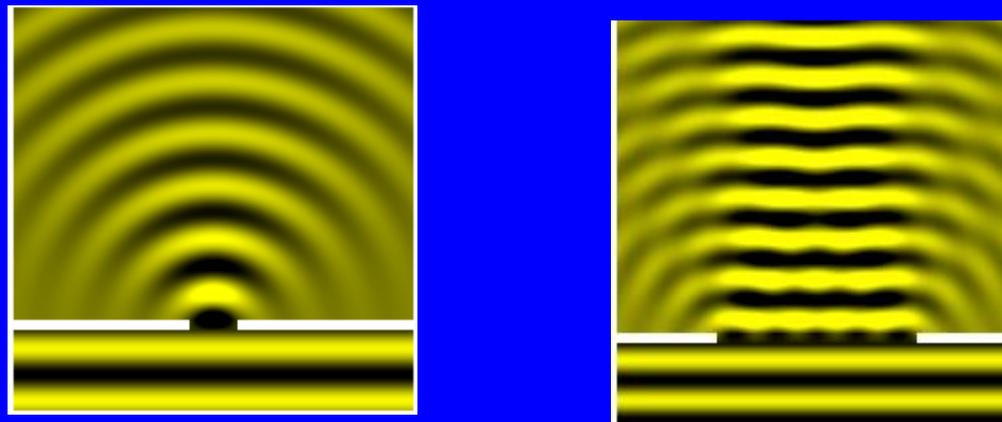


§5 光的衍射现象和惠更斯-菲涅耳原理

5.1 光的衍射现象

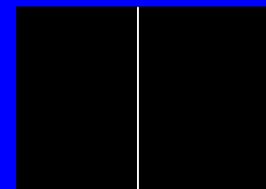
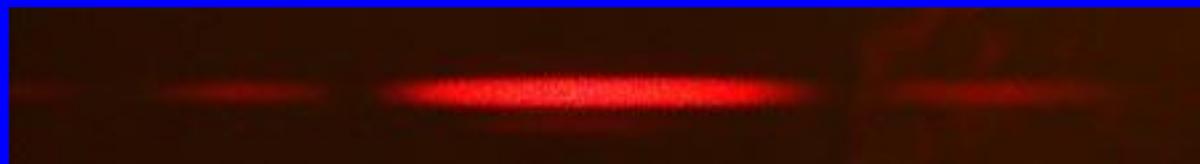
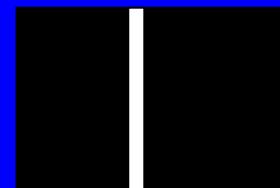
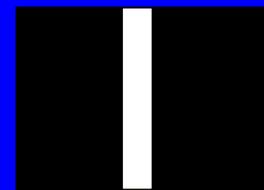
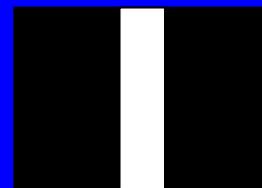
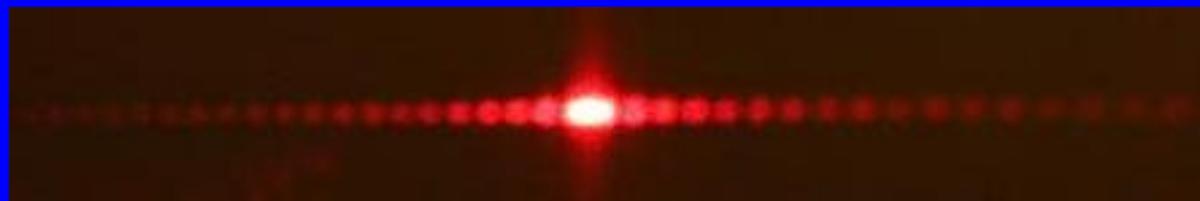
1) 衍射的定义:

