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双模量矩形板的大挠度弯曲计算分析

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摘 要:双模量矩形板在外载荷作用下,会形成各向同性的拉伸区和压缩区,把双模量矩形板看成两种各向同性 材料组成的层合板,采用弹性力学理论建立了双模量矩形板在外载荷作用下的静力平衡方程,利用静力平衡方程 确定了双模量矩形板的中性面位置,推导出了双模量矩形板的大挠度弯曲变形微分方程。用加权残值法求得了双 模量矩形板的大挠度弯曲变形时板中点挠度,把该方法计算结果与有限元计算结果进行了比较,说明了该计算方 法是可靠的,并讨论分析了双模量对矩形板大挠度弯曲变形的影响。

关键词: 双模量; 矩形板; 层合板; 中性面; 大挠度

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LARGE DEFLECTION BENDING CALCULATION AND ANALYSIS OF BIMODULOUS RECTANGULAR PLATE

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Abstract: A bimodulous rectangular plate could form an isotropic compression and a tensile area under external loads. Thusly, a bimodulous rectangular plate was regarded as a laminated plate composited of two kind of isotropic material. The static equilibrium equation of the bimodulous rectangular plate under the condition of external loads was established by using elastic mechanics theory. The location of the neutral plane in the bimodulous rectangular plate was determined by the utilization of static equilibrium equations. The large deflection bending deformation differential equations of the bimodulous rectangular plate was derived, and the middle point deflection of the bimodulous rectangular plate was gained with method of weighted residuals. Then the calculation results were compared with that obtained by finite element method, and it show that the method above is reliable. Meanwhile the effect of the bimodulous on the large deflection bending deformation of a rectangular plate was discussed and analyzed.

Key words: bimodulous; rectangular plate; laminated plate; neutral plane, large deflection

大量研究表明,土木工程中的混凝土、航天航 空及机械工程中的石墨、塑料、合金等许多材料都 具有拉压弹性模量不同的性质。尤其是近几年发展 起来的复合材料更具有明显的拉压弹性模量不同 的特性,如美国研制的具有世界一流水平的"超黑粉"纳米吸波材料中的黑粉,就是把石墨作为吸收剂制成石墨/环氧树脂复合材料,其中石墨材料的拉压弹性模量之比高达4倍。所以,用不同弹性模量

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本构关系对这些材料制成的结构进行计算分析已 备受关注^[1-2]。对于拉压不同弹性模量材料,弹性 系数不仅依赖于结构材料,还根据结构各点位移或 应力状态的不同而不同,亦即与结构材料、形状、 边界条件及外载荷有关。在梁、弹性平面等问题的 结构中,已开始考虑材料的不同弹性模量特性^[3-8]。 但是,对薄板及薄壳弯曲问题的分析与计算,仍延 用传统的同弹性模量弹性理论和默认中性面位置 与中心面位置重合^[9-10],且至今未见到有关双模量 矩形板大挠度弯曲变形计算分析的有关文献。本文 则把双模量矩形板看成两种各向同性材料组成的 层合板,利用静力平衡方程确定了双模量矩形板的 中性面位置,建立了双模量薄板弯曲变形的微分方 程,研究了双模量矩形薄板的大挠度弯曲变形问 题,讨论分析了拉压弹性模量对矩形薄板大挠度弯 曲变形的影响。

1 双模量矩形板中性面的确定

对于双模量矩形板的纯弯曲,其中性面不再位 于板厚的正中央,而是形成了各向同性的拉伸区和 压缩区。假设双模量矩形板变形前垂直于中面的直 线,变形后仍为垂直于中曲面的直线段,且长度不 变,即直法线假定。由弹性力学理论可知双模量矩 形板纯弯曲时的应力表达式为:

$$\begin{cases} \sigma_x = -\frac{E_1 z}{1 - \mu_1^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu_1 \frac{\partial^2 w}{\partial y^2} \right) \\ \sigma_y = -\frac{E_1 z}{1 - \mu_1^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu_1 \frac{\partial^2 w}{\partial x^2} \right), \quad z > 0 \quad (1) \\ \tau_{xy} = -\frac{E_1 z}{1 - \mu_1^2} \frac{\partial^2 w}{\partial x \partial y} \\ \sigma_x = -\frac{E_2 z}{1 - \mu_2^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu_2 \frac{\partial^2 w}{\partial y^2} \right) \\ \sigma_y = -\frac{E_2 z}{1 - \mu_2^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu_2 \frac{\partial^2 w}{\partial x^2} \right), \quad z < 0 \quad (2) \\ \tau_{xy} = -\frac{E_2 z}{1 + \mu_2} \frac{\partial^2 w}{\partial x \partial y} \end{cases}$$

