

On Rigorous Derivation of the Enskog Kinetic Equation

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Abstract. We develop a rigorous formalism for the description of the kinetic evolution of infinitely many hard spheres. On the basis of the kinetic cluster expansions of cumulants of groups of operators of finitely many hard spheres the nonlinear kinetic Enskog equation and its generalizations are justified. It is established that for initial states which are specified in terms of one-particle distribution functions the description of the evolution by the Cauchy problem of the BBGKY hierarchy and by the Cauchy problem of the generalized Enskog kinetic equation together with a sequence of explicitly defined functionals of a solution of stated kinetic equation is an equivalent. For the initial-value problem of the generalized Enskog equation the existence theorem is proved in the space of integrable functions.

Key words: BBGKY hierarchy; Enskog equation; kinetic cluster expansion; scattering cumulant; hard spheres.

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1 Introduction

It is well known the considerable advance [1–5] in the rigorous derivation of the Boltzmann kinetic equation in the Boltzmann-Grad scaling limit [6, 7]. The lack of similar progress for the Enskog kinetic equation [8–12] suggested by D. Enskog [13] as a generalization of the Boltzmann equation on dense gases, is conditioned by a priori stated collision integral of this kinetic equation. Thus, the rigorous derivation of the Enskog-type equations from dynamics of hard spheres [1], [14] remains an open problem so far.

In the paper we consider the problem of potentialities inherent in the description of the evolution of states of a hard sphere system in terms of a one-particle distribution function. In view of the generic initial data of hard spheres are determined by initial correlations connected with the forbidden configurations of particles, one of the problems consists in the derivation of

the kinetic equations from underlying dynamics in case of presence of correlations at initial time. We establish that in fact, if initial data is completely defined by a one-particle distribution function, then all possible states of infinitely many hard spheres at arbitrary moment of time can be described within the framework of a one-particle distribution function by the kinetic equation without any approximations.

We briefly outline the main results and structure of the paper. In section 2 we formulate some preliminary facts about dynamics of finitely many hard spheres and state the main result. In sections 3 and 4 the main results related to the origin of the Enskog kinetic evolution are proved. On the basis of the kinetic cluster expansions of cumulants of groups of operators of finitely many hard spheres which give the possibility to construct the kinetic equations in case of presence of correlations at initial time, in section 3 we derive the generalized Enskog equation. Then in section 4 we establish that all possible states are described by a sequence of explicitly defined functionals of a solution of constructed kinetic equation. Moreover, in section 5 we state the Markovian generalized Enskog equation and establish its links with the Enskog-type kinetic equations [8], [15–17]. In section 6 the existence theorem is proved for the abstract initial-value problem of the generalized Enskog equation in the space of integrable functions. Finally in section 7 we conclude with some observations and perspectives for future research.

2 On the origin of the Enskog kinetic evolution

2.1 Preliminaries: dynamics of finitely many hard spheres

We consider a system of identical particles of a unit mass interacting as hard spheres with a diameter $\sigma > 0$. Every particle is characterized by the phase coordinates $(q_i, p_i) \equiv x_i \in \mathbb{R}^3 \times \mathbb{R}^3$, $i \geq 1$. For configurations of such a system the following inequalities must be satisfied: $|q_i - q_j| \geq \sigma$, $i \neq j \geq 1$, i.e. the set $\mathbb{W}_n \equiv \{(q_1, \dots, q_n) \in \mathbb{R}^{3n} \mid |q_i - q_j| < \sigma \text{ for at least one pair } (i, j) : i \neq j \in (1, \dots, n)\}$, is the set of forbidden configurations of n hard spheres.

The evolution of all possible states is described by the sequence of marginal (s -particle) distribution functions $F_s(t, x_1, \dots, x_s)$, $s \geq 1$, that satisfy the initial-value problem of the BBGKY hierarchy [1]

$$\frac{\partial}{\partial t} F_s(t) = \mathcal{L}_s^* F_s(t) + \sum_{i=1}^s \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_{s+1} \mathcal{L}_{\text{int}}^*(i, s+1) F_{s+1}(t), \quad (1)$$

$$F_s(t)|_{t=0} = F_s^0, \quad s \geq 1. \quad (2)$$

If $t > 0$, in hierarchy of evolution equations (1) the operator \mathcal{L}_s^* is defined by the Poisson bracket of noninteracting particles with the corresponding boundary conditions on $\partial\mathbb{W}_s$ [1]

$$\mathcal{L}_s^* F_s(t) \doteq - \sum_{i=1}^s \left\langle p_i, \frac{\partial}{\partial q_i} \right\rangle_{\partial\mathbb{W}_s} F_s(t, x_1, \dots, x_s), \quad (3)$$

where $\langle \eta, (p_i - p_{s+1}) \rangle \doteq \sum_{\alpha=1}^3 \eta^\alpha (p_i^\alpha - p_{s+1}^\alpha)$ is a scalar product. Operator (3) is the generator of the Liouville equation for states and it is an adjoint operator to the generator \mathcal{L}_s of the

Liouville equation for observables [1]. The operator $\mathcal{L}_{\text{int}}^*(i, s+1)$ is defined by the expression

$$\begin{aligned} \sum_{i=1}^s \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_{s+1} \mathcal{L}_{\text{int}}^*(i, s+1) F_{s+1}(t) &= \sigma^2 \sum_{i=1}^s \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_{s+1} d\eta \langle \eta, (p_i - p_{s+1}) \rangle \times \\ &\times (F_{s+1}(t, x_1, \dots, q_i, p_i^*, \dots, x_s, q_i - \sigma\eta, p_{s+1}^*) - F_{s+1}(t, x_1, \dots, x_s, q_i + \sigma\eta, p_{s+1})), \end{aligned} \quad (4)$$

where $\mathbb{S}_+^2 \doteq \{\eta \in \mathbb{R}^3 \mid |\eta| = 1, \langle \eta, (p_i - p_{s+1}) \rangle > 0\}$ and the momenta p_i^*, p_{s+1}^* are defined by the following equalities

$$\begin{aligned} p_i^* &\doteq p_i - \eta \langle \eta, (p_i - p_{s+1}) \rangle, \\ p_{s+1}^* &\doteq p_{s+1} + \eta \langle \eta, (p_i - p_{s+1}) \rangle. \end{aligned} \quad (5)$$

If $t < 0$, the generator of the BBGKY hierarchy is defined by the corresponding expression [1].

We will consider initial data (2) satisfying the chaos condition [1], i.e. statistically independent hard spheres

$$F_s(t)|_{t=0} = \prod_{i=1}^s F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3s} \setminus \mathbb{W}_s}, \quad (6)$$

where $\mathcal{X}_{\mathbb{R}^{3s} \setminus \mathbb{W}_s} \equiv \mathcal{X}_s(q_1, \dots, q_s)$ is a characteristic function of allowed configurations $\mathbb{R}^{3s} \setminus \mathbb{W}_s$ of s hard spheres. Initial data (6) is intrinsic for the kinetic description of many-particle systems because in this case all possible states are characterized by means of a one-particle marginal distribution function.

To construct a solution of initial-value problem (1)-(2) we adduce some preliminaries about hard sphere dynamics. Let $L_n^1 \equiv L^1(\mathbb{R}^{3n} \times (\mathbb{R}^{3n} \setminus \mathbb{W}_n))$ be the space of integrable functions that are symmetric with respect to the permutations of the arguments x_1, \dots, x_n , equal to zero on the set of forbidden configurations \mathbb{W}_n and equipped with the norm: $\|f_n\|_{L^1(\mathbb{R}^{3n} \times \mathbb{R}^{3n})} = \int dx_1 \dots dx_n |f_n(x_1, \dots, x_n)|$. We denote by $L_{n,0}^1 \subset L_n^1$ the subspace of continuously differentiable functions with compact supports.

On the space L_n^1 the following one-parameter mapping: $\mathbb{R} \ni t \mapsto S_n(-t)f_n$ is defined

$$\begin{aligned} (S_n(-t)f_n)(x_1, \dots, x_n) &\equiv S_n(-t, 1, \dots, n)f_n(x_1, \dots, x_n) \doteq \\ &\doteq \begin{cases} f_n(X_1(-t, x_1, \dots, x_n), \dots, X_n(-t, x_1, \dots, x_n)), & \text{if } (x_1, \dots, x_n) \in (\mathbb{R}^{3n} \times (\mathbb{R}^{3n} \setminus \mathbb{W}_n)), \\ 0, & \text{if } (q_1, \dots, q_n) \in \mathbb{W}_n, \end{cases} \end{aligned} \quad (7)$$

where $X_i(-t)$ is a phase trajectory of i th particle constructed in [1]. We note that the phase trajectories of a hard sphere system are determined almost everywhere on the phase space $\mathbb{R}^{3n} \times (\mathbb{R}^{3n} \setminus \mathbb{W}_n)$ beyond the set \mathbb{M}_n^0 of the zero Lebesgue measure [1, 14]. The set \mathbb{M}_n^0 consists of the phase space points which are specified such initial data that during the evolution generate multiple collisions, i.e. collisions of more than two particles, more than one two-particle collision at the same instant and infinite number of collisions within a finite time interval.

On the space L_n^1 mapping (7) defines the isometric strong continuous group of operators [1], and on $L_{n,0}^1 \subset L_n^1$ the infinitesimal generator of group (7) is operator (3) which is the generator of the Liouville equation for hard spheres [1].

If $F_1^0 \in L_n^1$, then a nonperturbative solution of initial-value problem (1), (6) is a sequence of distribution functions $F_s(t, x_1, \dots, x_s)$, $s \geq 1$, represented by the following series [1], [18]

$$\begin{aligned} F_s(t, x_1, \dots, x_s) &= \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(s+n)} \setminus \mathbb{W}_{s+n}}, \end{aligned} \quad (8)$$

where the evolution operator $\mathfrak{A}_{1+n}(-t)$ is the $(n+1)$ th-order cumulant of groups of operators (7) defined by the expansion

$$\mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) = \sum_{\mathbb{P}: (\{Y\}, X \setminus Y) = \cup_i X_i} (-1)^{|\mathbb{P}|-1} (|\mathbb{P}|-1)! \prod_{X_i \subset \mathbb{P}} S_{|X_i|}(-t, X_i), \quad (9)$$

and the following notations are used: $\{Y\}$ is a set consisting of one element $Y \equiv (1, \dots, s)$, i.e. $|\{Y\}| = 1$, $\sum_{\mathbb{P}}$ is a sum over all possible partitions \mathbb{P} of the set $(\{Y\}, X \setminus Y) \equiv (\{Y\}, s+1, \dots, s+n)$ into $|\mathbb{P}|$ nonempty mutually disjoint subsets $X_i \in (\{Y\}, X \setminus Y)$. If $\|F_1^0\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < e^{-1}$, series (8) converges in the norm of the space L_n^1 for arbitrary $t \in \mathbb{R}$.

In case of initial data $F_1^0 \in L_{n,0}^1$ sequence of functions (8) is a strong solution of initial-value problem (1),(6) and for arbitrary initial data $F_1^0 \in L_n^1$ it is a weak solution [18].