波在传播过程遇到障碍物时偏离直线传播的现象



水波的衍射

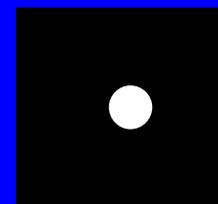
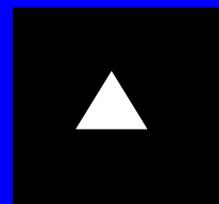
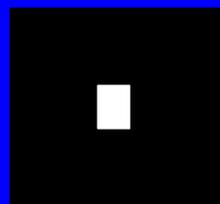
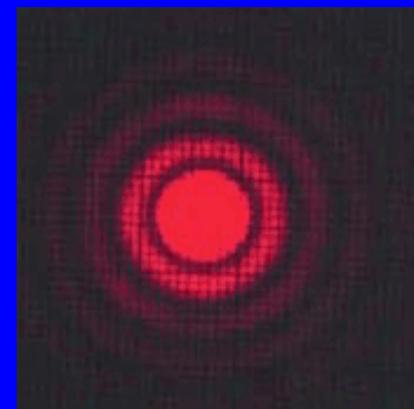
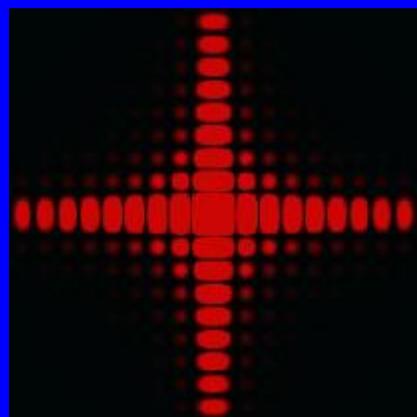
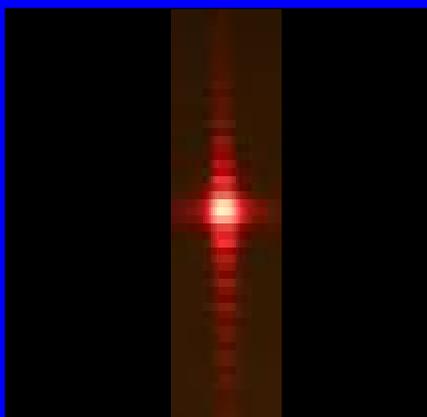
2) 光的衍射现象



单缝衍射图样

衍射屏

从矩孔到圆孔的衍射图样



3) 衍射现象的特点

(1) 在什么方向受限制，衍射图样就沿什么方向扩展

(2) 限制越厉害，衍射越强烈

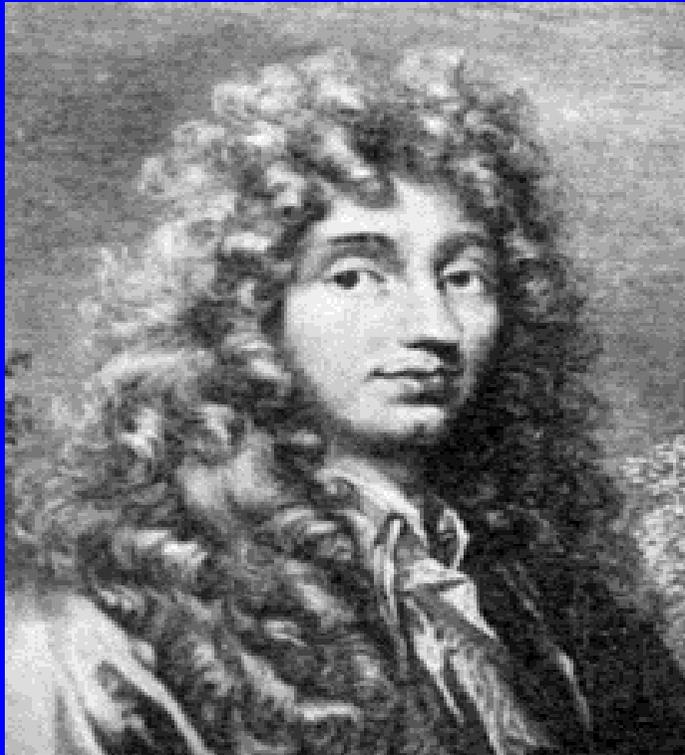
(3) 衍射强弱与障碍物尺寸的关系：

$r \sim 1000l$ 以上：衍射效应不明显

$r \sim 1000l - 10l$ ：衍射效应明显

$r \sim l$ ：向散射过渡

5.2 惠更斯-菲涅耳原理



惠更斯, C.

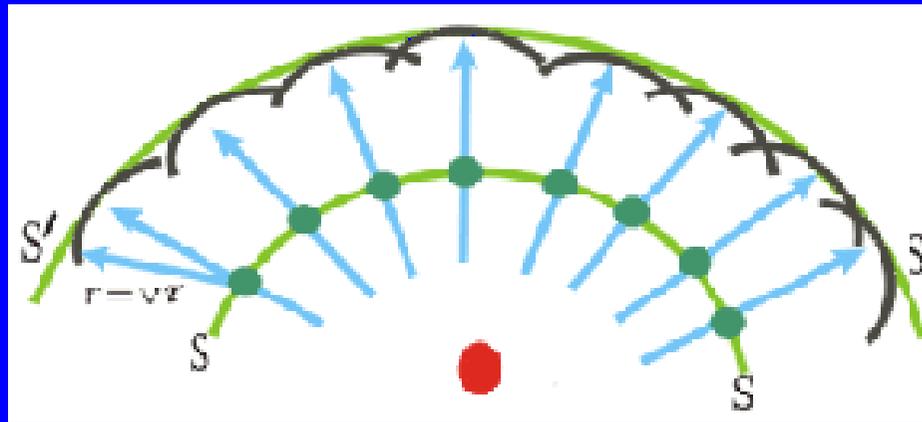


菲涅耳, A.-J.