式中: E_1 、 μ_1 分别为拉伸弹性模量及泊松比; E_2 、 μ_2 分别为压缩弹性模量及泊松比。 设*h*为板厚,受拉区高度为*h*₁,双模量矩形板 弯曲时横截面内力应满足以下关系:

$$\begin{cases} \frac{E_1}{1-\mu_1^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu_1 \frac{\partial^2 w}{\partial y^2} \right) \int_0^{h_1} z dz + \\ \frac{E_2}{1-\mu_2^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu_2 \frac{\partial^2 w}{\partial y^2} \right) \int_{h_1-h}^0 z dz = 0 \\ \frac{E_1}{1-\mu_1^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu_1 \frac{\partial^2 w}{\partial x^2} \right) \int_0^{h_1} z dz + \\ \frac{E_2}{1-\mu_2^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu_2 \frac{\partial^2 w}{\partial x^2} \right) \int_{h_1-h}^0 z dz = 0 \end{cases}$$
(3)

把式(3)中两式相加可得:

$$\frac{E_1}{1-\mu_1} \int_0^{h_1} z dz + \frac{E_2}{1-\mu_2} \int_{h_1-h}^0 z dz = 0$$
 (4)

由式(4)可以求得中性面的位置为:

$$h_1 = \frac{h\sqrt{E_2(1-\mu_1)}}{\sqrt{E_1(1-\mu_2)} + \sqrt{E_2(1-\mu_1)}}$$
(5)

2 薄板大挠度弯曲微分方程

由弹性力学得薄板微单元的平衡微分方程为:

$$\begin{cases} \frac{\partial \tau_{xz}}{\partial z} = -\frac{\partial \sigma_x}{\partial x} - \frac{\partial \tau_{xy}}{\partial y} \\ \frac{\partial \tau_{yz}}{\partial z} = -\frac{\partial \sigma_y}{\partial y} - \frac{\partial \tau_{xy}}{\partial x} \end{cases}$$
(6)

对于双模量矩形板可知其上下面边界条件为: $(\tau_{xz})|_{h_1-h}^{h_1}=0, \quad (\tau_{yz})|_{h_1-h}^{h_1}=0$ (7) 把式(6)代入式(7)中可得 τ_{xz} 、 τ_{yz} 的表达式为:

$$\begin{aligned} \tau_{xz} &= \frac{E_1}{2(1-\mu_1^2)} (z^2 - h_1^2) \frac{\partial}{\partial x} \nabla^2 w, \qquad 0 \leqslant z \leqslant h_1 \\ \tau_{xz} &= \frac{E_2}{2(1-\mu_2^2)} [z^2 - (h-h_1)^2] \frac{\partial}{\partial x} \nabla^2 w, \quad h_1 - h \leqslant z \leqslant 0 \end{aligned}$$

$$\tag{8}$$

$$\begin{cases} \tau_{yz} = \frac{E_1}{2(1-\mu_1^2)} (z^2 - h_1^2) \frac{\partial}{\partial y} \nabla^2 w, & 0 \le z \le h_1 \\ \tau_{yz} = \frac{E_2}{2(1-\mu_2^2)} [z^2 - (h-h_1)^2] \frac{\partial}{\partial y} \nabla^2 w, & h_1 - h \le z \le 0 \end{cases}$$
(9)

利用式(1)、式(2)、式(8)、式(9)可以得到下式:

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$$\begin{cases} M_x = \frac{E_1}{1-\mu_1^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu_1 \frac{\partial^2 w}{\partial y^2} \right) \int_{0}^{h} z^2 dz - \frac{E_2}{1-\mu_2^2} \left(\frac{\partial^2 w}{\partial x^2} + \mu_2 \frac{\partial^2 w}{\partial y^2} \right) \int_{h_1-h}^{0} z^2 dz = \left[\frac{E_2(h_1-h)^3}{3(1-\mu_2^2)} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \right] \frac{\partial^2 w}{\partial x^2} + \left[\frac{E_2(h_1-h)^3 \mu_2}{3(1-\mu_2^2) \mu_1} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \right] \mu_1 \frac{\partial^2 w}{\partial y^2} \\ M_y = -\frac{E_1}{1-\mu_1^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu_2 \frac{\partial^2 w}{\partial x^2} \right) \int_{0}^{h} z^2 dz - \frac{E_2}{1-\mu_2^2} \left(\frac{\partial^2 w}{\partial y^2} + \mu_2 \frac{\partial^2 w}{\partial x^2} \right) \int_{h_1-h}^{h} z^2 dz = (10) \\ \left[\frac{E_2(h_1-h)^3}{3(1-\mu_2^2)} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \right] \frac{\partial^2 w}{\partial y^2} + \left[\frac{E_2(h_1-h)^3 \mu_2}{3(1-\mu_2^2)} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \right] \frac{\partial^2 w}{\partial x^2} \\ M_{xy} = -\frac{E_1}{1+\mu_1} \frac{\partial^2 w}{\partial x \partial y} \int_{0}^{h} z^2 dz - \frac{E_2}{1+\mu_2} \frac{\partial^2 w}{\partial x \partial y} \int_{0}^{h} z^2 dz = \frac{E_2(h_1-h)^3}{3(1-\mu_1^2)} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \frac{\partial^2 w}{\partial x \partial y} \\ \left[\frac{Q_x}{2x} = \frac{\partial}{\partial x} \nabla^2 w \int_{0}^{h} \frac{E_1}{2(1-\mu_1^2)} (z^2-h_1^2) dz + \frac{\partial}{\partial x} \nabla^2 w \int_{0}^{h} \frac{E_1}{2(1-\mu_2^2)} (z^2-(h-h_1)^2) dz = \left[\frac{E_2(h_1-h)^3}{3(1-\mu_2^2)} - \frac{E_1h_1^3}{3(1-\mu_2^2)} \right] \frac{\partial}{\partial x} \nabla^2 w \\ \left\{ Q_y = \frac{\partial}{\partial y} \nabla^2 w \int_{0}^{h} \frac{E_1}{2(1-\mu_2^2)} (z^2-h_1^2) dz + \frac{\partial}{\partial y} \nabla^2 w \int_{0}^{h} \frac{E_1}{2(1-\mu_2^2)} (z^2-(h-h_1)^2) dz = \left[\frac{E_2(h_1-h)^3}{3(1-\mu_2^2)} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \right] \frac{\partial}{\partial y} \nabla^2 w \\ \left\{ \frac{\partial}{\partial y} \nabla^2 w \int_{0}^{h} \frac{E_1}{2(1-\mu_2^2)} (z^2-(h-h_1)^2) dz = \left[\frac{E_2(h_1-h)^3}{2(1-\mu_2^2)} - \frac{E_1h_1^3}{3(1-\mu_1^2)} \right] \frac{\partial}{\partial y} \nabla^2 w \right\} \right\} \right\}$$

由弹性力学理论可以知道,当双模量矩形板在 垂直于板面的外分布载荷 q(x, y) 作用下,外分布载 荷与双模量矩形板的内力应满足以下二式:

$$\begin{cases} \frac{\partial M_x}{\partial x} + \frac{\partial M_{xy}}{\partial y} = Q_x \\ \frac{\partial M_{xy}}{\partial x} + \frac{\partial M_y}{\partial y} = Q_y \\ \frac{\partial Q_x}{\partial x} + \frac{\partial Q_y}{\partial y} + N_x \frac{\partial^2 w}{\partial x^2} + 2N_{xy} \frac{\partial^2 w}{\partial x \partial y} + \\ N_y \frac{\partial^2 w}{\partial y^2} + q(x, y) = 0 \end{cases}$$
(13)

式中: $N_x \, \cdot \, N_y \, \cdot \, N_{xy}$ 是由横向分布载荷 q(x, y)引起的中面拉力。

由弹性力学理论可知还可以得到相容方程为:

$$\frac{\partial^2 \varepsilon_x}{\partial y^2} + \frac{\partial^2 \varepsilon_y}{\partial x^2} - \frac{\partial^2 \gamma_{xy}}{\partial x \partial y} = \left(\frac{\partial^2 w}{\partial x \partial y}\right)^2 - \frac{\partial^2 w}{\partial x^2} \frac{\partial^2 w}{\partial y^2} \quad (14)$$