In the next section we describe the evolution of states within the framework of the kinetic theory, i.e. in terms of a one-particle marginal distribution function.

2.2 The kinetic evolution of hard spheres: main results

In view of the fact that initial data (6) is completely determined by a one-particle marginal distribution function on allowed configurations, the states given in terms of the sequence $F(t) = (F_1(t, x_1), \dots, F_s(t, x_1, \dots, x_s), \dots)$ of marginal distribution functions (8) can be described within the framework of the sequence $F(t | F_1(t)) = (F_1(t, x_1), F_2(t, x_1, x_2 | F_1(t)), \dots, F_s(t, x_1, \dots, x_s | F_1(t)), \dots)$ of the marginal functionals of the state $F_s(t, x_1, \dots, x_s | F_1(t))$, $s \geq 2$ which are explicitly defined with respect to the solution $F_1(t, x_1)$ of the kinetic equation. We refer to such kinetic equation of a hard sphere system as the generalized Enskog kinetic equation.

Indeed for $t \geq 0$ the one-particle distribution function (8) is a solution of the following initial-value problem

$$\frac{\partial}{\partial t} F_1(t, x_1) = -\langle p_1, \frac{\partial}{\partial q_1} \rangle F_1(t, x_1) + \quad (10)$$

$$+ \sigma^2 \sum_{n=0}^{\infty} \frac{1}{n!} \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \langle \eta, (p_1 - p_2) \rangle \times$$

$$\times (\mathfrak{B}_{1+n}(t, \{1^*, 2_-^*\}, 3, \dots, n+2) F_1(t, q_1, p_1^*) F_1(t, q_1 - \sigma\eta, p_2^*) \prod_{i=3}^{n+2} F_1(t, x_i) -$$

$$- \mathfrak{B}_{1+n}(t, \{1, 2_+\}, 3, \dots, n+2) F_1(t, x_1) F_1(t, q_1 + \sigma\eta, p_2) \prod_{i=3}^{n+2} F_1(t, x_i)),$$

$$F_1(t, x_1)|_{t=0} = F_1^0(x_1). \quad (11)$$

In this evolution equation we use notations from definition (7) adopted to the conventional notation of the Enskog collision integral: indexes $(1^\sharp, 2^\sharp_\pm)$ denote that the evolution operator $\mathfrak{V}_{1+n}(t)$ acts on the corresponding phase points (q_1, p_1^\sharp) and $(q_1 \pm \sigma\eta, p_2^\sharp)$, and in general case the $(n+1)$ th-order evolution operator $\mathfrak{V}_{1+n}(t)$, $n \geq 0$, is defined by the expansion ($|Y| = 2$ in case of equation (10))

$$\begin{aligned} \mathfrak{V}_{1+n}(t, \{Y\}, X \setminus Y) &\doteq n! \sum_{k=0}^n (-1)^k \sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} \frac{1}{(n-m_1-\dots-m_k)!} \times \\ &\times \widehat{\mathfrak{A}}_{1+n-m_1-\dots-m_k}(t, \{Y\}, s+1, \dots, s+n-m_1-\dots-m_k) \times \\ &\times \prod_{j=1}^k \sum_{k_2^j=0}^{m_j} \dots \sum_{k_{n-m_1-\dots-m_j+s}^j}^{k_{n-m_1-\dots-m_j+s-1}^j} \prod_{i_j=1}^{s+n-m_1-\dots-m_j} \frac{1}{(k_{n-m_1-\dots-m_j+s+1-i_j}^j - k_{n-m_1-\dots-m_j+s+2-i_j}^j)!} \times \\ &\times \widehat{\mathfrak{A}}_{1+k_{n-m_1-\dots-m_j+s+1-i_j}^j - k_{n-m_1-\dots-m_j+s+2-i_j}^j}(t, i_j, s+n-m_1-\dots-m_j+1+ \\ &+ k_{s+n-m_1-\dots-m_j+2-i_j}^j, \dots, s+n-m_1-\dots-m_j+k_{s+n-m_1-\dots-m_j+1-i_j}^j), \end{aligned} \quad (12)$$

where we have in view that $k_1^j \equiv m_j$, $k_{n-m_1-\dots-m_j+s+1}^j \equiv 0$, and by the operator $\widehat{\mathfrak{A}}_{1+n}(t)$ we denote the $(1+n)$ th-order scattering cumulant

$$\widehat{\mathfrak{A}}_{1+n}(t, \{Y\}, X \setminus Y) \doteq \mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) \mathfrak{J}_{s+n}(X) \prod_{i=1}^{s+n} \mathfrak{A}_1(t, i). \quad (13)$$

In definition (13) the operator $\mathfrak{A}_{1+n}(-t)$ is the $(1+n)$ th-order cumulant (9) of groups (7) and the operator \mathfrak{J}_{s+n} is defined by the formula

$$\mathfrak{J}_{s+n}(X) f_{s+n} \doteq \mathcal{X}_{\mathbb{R}^{3(s+n)} \setminus \mathbb{W}_{s+n}} f_{s+n}, \quad (14)$$

where $\mathcal{X}_{\mathbb{R}^{3(s+n)} \setminus \mathbb{W}_{s+n}}$ is a characteristic function of allowed configurations of $s+n$ hard spheres. In section 4 we will give an illustration of evolution operators (12).

If $t < 0$, the generalized Enskog equation has the corresponding form

$$\begin{aligned} \frac{\partial}{\partial t} F_1(t, x_1) &= -\langle p_1, \frac{\partial}{\partial q_1} \rangle F_1(t, x_1) + \\ &+ \sigma^2 \sum_{n=0}^{\infty} \frac{1}{n!} \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \langle \eta, (p_1 - p_2) \rangle \times \\ &\times (\mathfrak{V}_{1+n}(t, \{1, 2_-\}, 3, \dots, n+2) F_1(t, q_1, p_1) F_1(t, q_1 - \sigma\eta, p_2) \prod_{i=3}^{n+2} F_1(t, x_i) - \\ &- \mathfrak{V}_{1+n}(t, \{1^*, 2_+^*\}, 3, \dots, n+2) F_1(t, q_1, p_1^*) F_1(t, q_1 + \sigma\eta, p_2^*) \prod_{i=3}^{n+2} F_1(t, x_i)). \end{aligned} \quad (15)$$

The marginal functionals of the state $F_s(t, x_1, \dots, x_s | F_1(t))$, $s \geq 2$, are represented by the

following expansions over products of the one-particle distribution function $F_1(t)$

$$\begin{aligned} F_s(t, x_1, \dots, x_s | F_1(t)) &\doteq \\ &\doteq \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{V}_{1+n}(t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1(t, x_i), \end{aligned} \quad (16)$$

where the $(n+1)$ th-order evolution operator $\mathfrak{V}_{1+n}(t)$, $n \geq 0$, is defined by expansion (12). If $\|F_1(t)\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < e^{-(3s+2)}$, series (16) converges in the norm of the space L_s^1 for arbitrary $t \in \mathbb{R}$.

We remark that in terms of marginal functional of the state (16) the collision integral of the generalized Enskog equation is represented by the expansion

$$\begin{aligned} \mathcal{I}_{GEE} &= \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_i - p_{s+1}) \rangle \times \\ &\times (F_2(t, q_1, p_1^*, q_1 - \sigma\eta, p_2^* | F_1(t)) - F_2(t, x_1, q_1 + \sigma\eta, p_2 | F_1(t))). \end{aligned}$$

In section 5 we establish the relationship of the generalized Enskog equation (10),(15) with the Markovian generalized Enskog equation and the revised Enskog equation [15, 16]. Here we only point out that first term of the collision integral expansion in kinetic equation (10) is exactly the collision integral of the Boltzmann-Enskog equation.

Thus, under the condition: $\|F_1^0\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < e^{-10}(1 + e^{-9})^{-1}$, initial-value problem of the BBGKY hierarchy (1),(6) is equivalent to initial-value problem of the generalized Enskog equation (10)-(11) and the sequence of marginal functionals of the state $F_s(t, x_1, \dots, x_s | F_1(t))$, $s \geq 2$, defined by expansions (16). The proof of this statement is the subject of the next two sections. The existence theorem of the Cauchy problem (10)-(11) in the space of integrable functions is stated in section 6.

Finally it should be noted that the possibility to describe the evolution of all possible states of hard spheres in an equivalent way by the restated Cauchy problem for the generalized Enskog equation and by a sequence of the marginal functionals of the state as compared with the corresponding Cauchy problem of the BBGKY hierarchy is an inherent property of the description of many-particle systems within the framework of the formalism of nonequilibrium grand canonical ensemble which adopted to the description of infinite-particle systems in suitable functional spaces [1].

2.3 A one-dimensional hard sphere system

We consider a one-dimensional system of hard spheres (hard rods of the length $\sigma > 0$). In view of that in this case for $t > 0$ operator (4) is defined by the expression [19]

$$\begin{aligned} &\sum_{i=1}^s \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_{s+1} \mathcal{L}_{\text{int}}^*(i, s+1) F_{s+1}(t) = \\ &\sum_{i=1}^s \int_0^{\infty} dPP (F_{s+1}(t, x_1, \dots, q_i, p_i - P, \dots, x_s, q_i - \sigma, p_i) - F_{s+1}(t, x_1, \dots, x_s, q_i - \sigma, p_i + P) + \\ &+ F_{s+1}(t, x_1, \dots, q_i, p_i + P, \dots, x_s, q_i + \sigma, p_i) - F_{s+1}(t, x_1, \dots, x_s, q_i + \sigma, p_i - P)), \end{aligned}$$

then for $t > 0$ the generalized Enskog kinetic equation for a one-dimensional hard sphere system has the form

$$\begin{aligned} \frac{\partial}{\partial t} F_1(t, x_1) &= -p_1 \frac{\partial}{\partial q_1} F_1(t, x_1) + \\ &+ \sum_{n=0}^{\infty} \frac{1}{n!} \int_0^{\infty} dPP(F_2(t, q_1, p_1 - P, q_1 - \sigma, p_1 | F_1(t)) - F_2(t, q_1, p_1, q_1 - \sigma, p_1 + P | F_1(t)) + \\ &+ F_2(t, q_1, p_1 + P, q_1 + \sigma, p_1 | F_1(t)) - F_2(t, q_1, p_1, q_1 + \sigma, p_1 - P | F_1(t))), \end{aligned}$$

where the marginal functional of the state $F_2(t, x_1, x_2 | F_1(t))$ is defined by formula (16). The series in right-hand side of the equality (17) converges provided that $\|F_1(t)\|_{L^1(\mathbb{R} \times \mathbb{R})} < e^{-8}$ (see section 4). If $t < 0$, the structure of the collision integral of the generalized Enskog equation induces by the structure of the corresponding term of the BBGKY hierarchy for a one-dimensional hard sphere system [19].

