\xi_{q}) = q^{-j_{q}/2}\Omega(|q^{j_{q}}a_{q}|_{q})\chi_{q}(q^{j_{q}}a_{q}\xi_{q})\Omega(\left|\frac{k_{q}}{q} + q^{j_{q}}\xi_{q}\right|_{q}), \quad j_{q} \in \mathbb{Z}_{+}.$$

Next, using (8.7) and repeating the above calculation almost word for word, we find that

(8.8)
$$\left(F[\zeta_q](\xi_q), F[\psi_{\alpha_q}^{(q)}\big|_{\mathbb{Z}_q}](\xi_q)\right) \neq 0,$$

only for finite number of indices j_q , $p < q \leq m$.

Thus $c_{\alpha} = c_{(\widehat{k};\widehat{j}|\widehat{a})} \neq 0$ for finite number of indices $j_q, 2 \leq q \leq m$. Moreover, for all wavelet functions $\widetilde{\Psi}_{\alpha}(x')$, in the product (8.2) we have m = P.

Now we consider the product (8.2) again

$$c_{\alpha} = c_{(\widehat{k};\widehat{j}\ \widehat{a})} = \left(\widetilde{\zeta}(x'), \widetilde{\Psi}_{\alpha}(x')\right) = \prod_{2 \le q \le m} \left(\zeta_q(x_q), \psi_{\alpha_q}^{(q)}(x_q)\right),$$

where according to the above calculation, m = P. Now let calculate the expression $(\zeta_q(x_q), \psi_{\alpha_q}^{(q)}(x_q))$. Here the function $\zeta_q(x_q) \in \Phi(\mathbb{Q}_q)$ belongs to one of the spaces $\mathcal{D}_{N_q}^{l_q}(\mathbb{Q}_q)$ and $\psi_{\alpha_q}^{(q)}(x_q)$ is given by (8.3). Using identity (8.6) we have

$$\Omega\left(|q^{-j_q}x_q - a_q|_q\right)\Omega\left(|q^{N_q}x_q|_q\right) = \begin{cases} \Omega\left(|q^{-j_q}x_q - a_q|_q\right)\Omega\left(|q^{N+j_q}a_q|_q\right), & -j_q \le N_q, \\ \Omega\left(|q^{N_q}x_q|_q\right)\Omega\left(|a_q|_q\right), & -j_q > N_q. \end{cases}$$

This relation implies that $(\zeta_q(x_q), \psi_{\alpha_q}^{(q)}(x_q)) \neq 0$ at least for $q^{N+j_q}a_q \in \mathbb{Z}_q$ $(-j_q \leq N_q)$ or $a_q = 0$ $(-j_q > N_q)$. Since the set of indices j_q is finite, the set of the above indices a_q is also finite. Thus products (8.5), (8.8) are nonzero only for finite number of $a_q \in I_q$. \Box

Corollary 8.1. Any test function $\zeta \in \Phi(\mathbb{A})$ can be represented in the form of a finite sum

(8.9)
$$\zeta(x) = \sum_{\alpha \in \widetilde{\Lambda}} c_{\alpha}(x_{\infty}) \widetilde{\Psi}_{\alpha}(x'), \quad x = (x_{\infty}, x') \in \mathbb{A},$$

where $\widetilde{\Lambda}$ is the set of indices (4.33); $c_{\alpha}(x_{\infty}) \in \Phi(\mathbb{R})$ are some Lizorkin test functions; $\widetilde{\Psi}_{\alpha}(x') = \widetilde{\Psi}_{(\widehat{k};\widehat{j}\ \widehat{a})}(x')$ are wavelet functions (5.6) which do not contain factors $\psi_{\alpha_q}^{(q)}(x_q) = \phi^{(q)}(x_q - a_q), a_p \in I_p$, (see (5.2), (5.3)) for $q \leq m$.

Using standard results from the book [56], we obtain the following assertion.

Proposition 8.1. Any distribution $f \in \Phi'(\mathbb{A})$ can be realized in the form of an infinite sum:

(8.10)
$$f(x) = \sum_{\alpha \in \widetilde{\Lambda}} b_{\alpha}(x_{\infty}) \widetilde{\Psi}_{\alpha}(x'), \quad x = (x_{\infty}, x') \in \mathbb{A},$$

where $b_{\alpha}(x_{\infty}) \in \Phi'(\mathbb{R})$ are some real Lizorkin distributions; $\widetilde{\Psi}_{\alpha}(x')$ are wavelet functions (5.6).

Here any distribution $f \in \Phi'(\mathbb{A})$ is associated with representation (8.10), where the coefficients

(8.11)
$$b_{\alpha}(x_{\infty}) \stackrel{def}{=} \langle f(x_{\infty}, x'), \widetilde{\Psi}_{\alpha}(x') \rangle, \quad \alpha \in \widetilde{\Lambda}.$$

And vice versa, taking into account Corollary 8.1 and the orthonormality of the wavelet basis (5.6), any infinite sum (8.10) is associated with the distribution $f \in \Phi'(\mathbb{A})$, whose action on a test function $\zeta(x) \in \Phi(\mathbb{A})$ is defined as

(8.12)
$$\left\langle f,\zeta\right\rangle = \left\langle \sum_{\beta\in\tilde{\Lambda}} b_{\beta}(x_{\infty})\widetilde{\Psi}_{\beta}(x'), \sum_{\alpha\in\tilde{\Lambda}} c_{\alpha}(x_{\infty})\widetilde{\Psi}_{\alpha}(x')\right\rangle$$
$$= \sum_{\alpha\in\tilde{\Lambda}} \left\langle b_{\alpha}(x_{\infty}), c_{\alpha}(x_{\infty})\right\rangle,$$

where the latter sum is finite.

Proposition 8.2. Any distribution $\tilde{f} \in \Phi'(\tilde{\mathbb{A}})$ can be realized as an infinite sum of the form

(8.13)
$$\widetilde{f}(x') = \sum_{p} \sum_{\alpha \in \widetilde{\Lambda}} b_{\alpha} \widetilde{\Psi}_{\alpha}(x'), \quad x' \in \widetilde{\mathbb{A}},$$

where b_{α} are constants; $\widetilde{\Psi}_{\alpha}(x')$ are wavelet functions (5.6).

It is clear that in Lemma 8.1, Corollary 8.1, and Propositions 8.1, 8.2 one can use wavelet functions (6.27) instead of wavelet functions (5.6).

9. PSEUDO-DIFFERENTIAL OPERATORS ON ADELES

9.1. Real and *p*-adic fractional operators. Let us introduce a distribution from $\mathcal{S}'(\mathbb{R})$ (we keep the notation $|\cdot| = |\cdot|_{\infty}$.)