因为 *N_x*、 *N_y*、 *N_{xy}*均为由横向分布载荷 *q*(*x*, *y*)引起的中面拉力,所以应变可表示为:

$$\begin{cases} \varepsilon_x = \frac{1}{E_1 h} (N_x - \mu_1 N_y) \\ \varepsilon_y = \frac{1}{E_1 h} (N_y - \mu_1 N_x) \\ \gamma_{xy} = \frac{2(1 + \mu_1)}{E_1 h} N_{xy} \end{cases}$$
(15)
$$\Leftrightarrow N_x = h \sigma_x = h \frac{\partial^2 \varphi}{\partial^2 y}, \quad N_y = h \sigma_y = h \frac{\partial^2 \varphi}{\partial^2 x},$$

$$N_{xy} = h\tau_{xy} = -h\frac{\partial^2 \varphi}{\partial x \partial y}$$
(16)

把式(10)-式(12)代入式(13)中,把式(15)-式(16)代入式(14)中,即可得到双模量薄板的大挠 度微分方程组为:

$$\begin{cases} D\nabla^4 w = h \left(\frac{\partial^2 \varphi}{\partial x^2} \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 \varphi}{\partial y^2} \frac{\partial^2 w}{\partial x^2} - 2 \frac{\partial^2 \varphi}{\partial x \partial y} \frac{\partial^2 w}{\partial x \partial y} \right) + \\ q(x, y) \\ \nabla^4 \varphi = E_1 \left[\left(\frac{\partial^2 w}{\partial x \partial y} \right)^2 - \frac{\partial^2 w}{\partial x^2} \frac{\partial^2 w}{\partial y^2} \right] \end{cases}$$

(17)

式中:

$$\nabla^4 = \frac{\partial^4}{\partial x^4} + 2\frac{\partial^4}{\partial x^2 \partial y^2} + \frac{\partial^4}{\partial y^4},$$
$$D = \frac{E_1 h_1^3}{3(1-\mu_1^2)} - \frac{E_2 (h_1 - h)^3}{3(1-\mu_2^2)}.$$

3 双模量矩形板大挠度的求解

当横向分布载荷作用在板上时,设挠度函数w 及平面力函数 φ 分别用梁函数表达式为:

 $w(x, y) = F_1 X_1 Y_1$, $\varphi(x, y) = F_2 X_2 Y_2$ (18) 式中: $X_1 \, \cdot \, Y_1 \, \cdot \, X_2 \, \cdot \, Y_2$ 皆为梁函数; $F_1 \, \cdot \, F_2$ 为 伽辽金常数。

把式(18)代入式(17)中利用伽辽金原理可得:

$$\begin{cases} DF_1 A_{11} = hF_2 A_{12} + A_{13} \\ F_2 A_{21} = E_1 F_1^2 A_{22} \end{cases}$$
(19)

式中:

由式(19)即可以得到方程:

$$DA_{11}F_1 = A_{13} + \frac{E_1hA_{22}A_{12}}{A_{21}}F_1^3$$
(20)

当双模量矩形板在集中载荷 *P* 作用时,由于梁 函数为三次多项式,有奇异性不能利用卡门方程, 但是可以利用泛函数:

$$U = 2(1 + \mu_1) \iint \left[\left(\frac{\partial^2 \varphi}{\partial x \partial y} \right)^2 - \frac{\partial^2 \varphi}{\partial x^2} \frac{\partial^2 \varphi}{\partial y^2} \right] dxdy - \frac{1}{2E_1 h} \iint \left(\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right)^2 dxdy + \frac{\partial^2 \varphi}{\partial y^2} = \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial y^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} \right]^2 dxdy + \frac{1}{2E_1 h} \left[\frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} + \frac{\partial^2 \varphi}{\partial x^2} \right]^2 dxdy + \frac{\partial^2 \varphi}{\partial x^2} \right]^2 dxdy + \frac{\partial^2 \varphi}{\partial x^2} +$$