1) 惠更斯原理：波面+次波

波阵面上各点都看成是子波波源

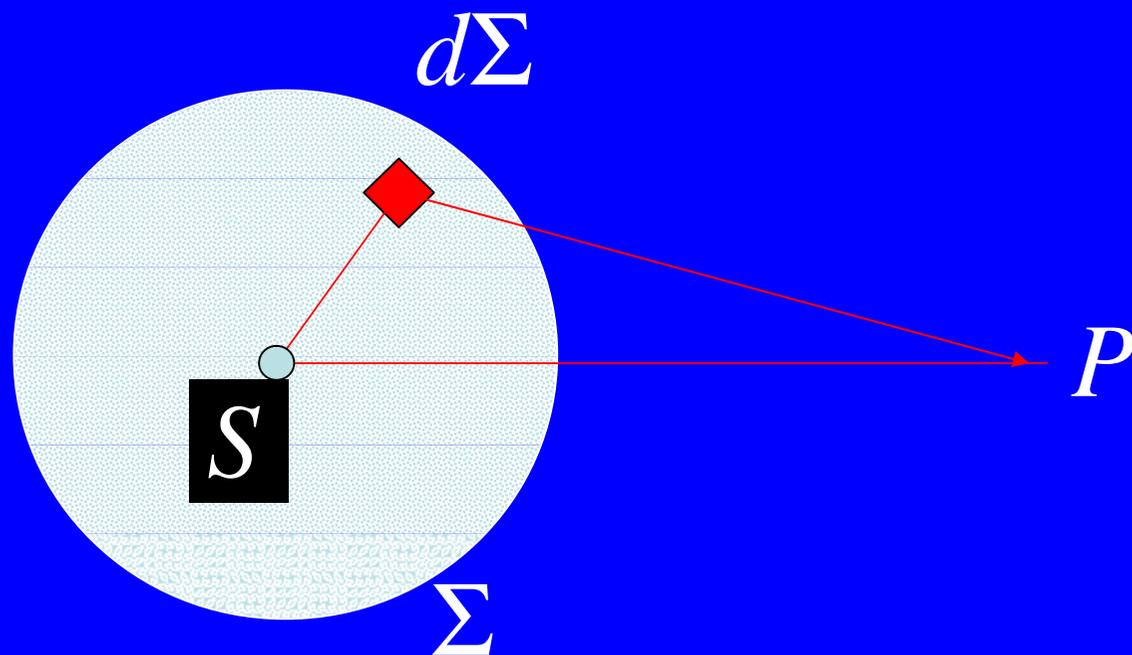
能定性解释光的传播方向问题



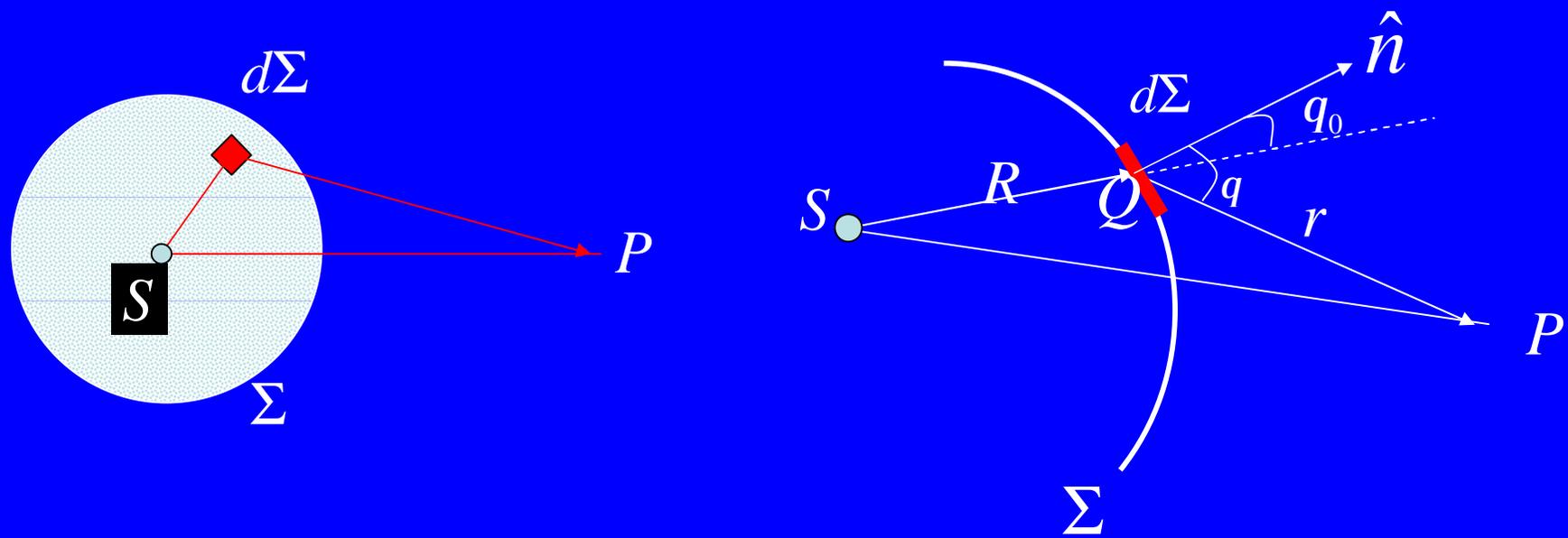
2) 菲涅耳原理：次波+相干叠加

波场中各点的强度由各子波在该点的相干叠加决定

能定量解释衍射图样中的强度分布



波前 S 上每个面元 dS 都可以看成是新的振动中心，它们发出次波。在空间某一点 P 的振动是所有这些次波在该点的相干迭加。

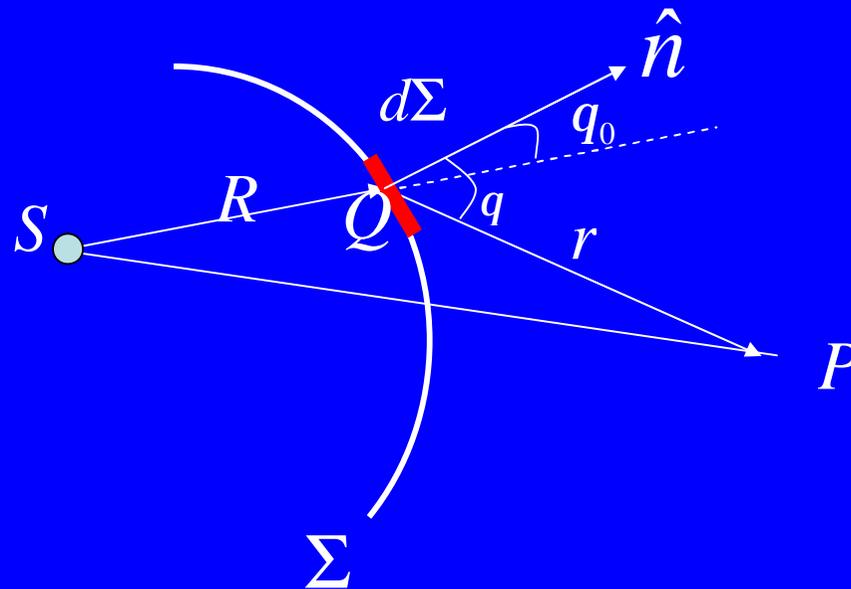


原理的数学表达式:

相干叠加----复振幅线性叠加

$$\tilde{U}(P) = \oiint_{(\Sigma)} d\tilde{U}$$

$$\begin{aligned}
 d\tilde{U}(P) &\propto d\Sigma \\
 &\propto \tilde{U}_0(Q) \\
 &\propto \frac{e^{ikr}}{r} \\
 &\propto F(q, q_0)
 \end{aligned}$$



$$\tilde{U}(P) = K \iint_{(\Sigma)} \tilde{U}_0(Q) F(q, q_0) \frac{e^{ikr}}{r} d\Sigma$$

-----Fresnel衍射积分公式

3) 菲涅耳-基尔霍夫衍射公式

$$\tilde{U}(P) = K \oiint_{(\Sigma)} \tilde{U}_0(Q) F(q, q_0) \frac{e^{ikr}}{r} d\Sigma$$

基尔霍夫推导出:

$$F(q, q_0) = \frac{1}{2} (\cos q + \cos q_0)$$

$$K = \frac{-i}{1} = \frac{e^{-ip/2}}{1}$$

$$\tilde{U}(P) = \frac{-i}{2} \oiint_{(\Sigma)} (\cos q + \cos q_0) \tilde{U}_0(Q) \frac{e^{ikr}}{r} d\Sigma$$