(9.1)
$$\kappa_{\alpha}^{(\infty)}(x) = \frac{|x|^{\alpha-1}}{\gamma_1(\alpha)}, \quad \alpha \neq -2s, \ \alpha \neq 1+2s, \quad s = 0, 1, 2, \dots, \quad x \in \mathbb{R},$$

called the *Riesz kernel*, where $|x|^{\alpha}$ is a homogeneous distribution of degree α defined in [54, Lemma 2.9.], [55, (25.19)], [23, Ch.I,§3.9.], and $\gamma_1(\alpha) = \frac{2^{\alpha} \pi^{\frac{1}{2}} \Gamma(\frac{\alpha}{2})}{\Gamma(\frac{1-\alpha}{2})}$. The Riesz kernel is an entire function of the complex variable α .

One can define the *Riesz kernel* (9.1) in the *real Lizorkin space of distributions* $\Phi'(\mathbb{R})$ (see [45], [54, Lemma 2.13.], [55, Lemma 25.2.]):

(9.2)
$$\kappa_{\alpha}^{(\infty)}(x) = \begin{cases} \frac{|x|^{\alpha-1}\Gamma(\frac{1-\alpha}{2})}{2^{\alpha}\sqrt{\pi}\Gamma(\frac{\alpha}{2})} = \frac{|x|^{\alpha-1}}{2\Gamma(\alpha)\cos(\frac{\pi\alpha}{2})}, & \alpha \neq -2s, \ \alpha \neq 1+2s, \\ (-1)^{s+1}\frac{|x|^{2s}\log|x|}{\pi(2s)!}, & \alpha = 1+2s, \\ (-1)^{s}\delta^{(2s)}(x), & \alpha = -2s, \ s \in \mathbb{Z}_{+}. \end{cases}$$

According to [23, Ch.II,§3.3.,(2)], [54, Lemma 2.13.], [55, Lemma 25.2.],

(9.3)
$$F[\kappa_{\alpha}^{(\infty)}(x)](\xi) = |x|^{-\alpha}, \quad x \in \mathbb{R}.$$

Define the real Riesz fractional operator D^{α}_{∞} on the Lizorkin space $\Phi(\mathbb{R})$ as a convolution

(9.4)
$$(D^{\alpha}_{\infty}\varphi)(x) = (\kappa^{(\infty)}_{-\alpha} * \varphi)(x) = \langle \kappa^{(\infty)}_{-\alpha}(\cdot), \varphi(x-\cdot) \rangle, \quad \varphi \in \Phi(\mathbb{R}),$$

where $\kappa_{\alpha}^{(\infty)}$ is given by (9.2). It is clear that [55, (25.2)]

(9.5)
$$(D^{\alpha}_{\infty}\varphi)(x) = F^{-1}[|\xi|^{\alpha}F[\varphi](\xi)](x), \quad \varphi \in \Phi(\mathbb{R}).$$

Lemma 9.1. The Lizorkin spaces of test functions $\Phi(\mathbb{R})$ and distributions $\Phi'(\mathbb{R})$ are invariant under the Riesz fractional operator (9.5) and $D^{\alpha}_{\infty}(\Phi'(\mathbb{R})) = \Phi'(\mathbb{R})$.

The *p*-adic fractional operator D_p^{α} was introduced on the space of distributions $\mathcal{D}'(\mathbb{Q}_p)$ in [58], [59, III.4.]. In [4], the fractional operator D_p^{α} was defined in the Lizorkin space of distributions $\Phi'(\mathbb{Q}_p)$ for all $\alpha \in \mathbb{C}$ by the following relations:

(9.6)
$$(D_p^{\alpha}f)(x) = F^{-1} [|\cdot|_p^{\alpha} F[f](\cdot)](x), \quad f \in \Phi'(\mathbb{Q}_p), \quad \alpha \in \mathbb{C}.$$

Representation (9.6) can be rewritten as a convolution

(9.7)
$$(D_p^{\alpha}f)(x) = (\kappa_{-\alpha}^{(p)} * f)(x) = \langle \kappa_{-\alpha}^{(p)}(\cdot), f(x-\cdot) \rangle, \quad f \in \Phi'(\mathbb{Q}_p), \quad \alpha \in \mathbb{C},$$

where (see [4])

(9.8)
$$\kappa_{\alpha}^{(p)}(x) = \begin{cases} \frac{|x|_{p}^{\alpha-1}}{\Gamma_{p}(\alpha)}, & \alpha \neq 0, \ 1, \\ \delta(x), & \alpha = 0, \\ -\frac{1-p^{-1}}{\log p} \log |x|_{p}, & \alpha = 1, \end{cases}$$

is the *Riesz kernel*, and

$$\Gamma_p(\alpha) \stackrel{def}{=} \int_{\mathbb{Q}_p} |x|_p^{\alpha-1} \chi_p(x) \, dx = \frac{1-p^{\alpha-1}}{1-p^{-\alpha}}$$

is the p-adic Γ -function (see [59, III, Theorem (4.2)], [64, VIII, (4.4)]).

Lemma 9.2. ([4]) The Lizorkin spaces of test functions $\Phi(\mathbb{Q}_p)$ and distributions $\Phi'(\mathbb{Q}_p)$ are invariant under the fractional operator (9.6) and $D_p^{\alpha}(\Phi'(\mathbb{Q}_p)) = \Phi'(\mathbb{Q}_p)$.

9.2. Adelic fractional operators on the Lizorkin spaces. Let us introduce on the space $\Phi(\mathbb{A})$ the *adelic fractional operator* $D^{\widehat{\gamma}}$ of order $\widehat{\gamma} = (\gamma_{\infty}, \gamma_2, \ldots, \gamma_p, \ldots) \in \mathbb{C}^{\infty}$, which is defined by its projection on any subspace of test functions $\Phi^{[m]}(\mathbb{A}) \subset \Phi(\mathbb{A})$ (for details, see Remark 4.1):

(9.9)
$$D^{\widehat{\gamma}}\Big|_{\Phi^{[m]}(\mathbb{A})} \stackrel{def}{=} D^{\widehat{\gamma}}_{(m)} = \bigotimes_{p \in \{\infty, 2, 3, \dots, m\}} D^{\gamma_p}_p \otimes \bigotimes_{p > m} Id_p, \quad m \in V_{\mathbb{Q}},$$

where $D_{\infty}^{\gamma_{\infty}}$ and $D_{p}^{\gamma_{p}}$ are defined by (9.4), (9.5) and (9.6), (9.7), respectively, Id_{p} is the identity operator in $\Phi(\mathbb{Q}_{p})$. Here the fractional operator $D_{(m)}^{\hat{\gamma}}$ is an infinite tensor product of one-dimensional operators.