$$\frac{D}{2} \iint \left[\left(\frac{\partial w}{\partial x} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right] dxdy + \\ \left[\frac{E_1 h_1^3}{3(1+\mu_1)} - \frac{E_2 (h_1 - h)^3}{3(1+\mu_2)} \right] \iint \left(\frac{\partial^2 w}{\partial x \partial y} \right)^2 dxdy + \\ 2\mu_1 \left[\frac{E_1 h_1^3}{3(1-\mu_1^2)} - \frac{E_2 (h_1 - h)^3 \mu_2}{3(1-\mu_1^2) \mu_1} \right] \iint \frac{\partial^2 w}{\partial x^2} \frac{\partial^2 w}{\partial y^2} dxdy + \\ \frac{1}{2} \iint \left[\frac{\partial^2 \varphi}{\partial x^2} \left(\frac{\partial w}{\partial y} \right)^2 + \frac{\partial^2 \varphi}{\partial y^2} \left(\frac{\partial w}{\partial x} \right)^2 - 2 \frac{\partial^2 \varphi}{\partial x \partial y} \frac{\partial w}{\partial x} \frac{\partial w}{\partial y} \right] dxdy - \\ P \iint w \delta(x - \xi) \delta(y - \eta) dxdy$$
(21)

式中*ξ*、η分别为集中载荷的横坐标和纵坐标。 把式(18)代入式(21)中,利用最小势能原理

$$\begin{split} \frac{\partial U}{\partial F_{1}} &= 0 \ , \ \ \frac{\partial U}{\partial F_{2}} = 0 \ \overline{P} / \frac{\partial \overline{P}}{\partial \overline{P}_{2}} = B_{13} \\ B_{11}F_{1} + B_{12}F_{1}F_{2} = B_{13} \\ B_{21}F_{2} = B_{22}F_{1}^{2} \end{split} \tag{22} \\ B_{11} &= D \int_{0}^{a} \int_{0}^{b} \left[\left(\frac{dX_{1}}{dx}Y_{1} \right)^{2} + \left(X_{1} \frac{dY_{1}}{dy} \right)^{2} \right] dxdy + \\ & 2 \left[\frac{E_{1}h_{1}^{3}}{3(1+\mu_{1})} - \frac{E_{2}(h_{1}-h)^{3}}{3(1+\mu_{2})} \right] \int_{0}^{a} \int_{0}^{b} \left(\frac{dX_{1}}{dx} \frac{dY_{1}}{dy} \right)^{2} dxdy + \\ & 4 \mu_{1} \left[\frac{E_{1}h_{1}^{3}}{3(1-\mu_{1}^{2})} - \frac{E_{2}(h_{1}-h)^{3}\mu_{2}}{3(1-\mu_{2}^{2})\mu_{1}} \right] \cdot \\ & \int_{0}^{a} \int_{0}^{b} \frac{d^{2}X_{1}}{dx^{2}} X_{1}Y_{1} \frac{d^{2}Y_{1}}{dy^{2}} dxdy, \\ B_{12} &= \int_{0}^{a} \int_{0}^{b} \left[\frac{d^{2}X_{2}}{dx^{2}} Y_{2} \left(X_{1} \frac{dY_{1}}{dy} \right)^{2} + X_{2} \frac{d^{2}Y_{2}}{dy^{2}} \left(\frac{dX_{1}}{dx} Y_{1} \right)^{2} - \\ & 2 \frac{dX_{2}}{dx} \frac{dY_{2}}{dy} \frac{dX_{1}}{dx} X_{1}Y_{1} \frac{dY_{1}}{dy} \right] dxdy, \\ B_{13} &= P \int_{0}^{a} \int_{0}^{b} \left[\frac{d^{2}X_{2}}{dx^{2}} Y_{2} + X_{2} \frac{d^{2}Y_{2}}{dy^{2}} \right] dxdy - \\ & 4(1+\mu_{1}) \int_{0}^{a} \int_{0}^{b} \left[\left(\frac{dX_{2}}{dx} \frac{dY_{2}}{dy} \right)^{2} - \frac{d^{2}X_{2}}{dx^{2}} X_{2}Y_{2} \frac{d^{2}Y_{2}}{dy^{2}} \right] dxdy, \\ B_{22} &= \frac{1}{2} \int_{0}^{a} \int_{0}^{b} \left[\frac{d^{2}X_{2}}{dx^{2}} Y_{2} \left(X_{1} \frac{dY_{1}}{dy} \right)^{2} + \\ & X_{2} \frac{d^{2}Y_{2}}{dy^{2}} \left(\frac{dX_{1}}{dx} Y_{1} \right)^{2} - 2 \frac{dX_{2}}{dx} \frac{dY_{2}}{dy} \frac{dX_{1}}{dx} X_{1}Y_{1} \frac{dY_{1}}{dy} \right] dxdy \end{split}$$