3 The generalized Enskog kinetic equation

3.1 Auxiliary facts: kinetic cluster expansions

To reformulate the Cauchy problem (1),(6) as the new Cauchy problem for a one-particle distribution function governed by the kinetic equation and the sequence of explicitly defined marginal functionals of the state which determined by a solution of such Cauchy problem, we introduce the following cluster expansions of cumulants (9) of groups of operators (7)

$$\begin{aligned} \mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) \mathfrak{J}_{s+n}(X) &= \tag{17} \\ &= \sum_{k_1=0}^n \frac{n!}{(n-k_1)!k_1!} \mathfrak{B}_{1+n-k_1}(t, \{Y\}, s+1, \dots, s+n-k_1) \times \\ &\times \sum_{k_2=0}^{k_1} \frac{k_1!}{k_2!(k_1-k_2)!} \cdots \sum_{k_{n-k_1+s}=0}^{k_{n-k_1+s-1}} \frac{k_{n-k_1+s-1}!}{k_{n-k_1+s}!(k_{n-k_1+s-1}-k_{n-k_1+s})!} \times \\ &\times \prod_{i=1}^{s+n-k_1} \mathfrak{A}_{1+k_{n-k_1+s+1-i}-k_{n-k_1+s+2-i}}(-t, i, s+n-k_1+1+k_{s+n-k_1+2-i}, \dots \\ &\dots, s+n-k_1+k_{s+n-k_1+1-i}) \mathfrak{J}_{1+k_{n-k_1+s+1-i}-k_{n-k_1+s+2-i}}(i, s+n-k_1+1+ \\ &+k_{s+n-k_1+2-i}, \dots, s+n-k_1+k_{s+n-k_1+1-i}), \quad n \geq 0, \end{aligned}$$

where the operator \mathfrak{J}_{s+n} is defined by formula (14) and the convention is assumed: $k_{s+1} \equiv 0$. In case of initial states (6) that involve correlations such type (17) cluster expansions (the kinetic cluster expansions) permits to take into consideration initial correlations in the kinetic equation.

We give a few examples of recurrence relations (17) in the terms of scattering cumulants (13). Acting on both sides of equality (17) by the evolution operators $\prod_{i=1}^{s+n} \mathfrak{A}_1(t, i)$, we obtain

$$\widehat{\mathfrak{A}}_1(t, \{Y\}) = \mathfrak{B}_1(t, \{Y\}),$$

$$\begin{aligned}
\widehat{\mathfrak{A}}_2(t, \{Y\}, s+1) &= \mathfrak{V}_2(t, \{Y\}, s+1) + \mathfrak{V}_1(t, \{Y\}) \sum_{i_1=1}^s \widehat{\mathfrak{A}}_2(t, i_1, s+1), \\
\widehat{\mathfrak{A}}_3(t, \{Y\}, s+1, s+2) &= \mathfrak{V}_3(t, \{Y\}, s+1, s+2) + \\
&+ 2! \mathfrak{V}_2(t, \{Y\}, s+1) \sum_{i_1=1}^{s+1} \widehat{\mathfrak{A}}_2(t, i_1, s+2) + \\
&+ \mathfrak{V}_1(t, \{Y\}) \left(\sum_{i_1=1}^s \widehat{\mathfrak{A}}_3(t, i_1, s+1, s+2) + 2! \sum_{1=i_1 < i_2}^s \widehat{\mathfrak{A}}_2(t, i_1, s+1) \widehat{\mathfrak{A}}_2(t, i_2, s+2) \right),
\end{aligned}$$

where $\widehat{\mathfrak{A}}_{1+n}(t)$ is the $(1+n)$ th-order ($n = 0, 1, 2$) scattering cumulant (13). It is easy to prove that solutions of these relations are given by the following expansions

$$\begin{aligned}
\mathfrak{V}_1(t, \{Y\}) &= \widehat{\mathfrak{A}}_1(t, \{Y\}) \doteq S_s(-t, 1, \dots, s) \mathfrak{I}_s(Y) \prod_{i=1}^s S_1(t, i), \quad (18) \\
\mathfrak{V}_2(t, \{Y\}, s+1) &= \widehat{\mathfrak{A}}_2(t, \{Y\}, s+1) - \widehat{\mathfrak{A}}_1(t, \{Y\}) \sum_{i_1=1}^s \widehat{\mathfrak{A}}_2(t, i_1, s+1), \\
\mathfrak{V}_3(t, \{Y\}, s+1, s+2) &= \widehat{\mathfrak{A}}_3(t, \{Y\}, s+1, s+2) - 2! \widehat{\mathfrak{A}}_2(t, \{Y\}, s+1) \times \\
&\times \sum_{i_1=1}^{s+1} \widehat{\mathfrak{A}}_2(t, i_1, s+2) - \widehat{\mathfrak{A}}_1(t, \{Y\}) \left(\sum_{i_1=1}^s \widehat{\mathfrak{A}}_3(t, i_1, s+1, s+2) - \right. \\
&\left. - 2! \sum_{i_1=1}^s \sum_{i_2=1}^{s+1} \widehat{\mathfrak{A}}_2(t, i_1, s+1) \widehat{\mathfrak{A}}_2(t, i_2, s+2) + 2! \sum_{1=i_1 < i_2}^s \widehat{\mathfrak{A}}_2(t, i_1, s+1) \widehat{\mathfrak{A}}_2(t, i_2, s+2) \right).
\end{aligned}$$

In the general case solutions of recurrence relations (17), i.e. the $(1+n)$ th-order ($n \geq 0$) evolution operators $\mathfrak{V}_{1+n}(t)$, are defined by the expansions (12)

$$\begin{aligned}
\mathfrak{V}_{1+n}(t, \{Y\}, X \setminus Y) &= n! \sum_{k=0}^n (-1)^k \sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} \frac{1}{(n-m_1-\dots-m_k)!} \times \quad (19) \\
&\times \widehat{\mathfrak{A}}_{1+n-m_1-\dots-m_k}(t, \{Y\}, s+1, \dots, s+n-m_1-\dots-m_k) \times \\
&\times \prod_{j=1}^k \sum_{k_2^j=0}^{m_j} \dots \sum_{k_{n-m_1-\dots-m_j+s}^j=0}^{k_{n-m_1-\dots-m_j+s-1}^j} \prod_{i_j=1}^{s+n-m_1-\dots-m_j} \frac{1}{(k_{n-m_1-\dots-m_j+s+1-i_j}^j - k_{n-m_1-\dots-m_j+s+2-i_j}^j)!} \times \\
&\times \widehat{\mathfrak{A}}_{1+k_{n-m_1-\dots-m_j+s+1-i_j}^j - k_{n-m_1-\dots-m_j+s+2-i_j}^j}(t, i_j, s+n-m_1-\dots-m_j+1+ \\
&+ k_{s+n-m_1-\dots-m_j+2-i_j}^j, \dots, s+n-m_1-\dots-m_j+k_{s+n-m_1-\dots-m_j+1-i_j}^j), \quad s \geq 1, n \geq 0.
\end{aligned}$$

This fact is verified as a result of the substitution of expressions (19) into recurrence relations (17). It should be emphasized that the correlations generated by the dynamics of hard spheres and initial correlations connected with the forbidden configurations are completely governed by evolution operators (19).

Let us indicate some properties of evolution operators (19). In the sense of the weak convergence of the space $L^1(\mathbb{R}^{3s} \times \mathbb{R}^{3s} \setminus \mathbb{W}_s)$ the infinitesimal generator of first-order evolution operator (18) is defined by operator (4)

$$\lim_{t \rightarrow 0} \frac{1}{t} (\mathfrak{V}_1(t, \{Y\}) - I) f_s = \sum_{i < j=1}^s \mathcal{L}_{\text{int}}^*(i, j) f_s.$$

In the general case, i.e. $n \geq 1$, in the same sense as for evolution operators (18) it holds

$$\lim_{t \rightarrow 0} \frac{1}{t} \mathfrak{V}_{1+n}(t, \{Y\}, Y \setminus X) f_{s+n} = 0.$$

Since in case of a system of non-interacting particles for scattering cumulants (13) the equalities hold: $\widehat{\mathfrak{A}}_{1+n}(t, \{Y\}, X \setminus Y) = \prod_{i=1}^s S_1(-t, i) \mathfrak{J}_s(Y) \prod_{j=1}^s S_1(t, j) \delta_{n,0}$, $n \geq 0$, then we have

$$\mathfrak{V}_{1+n}(t, \{Y\}, X \setminus Y) = \prod_{i=1}^s S_1(-t, i) \mathfrak{J}_s(Y) \prod_{j=1}^s S_1(t, j) \delta_{n,0},$$

where $\delta_{n,0}$ is a Kronecker symbol.

Similarly, at initial time $t = 0$ it is true: $\mathfrak{V}_{1+n}(0, \{Y\}, X \setminus Y) = \mathfrak{J}_s(Y) \delta_{n,0}$.

3.2 The derivation of the generalized Enskog equation

We establish that the one-particle distribution function defined by expansion (8),(9) in case of $s = 1$, i.e.

$$F_1(t, x_1) = \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_2 \dots dx_{n+1} \mathfrak{A}_{1+n}(-t, x_1, \dots, x_{n+1}) \prod_{i=1}^{n+1} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(1+n)} \setminus \mathbb{W}_{1+n}}, \quad (20)$$

is governed by kinetic equation (10) or (15).

In view of the validity of the following equalities for cumulants (9) of groups (7) for $f \in L_0^1$ in the sense of the norm convergence [1], [4]

$$\begin{aligned} \lim_{t \rightarrow 0} \frac{1}{t} \mathfrak{A}_1(-t, 1) f_1(x_1) &= \mathcal{L}_1^*(1) f_1(x_1), \\ \lim_{t \rightarrow 0} \frac{1}{t} \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_2 \mathfrak{A}_2(-t, 1, 2) f_2(x_1, x_2) &= \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_2 \mathcal{L}_{\text{int}}^*(1, 2) f_2(x_1, x_2), \\ \lim_{t \rightarrow 0} \frac{1}{t} \int_{\mathbb{R}^{3n} \times (\mathbb{R}^{3n} \setminus \mathbb{W}_n)} dx_2 \dots dx_{n+1} \mathfrak{A}_{1+n}(-t, 1, \dots, n+1) f_{n+1} &= 0, \quad n \geq 2, \end{aligned}$$

where the operators $\mathcal{L}_1^*(1)$ and $\mathcal{L}_{\text{int}}^*(1, 2)$ are defined by formulas (3) and (4) respectively, as a result of the differentiation over the time variable of expression (20) in the sense of pointwise convergence on the space $L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ we obtain [1], [14]

$$\begin{aligned} \frac{\partial}{\partial t} F_1(t, x_1) &= -\langle p_1, \frac{\partial}{\partial q_1} \rangle F_1(t, x_1) + \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_2 \mathcal{L}_{\text{int}}^*(1, 2) \sum_{n=0}^{\infty} \frac{1}{n!} \times \\ &\times \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \mathfrak{A}_{1+n}(-t, \{1, 2\}, 3, \dots, n+2) \prod_{i=1}^{n+2} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(n+2)} \setminus \mathbb{W}_{n+2}}. \end{aligned} \quad (21)$$