If

$$\zeta(x) = \zeta_{\infty}(x_{\infty}) \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \zeta_p(x_p) \in \Phi^{[m]}(\mathbb{A})$$

is an elementary function, i.e., $\zeta_p(x_p) = \phi^{(p)}(\xi_p) = \Omega(|x_p|_p)$ for all $p > m(\zeta)$, then taking into account Definition (9.9) and (9.5), (9.6), we obtain

(9.10)
$$\left(D^{\widehat{\gamma}} \Big|_{\Phi^{[m]}(\mathbb{A})} \zeta \right)(x) = \left(D^{\widehat{\gamma}}_{(m)} \zeta \right)(x) = F^{-1} \left[\left| \cdot \right|^{\widehat{\gamma}}_{(m)} F[\zeta](\cdot) \right](x),$$

where

(9.11)
$$|\xi|_{(m)}^{\widehat{\gamma}} = \prod_{p \in \{\infty, 2, 3, \dots, m\}} |\xi_p|_p^{\gamma_p} \prod_{p > m} \phi^{(p)}(\xi_p), \quad \xi = (\xi_{\infty}, \xi_2, \dots, \xi_p, \dots) \in \mathbb{A}$$

is a symbol of the operator $D_{(m)}^{\hat{\gamma}}, m \in V_{\mathbb{Q}}$.

It is clear that $|\xi|^{\widehat{\gamma}}_{(m)}$ for all $\widehat{\gamma} = (\gamma_{\infty}, \gamma_2, \dots, \gamma_p, \dots) \in \mathbb{C}^{\infty}$ such that for all p > m, $\gamma_p = 0$ is a multiplier in the space of test functions $\Psi^{[m]}(\mathbb{Q}_p), m \in V_{\mathbb{Q}}$.

According to (9.4), (9.7), (2.10), relation (9.10) can be rewritten as

(9.12)
$$(D_{(m)}^{\widehat{\gamma}}\zeta)(x) = (\kappa_{-\widehat{\gamma};(m)} * \zeta)(x) = \langle \kappa_{-\widehat{\gamma};(m)}(\cdot), \zeta(x-\cdot) \rangle, \quad \zeta \in \Phi(\mathbb{A}),$$

where

$$\kappa_{\widehat{\gamma};(m)}(\xi) = \prod_{p \in \{\infty, 2, 3, \dots, m\}} \kappa_{\gamma_p}^{(p)}(\xi_p) \prod_{p > m} \phi^{(p)}(\xi_p)$$

is the *adelic Riesz kernel*, and $\kappa_{\gamma_{\infty}}^{(\infty)}$ and $\kappa_{\gamma_{p}}^{(p)}$ are given by (9.2) and (9.8), respectively. For $f \in \Phi'(\mathbb{A})$ we define the distribution $D^{\widehat{\gamma}} f$ by the relation

(9.13)
$$\langle D^{\widehat{\gamma}}f,\zeta\rangle \stackrel{def}{=} \langle f,D_{(m)}^{\widehat{\gamma}}\zeta\rangle, \quad \forall \zeta \in \Phi^{[m]}(\mathbb{A}), \quad m \in V_{\mathbb{Q}} \setminus \infty.$$

It is easy to see that Lemmas 9.1, 9.2 imply the following statements.

Lemma 9.3. The Lizorkin spaces of test functions $\Phi(\mathbb{A})$ and distributions $\Phi'(\mathbb{A})$ are invariant under the fractional operator $D^{\hat{\gamma}}_{(m)}$ and $D^{\hat{\gamma}}_{(m)}(\Phi'(\mathbb{A})) = \Phi'(\mathbb{A})$.

Proposition 9.1. The family of operators $\{D^{\widehat{\gamma}} : \widehat{\gamma} \in \mathbb{C}^{\infty}\}$ on the space of distributions $\Phi'(\mathbb{A})$ forms an abelian group: if $f \in \Phi'(\mathbb{A})$ then

$$\begin{split} D^{\widehat{\gamma}}D^{\beta}f &= D^{\widehat{\beta}}D^{\widehat{\gamma}}f = D^{\widehat{\gamma}+\widehat{\beta}}f, \\ D^{\widehat{\gamma}}D^{-\widehat{\gamma}}f &= f, \qquad \widehat{\gamma}, \ \widehat{\beta} \in \mathbb{C}^{\infty}. \end{split}$$

For the case $\hat{\gamma} = (\gamma_{\infty}, \gamma_2, \dots, \gamma_p, \dots) \in \mathbb{C}^{\infty}$ such that $\gamma_{\infty} = \gamma_2 = \dots = \gamma_p = \dots = \gamma \in \mathbb{C}$ we shall write D^{γ} instead of $D^{\hat{\gamma}}$. Similarly to (9.9), the *adelic fractional operator* D^{γ} of order $\gamma \in \mathbb{C}^{\infty}$ is defined by its projection on any subspace $\Phi^{[m]}(\mathbb{A}) \subset \Phi(\mathbb{A})$ (see (7.5)):

(9.14)
$$D^{\gamma}\Big|_{\Phi^{[m]}(\mathbb{A})} \stackrel{def}{=} D^{\gamma}_{(m)} = \bigotimes_{p \in \{\infty, 2, 3, \dots, m\}} D^{\gamma}_{p} \otimes \bigotimes_{p > m} Id_{p}, \quad m \in V_{\mathbb{Q}},$$

where D_{∞}^{γ} and D_{p}^{γ} are defined by (9.4), (9.5) and (9.6), (9.7), respectively, Id_{p} is the identity operator in $\Phi(\mathbb{Q}_{p})$. If

$$\zeta(x) = \zeta_{\infty}(x_{\infty}) \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \zeta_p(x_p) \in \Phi^{[m]}(\mathbb{A})$$

is an elementary function, i.e., $\zeta_p(x_p) = \phi^{(p)}(\xi_p) = \Omega(|x_p|_p)$ for all $p > m(\zeta)$, then taking into account Definition (9.14) and (9.5), (9.6), we obtain

(9.15)
$$\left(D^{\gamma}\Big|_{\Phi^{[m]}(\mathbb{A})}\zeta\right)(x) = \left(D^{\gamma}_{(m)}\zeta\right)(x) = F^{-1}\left[|\cdot|^{\gamma}_{(m)}F[\zeta](\cdot)\right](x),$$

where

(9.16)
$$|\xi|_{(m)}^{\gamma} = \prod_{p \in \{\infty, 2, 3, \dots, m\}} |\xi_p|_p^{\gamma} \prod_{p > m} \phi^{(p)}(\xi_p), \quad \xi = (\xi_{\infty}, \xi_2, \dots, \xi_p, \dots) \in \mathbb{A}.$$

is the symbol of the operator $D_{(m)}^{\gamma}$, $m \in V_{\mathbb{Q}}$.