由式(22)可以得到方程:

$$B_{11}F_1 + \frac{B_{12}B_{22}}{B_{21}}F_1^3 = B_{13}$$
(23)

求解式(20)、式(23)即可得到双模量矩形板在外 载荷作用下大挠度变形时板的中点挠度。

对于在均布载荷q作用下的四边简支方薄板 可设w及φ的表达式为:

$$\begin{cases} w = F_1 \left(\frac{x^4}{a^4} - 2\frac{x^3}{a^3} + \frac{x}{a} \right) \left(\frac{y^4}{a^4} - 2\frac{y^3}{a^3} + \frac{y}{a} \right) \\ \varphi = F_2 \left(\frac{x^4}{a^4} - 2\frac{x^3}{a^3} + \frac{x^2}{a} \right) \left(\frac{y^4}{a^4} - 2\frac{y^3}{a^3} + \frac{y^2}{a} \right) \end{cases}$$
(24)

把式(24)代入式(20)中且设方薄板中点挠度为 w₀可以得到:

$$qa^4 = 241.7Dw_0 + 3.6502E_1hw_0^3 \tag{25}$$

当 $E_1 = E_2 = E$, $\mu_1 = \mu_2 = 0.316$ 时可把式(25) 化为:

$$\frac{qa^4}{Eh^4} = 22.3759\frac{w_0}{h} + \frac{3.6502w_0^3}{h^3}$$
(26)

对于在均布载荷 q 作用下的薄方板,当 $E_1 = E_2 = E$, $\mu_1 = \mu_2 = 0.316$ 时,贾春元则给出了 载荷与板中点挠度的关系式为^[11]:

$$\frac{qa^4}{Eh^4} = 22.25\frac{w_0}{h} + \frac{3.9008w_0^3}{h^3}$$
(27)

 $\frac{w_0}{h}$ 分别取 1、1.2、1.4、1.6,式(26)与式(27)

的计算结果误差分别为 0.4768%、 0.8432%、 1.2218%、 1.5996%, 说明本文方法是可靠的。

4 算例分析

为了具体分析拉压弹性模量对双模量矩形板 大挠度弯曲变形的影响,以表1所示材料四边简支 方板作用均布载荷为例进行计算。其中材料1为玻 璃丝织物,材料2为硅铝合金^[10]。材料3-材料6 为假想的分别与材料1和材料2拉区材料参数(或压 区材料参数)相同的单模量材料。

	表1 材料特性参数
Table1	Material characteristic parameters

			-	
材料	E_1 /GPa	E_2/GPa	μ_1	μ_2
材料1	30.38	16.17	0.35	0.19
材料 2	66.93	73.42	0.34	0.39
材料 3	30.38	30.38	0.35	0.35
材料 4	16.17	16.17	0.19	0.19
材料 5	66.93	66.93	0.34	0.34
材料 6	73.42	73.42	0.39	0.39
111110	13.72	13.72	0.57	0.57

为了验证本文计算方法正确性,分别用 ANSYS7.0 和本文方法(即式(25))计算了由材料1或 材料3组成的均布载荷作用下四边简支方板中点挠 度(w₀)。板边长 a=2000mm, h=100mm。由材料 1 组成的板有限元计算时采用 8 节点 SOLID185 单 元,该单元具有大变形,大应变能力。由式(5)计算 出受拉区高度 h_1 =60mm,上层 60mm 材料参数为 E=16.17GPa, µ=0.19, 下层 40mm 材料参数为 *E*=30.38, μ=0.35, 厚度方向单元尺寸为 20mm, 边 长方向单元尺寸为 40mm, 采用 Large Displacement static analysis 求解。由材料3组成的板有限元计算 时采用采用 8 节点 SOLID185 单元,该单元具有大 变形,大应变能力。材料参数为 E=30.38, μ=0.35, 厚度方向单元尺寸为 20mm, 边长方向单元尺寸为 40mm。采用 Large Displacement static analysis 求解。 $qa^4/h^4 = 100$ GPa 时材料 1 和材料 3 组成四边简支 方板时节点平面外位移如图1和图2所示。计算结 果及结果比较如表2所示。结果比较数值为本文结 果与ANSYS7.0计算结果差值除以ANSYS7.0计算 结果。