----菲涅耳-基尔霍夫衍射公式

衍射屏情形：

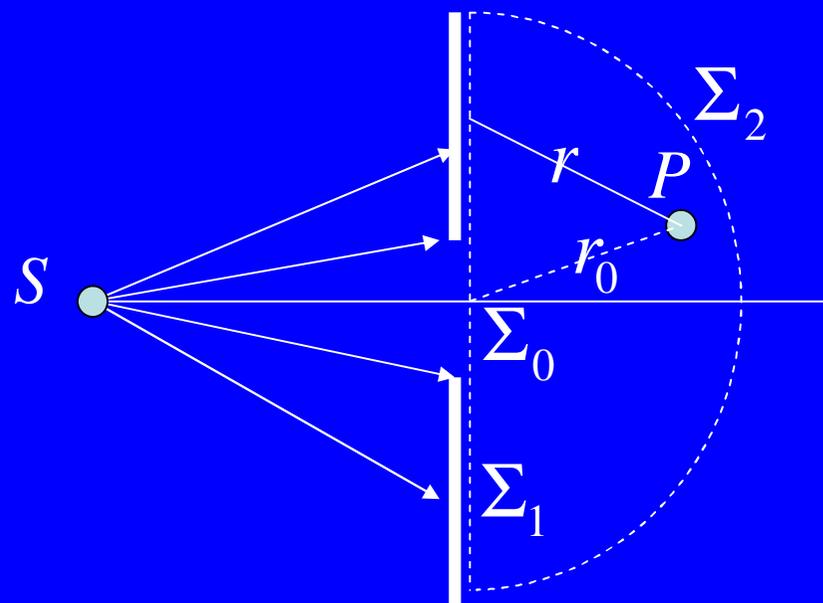
取波前为如下闭合曲面

$$\Sigma = \Sigma_0 + \Sigma_1 + \Sigma_2$$

则

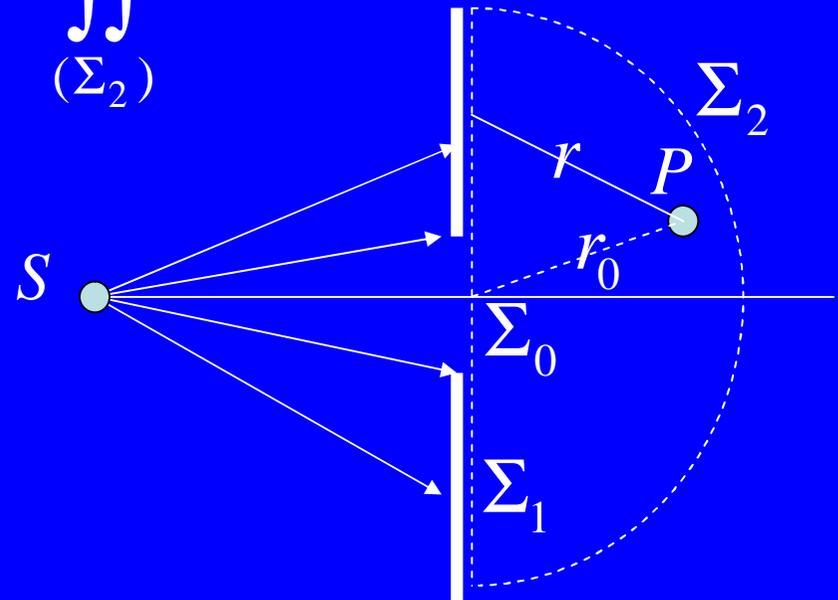
$$\tilde{U}(P) = \oiint_{(\Sigma)} = \iint_{(\Sigma_0)} + \iint_{(\Sigma_1)} + \iint_{(\Sigma_2)}$$

$$\tilde{U}_0(Q) = \begin{cases} \tilde{U}_0(Q) & \Sigma_0 \\ 0 & \Sigma_1 \end{cases} \quad \text{--- 基尔霍夫边界条件}$$



$$\tilde{U}(P) = \oiint_{(\Sigma)} = \iint_{(\Sigma_0)} + \iint_{(\Sigma_1)} + \iint_{(\Sigma_2)}$$

$$\tilde{U}_0(Q) = \begin{cases} \tilde{U}_0(Q) & \Sigma_0 \\ 0 & \Sigma_1 \end{cases}$$



可证明： $\iint_{(\Sigma_2)} = 0$

有：

$$\tilde{U}(P) = \frac{-i}{2I} \iint_{(\Sigma_0)} (\cos q + \cos q_0) \tilde{U}_0(Q) \frac{e^{ikr}}{r} d\Sigma$$

$$\tilde{U}(P) = \frac{-i}{2I} \iint_{(\Sigma_0)} (\cos q + \cos q_0) \tilde{U}_0(Q) \frac{e^{ikr}}{r} d\Sigma$$