9.3. One class of adelic pseudo-differential operators. On the adelic Lizorkin space of distributions $\Phi'(\mathbb{A})$ we introduce a pseudo-differential operators A with the symbol

$$\mathcal{A}(\xi) = \prod_{p \in V_{\mathbb{Q}}} \mathcal{A}_p(\xi_p), \quad \xi \in \mathbb{A},$$

where $\mathcal{A}_{\infty}(\xi_{\infty}) \in C^{\infty}(\mathbb{R} \setminus \{0\})$ and $\left|\frac{d^s}{d\xi_{\infty}^s} \mathcal{A}_{\infty}(\xi_{\infty})\right| \leq M_s |\xi_{\infty}|^{-m_s}$ $(m_s \geq 0)$ in a neighborhood of $\xi_{\infty} = 0$; $\mathcal{A}_p(\xi_p) \in \mathcal{E}(\mathbb{Q}_p \setminus \{0\}), p = 2, 3, \dots$

The adelic pseudo-differential operator A is defined by its projection on any subspace of test functions $\Phi^{[m]}(\mathbb{A}) \subset \Phi(\mathbb{A})$ (see (7.5)):

(9.17)
$$A\big|_{\Phi^{[m]}(\mathbb{A})} \stackrel{def}{=} A_{(m)} = F^{-1} \mathcal{A}_{(m)} F, \quad m \in V_{\mathbb{Q}},$$

where

(9.18)
$$\mathcal{A}_{(m)}(\xi) = \mathcal{A}_{\infty}(\xi_{\infty}) \prod_{p \in \{2,3,\dots,m\}} \mathcal{A}_{p}(\xi_{p}) \prod_{p > m} \phi^{(p)}(\xi_{p}), \quad \xi \in \mathbb{A},$$

is the symbol of the pseudo-differential operator $A_{(m)}$. If

$$\zeta(x) = \zeta_{\infty}(x_{\infty}) \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \zeta_p(x_p) \in \Phi^{[m]}(\mathbb{A})$$

is an elementary function, i.e., $\zeta_p(x_p) = \phi^{(p)}(\xi_p) = \Omega(|x_p|_p)$ for all $p > m(\zeta)$, then taking into account Definition (9.17), (9.18), we obtain

(9.19)
$$A|_{\Phi^{[m]}(\mathbb{A})}\zeta = A_{(m)}\zeta = \left(F^{-1}\mathcal{A}_{(m)}F\right)\zeta,$$

where the symbol $\mathcal{A}_{(m)}$ is defined by (9.18). For any test function $\zeta = \sum_{k=1}^{r} \zeta^{r} \in \Phi^{[m]}(\mathbb{A})$, where ζ^{r} are elementary functions, the operator A is defined by the relation

$$(A|_{\Phi^{[m]}(\mathbb{A})}\zeta)(x) = (A_{(m)}\zeta)(x) = \sum_{k=1}^{r} (A_{(m)}\zeta^{r})(x).$$

Now we define a conjugate pseudo-differential operator A^T on $\Phi(\mathbb{A})$. For any function $\zeta \in \Phi^{[m]}(\mathbb{A})$:

$$(A^T|_{\Phi^{[m]}(\mathbb{A})}\zeta)(x) \stackrel{def}{=} (A^T_{(m)}\zeta)(x) = F^{-1}\left[\overline{\mathcal{A}_{(m)}(-\xi)} F[\zeta](\xi)\right](x).$$

Then the operator A in the Lizorkin space of distributions $\Phi'(\mathbb{A})$ is defined in the usual way: for $f \in \Phi'(\mathbb{A})$ we have

(9.20)
$$\langle Af, \zeta \rangle = \langle f, A_{(m)}^T \zeta \rangle, \quad \forall \zeta \in \Phi^{[m]}(\mathbb{A}), \quad m \in V_{\mathbb{Q}} \setminus \infty.$$

It follows from the latter relation and (2.8) that

(9.21)
$$Af = \left(F^{-1}\mathcal{A}F\right)f, \qquad f \in \Phi'(\mathbb{A}).$$

Lemma 9.4. The Lizorkin spaces $\Phi(\mathbb{A})$ and $\Phi'(\mathbb{A})$ are invariant under the operator (9.21).

Proof. It is sufficient to prove this lemma for elementary functions. Let $\zeta \in \Phi^{[m]}(\mathbb{A})$ be an elementary function. Since the function $\mathcal{A}_{(m)}(\xi)$ defined by (9.18) is a multiplier in the space $\Psi^{[m]}(\mathbb{Q}_p)$, $m \in V_{\mathbb{Q}}$, both functions $F[\zeta](\xi)$ and $\mathcal{A}_{(m)}(\xi)F[\zeta](\xi)$ belong to $\Psi^{[m]}(\mathbb{A})$ and, consequently, $A\zeta \in \Phi^{[m]}(\mathbb{A})$, $m \in V_{\mathbb{Q}}$. Thus the pseudo-differential operator A (see (9.17), (9.19)) is well defined and the Lizorkin space $\Phi(\mathbb{A})$ is invariant under its action. Therefore, if $f \in \Phi'(\mathbb{A})$, then according to (9.20), (9.21), $Af = (F^{-1}\mathcal{A}F)f \in \Phi'(\mathbb{A})$, i.e., the Lizorkin space of distributions $\Phi'(\mathbb{A})$ is invariant under A.

The fractional operators (9.10) and (9.15) are particular cases of operator (9.19).

Similarly to the above constructions, we introduce the corresponding operators on the Lizorkin spaces $\Phi(\widetilde{\mathbb{A}})$ and $\Phi'(\widetilde{\mathbb{A}})$: the pseudo-differential operator

(9.22)
$$A_0\zeta = \left(F^{-1}\mathcal{A}_0\right)F\right)\zeta, \quad \zeta \in \Phi(\widetilde{\mathbb{A}}),$$

with the symbol

(9.23)
$$\mathcal{A}_0(\xi') = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} \mathcal{A}_p(\xi_p), \quad \xi' = (\xi_2, \dots, \xi_p, \dots) \in \widetilde{\mathbb{A}};$$

the fractional operator

(9.24)
$$(D_0^{\widehat{\gamma}}\zeta)(x') = F^{-1} [|\cdot|_0^{\widehat{\gamma}} F[\zeta](\cdot)](x'), \quad \zeta \in \Phi(\widetilde{\mathbb{A}}),$$

with the symbol

(9.25)
$$|\xi'|_0^{\widehat{\gamma}} = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} |\xi_p|_p^{\gamma_p}, \quad \xi' = (\xi_2, \dots, \xi_p, \dots) \in \widetilde{\mathbb{A}};$$

and the fractional operator

(9.26)
$$(D_0^{\gamma}\zeta)(x') = F^{-1} [|\cdot|_0^{\gamma} F[\zeta](\cdot)](x'), \quad \zeta \in \Phi(\widetilde{\mathbb{A}}),$$

with the symbol

(9.27)
$$|\xi'|_0^{\gamma} = \prod_{p \in V_{\mathbb{Q}} \setminus \infty} |\xi_p|_p^{\gamma}, \quad \xi' = (\xi_2, \dots, \xi_p, \dots) \in \widetilde{\mathbb{A}}.$$

9.4. On the eigenfunctions of adelic pseudo-differential operators.

Theorem 9.1. Let A_0 be a pseudo-differential operator (9.22) with symbol (9.23), and

$$\widetilde{\Psi}_{\alpha}(x') = \widetilde{\Psi}_{(\widehat{k};\widehat{j}\ \widehat{a})}(x') = \bigotimes_{2 \le q \le m} \psi_{\alpha_q}^{(q)}(x_q) \otimes \bigotimes_{m < q} \phi^{(q)}(x_q), \quad \alpha \in \widetilde{\Lambda}, \quad x' \in \widetilde{\mathbb{A}},$$

be wavelet functions (5.6) which do not contain factors $\psi_{\alpha_q}^{(q)}(x_q) = \phi^{(q)}(x_q - a_q), a_p \in I_p,$ (see (5.2), (5.3)) for $q \leq m$; $\tilde{\Lambda}$ is the indexes set (4.33). Then $\tilde{\Psi}_{\alpha}$ is an eigenfunction of the operator A_0 if and only if

(9.28)
$$\prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \left(q^{-j_q} (-q^{-1}k_q + \eta_q) \right) = \prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \left(-q^{-j_q-1}k_q \right), \quad j_q \in \mathbb{Z}_+,$$

holds for all $\eta_q \in \mathbb{Z}_q$, $q = 2, 3, \ldots, m$. The corresponding eigenvalue is the following $\lambda = \prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \left(-q^{j_q-1}k_q \right), \ i.e.,$

$$A_0 \Big|_{\Phi^{[m]}(\widetilde{\mathbb{A}})} \widetilde{\Psi}_{\alpha}(x') \stackrel{def}{=} A_{0(m)} \widetilde{\Psi}_{\alpha}(x') \\ = \Big(\prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \big(-q^{j_q-1}k_q \big) \Big) \widetilde{\Psi}_{\alpha}(x'), \quad \alpha \in \widetilde{\Lambda}, \quad x' \in \widetilde{\mathbb{A}}.$$

Proof. In view of (9.23) and (2.10), we have

$$A_{0}|_{\Phi^{[m]}(\widetilde{\mathbb{A}})}\widetilde{\Psi}_{\alpha}(x') \stackrel{def}{=} A_{0(m)}\widetilde{\Psi}_{\alpha}(x') = F^{-1} \Big[\mathcal{A}_{0(m)}(\xi') F[\widetilde{\Psi}_{\alpha}(\xi')] \Big](x')$$

$$= \prod_{2 \le q \le p} F^{-1} \Big[\mathcal{A}_{q} F[\psi_{\alpha_{q}}^{(q)}] \Big](x_{q}) \prod_{p < q \le m} F^{-1} \Big[\mathcal{A}_{q} F[\psi_{\alpha_{q}}^{(q)}] \Big](x_{q}) \prod_{q > m} \phi^{(q)}(x_{q}),$$

$$(\psi_{\alpha_{q}}^{(q)}) = \psi_{\alpha_{q}}^{(q)} (x_{q}) - \psi_{\alpha_{q}}^{(q)} (x_{q}) = (x_{q}) \sum_{p < q \le m} F^{-1} \Big[\mathcal{A}_{q} F[\psi_{\alpha_{q}}^{(q)}] \Big](x_{q}) \prod_{q > m} \phi^{(q)}(x_{q}),$$

where $\psi_{\alpha_q}^{(q)}(x_q) = \psi_{k_q; j_q a_q}^{(q)}(x_q) = q^{j_q/2} \chi_q \left(\frac{k_q}{q} (q^{-j_q} x_q - a_q)\right) \Omega \left(|q^{-j_q} x_q - a_q|_q\right)$ (see (5.2), (5.3)), and $x_q \in \mathbb{Q}_q$, $2 \le q \le p$; $x_q \in \mathbb{Z}_q$, $p < q \le m$. n

If condition
$$(9.28)$$
 holds, then using formulas (8.4) , (8.7) , we obtain

$$\begin{split} A_0 \big|_{\Phi^{[m]}(\widetilde{\mathbb{A}})} \widetilde{\Psi}_{\alpha}(x') &= A_{0(m)} \widetilde{\Psi}_{\alpha}(x') \\ &= \prod_{2 \leq q \leq p} q^{-j_q/2} F^{-1} \Big[\mathcal{A}_q(\xi_q) \chi_q \big(q^{j_q} a_q \xi_q \big) \Omega \Big(\Big| \frac{k_q}{q} + q^{j_q} \xi_q \Big|_q \Big) \Big](x_q) \times \\ &\times \prod_{p < q \leq m} q^{-j_q/2} F^{-1} \Big[\mathcal{A}_q(\xi_q) \Omega(|q^{j_q} a_q|_q) \chi_q \big(q^{j_q} a_q \xi_q \big) \Omega \Big(\Big| \frac{k_q}{q} + q^{j_q} \xi_q \Big|_q \Big) \Big](x_q) \times \\ &\times \prod_{q > m} \phi^{(q)}(x_q). \end{split}$$

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Making the change of variables $\xi_q = q^{-j_q}(\eta_q - q^{-1}k_q)$ and using formulas (2.8), (2.10), we have for $2 \le q \le p$

$$q^{-j_q/2}F^{-1}\Big[\mathcal{A}_q(\xi_q)\chi_q(q^{j_q}a_q\xi_q)\Omega\Big(\Big|\frac{k_q}{q} + q^{j_q}\xi_q\Big|_q\Big)\Big](x_q)$$

= $\mathcal{A}_q\Big(-q^{j_q-1}k_q\Big)q^{j_q/2}\int_{\mathbb{Q}_q}\chi_q\Big((-q^{-1}k_q + \eta_q)(-q^{-j_q}x_q + a_q)\Big)\Omega(|\eta_q|_q)\,d\eta_q$
= $\mathcal{A}_q\Big(-q^{j_q-1}k_q\Big)\psi_{\alpha_q}^{(q)}(x_q), \quad x_q \in \mathbb{Q}_q.$

In the same way, we obtain for $p < q \le m$

$$\begin{aligned} q^{-j_q/2} F^{-1} \Big[\mathcal{A}_q(\xi_q) \Omega(|q^{j_q} a_q|_q) \chi_q(q^{j_q} a_q \xi_q) \Omega\Big(\Big| \frac{k_q}{q} + q^{j_q} \xi_q \Big|_q \Big) \Big](x_q) \\ &= \mathcal{A}_q \Big(-q^{j_q-1} k_q \Big) \Omega(|q^{j_q} a_q|_q) \psi^{(q)}_{\alpha_q}(x_q) \\ &= \mathcal{A}_q \Big(-q^{j_q-1} k_q \Big) \Omega(|q^{j_q} a_q|_q) q^{j_q/2} \chi_q \Big(\frac{k_q}{q} (q^{-j_q} x_q - a_q) \Big) \Omega\Big(|q^{-j_q} x_q - a_q|_q \Big). \end{aligned}$$

