

# Inverse Filtering for Compressible Flows

ACL Report 2013-1

Antony Jameson

September, 2013

The instantaneous equations are :

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j}(\rho u_j) = 0 \quad (1)$$

$$\frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_j}(\rho u_i u_j) + \frac{\partial p}{\partial x_i} = \frac{\partial}{\partial x_j}(\sigma_{ij}) \quad (2)$$

$$\frac{\partial}{\partial t}(\rho E) + \frac{\partial}{\partial x_j}(\rho u_j H) = \frac{\partial}{\partial x_j}(u_i \sigma_{ij} + k \frac{\partial T}{\partial x_j}) \quad (3)$$

Let P be an invertible filter with inverse Q such that,

Commute with derivatives,

$$P \frac{\partial}{\partial x_i} = \frac{\partial P}{\partial x_i}, \quad Q \frac{\partial}{\partial x_i} = \frac{\partial Q}{\partial x_i} \quad (4)$$

Note that P is not a projector such that,

$$P^2 = P$$

If this were so, then

$$P^2 Q = PQ = I$$

and

$$P^2 Q = P(PQ) = P$$

so,

$$P = I$$

Define filtered quantities,

$$\bar{\rho} = P\rho, \rho = Q\bar{\rho}, \bar{p} = Pp, p = Q\bar{p} \quad (5)$$

and mass weighted filtered velocity

$$\tilde{u}_i = \frac{P(\rho u_i)}{P\rho} = \frac{P(\rho u_i)}{\bar{\rho}} \quad (6)$$

$$\bar{\rho}\tilde{u}_i = P(\rho u_i), \rho u_i = Q(\bar{\rho}\tilde{u}_i) \quad (7)$$

Then substituting for  $\rho$  and  $\rho u_i$ , the mass conservation equation becomes,

$$\frac{\partial}{\partial t}Q\bar{\rho} + \frac{\partial}{\partial x_j}Q(\bar{\rho}\tilde{u}_i) = 0 \quad (8)$$

and multiplying by  $P = Q^{-1}$ ,

$$\frac{\partial\bar{\rho}}{\partial t} + \frac{\partial}{\partial x_j}(\bar{\rho}\tilde{u}_i) = 0 \quad (9)$$

The filtered Momentum equations are,

$$\frac{\partial}{\partial t}P(\rho u_i) + \frac{\partial}{\partial x_j}P(\rho u_i u_j) + \frac{\partial}{\partial x_i}Pp = \frac{\partial}{\partial x_j}P(\sigma_{ij}) \quad (10)$$

substituting for  $\rho u_i$  and  $p$ ,

$$\frac{\partial}{\partial t}(\bar{\rho}\tilde{u}_i) + \frac{\partial}{\partial x_j}P(\rho u_i u_j) + \frac{\partial}{\partial x_i}\bar{p} = \frac{\partial}{\partial x_j}P(\bar{\sigma}_{ij}) \quad (11)$$

where  $\bar{\sigma}_j$  is evaluated from  $\bar{u} = Pu_i$ .

To reduce this to the usual form,

$$\rho u_i u_j = \frac{Q(\bar{\rho}\tilde{u}_i)Q(\bar{\rho}\tilde{u}_j)}{Q\bar{\rho}} \quad (12)$$

so,

$$\frac{\partial}{\partial x_j}P(\rho u_i u_j) = \frac{\partial}{\partial x_j}P(\bar{\rho}\tilde{u}_i\tilde{u}_j) + \frac{\partial}{\partial x_j}\tilde{\tau}_{ij} \quad (13)$$

where,

$$\tilde{\tau}_{ij} = Q^{-1} \left[ \frac{Q(\bar{\rho}\tilde{u}_i)Q(\bar{\rho}\tilde{u}_j)}{Q\bar{\rho}} - Q(\bar{\rho}\tilde{u}_i\tilde{u}_j) \right] \quad (14)$$

Also,

$$\frac{\partial}{\partial x_j}\bar{\sigma}_{ij} = \frac{\partial}{\partial x_j}(\tilde{\sigma}_{ij} + \eta_{ij}) \quad (15)$$

where,

$$\tilde{\sigma}_{ij} = \mu \left( \frac{\partial \tilde{u}_i}{\partial x_j} + \frac{\partial \tilde{u}_j}{\partial x_i} \right) + \lambda \delta_{ij} \frac{\partial \tilde{u}_k}{\partial x_k} \quad (16)$$

and  $\eta_{ij}$  is the same quantity evaluated with,

$$\tilde{u}_i - \bar{u}_i = \frac{\bar{m}_i}{\bar{\rho}} - \bar{u}_i \quad (17)$$

where,

$$m_i = \rho u_i, \quad \bar{m}_i = P(\rho u_i) \quad (18)$$

then,

$$\bar{u}_i = P \frac{m_i}{\rho} = Q^{-1} \left( \frac{Q \bar{m}_i}{Q \bar{\rho}} \right) \quad (19)$$

so,

$$\tilde{u}_i - \bar{u}_i = Q^{-1} \left[ Q \left( \frac{\bar{m}_i}{\bar{\rho}} \right) - \frac{Q \bar{m}_i}{Q \bar{\rho}} \right] \quad (20)$$