To represent the second term of the right-hand side in this equality in the terms of one-particle distribution function (20) we expand cumulants (9) of groups (7) into kinetic cluster expansions (17) for the case $s = 2$. Then we transform the series over the summation index n and the sum over the index k_1 to the two-fold series. As a result it holds

$$\begin{aligned}
& \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \mathfrak{A}_{1+n}(-t, \{1, 2\}, 3, \dots, n+2) \mathfrak{J}_{n+2}(1, \dots, n+2) \prod_{i=1}^{n+2} F_1^0(x_i) \quad (22) \\
&= \sum_{n=0}^{\infty} \frac{1}{n!} \sum_{k_1=0}^{\infty} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{n+k_1}} dx_3 \dots dx_{n+2+k_1} \mathfrak{Y}_{1+n}(t, \{1, 2\}, 3, \dots, n+2) \sum_{k_2=0}^{k_1} \dots \\
&\dots \sum_{k_{n+2}=0}^{k_{n+1}} \frac{1}{k_{n+2}!(k_{n+1} - k_{n+2})! \dots (k_1 - k_2)!} \prod_{i=1}^{n+2} \mathfrak{A}_{1+k_{n+3-i}-k_{n+4-i}}(-t, i, n+3+k_{n+4-i}, \\
&\dots, n+2+k_{n+3-i}) \mathfrak{X}_{1+k_{n+3-i}-k_{n+4-i}}(q_i, q_{n+3+k_{n+4-i}}, \dots, q_{n+2+k_{n+3-i}}) \prod_{j=1}^{n+2+k_1} F_1^0(x_j).
\end{aligned}$$

The last series can be expressed in terms of one-particle distribution function (20) according to the following product formula for series (20)

$$\begin{aligned}
\prod_{i=1}^{n+2} F_1(t, x_i) &= \sum_{k_1=0}^{\infty} \sum_{k_2=0}^{k_1} \dots \sum_{k_{n+2}=0}^{k_{n+1}} \frac{1}{k_{n+2}!(k_{n+1} - k_{n+2})! \dots (k_1 - k_2)!} \times \quad (23) \\
&\times \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{k_1}} dx_{n+3} \dots dx_{n+2+k_1} \prod_{i=1}^{n+2} \mathfrak{A}_{1+k_{n+3-i}-k_{n+4-i}}(-t, i, n+3+k_{n+4-i}, \\
&\dots, n+2+k_{n+3-i}) \mathfrak{X}_{1+k_{n+3-i}-k_{n+4-i}}(q_i, q_{n+3+k_{n+4-i}}, \dots, q_{n+2+k_{n+3-i}}) \prod_{j=1}^{n+2+k_1} F_1^0(x_j).
\end{aligned}$$

According to equalities (22) and (23), and taking into account definition (4) of the operator $\mathcal{L}_{\text{int}}^*(1, 2)$ for $t > 0$, from equality (21) we finally derive

$$\begin{aligned}
\frac{\partial}{\partial t} F_1(t, x_1) &= -\langle p_1, \frac{\partial}{\partial q_1} \rangle F_1(t, x_1) + \\
&+ \sigma^2 \sum_{n=0}^{\infty} \frac{1}{n!} \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \langle \eta, (p_1 - p_2) \rangle \times \\
&\times (\mathfrak{Y}_{1+n}(t, \{1^*, 2^*\}, 3, \dots, n+2) F_1(t, q_1, p_1^*) F_1(t, q_1 - \sigma\eta, p_2^*) \prod_{i=3}^{n+2} F_1(t, x_i) - \\
&- \mathfrak{Y}_{1+n}(t, \{1, 2_+\}, 3, \dots, n+2) F_1(t, x_1) F_1(t, q_1 + \sigma\eta, p_2) \prod_{i=3}^{n+2} F_1(t, x_i)),
\end{aligned}$$

where we used notations accepted in equation (10). We treat the constructed identity for the one-particle distribution function as the kinetic equation for a hard sphere system. We refer to

this evolution equation as the generalized Enskog kinetic equation. The collision integral series converges (see next section in the general case) under the condition: $\|F_1(t)\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < e^{-8}$.

We emphasize that coefficients of the derived kinetic equation (10) are determined by initial correlations connected with the forbidden configurations of particles.

For quantum many-particle system the generalized kinetic equation has been formulated in [20] (see also a review [21]) and for systems of classical particles interacting via smooth potential it was stated by other method in [1], [22, 23] and for discrete velocity models in [24].

4 Marginal functionals of the state

4.1 A mean value functional within the framework of the kinetic evolution

By means of kinetic cluster expansions (17) in case of $s \geq 2$, solution expansions (8) of the BBGKY hierarchy (1) for hard spheres can be represented in the form of the expansions with respect to one-particle distribution function (20) which is governed by the generalized Enskog equation (10).

Indeed, we expand cumulants (9) of groups (7) from solution expansions (8) into kinetic cluster expansions (17) in case of $s \geq 2$. Then, transforming the series over the summation index n and the sum over the index k_1 to the two-fold series, as a result we obtain

$$\begin{aligned}
F_s(t, Y) &= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) \mathfrak{J}_{s+n}(X) \prod_{i=1}^{s+n} F_1^0(x_i) = \\
&= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{B}_{1+n}(t, \{Y\}, X \setminus Y) \sum_{k_1=0}^{\infty} \sum_{k_2=0}^{k_1} \dots \\
&\dots \sum_{k_{n+s}=0}^{k_{n+s-1}} \frac{1}{k_{n+s}!(k_{n+s-1} - k_{n+s})! \dots (k_1 - k_2)!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{k_1}} dx_{n+s+1} \dots \\
&\dots dx_{n+s+k_1} \prod_{i=1}^{n+s} \mathfrak{A}_{1+k_{n+s+1-i} - k_{n+s+2-i}}(-t, i, n + s + 1 + k_{n+s+2-i}, \dots, n + s + k_{n+s+1-i}) \\
&\times \prod_{j=1}^{n+s+k_1} F_1^0(x_j) \mathfrak{X}_{1+k_{n+s+1-i} - k_{n+s+2-i}}(q_i, q_{n+s+1+k_{n+s+2-i}}, \dots, q_{n+s+k_{n+s+1-i}}),
\end{aligned}$$

where we use notations accepted above. According to the validity of product formula (23) in case of the arbitrary index $s \geq 2$, in obtained expansion the series over the index k_1 can be expressed in terms of one-particle distribution function (20). Hence the following equality holds

$$\begin{aligned}
&\sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) \mathfrak{J}_{s+n}(X) \prod_{i=1}^{s+n} F_1^0(x_i) = \\
&= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{B}_{1+n}(t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1(t, x_i),
\end{aligned}$$

where the $(n + 1)$ th-order evolution operator $\mathfrak{V}_{1+n}(t)$ is defined by formula (12).

Thus, if initial data is completely defined by a one-particle distribution function on allowed configurations, then the evolution of states governed by the BBGKY hierarchy (1) can be completely described by the generalized kinetic equation (10) and the sequence of marginal functionals of the state $F_s(t, x_1, \dots, x_s | F_1(t))$, $s \geq 2$, defined by expansions (16). In other words, for mentioned initial states all possible states of a hard sphere system at arbitrary moment of time can be described within the framework of a one-particle distribution function without any approximations. Therefore the statement of section 2.2 about an equivalence of stated Cauchy problems is valid.

Typical properties for the kinetic description of the evolution of constructed marginal functionals of the state (16) are induced by the properties of evolution operators (12).

We indicate that expansions (16) of marginal functionals of the state are nonequilibrium analog of the Mayer-Ursell expansions over powers of the density of equilibrium distribution functions. In particular, if the one-particle distribution function is the Maxwellian distribution function, then marginal functionals of the state (16) is transformed into the Mayer-Ursell expansions [17], [25, 26].

Owing that a sequence of solutions of the BBGKY hierarchy is represented in terms of a solution of the generalized Enskog equation by the marginal functionals of the state (16), the average value of the s -ary marginal observable $B^{(s)} = (0, \dots, 0, a_s(x_1, \dots, x_s), 0, \dots)$ [21] is defined in terms of the functional $F_s(t, x_1, \dots, x_s | F_1(t))$, by the formula

$$\langle B^{(s)} \rangle(t) = \frac{1}{s!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^s} dx_1 \dots dx_s a_s(x_1, \dots, x_s) F_s(t, x_1, \dots, x_s | F_1(t)),$$

and, in particular the average value of the additive marginal observable $B^{(1)} = (0, a_1(x_1), 0, \dots)$ is determined by

$$\langle B^{(1)} \rangle(t) = \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 a_1(x_1) F_1(t, x_1), \quad (24)$$

where the function $F_1(t, x_1)$ is a solution of the Cauchy problem (10)-(11).

In fact functionals (16) characterize the correlations of states of a hard sphere system. We illustrate this statement by giving an example of the marginal functional from the collision integral of the generalized Enskog equation. In this case the marginal correlation functional $G_2(t, x_1, x_2 | F_1(t))$ is defined by the cluster expansion

$$F_2(t, x_1, x_2 | F_1(t)) = F_1(t, x_1)F_1(t, x_2) + G_2(t, x_1, x_2 | F_1(t)),$$

and it is represented by the series similar to (16)

$$\begin{aligned} G_2(t, x_1, x_2 | F_1(t)) &= (\widehat{\mathfrak{A}}_1(t, \{1, 2\}) - I)F_1(t, x_1)F_1(t, x_2) + \\ &+ \sum_{n=1}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \mathfrak{V}_{1+n}(t, \{1, 2\}, 3, \dots, n+2) \prod_{i=1}^{n+2} F_1(t, x_i). \end{aligned} \quad (25)$$

Basically this cluster expansion gives the classification of all possible currently in use scaling limits [3], [6, 7] since it characterizes the deviation of the state from uncorrelated state, i.e. the

state satisfying a chaos property. In the scaling limits it is assumed that chaos property of initial state preserves in time, and hence the scaling limit means such limit of dimensionless parameters of a system in which the marginal correlation functional $G_2(t, 1, 2 | F_1(t))$ vanishes. According to definition (25), it is possible, if particles of every finite particle cluster move without collisions. In consequence of the definition of a dispersion, for example, of an additive-type observable

$$\begin{aligned} & \langle (B^{(1)} - \langle B^{(1)} \rangle(t))^2 \rangle(t) = \\ & = \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 (a_1^2(x_1) - \langle B^{(1)} \rangle^2(t)) F_1(t, x_1) + \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^2} dx_1 dx_2 a_1(x_1) a_1(x_2) G_2(t, x_1, x_2 | F_1(t)), \end{aligned}$$

where the functional $\langle B^{(1)} \rangle(t)$ is determined by expression (24), the value of the dispersion is minimal for states characterized by marginal correlation functionals (25) equal to zero, i.e. from macroscopic point of view the scaling evolution of hard sphere states is the evolution with minimal dispersion.

Thus, in case of the Boltzmann-Grad scaling limit [1–6] the limit of solution (20) of the generalized Enskog equation is governed by the Boltzmann equation for hard spheres and the limit of marginal functionals (16) are products of the limit one-particle distribution functions that means the propagation of a chaos.