The latter expression is nonzero only if $q^{j_q}a_q \in \mathbb{Z}_q$ and $|x_q - q^{j_q}a_q|_q \leq q^{-j_q}$. This implies that $x_q \in \mathbb{Z}_q$. Thus,

$$q^{-j_q/2}F^{-1}\Big[\mathcal{A}_q(\xi_q)\Omega(|q^{j_q}a_q|_q)\chi_q(q^{j_q}a_q\xi_q)\Omega\Big(\Big|\frac{k_q}{q}+q^{j_q}\xi_q\Big|_q\Big)\Big](x_q)$$
$$=\mathcal{A}_q\Big(-q^{j_q-1}k_q\Big)\psi_{\alpha_q}^{(q)}(x_q), \quad x_q\in\mathbb{Z}_q.$$

Consequently, $A_0 \widetilde{\Psi}_{\alpha}(x') = A_{0(m)} \widetilde{\Psi}_{\alpha}(x') = \lambda \widetilde{\Psi}_{\alpha}(x')$, where

$$\lambda = \prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \left(-q^{j_q - 1} k_q \right), \quad x' \in \widetilde{\mathbb{A}}.$$

Conversely, if $A_0|_{\Phi^{[m]}(\widetilde{\mathbb{A}})} \widetilde{\Psi}_{\alpha}(x') = A_{0(m)} \widetilde{\Psi}_{\alpha}(x') = \lambda \widetilde{\Psi}_{\alpha}(x')$, then taking the Fourier transform of both left- and right-hand sides of this identity and using formulas (8.4), (8.7), we obtain

$$\left(\mathcal{A}_0(\xi') - \lambda\right) \prod_{2 \le q \le m} \Omega\left(\left|\frac{k_q}{q} + q^{j_q}\xi_q\right|_q\right) = 0, \quad \xi' \in \widetilde{\mathbb{A}}.$$

If now $\frac{k_q}{q} + q^{j_q}\xi_q = \eta_q$, $\eta_q \in \mathbb{Z}_q$, then $\xi_q = \left(\eta_q - \frac{k_q}{q}\right)q^{-j_q}$ for $2 \le q \le m$. Thus,

$$\lambda = \mathcal{A}_{0(m)} \Big(\big(\eta_2 - 2^{-1} k_2 \big) 2^{-j_2}, \dots, \big(\eta_m - m^{-1} k_m \big) m^{-j_m} \Big), \quad \forall \, \eta_q \in \mathbb{Z}_q, \, 2 \le q \le m.$$

In particular, $\lambda = \mathcal{A}_{0(m)} \left(2^{-j_2 - 1} k_2, \dots, m^{-j_m - 1} k_m \right)$ and consequently (9.28) holds. \Box

According to (9.24), (9.25), the adelic fractional operator $D_{0(m)}^{\widehat{\gamma}}$ has the symbol $\mathcal{A}_{0(m)}(\xi') = |\xi'|_{0(m)}^{\widehat{\gamma}} = \prod_{q \in \{2,3,\dots,m\}} |\xi_q|_p^{\gamma_q} \prod_{q>m} \phi^{(q)}(\xi_q), \quad \xi' = (\xi_2,\dots,\xi_p,\dots) \in \widetilde{\mathbb{A}}.$ It is easy to see that the symbol $\mathcal{A}_{0(m)}(\xi') = |\xi'|_{0(m)}^{\widehat{\gamma}}$ satisfies condition (9.28):

$$\prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \left(q^{-j_q} (-q^{-1}k_q + \eta_q) \right) = \prod_{q \in \{2,3,\dots,m\}} \left| q^{-j_q} \left(-q^{-1}k_q + \eta_q \right) \right|_q^{\gamma_q}$$
$$= \prod_{q \in \{2,3,\dots,m\}} q^{j_q} \left| -q^{-1}k_q + \eta_q \right|_q^{\gamma_q} = \prod_{q \in \{2,3,\dots,m\}} q^{\gamma_q(j_q+1)}$$
$$= \prod_{q \in \{2,3,\dots,m\}} \mathcal{A}_q \left(-q^{-j_q-1}k_q \right), \quad \forall \eta_q \in \mathbb{Z}_q, \quad 2 \le q \le m.$$

Thus according to Theorem 9.1, we have

Corollary 9.1. The wavelet function (5.6)

$$\widetilde{\Psi}_{\alpha}(x') = \widetilde{\Psi}_{(\widehat{k};\widehat{j}\ \widehat{a})}(x') = \bigotimes_{2 \le q \le m} \psi_{\alpha_q}^{(q)}(x_q) \otimes \bigotimes_{m < q} \phi^{(q)}(x_q), \quad x' \in \widetilde{\mathbb{A}}, \quad \alpha \in \widetilde{\Lambda} = \bigcup_{m \in V_{\mathbb{Q}} \setminus \infty} \widetilde{\Lambda}_m.$$

which does not contain the factors $\psi_{\alpha_q}^{(q)}(x_q) = \phi^{(q)}(x_q - a_q), a_p \in I_p \text{ for } q \leq m$, is an eigenfunction of the adelic fractional operator $D_0^{\widehat{\gamma}}$:

$$(9.29) \qquad D_0^{\widehat{\gamma}}\big|_{\Phi^{[m]}(\widetilde{\mathbb{A}})} \widetilde{\Psi}_{\alpha}(x') = D_{0(m)}^{\widehat{\gamma}} \widetilde{\Psi}_{\alpha}(x') = \Big(\prod_{q \in \{2,3,\dots,m\}} q^{\gamma_q(j_q+1)}\Big) \widetilde{\Psi}_{\alpha}(x'), \quad x' \in \widetilde{\mathbb{A}}.$$

Corollary 9.2. The wavelet function (5.6)

$$\widetilde{\Psi}_{\alpha}(x') = \widetilde{\Psi}_{(\widehat{k};\widehat{j}\ \widehat{a})}(x') = \bigotimes_{2 \le q \le m} \psi_{\alpha_q}^{(q)}(x_q) \otimes \bigotimes_{m < q} \phi^{(q)}(x_q), \quad x' \in \widetilde{\mathbb{A}}, \quad \alpha \in \widetilde{\Lambda} = \bigcup_{m \in V_{\mathbb{Q}} \setminus \infty} \widetilde{\Lambda}_m.$$

which does not contain the factors $\psi_{\alpha_q}^{(q)}(x_q) = \phi^{(q)}(x_q - a_q), a_p \in I_p$, for $q \leq m$, is an eigenfunction of the adelic fractional operator D_0^{γ} :

$$(9.30) \qquad D_0^{\gamma}\big|_{\Phi^{[m]}(\widetilde{\mathbb{A}})}\widetilde{\Psi}_{\alpha}(x') = D_{0(m)}^{\gamma}\widetilde{\Psi}_{\alpha}(x') = \Big(\prod_{q \in \{2,3,\dots,m\}} q^{j_q+1}\Big)^{\gamma}\widetilde{\Psi}_{\alpha}(x'), \quad x' \in \widetilde{\mathbb{A}}.$$