图 1 节点平面外位移(材料 1)





图 2 节点平面外位移(材料 3)

Fig.2 The nodes displacement out of plane (material 3)

表 2 本文结果与 ANSYS 7.0 结果比较

Table 2The results comparison between this paper and
ANSYS 7.0

材料	$qa^4/h^4/(\times 10^{10})$		10	20	30	40	50	60	70
材料1	1	本文方法	21.0	40.8	59.8	75.0	89.3	102.2	113.8
	w_0/mm	ANSYS 7.0	20.4	40.2	58.7	75.3	90.6	104.5	116.4
	误差/(%)		2.94	1.49	1.87	-0.40	-1.43	-2.20	-2.23
材料3	w ₀ /mm	本文方法	14.3	28.3	41.9	54.8	66.9	78.4	89.1
		ANSYS 7.0	14.1	27.6	40.9	54.3	65.7	78.1	88.5
	误差/(%)		1.42	2.54	2.44	0.92	1.83	0.38	0.68

为了分析均布载荷作用下四边简支方板大挠 度弯曲时载荷与挠度的非线性关系,采用式(25)计 算了均布载荷作用下四边简支方板 ^w₀ 分别取 0.2,

0.4,	,	1.6 时 $\frac{qa^{+}}{h^4}$ 的计算结果,	如表3所示。
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表 3 载荷与挠度的非线性关系 /(×10¹⁰) Table 3 The nonlinear relation between load and n deflection

W ₍	h/h	0.2	0.4	0.6	0.8	1.0	1.2	1.4	1.6
	材料1	9.51	19.56	30.66	43.37	58.21	75.70	96.39	120.81
aa^4	材料 2	32.78	66.73	103.02	142.83	187.34	237.70	295.10	360.72
$\frac{qu}{4}$	材料3	14.04	28.60	44.23	61.46	80.82	102.84	128.05	156.99
h^{-}	材料4	6.80	13.89	21.55	30.05	39.69	50.75	63.50	78.24
(N/m^2)	材料5	30.68	62.54	96.73	134.45	176.86	225.13	280.44	343.95
	材料6	35.10	71.48	110.43	153.25	201.21	255.60	317.71	388.82

为了分析双模量对均布载荷作用下四边简支 方板中点挠度影响,比较了双模量四边简支方板和 分别与该双模量材料受拉区材料或受压区材料相 同的单模量四边简支方板的中点挠度相差大小,比 较时将单模量四边简支方板中点挠度相差大小,比 较时将单模量四边简支方板中点挠度与双模量四 边简支方板中点挠度除以双模量四边简支方板中 点挠度,计算结果如表4所示。

表 4 不同材料 qa^4/h^4 计算结果比较 /(%) Table 4 The results comparison of qa^4/h^4 between different materials

w_0/h	0.2	0.4	0.6	0.8	1.0	1.2	1.4	1.6
材料1和材料3	3 47.63	46.26	44.25	41.72	38.86	35.85	32.85	29.95
材料1和材料4	4-28.50	-28.96	-29.73	-30.71	-31.81	-32.97	-34.12	-35.24
材料2和材料5	5 -6.41	-6.28	-6.10	-5.87	-5.59	-5.29	-4.97	-4.65
材料2和材料6	5 7.08	7.12	7.20	7.29	7.40	7.53	7.66	7.79

由表 2 可以看出,采用 ANSYS7.0 和本文方法 计算的结果比较相近,二者相差都小于工程上所充 许的误差(5%),这验证了本文方法的可靠性。

对表 3、表 4 的计算结果进行分析可知,对于 拉压弹性模量不同材料的薄板大挠度弯曲变形,不 考虑拉压弹性模量相异时其计算结果与实际情况 相差较大,超过了工程上所允许的计算误差。所以, 对于拉压弹性模量相差较大的薄板大挠度弯曲变 形,其中点挠度的计算不宜采用相同弹性模量经典 薄板理论,而应该采用双模量薄板大挠度理论。

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