4.2 On the convergence of a series for marginal functionals

We prove the existence of marginal functionals of the state (16) for $F_1(t) \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$. In case of smooth interaction potentials this fact was established in [22, 23] by other method.

Owing to the fact that operator (7) is the isometric operator and the validity for $|X \setminus Y| > 1$ the identity: $\sum_{P: (X \setminus Y) = \cup_i X_i} (-1)^{|P|-1} (|P| - 1)! = 0$, for cumulants (9) of groups (7) the following identity holds

$$\int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \mathfrak{A}_n(t, s+1, \dots, s+n) \prod_{i=1}^{s+n} F_1(t, x_i) = 0, \quad n > 1.$$

As a result of the validity of these identities we can represent marginal functionals of the state (16) by the following renormalized series

$$F_s(t, x_1, \dots, x_s | F_1(t)) = \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_{s+1} \dots dx_{s+n} \tilde{\mathfrak{A}}_{1+n}(t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1(t, x_i), \quad (26)$$

where the $(1+n)th$ -order renormalized evolution operator $\tilde{\mathfrak{A}}_{1+n}(t)$ is defined by the expansion

$$\begin{aligned} \tilde{\mathfrak{A}}_{1+n}(t, \{Y\}, X \setminus Y) &= n! \sum_{k=0}^n (-1)^k \sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} \frac{1}{(n-m_1-\dots-m_k)!} \times \\ &\times \hat{\mathfrak{A}}_{1+n-m_1-\dots-m_k}(t, \{Y\}, s+1, \dots, s+n-m_1-\dots-m_k) \times \end{aligned} \quad (27)$$

$$\begin{aligned} & \times \prod_{j=1}^k \sum_{k_2^j=0}^{m_j} \cdots \sum_{k_s^j=0}^{s-1} \prod_{i_j=1}^s \frac{1}{(k_{s+1-i_j}^j - k_{s+2-i_j}^j)!} \widehat{\mathfrak{A}}_{1+k_{s+1-i_j}^j - k_{s+2-i_j}^j}(t, i_j, \\ & s+n-m_1-\dots-m_j+1+k_{s+2-i_j}^j, \dots, s+n-m_1-\dots-m_j+k_{s+1-i_j}^j). \end{aligned}$$

and as above in (12) the conventions are assumed: $k_1^j \equiv m_j$, $k_{s+1}^j \equiv 0$.

To prove the convergence of the series of marginal functional (16) we use the estimate of the cumulants $\mathfrak{A}_{1+n}(t)$ of groups of operators (7) [18]

$$\int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{s+n}} dx_1 \dots dx_{s+n} |\mathfrak{A}_{1+n}(-t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1(t, x_i)| \leq n! e^{n+2} \|F_1(t)\|_{L_1^1}^{s+n}.$$

In consequence of this estimate, for $(1+n)$ th-order renormalized evolution operator (27) the inequality holds

$$\begin{aligned} & \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{s+n}} dx_1 \dots dx_{s+n} |\widetilde{\mathfrak{A}}_{1+n}(t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1(t, x_i)| \leq n! \|F_1(t)\|_{L_1^1}^{s+n} \times \quad (28) \\ & \times \sum_{k=0}^n \sum_{m_1=1}^n \cdots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} e^{n-m_1-\dots-m_k+2} \prod_{j=1}^k \sum_{k_2^j=0}^{m_j} \cdots \sum_{k_s^j=0}^{k_{s-1}^j} \prod_{i_j=1}^s e^{k_{s+1-i_j}^j - k_{s+2-i_j}^j + 2}. \end{aligned}$$

So far as it is redefined that $k_1^j \equiv m_j$, $k_{s+1}^j \equiv 0$, we have

$$\prod_{i_j=1}^s e^{k_{s+1-i_j}^j - k_{s+2-i_j}^j + 2} = e^{m_j+2s}. \quad (29)$$

Then the following equality is true

$$\sum_{k_2^j=0}^{m_j} \cdots \sum_{k_s^j=0}^{k_{s-1}^j} 1 = \frac{1}{(s-1)!} (m_j+1)(m_j+2)\dots(m_j+s-1). \quad (30)$$

Indeed, using the equality (in case of $m=0$ it is the sum of the arithmetical progression)

$$\sum_{k=1}^n k(k+1)\dots(k+m) = \frac{1}{m+2} n(n+1)\dots(n+m+1), \quad (31)$$

as a result of the sequential computation of sums we obtain

$$\begin{aligned} & \sum_{k_2^j=0}^{m_j} \cdots \sum_{k_s^j=0}^{k_{s-1}^j} 1 = \sum_{k_2^j=0}^{m_j} \cdots \sum_{k_{s-1}^j=0}^{k_{s-2}^j} (k_{s-1}^j + 1) = \sum_{k_2^j=0}^{m_j} \cdots \sum_{k_{s-2}^j=0}^{k_{s-3}^j} \frac{(k_{s-2}^j + 1)(k_{s-2}^j + 2)}{2} = \\ & = \dots = \frac{1}{(s-1)!} (m_j+1)(m_j+2)\dots(m_j+s-1). \end{aligned}$$

According to estimates (29) and (30), we have

$$\begin{aligned} & \prod_{j=1}^k \sum_{k_2^j=0}^{m_j} \dots \sum_{k_s^j=0}^{k_{s-1}^j} \prod_{i_j=1}^s e^{k_{s+1}^j - i_j - k_{s+2}^j + 2} = \prod_{j=1}^k e^{m_j + 2s} \frac{(m_j + 1)(m_j + 2) \dots (m_j + s - 1)}{(s - 1)!} \leq \quad (32) \\ & \leq \prod_{j=1}^k e^{m_j + 2s} \frac{(m_j + s - 1)^{s-1}}{(s - 1)!} \leq \prod_{j=1}^k e^{2m_j + 3s - 1}. \end{aligned}$$

Thus, in consequence of estimate (32) inequality (28) take the form

$$\begin{aligned} & \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{s+n}} dx_1 \dots dx_{s+n} |\tilde{\mathfrak{V}}_{1+n}(t, \{Y\}, X \setminus Y) \prod_{i=1}^{s+n} F_1(t, x_i)| \leq \quad (33) \\ & \leq n! \|F_1(t)\|_{L_1^1}^{s+n} e^{n+2} \sum_{k=0}^n e^{k(3s-1)} \sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} e^{m_1+\dots+m_k}. \end{aligned}$$

Let us estimate the sums $\sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} e^{m_1+\dots+m_k}$. Calculating the sum over index m_k as the sum of a geometric progression, we obtain

$$\sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} e^{m_1+\dots+m_k} \leq e^{n+1} \sum_{m_1=1}^n \dots \sum_{m_{k-1}=1}^{n-m_1-\dots-m_{k-2}} 1.$$

Then, applying equality (31), as a result of the sequential computation of sums the following inequality holds

$$\begin{aligned} & \sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} e^{m_1+\dots+m_k} \leq e^{n+1} \sum_{m_1=1}^n \dots \sum_{m_{k-2}=1}^{n-m_1-\dots-m_{k-3}} (n - m_1 - \dots - m_{k-2}) = \\ & = \dots = e^{n+1} \frac{(n - k + 2) \dots (n - 1)n}{(k - 1)!}. \end{aligned}$$

According the last inequality, for the majorant of estimate (33) we get

$$\begin{aligned} & \sum_{k=0}^n e^{k(3s-1)} \sum_{m_1=1}^n \dots \sum_{m_k=1}^{n-m_1-\dots-m_{k-1}} e^{m_1+\dots+m_k} \leq \quad (34) \\ & \leq 1 + e^{n+3s} + e^{n+1} \sum_{k=2}^n e^{k(3s-1)} \frac{(n - k + 2) \dots (n - 1)n}{(k - 1)!} \leq \\ & \leq 1 + e^{n+3s} + e^{n+1} \sum_{k=2}^n e^{k(3s-1)} \frac{n^{k-1}}{(k - 1)!} \leq 1 + e^{2n+1} \sum_{k=1}^n e^{k(3s-1)}. \end{aligned}$$

Thus, in consequence of inequalities (33) and (34) for marginal functionals of the state (26)

the following estimate holds

$$\begin{aligned} \|F_s(t | F_1(t))\|_{L^1_s} &\leq \sum_{n=0}^{\infty} \|F_1(t)\|_{L^1_1}^{s+n} e^{n+2} + \sum_{n=1}^{\infty} \|F_1(t)\|_{L^1_1}^{s+n} e^{3n+3} \sum_{k=1}^n e^{k(3s-1)} \leq \\ &\leq \|F_1(t)\|_{L^1_1}^s e^2 \sum_{n=0}^{\infty} \|F_1(t)\|_{L^1_1}^n e^n + \|F_1(t)\|_{L^1_1}^s e^{3s+2} \sum_{n=1}^{\infty} \|F_1(t)\|_{L^1_1}^n e^{(3s+2)n}. \end{aligned} \quad (35)$$

Hence the sequence of marginal functionals of the state $F_s(t, x_1, \dots, x_s | F_1(t))$, $s \geq 2$, exists and represents by converged series (16) provided that

$$\|F_1(t)\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < e^{-(3s+2)}. \quad (36)$$

On the other hand, owing the estimate for series (20)

$$\|F_1(t)\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} \leq \|F_1^0\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} e^2 \sum_{n=0}^{\infty} e^n \|F_1^0\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)}^n, \quad (37)$$

and according estimates (36) and (37), we conclude that the Enskog collision integral exists under the condition on initial data

$$\|F_1^0\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < \frac{e^{-10}}{1 + e^{-9}}. \quad (38)$$

5 Some properties of the collision integral

5.1 The Markovian generalized Enskog equation

To state the Markovian generalized Enskog equation we first represent the collision integral \mathcal{I}_{GEE} of the generalized Enskog equation (10) as an expansion with respect to the Boltzman-Enskog collision integral (in case of $t > 0$)

$$\mathcal{I}_{BEE} \equiv \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle (F_1(t, q_1, p_1^*) F_1(t, q_1 - \sigma\eta, p_2^*) - F_1(t, x_1) F_1(q_1 + \sigma\eta, p_2)),$$

where the momenta p_i^*, p_{s+1}^* are defined by equalities (5). We observe that such expansion of the collision integral \mathcal{I}_{GEE} is given in terms of marginal correlation functionals (25). Indeed, in view that it holds: $(\widehat{\mathfrak{A}}_1(t, \{1^\sharp, 2^\sharp_\pm\}) - I) F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) = 0$, we have

$$\begin{aligned} \mathcal{I}_{GEE} &= \mathcal{I}_{BEE} + \sum_{n=1}^{\infty} \mathcal{I}_{GEE}^{(n)} \equiv \\ &\equiv \mathcal{I}_{BEE} + \sigma^2 \sum_{n=1}^{\infty} \frac{1}{n!} \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \langle \eta, (p_1 - p_2) \rangle \times \\ &\times (\mathfrak{Y}_{1+n}(t, \{1^*, 2^*_-\}, 3, \dots, n+2) F_1(t, q_1, p_1^*) F_1(t, q_1 - \sigma\eta, p_2^*) \prod_{i=3}^{n+2} F_1(t, x_i) - \\ &- \mathfrak{Y}_{1+n}(t, \{1, 2_+\}, 3, \dots, n+2) F_1(t, x_1) F_1(t, q_1 + \sigma\eta, p_2) \prod_{i=3}^{n+2} F_1(t, x_i)), \end{aligned} \quad (39)$$

where notations accepted in equation (10) are used.