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References

- S. Albeverio, S. Evdokimov, M. Skopina, p-Adic multiresolution analysis and wavelet frames, Journal of Fourier Analysis and Applications, 16, no. 5, (2010), 693–714.
- [2] S. Albeverio, S. Evdokimov, M. Skopina, *p-Adic non-orthogonal wavelet bases*, Proc. Steklov Inst. Math., v. 265, Moscow, 2009, 1–12.
- [3] S. Albeverio, A. Yu. Khrennikov, V. M. Shelkovich, *Theory of p-adic distributions: Linear and nonolinear models*, London Mathematical Society Lecture Note Series (No. 370), Cambridge University Press, Cambridge, New York, Melbourne, Madrid, Cape Town, Singapore, Sao Paulo, Delhi, Dubai, Tokyo, 2010.
- [4] S. Albeverio, A. Yu. Khrennikov, V. M. Shelkovich, Harmonic analysis in the p-adic Lizorkin spaces: fractional operators, pseudo-differential equations, p-adic wavelets, Tauberian theorems, Journal of Fourier Analysis and Applications, Vol. 12, Issue 4, (2006), 393–425.
- [5] S. Albeverio, A. Yu. Khrennikov, and V. M. Shelkovich, p-Adic semi-linear evolutionary pseudodifferential equations in the Lizorkin space Dokl. Ross. Akad. Nauk, 415, no. 3, (2007), 295–299. English transl. in Russian Doklady Mathematics, 76, no. 1, (2007), 539–543.
- [6] S. Albeverio, A. Yu. Khrennikov, V. M. Shelkovich, The Cauchy problems for evolutionary pseudodifferential equations over p-adic field and the wavelet theory, Journal of Mathematical Analysis and Applications, 375, (2011), 82–98.
- [7] S. Albeverio, S. V. Kozyrev, Multidimensional ultrametric pseudodifferential equations, Proc. Steklov Inst. Math., v. 265, Moscow, 2009, 13–29.
- [8] V. Anashin and A. Yu. Khrennikov, Applied algebraic dynamics, De Gruyter, Berlin, New York, 2009.
- [9] I. Ya. Aref'eva, B. G. Dragovic, and I. V. Volovich On the adelic string amplitudes, Phys. Lett. B, 209, no. 4, (1998), 445–450.
- [10] J. J. Benedetto, and R. L. Benedetto, A wavelet theory for local fields and related groups, The Journal of Geometric Analysis, 3, (2004), 423–456.
- [11] R. L. Benedetto, *Examples of wavelets for local fields*, Wavelets, Frames, and operator Theory, (College Park, MD, 2003), Am. Math. Soc., Providence, RI, (2004), 27–47.

- [12] Yu.M. Berezanskii, Selfadjoint operators in spaces of functions of infinitely many variables, Translated from the Russian by H. H. McFaden. Translation edited by Ben Silver. Translations of Mathematical Monographs, 63. American Mathematical Society, Providence, RI, 1986.
- [13] A. Borel, Some finiteness properties of adele groups over number fields, Publications mathématiques de l'I.H.É.S., tome 16, (1963), 5–30.
- [14] L. Brekke and P. G. O. Freund, *p*-Adic numbers in physics, Phys. Rep., 233, 1, (1993).
- [15] J. W. S. Cassels and A. Fröhlich, Algebraic Number Theory, Academic Press, 1967.
- [16] C. Chevalley, Généralisation de la théorie du corps de classes pour les extensions infinies, J. Math. Pures Appl., 15, (1936), 359–371.
- [17] K. Conrad, The classical and adelic point of view in number theory, (2010), http://www.mccme.ru/ium/postscript/f10/conrad-lecturenotes.pdf
- [18] I. Daubechies, Ten Lectures on wavelets, CBMS-NSR Series in Appl. Math., SIAM, 1992.
- [19] B. Dragovich, On Generalized Functions in Adelic Quantum Mechanics, Integral Transform. Spec. Funct. 6, (1998), 197–203.
- [20] B. Dragovich, Adelic model of harmonic oscillator, (Russian), Teoret. Mat. Fiz., 101, (1994), 349–359; translation in Theoret. and Math. Phys., 101, no. 3, (1994), 1404–1412.
- [21] B. Dragovich, p-Adic and adelic quantum mechanics, Proc. V.A. Steklov Inst. Math. 245, (2004), 72–85.
- [22] B. Dragovich, Ya. V. Radyno, and A. Khrennikov, *Distributions on adeles*, Journal of Mathematical Sciences, 142, no. 3, (2007), 2105–2112.
- [23] I. M. Gel'fand and G. E. Shilov, *Generalized Functions*, vol. 1, Properties and Operations. New York, Acad. Press, 1964.
- [24] I. M. Gel'fand, M. I. Graev, and I. I. Piatetskii-Shapiro, Generalized functions. vol 6: Representation theory and automorphic functions. Nauka, Moscow, 1966. Translated from the Russian by K.A. Hirsch, Published in 1990, Academic Press (Boston).
- [25] W. Karwowski, M. Vilela, Hierarchical structures and asymmetric stochastic processes on p-adics and adeles, Journal of Mathematical Physics, Volume 35, Issue 9, (1994), 4637–4650.
- [26] A. Khrennikov, p-Adic valued distributions in mathematical physics, Kluwer Academic Publ., Dordrecht, 1994.
- [27] A. Khrennikov, Non-archimedean analysis: quantum paradoxes, dynamical systems and biological models, Kluwer Academic Publ., Dordrecht, 1997.
- [28] A. Yu. Khrennikov, Information dynamics in cognitive, psychological, social and anomalous phenomena, Kluwer Academic Publ., Dordrecht, 2004.
- [29] A. Yu. Khrennikov, and S. V. Kozyrev, Wavelets on ultrametric spaces, Applied and Computational Harmonic Analysis, 19, (2005), 61–76.
- [30] A. Yu. Khrennikov and M. Nilsson, p-Adic Deterministic and Random Dynamics, Kluwer, Dordrecht, 2004.
- [31] A. Khrennikov, Ya. V. Radyno, On adelic analog of Laplacian, Proceedings of Jangjeon Mathematical Society, 6, no. 1, (2003), 1–18.
- [32] A.Yu. Khrennikov, V.M. Shelkovich, p-Adic multidimensional wavelets and their application to p-adic pseudo-differential operators, (2006), Preprint at the url: http://arxiv.org/abs/math-ph/0612049
- [33] A. Yu. Khrennikov, V. M. Shelkovich, Non-Haar p-adic wavelets and their application to pseudodifferential operators and equations, Applied and Computational Harmonic Analysis, 28, (2010), 1–23.
- [34] A.Yu. Khrennikov, V.M. Shelkovich, M. Skopina, p-Adic refinable functions and MRA-based wavelets, Journal of Approximation Theory, 161, (2009), 226–238.
- [35] A.Yu. Khrennikov, V.M. Shelkovich, M. Skopina, *p-Adic orthogonal wavelet bases*, *p-Adic Numbers*, Ultrametric Analysis and Applications, 1, no. 2, (2009), 145–156.
- [36] A. N. Kochubei, Parabolic equations over the field of p-adic numbers, Izv. Akad. Nauk SSSR Ser. Mat., 55, no. 6 (1991), 1312–1330. English transl. in Math. USSR Izv., 39, (1992), 1263–1280
- [37] A. N. Kochubei, Schrödinger-type operator over the p-adic number field, Teoret. Mat. Fiz., 86, (1991), no. 3, 323–333; translation in Theoret. and Math. Phys., 86, (1991), no. 3, 221–228.
- [38] A. N. Kochubei, Pseudo-differential equations and stochastics over non-archimedean fields, Marcel Dekker. Inc. New York, Basel, 2001.

- [39] S. V. Kozyrev, Wavelet analysis as a p-adic spectral analysis, Izvestia Ross. Akademii Nauk, Seria Math., 66 no. 2, 149–158; English transl. in Izvestiya: Mathematics, 66 no. 2, 367–376.
- [40] S. V. Kozyrev, Methods and applications of ultrametric and p-adic analysis: from wavelet theory to biophysics, Sovrem. Probl. Mat. Issue 12. Steklov Inst. Math., Moscow, 2008. http://www.mi.ras.ru/spm/pdf/012.pdf.
- [41] S. V. Kozyrev, p-Adic pseudodifferential operators and p-adic wavelets, Theor. Math. Physics, 138, no. 3, 1–42; English transl. in Theoretical and Mathematical Physics, 138, no. 3, 322–332.
- [42] A.V. Kosyak, Regular, quasiregular and induces representations of infinite-dimensional groups, EMS, 453 p. (in preparation).
- [43] P. I. Lizorkin, Generalized Liouville differentiation and the functional spaces $L_p^r(E_n)$. Imbedding theorems, (Russian) Mat. Sb. (N.S.) **60**(102), (1963), 325–353.
- [44] P. I. Lizorkin, Generalized Liouville differentiation and the multiplier method in the theory of imbeddings of classes of differentiable functions, (Russian) Trudy Mat. Inst. Steklov. Vol. 105, 1969 89–167.
- [45] P. I. Lizorkin, Operators connected with fractional differentiation, and classes of differentiable functions, (Russian) Studies in the theory of differentiable functions of several variables and its applications, IV. Trudy Mat. Inst. Steklov. Vol. 117, (1972), 212–243.
- [46] S. Mallat, Multiresolution representation and wavelets, Ph. D. Thesis, University of Pennsylvania, Philadelphia, PA. 1988.
- [47] S. Mallat, An efficient image representation for multiscale analysis, In: Proc. of Machine Vision Conference, Lake Taho. 1987.
- [48] Ju. I. Manin, Non-Archimedean integration and p-adic Jacquet-Langlands L-functions, (Russian) Uspehi Mat. Nauk, **31**(187), no. 1, (1976), 5–54. Russian Mathematical Surveys, **31**:1, (1976), 5–57
- [49] Yu. I. Manin and A. A. Panchishkin, Introduction to modern number theory fundamental problems, Ideas and theories, 2nd Edition, Spinger-Verlag, Berlin Heidelberg, 2005.
- [50] Y. Meyer, Ondelettes and fonctions splines, Séminaire EDP. Paris. December 1986.
- [51] Y. Meyer, Principe d'incertitude, bases hilbertiennes et algèbres d'opérateurs, Séminaire N. Bourbaki, 1985–1986, exp. no. 662, p.209–223.
- [52] J. von Neumann, On infinite direct products, Compositio mathematica, Vol. 15, no. 1, (1939), 1–77.
- [53] I. Novikov, V. Protasov, and M. Skopina, *Wavelet Theory*, Translations of Mathematical Monographs, vol. 239, AMS, 2011.
- [54] S. G. Samko, Hypersingular integrals and their applications. Taylor & Francis, New York, 2002.
- [55] S. G. Samko, A. A. Kilbas, and O. I. Marichev, Fractional Integrals and Derivatives and Some of Their Applications. Minsk, Nauka i Tekhnika, 1987 (in Russian).
- [56] H. H. Schaefer, *Topological vector spaces*, The Macmillan Co., New York; Collier-Macmillan Ltd., London, 1966.
- [57] V. M. Shelkovich, M. Skopina, p-Adic Haar multiresolution analysis and pseudo-differential operators, Journal of Fourier Analysis and Applications, Vol. 15, Issue 3, (2009), 366–393.
- [58] M. H. Taibleson, Harmonic analysis on n-dimensional vector spaces over local fields. I. Basic results on fractional integration, Math. Annalen, 176, (1968), 191–207.
- [59] M. H. Taibleson, Fourier analysis on local fields, Princeton University Press, Princeton, 1975.
- [60] V. S. Vladimirov, and T. M. Zapuzhak, Adelic formulas for string amplitudes in fields of algebraic numbers, Lett. Math. Phys., 37, (1996), 232–242.
- [61] V. S. Vladimirov, Adelic Formulas for Gamma and Beta Functions of Completions of Algebraic Number Fields and Their Applications to String Amplitudes, Izv. Ross. Akad. Nauk, Ser. Mat., 60, (1996), no. 1, 63–86 (in Russian); English translation in Izvestiya: Mathematics, 60, (1996), no. 1, 67–90.
- [62] V. S. Vladimirov, p-Adic analysis and p-adic quantum mechanics, Ann. of the NY Ac. Sci.: Symposium in Frontiers of Math., 1988.
- [63] V. S. Vladimirov, On the spectrum of some pseudo-differential operators over p-adic number field, Algebra and analysis, 2, no. 6, 1990, 107–124.
- [64] V. S. Vladimirov, I. V. Volovich, and E. I. Zelenov, p-Adic analysis and mathematical physics. World Scientific Publishing, River Edge, NJ, 1994.
- [65] A. Weil, Adeles and algebraic groups, Progress in Mathematics, vol. 23, Birkhäuser Boston, Mass., 1982, With appendices by M. Demazure and Takashi Ono.

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[66] K. Yasuda, Markov Processes on the Adeles and Representations of Euler Products, J. Theoret. Probab. 23, (2010), 748–769.

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