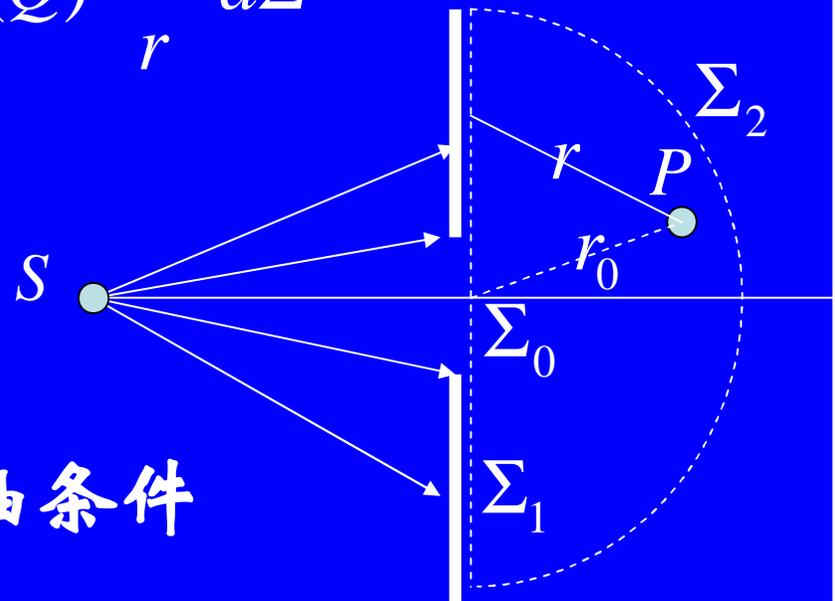
公式简化:

光孔和接收范围满足傍轴条件

$$q \approx q_0 \approx 0 \quad r \approx r_0$$

有:

$$\tilde{U}(P) = \frac{-i}{I r_0} \iint_{(\Sigma_0)} \tilde{U}_0(Q) e^{ikr} d\Sigma$$