On the kinetic (macroscopic) scale the typical length for the kinetic phenomena is the mean free pass. Then, observing that in the kinetic scale of the variation of variables [2], [7] the groups of operators (7) of finitely many hard spheres depend on microscopic time variable $\varepsilon^{-1}t$, where $\varepsilon \geq 0$ is a scale parameter, the dimensionless marginal functionals of the state are represented in the form: $F_s(\varepsilon^{-1}t, x_1, \dots, x_s | F_1(t))$, $s \geq 2$. In the formal limit (the Markovian limit) $\varepsilon \rightarrow 0$, the limit marginal functional of the state $F_s(x_1, \dots, x_s | F_1(t))$ is represented by expansion (16) (and also (25)) with the limit generating evolution operators $\lim_{\varepsilon \rightarrow 0} \mathfrak{V}_{1+n}(\varepsilon^{-1}t)$, $n \geq 0$, for example,

$$\lim_{\varepsilon \rightarrow 0} \mathfrak{V}_1(\varepsilon^{-1}t, \{Y\}) = \lim_{\varepsilon \rightarrow 0} \widehat{\mathfrak{A}}_1(\varepsilon^{-1}t, \{Y\}).$$

We note that the limit of the first two terms of dimensionless marginal functional expansions (16) coincide with corresponding terms of the Markovian functionals constructed by the perturbation method with the use of the weakening of correlation condition by N.N. Bogolyubov [1], [17], [25].

The Markovian generalized Enskog equation is derived in the formal Markovian limit $\varepsilon \rightarrow 0$ from the generalized Enskog equation (10) and it has the form

$$\begin{aligned} \frac{\partial}{\partial t} F_1(t, x_1) &= -\langle p_1, \frac{\partial}{\partial q_1} \rangle F_1(t, x_1) + \mathcal{I}_{BEE} + \\ &+ \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \sum_{n=1}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} (\mathfrak{V}_{1+n}(\{1^*, 2^*\}, 3, \dots \\ &\dots, n+2) F_1(t, q_1, p_1^*) F_1(t, q_1 - \sigma\eta, p_2^*) \prod_{i=3}^{n+2} F_1(t, x_i) - \\ &- \mathfrak{V}_{1+n}(\{1, 2_+\}, 3, \dots, n+2) F_1(t, x_1) F_1(t, q_1 + \sigma\eta, p_2) \prod_{i=3}^{n+2} F_1(t, x_i)), \end{aligned} \quad (40)$$

where by $\mathfrak{V}_{1+n}(\{1^\sharp, 2_\pm^\sharp\}, 3, \dots, n+2)$ we denote as above the limit generating evolution operators: $\mathfrak{V}_{1+n} = \lim_{\varepsilon \rightarrow 0} \mathfrak{V}_{1+n}(\varepsilon^{-1}t)$, $n \geq 0$.

Thus, in the kinetic scale of the variation of variables the structure of the collision integral of the generalized Enskog equation (40) describe all possible correlations which are created by hard sphere dynamics and by the propagation of initial correlations stipulated by forbidden configurations.

We remark that for the first time two terms of the Markovian collision integral in case of smooth interaction potentials was established by N.N. Bogolyubov [17] within the framework of the perturbation theory and, using an analogy with the virial equilibrium expansions, these terms of the collision integral in non-perturbative form were formulated by M.S. Green and R.A. Piccirelli [26], [28] and by E.G.D. Cohen [29].

To derive the Bogolyubov collision integral for hard spheres from (40) we apply analogs of the Duhamel equations [27] to generating operators of marginal correlation functionals (25). Indeed, if $f_3 \in L^1(\mathbb{R}^9 \times (\mathbb{R}^9 \setminus \mathbb{W}_3))$, for the scattering cumulant $\widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, 1, 2, 3)$ (13) an analog

of the Duhamel equation holds [4]

$$\begin{aligned} \widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, \{1, 2\}, 3)f_3(x_1, x_2, x_3) &= \int_0^{\varepsilon^{-1}t} d\tau S_2(-\tau, 1, 2)S_1(-\tau, 3) \sum_{i_1=1}^2 \mathcal{L}_{\text{int}}^*(i_1, 3) \times \\ &\times S_3(-\varepsilon^{-1}t + \tau, 1, 2, 3) \mathcal{X}_{\mathbb{R}^9 \setminus \mathbb{W}_3} \prod_{i_2=1}^3 S_1(\varepsilon^{-1}t, i_2) f_3(x_1, x_2, x_3), \end{aligned}$$

and, consequently, for the first term of the collision integral expansion we obtain

$$\mathcal{I}_{GEE}^{(1)} = \mathcal{I}_{GEE,+}^{(1)} - \mathcal{I}_{GEE,-}^{(1)},$$

where we denote by $\mathcal{I}_{GEE,\pm}^{(1)}$ the following parts of the collision integral $\mathcal{I}_{GEE}^{(1)}$

$$\begin{aligned} \mathcal{I}_{GEE,\pm}^{(1)} &= \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_\pm^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_3 \int_0^{\varepsilon^{-1}t} d\tau S_2(-\tau, 1^\sharp, 2^\sharp_\pm) ((\mathcal{L}_{\text{int}}^*(1^\sharp, 3) + \\ &+ \mathcal{L}_{\text{int}}^*(2^\sharp_\pm, 3)) S_3(-\varepsilon^{-1}t + \tau, 1^\sharp, 2^\sharp_\pm, 3) \mathcal{X}_2(q_1, q_3) \mathcal{X}_2(q_1 \pm \sigma\eta, q_3) S_1(\varepsilon^{-1}t - \tau, 1^\sharp) \times \\ &S_1(\varepsilon^{-1}t - \tau, 2^\sharp_\pm) S_1(\varepsilon^{-1}t - \tau, 3) - S_2(-\varepsilon^{-1}t + \tau, 1^\sharp, 2^\sharp_\pm) S_1(\varepsilon^{-1}t - \tau, 1^\sharp) S_1(\varepsilon^{-1}t - \tau, 2^\sharp_\pm) \times \\ &\times (\mathcal{L}_{\text{int}}^*(1^\sharp, 3) S_2(-\varepsilon^{-1}t + \tau, 1^\sharp, 3) + \mathcal{L}_{\text{int}}^*(2^\sharp_\pm, 3) S_2(-\varepsilon^{-1}t + \tau, 2^\sharp_\pm, 3))) \times \\ &\times S_1(\tau, 1^\sharp) S_1(\tau, 2^\sharp_\pm) S_1(\tau, 3) F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) F_1(t, x_3). \end{aligned}$$

In the Markovian limit $\varepsilon \rightarrow 0$ this expression coincide with the corresponding correction to the Boltzmann-Enskog collision integral constructed by the perturbation method with the use of the weakening of correlation condition in [17], [25].

5.2 To the classification of the Enskog-type kinetic equations

We establish links of the Markovian generalized Enskog equation (40) with the revised Enskog equation [8–10]. The classification and links of the Enskog-type kinetic equations derived upon some heuristic arguments is given in [8], [12].

The collision integral of the revised Enskog equation has the form [8–10]

$$\begin{aligned} \mathcal{I}_{REE} &= \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_\pm^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle (g_2(q_1, q_1 - \sigma\eta | n(t)) F_1(t, q_1, p_1^*) F_1(t, q_1 - \sigma\eta, p_2^*) - \\ &g_2(q_1, q_1 + \sigma\eta | n(t)) F_1(t, q_1, p_1) F_1(t, q_1 + \sigma\eta, p_2)), \end{aligned}$$

where $n(t) \doteq \int_{\mathbb{R}^3} dp_1 F_1(t, q_1, p_1)$ and the functional $g_2(q_1, q_2 | n(t))$ is defined in terms of the formal Mayer cluster expansion

$$\begin{aligned} g_2(q_1, q_2 | n(t)) &= \\ &= \mathcal{X}_2(q_1, q_2) \left(1 + \sum_{k=1}^{\infty} \frac{1}{k!} \int_{\mathbb{R}^{3k}} dq_3 \dots dq_{k+2} V(\{q_1, q_2\}, q_3, \dots, q_{k+2}) \prod_{i=2}^{k+2} n(t, q_i) \right), \end{aligned} \tag{41}$$

and $V(\{q_1, q_2\}, q_3, \dots, q_{k+2})$ is the sum of all graphs of k labeled points which are biconnected when the Mayer function $\mathcal{X}_2(q_1, q_2) - 1$ is added, $\mathcal{X}_2(q_1, q_2) \equiv \mathcal{X}_{\mathbb{R}^6 \setminus \mathbb{W}_2}$ is a characteristic function

of allowed configurations of two hard spheres. Taking into account that the Boltzmann-Enskog collision integral is defined by the first term of expansion (41), we denote the collision integral of the revised Enskog equation in the form of the corresponding expansion

$$\mathcal{I}_{REE} \equiv \mathcal{I}_{BEE} + \sum_{k=1}^{\infty} (\mathcal{I}_{REE,+}^{(k)} - \mathcal{I}_{REE,-}^{(k)}), \quad (42)$$

where

$$\begin{aligned} \mathcal{I}_{REE,\pm}^{(k)} &\doteq \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \int_{\mathbb{R}^{3k}} dq_3 \dots dq_{k+2} V(\{q_1, q_1 \pm \sigma\eta\}, q_3, \dots, q_{k+2}) \times \\ &\times F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) \prod_{i=2}^{k+2} n(t, q_i). \end{aligned}$$

We compare term by term collision integral (42) of the revised Enskog equation and the collision integral (39) of the Markovian generalized Enskog equation (40) by the example of the first terms of these expansions. In case of collision integral (42) this term has the form

$$\begin{aligned} \mathcal{I}_{REE,\pm}^{(1)} &= \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_3 (\mathcal{X}_2(q_1, q_3) \mathcal{X}_2(q_1 \pm \sigma\eta, q_3) - \\ &- \mathcal{X}_2(q_1, q_3) - \mathcal{X}_2(q_1 \pm \sigma\eta, q_3) + 1) F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) F_1(t, x_3), \end{aligned} \quad (43)$$

and in case of collision integral (39) the corresponding term has the following form

$$\begin{aligned} \mathcal{I}_{GEE,\pm}^{(1)} &= \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_3 \lim_{\varepsilon \rightarrow 0} (\widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, \{1^\sharp, 2_\pm^\sharp\}, 3) - \\ &- \widehat{\mathfrak{A}}_1(\varepsilon^{-1}t, \{1^\sharp, 2_\pm^\sharp\}) \widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, 1^\sharp, 3) - \widehat{\mathfrak{A}}_1(\varepsilon^{-1}t, \{1^\sharp, 2_\pm^\sharp\}) \widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, 2_\pm^\sharp, 3)) \times \\ &\times F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) F_1(t, x_3), \end{aligned} \quad (44)$$

where we use notations accepted in equation (10) and the evolution operator $\widehat{\mathfrak{A}}_n(\varepsilon^{-1}t)$ is a corresponding order scattering cumulant (13).

We transform expression (44) to the form adopted to the structure of expression (43). With that end in view we introduce the notion of the scattering operator

$$\widehat{S}_n(1, \dots, n) = \lim_{\varepsilon \rightarrow 0} S_n(-\varepsilon^{-1}t, 1, \dots, n) \prod_{i=1}^n S_1(\varepsilon^{-1}t, i), \quad n \geq 1,$$

where the evolution operator $S_n(-\varepsilon^{-1}t)$ is defined by (7). Then, according to definition (13) of scattering cumulants, for the first summand in (44) the following equality holds

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, \{1^\sharp, 2_\pm^\sharp\}, 3) &= \mathfrak{I}_2(1, 3) \mathfrak{I}_2(2_\pm, 3) \widehat{S}_3(1^\sharp, 2_\pm^\sharp, 3) - \\ &- \lim_{\varepsilon \rightarrow 0} S_2(-\varepsilon^{-1}t, 1^\sharp, 2_\pm^\sharp) S_1(-\varepsilon^{-1}t, 3) \mathfrak{I}_2(1, 3) \mathfrak{I}_2(2_\pm, 3) \times \\ &\times S_1(\varepsilon^{-1}t, 1^\sharp) S_1(\varepsilon^{-1}t, 2_\pm^\sharp) S_1(\varepsilon^{-1}t, 3), \end{aligned}$$

and for other summands in (44), respectively

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \widehat{\mathfrak{A}}_1(\varepsilon^{-1}t, \{1^\sharp, 2^\sharp_\pm\}) \widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, 1^\sharp, 3) &= \widehat{S}_2(1^\sharp, 2^\sharp_\pm) \mathfrak{I}_2(1, 3) \widehat{S}_2(1^\sharp, 3) - \\ &- \lim_{\varepsilon \rightarrow 0} S_2(-\varepsilon^{-1}t, 1^\sharp, 2^\sharp_\pm) S_1(-\varepsilon^{-1}t, 3) \mathfrak{I}_2(1, 3) S_1(\varepsilon^{-1}t, 1^\sharp) S_1(\varepsilon^{-1}t, 2^\sharp_\pm) S_1(\varepsilon^{-1}t, 3), \\ \lim_{\varepsilon \rightarrow 0} \widehat{\mathfrak{A}}_1(\varepsilon^{-1}t, \{1^\sharp, 2^\sharp_\pm\}) \widehat{\mathfrak{A}}_2(\varepsilon^{-1}t, 2^\sharp_\pm, 3) &= \widehat{S}_2(1^\sharp, 2^\sharp_\pm) \mathfrak{I}_2(2_\pm, 3) \widehat{S}_2(2^\sharp_\pm, 3) - \\ &- \lim_{\varepsilon \rightarrow 0} S_2(-\varepsilon^{-1}t, 1^\sharp, 2^\sharp_\pm) S_1(-\varepsilon^{-1}t, 3) \mathfrak{I}_2(2_\pm, 3) S_1(\varepsilon^{-1}t, 1^\sharp) S_1(\varepsilon^{-1}t, 2^\sharp_\pm) S_1(\varepsilon^{-1}t, 3), \end{aligned}$$

where the operator $\mathfrak{I}_2(i, j)$ is defined by (14).

As a result of the validity of these equalities, taking into account that $\widehat{S}_2(1^\sharp, 2^\sharp_\pm) = I$ is the identity operator and, owing to the following equality

$$\begin{aligned} \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_3 \lim_{\varepsilon \rightarrow 0} S_2(-\varepsilon^{-1}t, 1^\sharp, 2^\sharp_\pm) S_1(-\varepsilon^{-1}t, 3) (\mathfrak{I}_2(1, 3) - I) (\mathfrak{I}_2(2_\pm, 3) - I) \times \\ \times S_1(\varepsilon^{-1}t, 1^\sharp) S_1(\varepsilon^{-1}t, 2^\sharp_\pm) S_1(\varepsilon^{-1}t, 3) F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) F_1(t, x_3) = 0, \end{aligned}$$

expression (44) takes the form

$$\begin{aligned} \mathcal{I}_{GEE, \pm}^{(1)} &= \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_3 (\mathcal{X}_2(q_1, q_3) \mathcal{X}_2(q_1 \pm \sigma\eta, q_3) \times \\ &\times \widehat{S}_3(1^\sharp, 2^\sharp_\pm, 3) - \mathcal{X}_2(q_1, q_3) \widehat{S}_2(1^\sharp, 3) - \mathcal{X}_2(q_1 \pm \sigma\eta, q_3) \widehat{S}_2(2^\sharp_\pm, 3) + I) \times \\ &\times F_1(t, q_1, p_1^\sharp) F_1(t, q_1 \pm \sigma\eta, p_2^\sharp) F_1(t, x_3). \end{aligned} \quad (45)$$

It is clear that expression (45) coincides with corresponding term (43) of the collision integral expansion of the revised Enskog equation, if the correlations created in the process of evolution by pair interactions of hard sphere can be neglected.

As we note above in section 4.1 expansions of marginal functionals of the state (16) are nonequilibrium analog of the Mayer-Ursell expansions over powers of the density of equilibrium distribution functions [17], [25, 26]. Hence the structures of expansions of collision integral (42) of the revised Enskog equation and collision integral (39) of the Markovian generalized Enskog equation (40) are the same, so observations given for the first terms of these collision integral expansions take place for every term. In fact every term of collision integral (39) give us the classification of correlations generated by pair collisions of hard spheres.

Thus, in case of neglecting of correlations which are generated by hard sphere dynamics in particle clusters of more than two particles, the collision integral of the generalized Enskog kinetic equation (40) transforms to collision integral (42) of the revised Enskog equation. The kinetic evolution of hard spheres described by the revised Enskog equation makes it possible to take into consideration only the initial correlations specified by the forbidden configurations of particles.

6 The Cauchy problem of the generalized Enskog equation

6.1 An existence theorem

We consider the abstract Cauchy problem (10)-(11) ((15),(11)) in the space of integrable functions $L^1(\mathbb{R}^3 \times \mathbb{R}^3)$. The following statement is true.

Theorem 1. *A global in time solution of the Cauchy problem of the generalized Enskog equation (10)-(11) ((15),(11)) is determined by the expansion*

$$F_1(t, x_1) = \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_2 \dots dx_{n+1} \mathfrak{A}_{1+n}(-t, 1, \dots, n+1) \prod_{i=1}^{n+1} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^3(1+n) \setminus \mathbb{W}_{1+n}}, \quad (46)$$

where the cumulants $\mathfrak{A}_{1+n}(-t)$, $n \geq 0$, of groups of operators (7) are defined by formula (9). If $\|F_1^0\|_{L^1(\mathbb{R}^3 \times \mathbb{R}^3)} < e^{-10}(1 + e^{-9})^{-1}$, then for $F_1^0 \in L_0^1(\mathbb{R}^3 \times \mathbb{R}^3)$ it is a strong (classical) solution and for an arbitrary initial data $F_1^0 \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ it is a weak (generalized) solution.

Proof. Let $F_1^0 \in L_0^1(\mathbb{R}^3 \times \mathbb{R}^3)$. Series (46) converges in the norm of the space $L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ provided that (38). We use the result of the differentiation over the time variable of expression (46) in the sense of pointwise convergence on the space $L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ from section 3.3. Then for $F_1^0 \in L_0^1(\mathbb{R}^3 \times \mathbb{R}^3)$ it holds

$$\begin{aligned} & \lim_{\Delta t \rightarrow 0} \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 \left| \frac{1}{\Delta t} (F_1(t + \Delta t, x_1) - F_1(t, x_1)) - \right. \\ & - \left(- \langle p_1, \frac{\partial}{\partial q_1} \rangle F_1(t, x_1) + \sum_{n=0}^{\infty} \frac{1}{n!} \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_2 \mathcal{L}_{\text{int}}^*(1, 2) \times \right. \\ & \left. \left. \times \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \mathfrak{A}_{1+n}(t, \{1, 2\}, 3, \dots, n+2) \prod_{i=1}^{n+2} F_1(t, x_i) \right) \right| = 0, \end{aligned} \quad (47)$$

where the operator $\mathcal{L}_{\text{int}}^*(1, 2)$ is defined by formula (4) and abridged notations are used: the symbols $F_1(t, x_1)$ and $\prod_{i=1}^{n+2} F_1(t, x_i)$ are implied series (46) and for $s = 2$ series (23), respectively.

The validity of equality (47) means that a strong solution of initial-value problem (10)-(11) is given by series (46).

Let us prove that in case of initial data $F_1^0 \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ a weak solution of the initial-value problem of the generalized Enskog equation (10) is represented by series (46). With this purpose we introduce the functional

$$(\varphi_1, F_1(t)) \doteq \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 \varphi_1(x_1) F_1(t, x_1), \quad (48)$$

where $\varphi_1 \in L_0(\mathbb{R}^3 \times \mathbb{R}^3) \subset L(\mathbb{R}^3 \times \mathbb{R}^3)$ is continuously differentiable bounded function with compact support and the function $F_1(t)$ is defined by series (46) which is convergent series in the norm of the space $L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ under condition (38). This functional exists since the function $F_1(t)$ is integrable and φ_1 is a bounded function.

Using expansion (46), we transform functional (48) as follows

$$\begin{aligned}
 (\varphi_1, F_1(t)) &= \\
 &= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{n+1}} dx_1 \dots dx_{n+1} \varphi_1(x_1) \mathfrak{A}_{1+n}(-t, 1, \dots, n+1) \prod_{i=1}^{n+1} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(1+n)} \setminus \mathbb{W}_{1+n}} = \\
 &= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{n+1}} dx_1 \dots dx_{n+1} \mathfrak{A}_{1+n}(t, 1, \dots, n+1) \varphi_1(x_1) \prod_{i=1}^{n+1} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(1+n)} \setminus \mathbb{W}_{1+n}},
 \end{aligned}$$

where the cumulant $\mathfrak{A}_{1+n}(t)$, $n \geq 0$ is adjoint to the cumulant $\mathfrak{A}_{1+n}(-t)$, $n \geq 0$, in the sense of functional (48). For $F_1^0 \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ and $\varphi_1 \in L_0(\mathbb{R}^3 \times \mathbb{R}^3)$ in the sense of the $*$ -weak convergence on the space $L(\mathbb{R}^3 \times \mathbb{R}^3)$ the following equality holds

$$\begin{aligned}
 \lim_{\Delta t \rightarrow 0} \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{n+1}} dx_1 \dots dx_{n+1} \frac{1}{\Delta t} (\mathfrak{A}_{1+n}(-t - \Delta t) - \\
 - \mathfrak{A}_{1+n}(-t)) \varphi_1(x_1) \prod_{i=1}^{n+1} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(1+n)} \setminus \mathbb{W}_{1+n}} = \\
 = (\mathcal{L}_1 \varphi_1, F_1(t)) + \sum_{n=1}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{n+2}} dx_1 \dots dx_{n+2} \mathcal{L}_{\text{int}}(1, 2) \varphi_1(x_1) \times \\
 \times \mathfrak{A}_{1+k}(-t, \{1, 2\}, 3, \dots, n+2) \prod_{i=1}^{n+2} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(n+2)} \setminus \mathbb{W}_{n+2}}.
 \end{aligned} \tag{49}$$

In obtained expression (49) the operator $\mathcal{L}_{\text{int}}(1, 2)$ is adjoint to the operator $\mathcal{L}_{\text{int}}^*(1, 2)$ and it is defined on $L_0(\mathbb{R}^6 \times (\mathbb{R}^6 \setminus \mathbb{W}_3))$ by the formula

$$\begin{aligned}
 \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^2} dx_1 dx_2 \mathcal{L}_{\text{int}}(1, 2) \varphi_2(x_1, x_2) \doteq \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, (p_1 - p_2) \rangle \times \\
 \times (\varphi_2(q_1, p_1^*, q_1 + \sigma\eta, p_2^*) - \varphi_2(q_1, p_1, q_1 + \sigma\eta, p_2)).
 \end{aligned} \tag{50}$$

For $F_1^0 \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ and $\varphi_1 \in L_0(\mathbb{R}^3 \times \mathbb{R}^3)$ the limit functionals (49) exist. Using equality (23), we transform the second functional in (49) to the form

$$\begin{aligned}
 \sum_{n=1}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^{n+2}} dx_1 \dots dx_{n+2} \mathcal{L}_{\text{int}}(1, 2) \varphi_1(x_1) \mathfrak{A}_{1+k}(-t, \{1, 2\}, 3, \dots, n+2) \times \\
 \times \prod_{i=1}^{n+2} F_1^0(x_i) \mathcal{X}_{\mathbb{R}^{3(n+2)} \setminus \mathbb{W}_{n+2}} = \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^2} dx_1 dx_2 \mathcal{L}_{\text{int}}(1, 2) \varphi_1(x_1) \times \\
 \times \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \mathfrak{A}_{1+n}(t, \{1, 2\}, 3, \dots, n+2) \prod_{i=1}^{n+2} F_1(t, x_i).
 \end{aligned} \tag{51}$$

Therefore, in consequence of equalities (49),(51), for functional (48) we obtain the equality

$$\begin{aligned} \frac{d}{dt}(\varphi_1, F_1(t)) &= (\mathcal{L}_1\varphi_1, F_1(t)) + \\ &+ \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^2} dx_1 dx_2 \mathcal{L}_{\text{int}}(1, 2)\varphi_1(x_1) \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots dx_{n+2} \mathfrak{V}_{1+n}(t) \prod_{i=1}^{n+2} F_1(t, i), \end{aligned}$$

or in the explicit form

$$\begin{aligned} \frac{d}{dt}(\varphi_1, F_1(t)) &= \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 \langle p_1, \frac{\partial}{\partial q_1} \rangle \varphi_1(x_1) F_1(t, x_1) + \\ &+ \sigma^2 \int_{\mathbb{R}^3 \times \mathbb{R}^3} dx_1 \int_{\mathbb{R}^3 \times \mathbb{S}_+^2} dp_2 d\eta \langle \eta, p_1 - p_2 \rangle (\varphi_1(q_1, p_1^*) - \varphi_1(q_1, p_1)) \sum_{n=0}^{\infty} \frac{1}{n!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^n} dx_3 \dots \\ &\dots dx_{n+2} \mathfrak{V}_{1+n}(t, \{1, 2_+\}, 3, \dots, n+2) F_1(t, x_1) F_1(t, q_1 + \sigma\eta, p_2) \prod_{i=3}^{n+2} F_1(t, x_i). \end{aligned} \quad (52)$$

This equality (52) means that series (46) for arbitrary initial data $F_1^0 \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ is a weak solution of the Cauchy problem of the generalized Enskog equation (10)-(11). \square

6.2 A weak solution in the extended sense

For the Cauchy problem (10)-(11) we introduced the notion of a weak solution in the extended sense. Let $(\varphi, F(t | F_1(t)))$ be the functional

$$(\varphi, F(t | F_1(t))) \doteq \sum_{s=0}^{\infty} \frac{1}{s!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^s} dx_1 \dots dx_s \varphi_s(x_1, \dots, x_s) F_s(t, x_1, \dots, x_s | F_1(t)), \quad (53)$$

where $\varphi = (\varphi_0, \varphi_1(x_1), \dots, \varphi_s(x_1, \dots, x_s), \dots)$ is a finite sequence of bounded infinitely times differentiable functions with compact supports $\varphi_s \in L_0(\mathbb{R}^{3s} \times (\mathbb{R}^{3s} \setminus \mathbb{W}_s)) \in L(\mathbb{R}^{3s} \times (\mathbb{R}^{3s} \setminus \mathbb{W}_s))$, $s \geq 0$, and elements of the sequence $F(t, | F_1(t)) \doteq (1, F_1(t, x_1), F_2(t, x_1, x_2 | F_1(t)), \dots, F_s(t, x_1, \dots, x_s | F_1(t)), \dots)$ are defined by formulas (46) and (16) for the first and other elements respectively. Owing to estimates (37) and (35) this functional exists provided that (36).

On functions $\varphi_s \in L_0(\mathbb{R}^{3s} \times (\mathbb{R}^{3s} \setminus \mathbb{W}_s))$, $s \geq 1$, it is defined the operator \mathcal{B}^+ which is adjoint to the generator of the BBGKY hierarchy of hard spheres (1) for $t > 0$,

$$(\mathcal{B}^+\varphi)_s(x_1, \dots, x_s) \doteq \mathcal{L}_s(Y)\varphi_s(x_1, \dots, x_s) + \sum_{j_1 \neq j_2=1}^s \mathcal{L}_{\text{int}}(j_1, j_2)\varphi_{s-1}((x_1, \dots, x_s) \setminus (x_{j_1})),$$

where the operator $\mathcal{L}_{\text{int}}(j_1, j_2)$ is defined by (50) and the Liouville operator $\mathcal{L}_s(Y)$ is adjoint to operator (3) in the sense of functional (53) [1].

We are said to be that expansion (46) is a weak solution of the Cauchy problem (10)-(11) in the extended sense, if for functional (53) the following equality is true

$$\frac{d}{dt}(\varphi, F(t | F_1(t))) = (\mathcal{B}^+ \varphi, F(t | F_1(t))). \quad (54)$$

In view of estimates (35) the functional on the right-hand side of this equation exists provided that (36).

To verify this definition we transform functional (53) as follows [21]

$$\begin{aligned} & (\varphi, F(t | F_1(t))) = \\ & = \sum_{s=0}^{\infty} \frac{1}{s!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^s} dx_1 \dots dx_s \sum_{n=0}^s \frac{1}{(s-n)!} \sum_{j_1 \neq \dots \neq j_{s-n}=1}^s \sum_{Z \subset Y \setminus (j_1, \dots, j_{s-n})} (-1)^{|Y \setminus (j_1, \dots, j_{s-n}) \setminus Z|} \times \\ & \times S_{s-n+|Z|}(t, (j_1, \dots, j_{s-n}) \cup Z) \varphi_{s-n}(x_{j_1}, \dots, x_{j_{s-n}}) \prod_{i=1}^s F_1^0(x_i), \end{aligned}$$

where $\sum_{Z \subset Y \setminus (j_1, \dots, j_{s-n})}$ is a sum over all subsets $Z \subset Y \setminus (j_1, \dots, j_{s-n})$ of the set $Y \setminus (j_1, \dots, j_{s-n}) \subset Y \equiv (1, \dots, s)$. For $F_1^0 \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ and $\varphi_s \in L_0(\mathbb{R}^{3s} \times (\mathbb{R}^{3s} \setminus \mathbb{W}_s))$, $s \geq 1$, this functional exists.

Skipping the details, as a result for $\varphi \in L_0(\mathbb{R}^{3s} \times (\mathbb{R}^{3s} \setminus \mathbb{W}_s))$, $s \geq 1$, the derivative of functional (53) over the time variable in the sense of the *-weak convergence on the space of sequences of bounded functions takes the form

$$\begin{aligned} \frac{d}{dt}(\varphi, F(t | F_1(t))) & = \sum_{s=0}^{\infty} \frac{1}{s!} \int_{(\mathbb{R}^3 \times \mathbb{R}^3)^s} dx_1 \dots dx_s (\mathcal{L}_s(Y) \varphi_s(x_1, \dots, x_s) + \\ & + \sum_{j_1 \neq j_2=1}^s \mathcal{L}_{\text{int}}(j_1, j_2) \varphi_{s-1}((x_1, \dots, x_s) \setminus (x_{j_1})) F_s(t, x_1, \dots, x_s | F_1(t)). \end{aligned} \quad (55)$$

According to the notion of a weak solution in the extended sense (54), equality (55) means that for arbitrary initial data $F_1^0 \in L_1^1$ a weak solution of initial-value problem of the generalized Enskog equation (10) is determined by series (46). In fact, we have also proved that marginal functionals of the state (16) are weak solutions of the BBGKY hierarchy of hard spheres.

7 Conclusion

Thus, in case of initial data (6) solution (8) of the Cauchy problem (1)-(2) of the BBGKY hierarchy for hard spheres and a solution of the Cauchy problem of the generalized Enskog equation (10)-(11) together with marginal functionals of the state (16) give two equivalent approaches to the description of the evolution of a hard sphere system. We emphasize that in addition the coefficients of the generalized Enskog equation (10) and marginal functionals of the state (16) are determined by the operators of initial correlations specified by the forbidden configurations of hard spheres.

The structure of the collision integral expansion of the generalized Enskog equation (10) is such that the first term of this expansion is the Boltzman-Enskog collision integral and the next terms describe all possible correlations which are created by hard sphere dynamics and by the propagation of initial correlations connected with the forbidden configurations.

On the basis of the derived generalized Enskog equation (10) we state the Markovian generalized Enskog equation (40). In case of absence of correlations which generated by hard sphere dynamics the collision integral of this kinetic equation transforms to the collision integral of the revised Enskog equation [8]. Therefore the kinetic evolution of hard spheres described by the revised Enskog equation makes it possible to take into consideration only the initial correlations specified by the forbidden configurations of particles.

In the end it should be emphasized that the kinetic evolution is an inherent property of infinite-particle systems. In spite of the fact that in terms of a one-particle marginal distribution function from the space of integrable functions can be described a hard sphere system with the finite average number of particles, the Enskog kinetic equation has been derived on the basis of the formalism of nonequilibrium grand canonical ensemble since its framework is adopted to the description of infinite-particle systems in suitable functional spaces [1] as well.

We remark also that developed approach is also related to the problem of a rigorous derivation of the non-Markovian kinetic-type equations from underlying many-particle dynamics which make possible to describe the memory effects of particle and energy transport, for example, the anomalous transport in turbulent plasma, the Brownian motion of macroparticles in complex fluids.

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