Finally, the momentum equation is,

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{u}_i) + \frac{\partial}{\partial x_j} (\bar{\rho} \tilde{u}_i \tilde{u}_j) + \frac{\partial \bar{p}}{\partial x_i} = \frac{\partial}{\partial x_j} \sigma_{ij} - \frac{\partial}{\partial x_j} (\tilde{\tau}_{ij} + \eta_{ij}) \quad (21)$$

Here,

$$\tilde{\tau}_{ij} = Q^{-1} \left[ \frac{Q(\bar{\rho} \tilde{u}_i) Q(\bar{\rho} \tilde{u}_j)}{Q \bar{\rho}} - Q(\bar{\rho} \tilde{u}_i \tilde{u}_j) \right] \quad (22)$$

$$\tilde{\tau}_{ij} = Q^{-1} \left[ \frac{Q \bar{m}_i Q \bar{m}_j}{Q \bar{\rho}} - Q \left( \frac{\bar{m}_i \bar{m}_j}{\bar{\rho}} \right) \right] \quad (23)$$

The filtered energy equations are,

$$\frac{\partial}{\partial t} P(\rho E) + \frac{\partial}{\partial x_j} [P(\rho u_j H)] = \frac{\partial}{\partial x_j} [P(u_i \sigma_{ij} + k \frac{\partial T}{\partial x_j})] \quad (24)$$

Define,

$$\tilde{E} = \frac{P(\rho E)}{P \rho} = \frac{P(\rho E)}{\bar{\rho}} \quad (25)$$

$$\bar{\rho} \tilde{E} = P(\rho E), \quad \rho E = Q(\bar{\rho} \tilde{E}) \quad (26)$$

$$\bar{\rho} \tilde{H} = \bar{\rho} \tilde{E} + \bar{p}, \quad Q(\bar{\rho} \tilde{H}) = \rho E + p \quad (27)$$

Then,

$$\rho u_j H = \frac{\rho u_j \rho H}{\rho} = \frac{Q(\bar{\rho} \tilde{u}_i) Q(\bar{\rho} \tilde{H})}{Q \bar{\rho}} \quad (28)$$

$$u_i \sigma_{ij} = \frac{\rho u_i \sigma_{ij}}{\rho} = \frac{Q(\bar{\rho} \tilde{u}_i) Q \sigma_{ij}}{Q \bar{\rho}} = \frac{Q(\bar{\rho} \tilde{u}_i) Q(\tilde{\tau}_{ij} + \eta_{ij})}{Q \bar{\rho}} \quad (29)$$

$$P(\rho u_j H) = Q^{-1} \left[ \frac{Q(\bar{\rho} \tilde{u}_j) Q(\bar{\rho} \tilde{H})}{Q \bar{\rho}} \right] \quad (30)$$

$$P(u_i \sigma_{ij}) = Q^{-1} \left[ \frac{Q(\bar{\rho} \tilde{u}_i) Q(\tilde{\sigma}_{ij} + \eta_{ij})}{Q \bar{\rho}} \right] \quad (31)$$

Thus, the energy equations can be written as,

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{E}) + \frac{\partial}{\partial x_j} (\bar{\rho} \tilde{u}_j \tilde{H}) = \frac{\partial}{\partial x_j} (\tilde{u}_i \tilde{\sigma}_{ij} + k \frac{\partial \bar{T}}{\partial x_j}) - \frac{\partial}{\partial x_j} (\alpha_{ij} + \beta_{ij}) \quad (32)$$

Where,

$$\alpha_{ij} = Q^{-1} \left[ \frac{Q(\bar{\rho} \tilde{u}_j) Q(\bar{\rho} \tilde{H})}{Q \bar{\rho}} - Q(\bar{\rho} \tilde{u}_j \tilde{H}) \right] \quad (33)$$

and

$$\beta_{ij} = Q^{-1} \left[ \frac{Q(\bar{\rho} \tilde{u}_i) Q(\tilde{\sigma}_{ij} + \eta_{ij})}{Q \bar{\rho}} - Q(\tilde{u}_i \tilde{\sigma}_{ij}) \right] \quad (34)$$