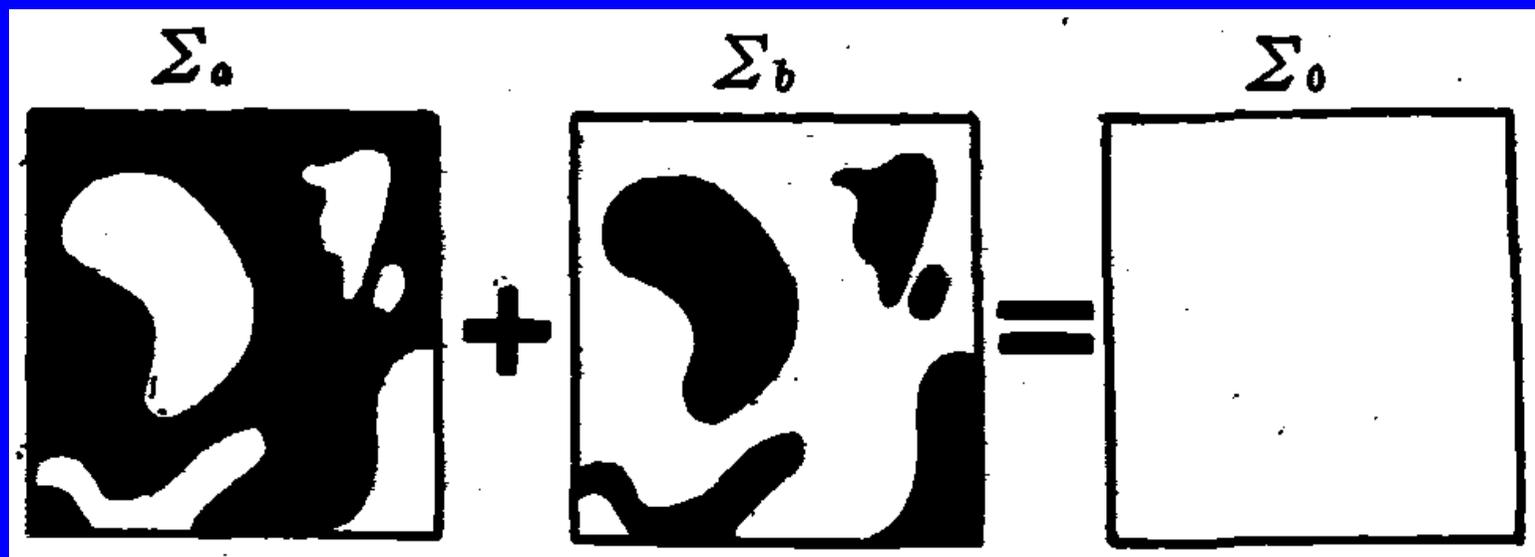


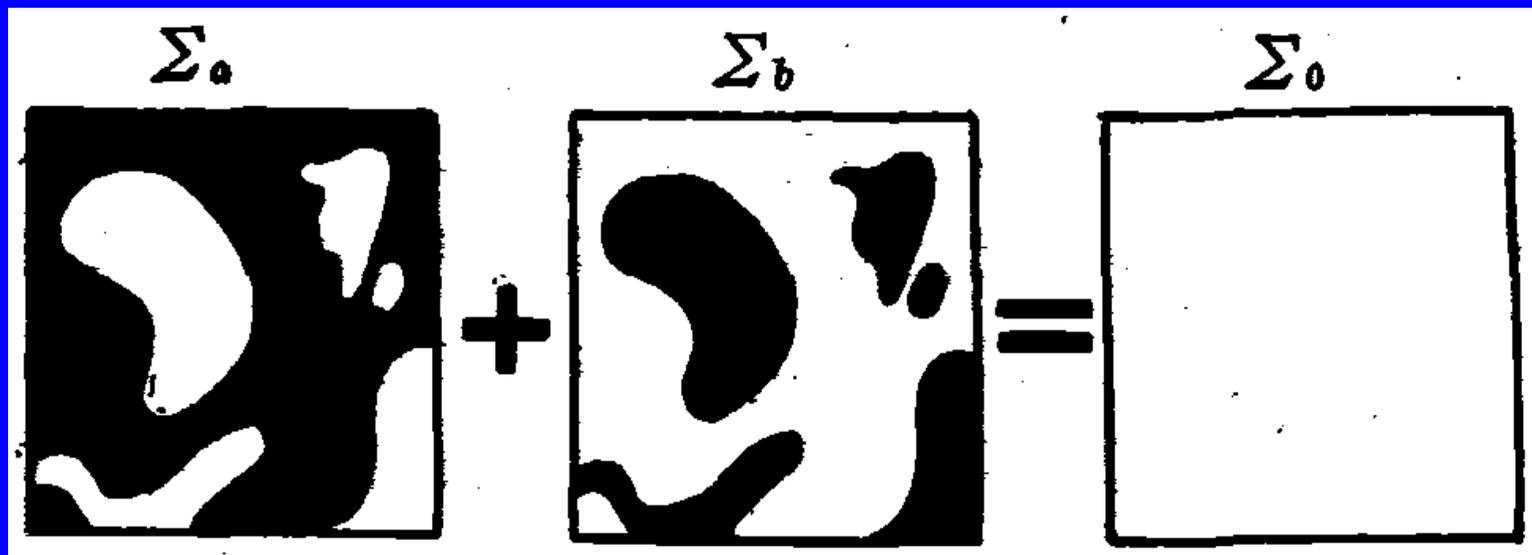
5.3 巴俾涅原理

菲涅耳-基尔霍夫衍射公式

$$\tilde{U}(P) = \frac{-i}{2l} \iint_{(\Sigma_0)} (\cos q + \cos q_0) \tilde{U}_0(Q) \frac{e^{ikr}}{r} d\Sigma$$

设有衍射屏： $\Sigma_0 = \Sigma_a + \Sigma_b$ 两屏互补





$$\tilde{U}(P) = \iint_{(\Sigma_0)} = \iint_{(\Sigma_a)} + \iint_{(\Sigma_b)} = \tilde{U}_a(P) + \tilde{U}_b(P)$$

当 $\Sigma_0 \rightarrow \infty$

则 $\tilde{U}(P) = \tilde{U}_0(P)$ --- 自由波场在P点的复振幅

$$\tilde{U}_0(P) = \tilde{U}_a(P) + \tilde{U}_b(P)$$

利用巴俾涅原理，已知一种衍射屏形成的波场的复振幅或光强，往往可以计算出其互补屏的复振幅和光强。

5.4 衍射的分类

1) 衍射系统 光源，衍射屏，接收屏

2) 衍射分类

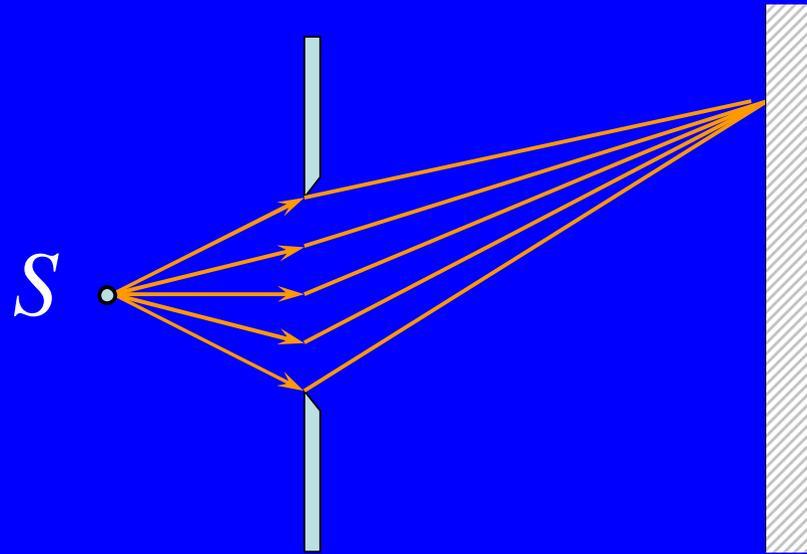
按光源和考察点（光屏）到障碍物距离的不同进行分类

(1) 菲涅耳衍射

(2) 夫琅和费衍射

(1) 菲涅耳衍射 (近场衍射)

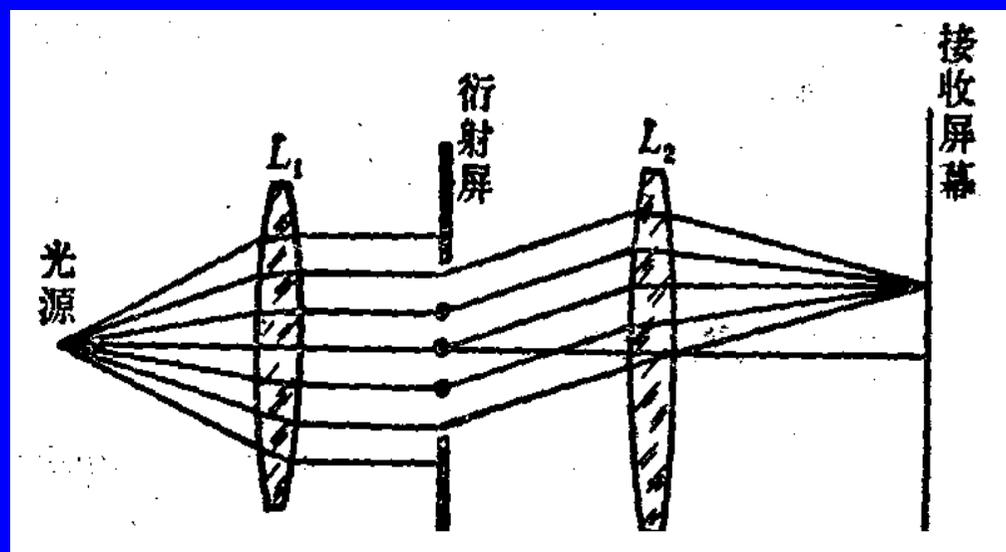
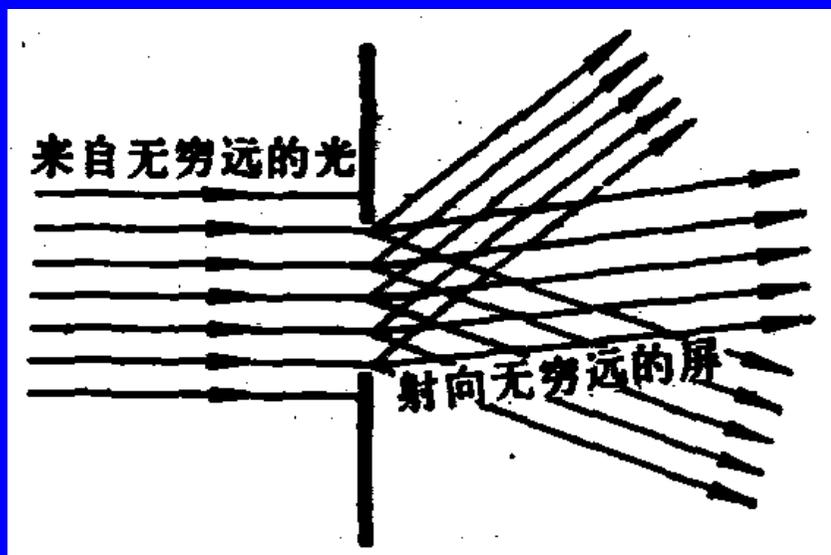
障碍物 (孔隙) 距光源和光屏的距离都是有限的, 或其中之一是有限的。



观察比较方便, 但定量计算却很复杂 (需完成复杂的Fresnel积分)。

(2) 夫琅和费衍射 (远场衍射)

光源和接收屏幕距离衍射屏幕无限远



等效形式

作业:

习题: